

LIA TYL ILC-top

(one aspect of it)

$$e^+ e^- \rightarrow t \bar{t} \rightarrow \mu^+ \mu^- b \bar{b} \nu_\mu \bar{\nu}_\mu$$

The Matrix Element Method
Conjugate (optimal) variables
Results
Kinematics
Conclusion

The Matrix Element Method

$$| \mathcal{M} |^2$$

All information per event is used
(Petra->LEP->Tevatron->LHC)

H.J. Behrends et al., CELLO Collab. Z. Phys. C43 (1989)

Here, we are dealing with a rich 9-dimension Phase Space:

$$dLips = d \cos \theta_t d \cos \theta_b d\phi_b d \cos \theta_{\bar{b}} d\phi_{\bar{b}} d \cos \theta_{\mu^+} d\phi_{\mu^+} d \cos \theta_{\mu^-} d\phi_{\mu^-}$$

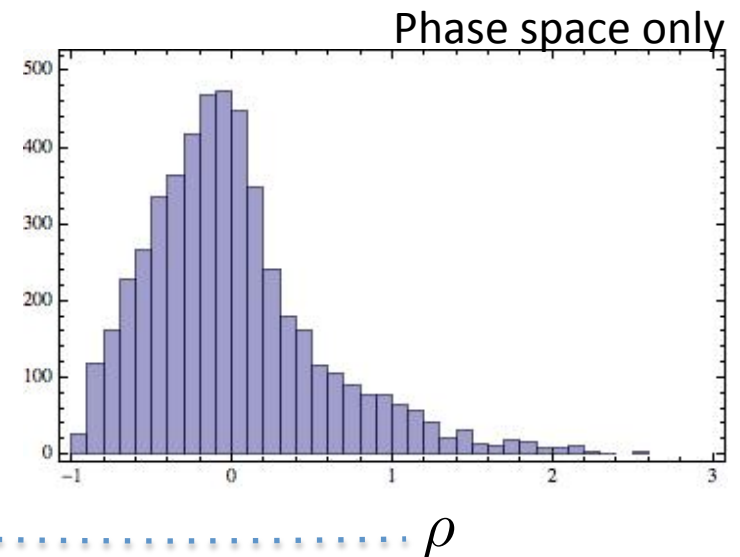
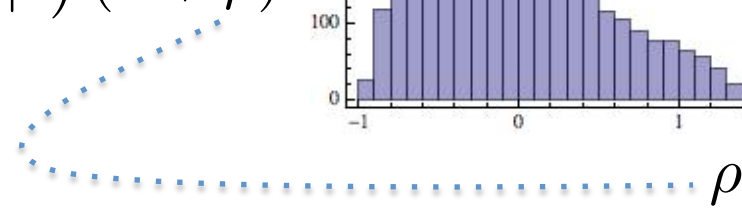
and 16^2 amplitudes²

Example :

Correction term due to
interferences between
the helicity amplitudes

$$| t^{\uparrow} \bar{t}^{\uparrow} + t^{\uparrow} \bar{t}^{\downarrow} + t^{\downarrow} \bar{t}^{\uparrow} + t^{\downarrow} \bar{t}^{\downarrow} |^2 =$$

$$\left(| t^{\uparrow} \bar{t}^{\uparrow} | + | t^{\uparrow} \bar{t}^{\downarrow} | + | t^{\downarrow} \bar{t}^{\uparrow} | + | t^{\downarrow} \bar{t}^{\downarrow} | \right)^2 (1 + \rho)$$



$\alpha \leftarrow$ theoretical parameters

$\mathcal{L}(\alpha)$

A likelihood analysis of events easily handle the 9-dimension analysis

It can be shown that, in effect, the Likelihood analysis implicitly makes use of “conjugate” kinematical variables defined by:

$$\omega_i = \frac{\partial |\mathcal{M}|^2(\alpha)}{\partial \alpha_i} \bigg|_{\alpha^0} \frac{1}{|\mathcal{M}|^2(\alpha^0)}$$
$$\Omega_i = \frac{\partial N(\alpha)}{\partial \alpha_i} \bigg|_{\alpha^0} \frac{1}{N(\alpha^0)}$$

Using only the distribution of events in Phase Space:

$$V_{ij}^{(-1)} \equiv \Lambda_{ij} = N \langle (\omega_i - \Omega_i)(\omega_j - \Omega_j) \rangle_0$$

Using also the yields:

$$V_{ij}^{(-1)} \equiv \Lambda_{ij} = N \langle \omega_i \omega_j \rangle_0$$

Simultaneous 10-parameter fit

$\mathcal{R}e \delta\tilde{F}_{1V}^\gamma$	$\mathcal{R}e \delta\tilde{F}_{1V}^Z$	$\mathcal{R}e \delta\tilde{F}_{1A}^\gamma$	$\mathcal{R}e \delta\tilde{F}_{1A}^Z$	$\mathcal{R}e \delta\tilde{F}_{2V}^\gamma$	$\mathcal{R}e \delta\tilde{F}_{2V}^Z$	$\mathcal{R}e \delta\tilde{F}_{2A}^\gamma$	$\mathcal{R}e \delta\tilde{F}_{2A}^Z$	$\mathcal{I}m \delta\tilde{F}_{2A}^\gamma$	$\mathcal{I}m \delta\tilde{F}_{2A}^Z$
0.0037	-0.18	-0.09	+0.14	+0.62	-0.15	0	0	0	0
	0.0063	+0.14	-0.06	-0.13	+0.61	0	0	0	0
		0.0053	-0.15	-0.05	+0.09	0	0	0	0
			0.0083	+0.06	-0.04	0	0	0	0
				0.0105	-0.19	0	0	0	0
					0.0169	0	0	0	0
						0.0068	-0.15	0	0
							0.0118	0	0
								0.0069	-0.17
									0.0100

0.4 \rightarrow 1.7%

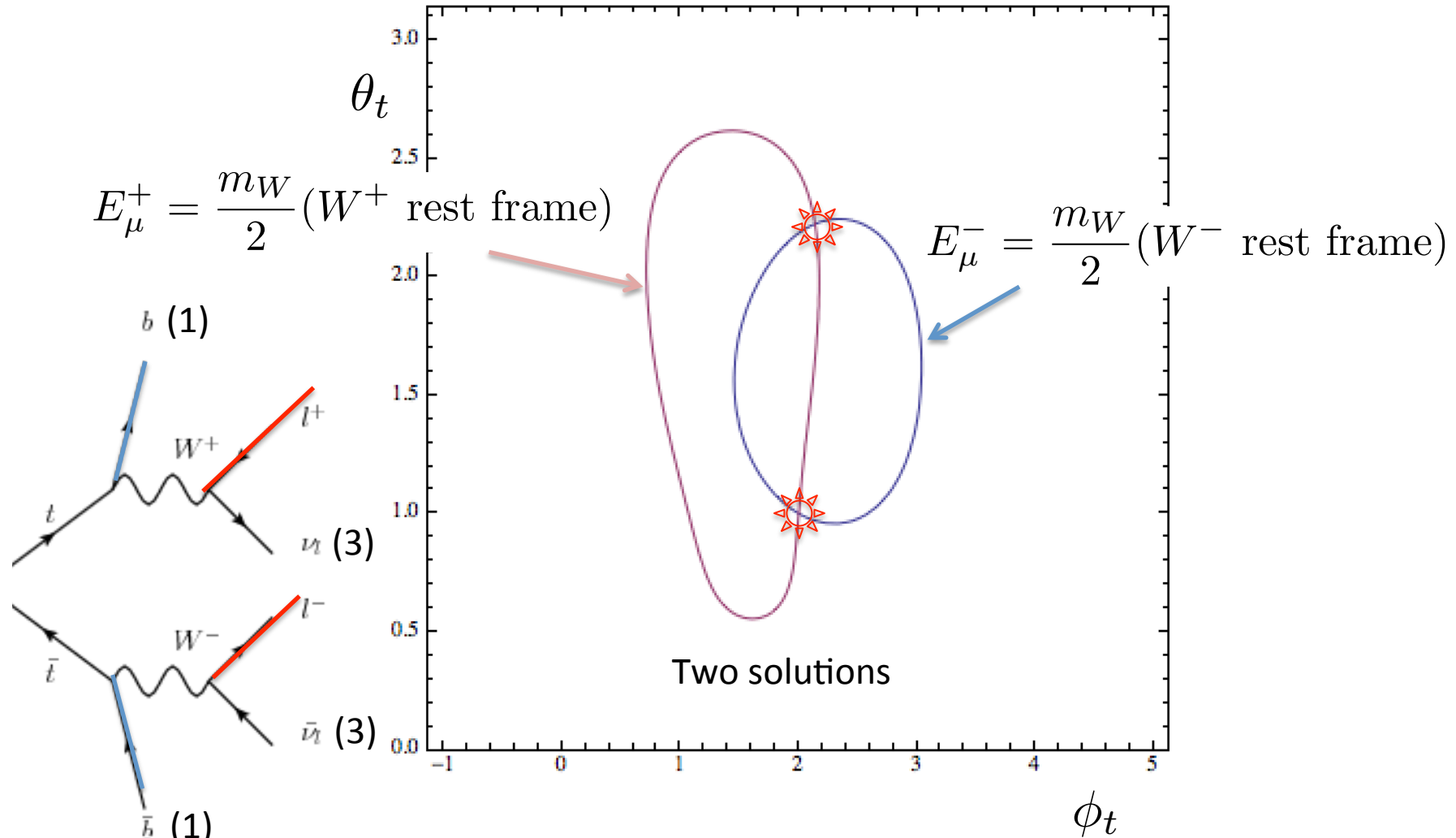
500 GeV & 500 fb⁻¹

Luminosity split 50/50 between $\pm 80\% \mp 30\%$ beam polarizations

17 500 events produced

(no reconstruction & detector & background & Physics effects taken care of)

Kinematical reconstruction

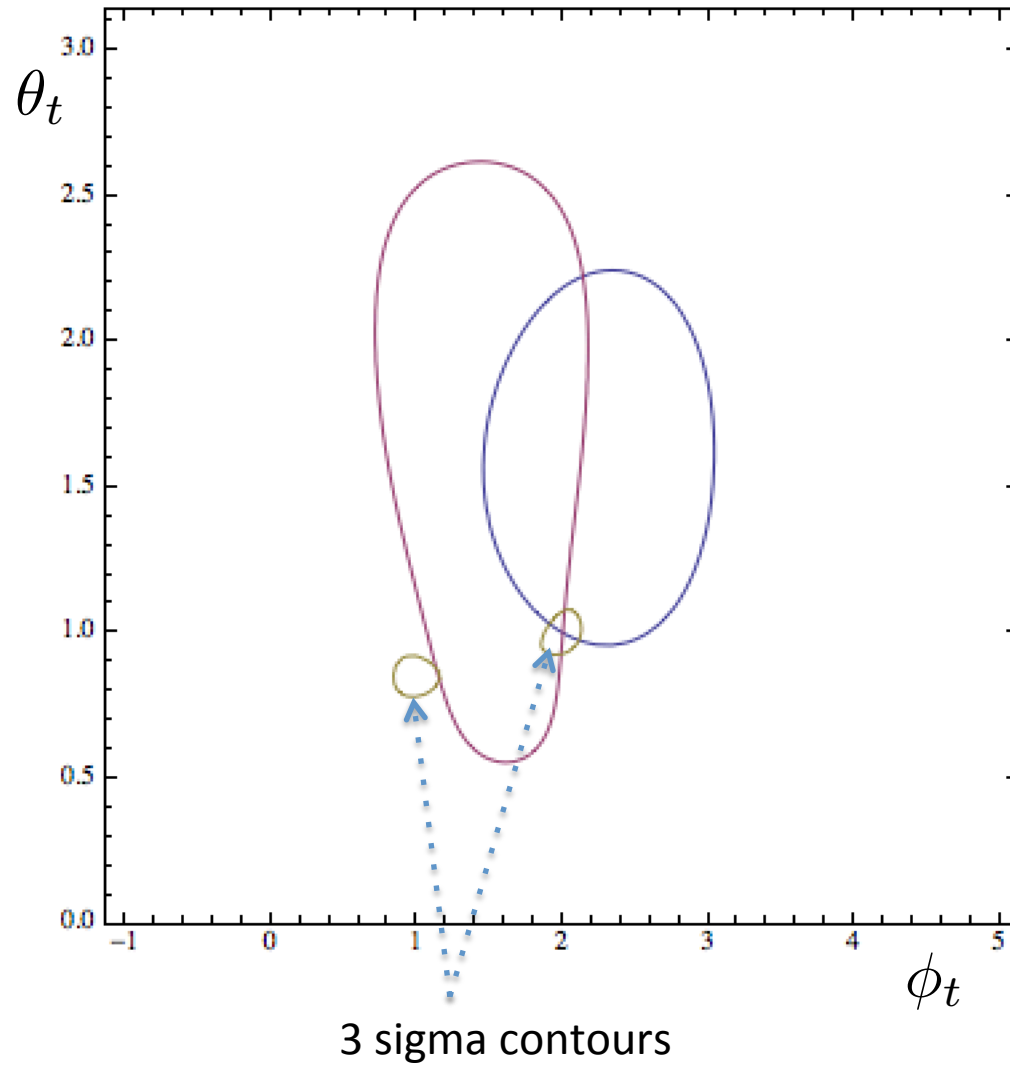


$$1+3+3+1=4+2+2$$

$$m_t \quad (2)(2) \quad m_W$$

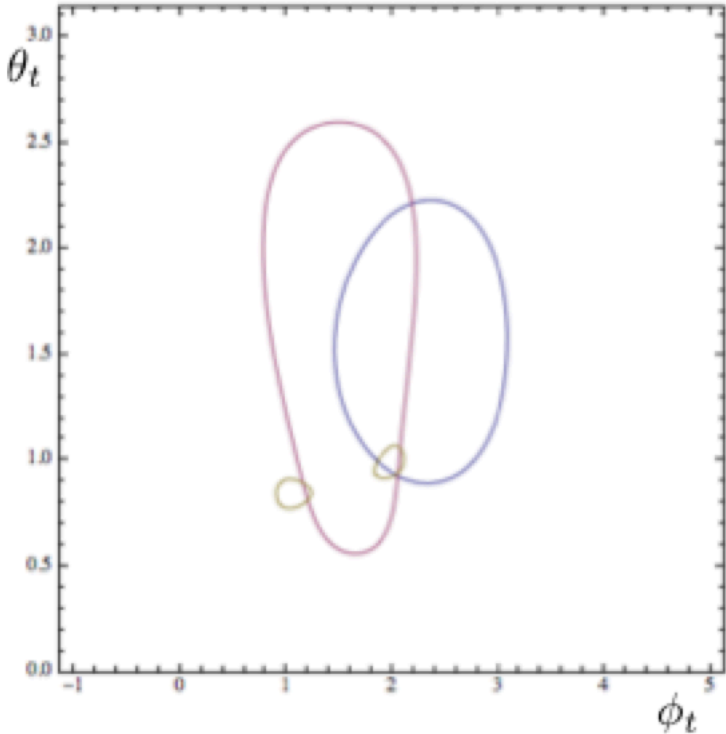
- directions of the b jets
- direction and energies of the e/mu

The energies of the b jets are used to select the right solution



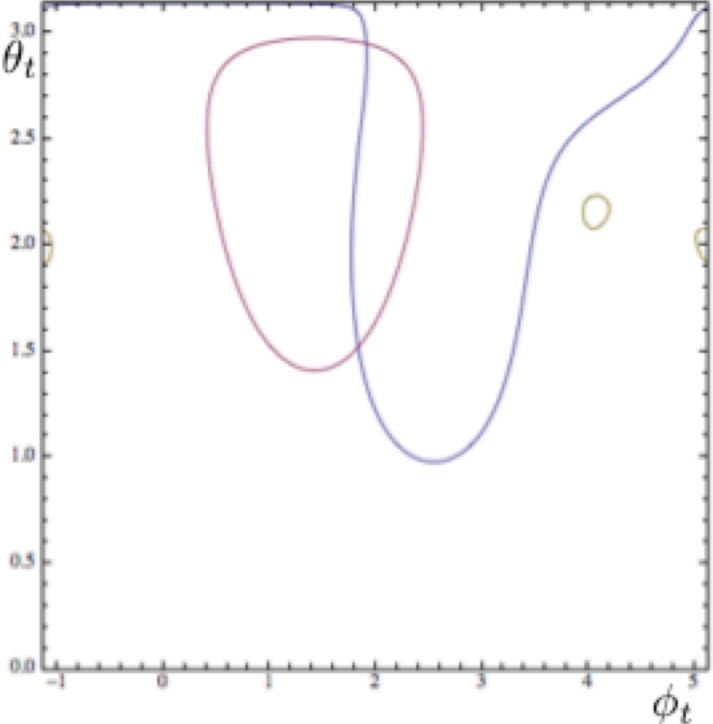
... and some complications.

The W and top widths do not help : remain manageable

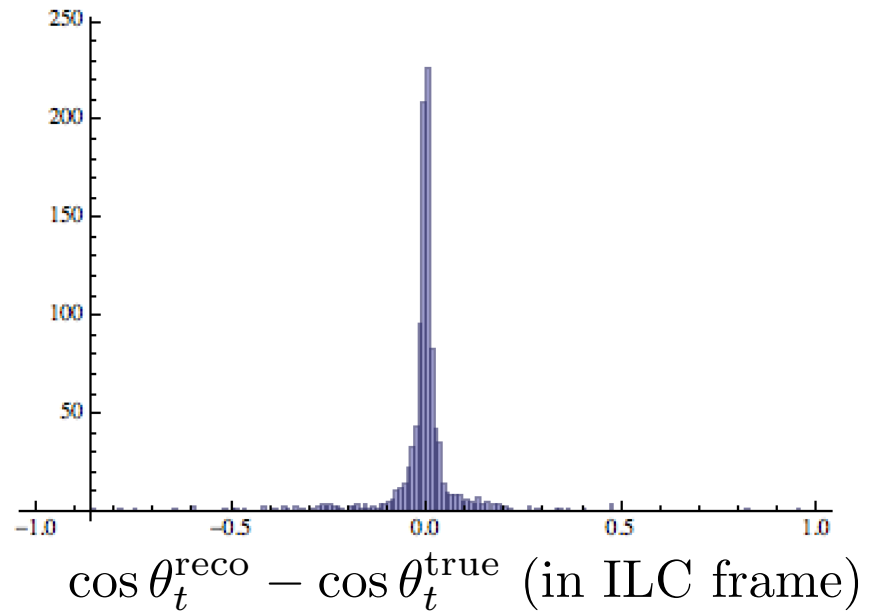


Picking up the best solution (unperfect) works

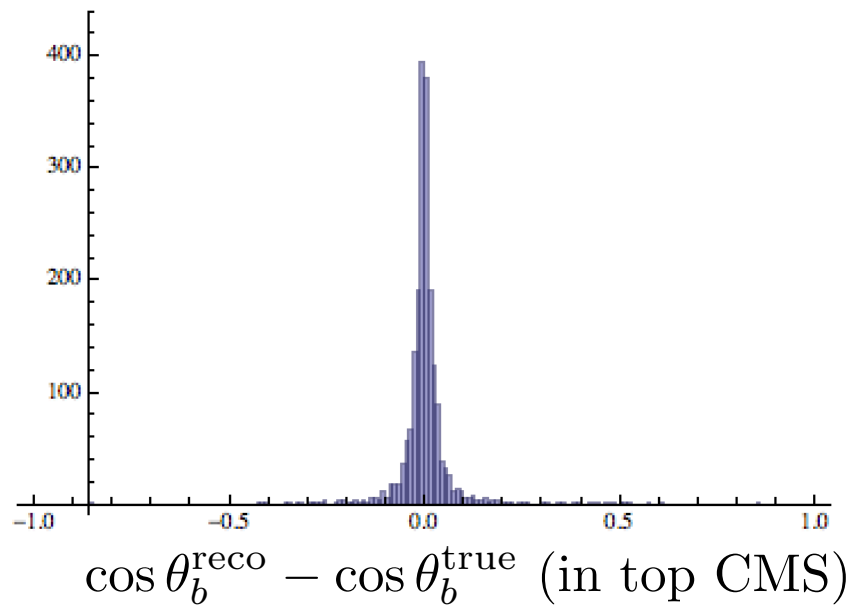
If one select the wrong b assignment : no solution



... in fact, 5% about manage to pass through



Using GRACE events, and letting the masses free to vary



Conclusion :

The final state

$$e^+e^- \rightarrow t\bar{t} \rightarrow \mu^+\mu^-b\bar{b}\nu_\mu\bar{\nu}_\mu$$

Appears promising for top studies :
The sensitivity to top couplings are
similar to the semi-leptonic mode,
but with different limitations.

A lot of work ahead