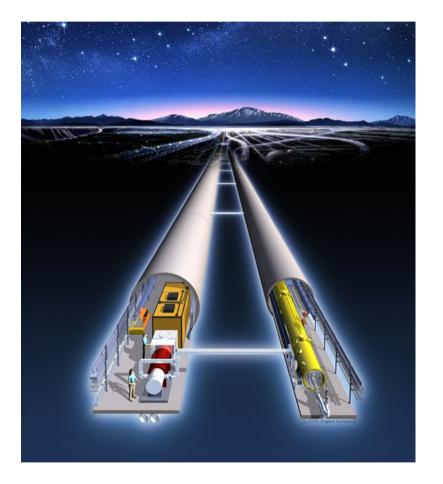


Top quark coupling at LC - Experimental issues -





Roman Pöschl



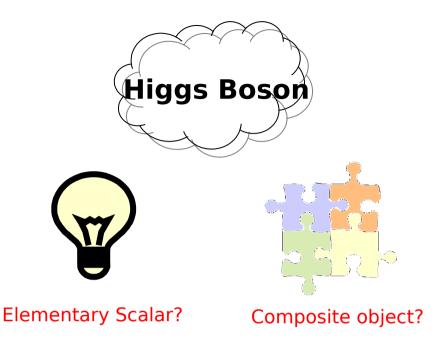


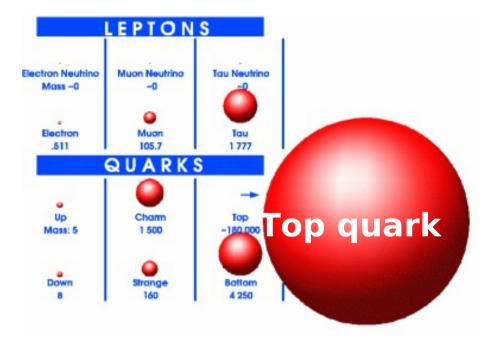
Comprendre le monde, construire l'avenir®

Based on collaboration With IFIC Valencia

3rd Journées Collisionneur Linéaire – LPSC/Grenoble Dec. 2014

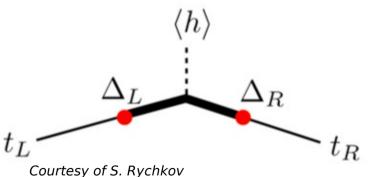
An enigmatic couple ...





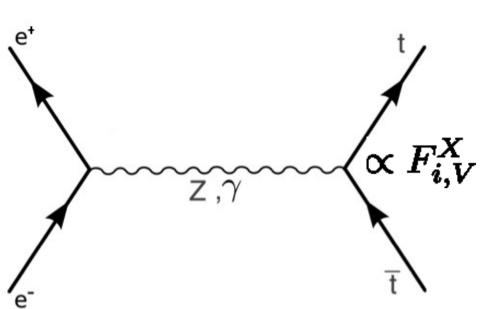
- Higgs and top quark are intimately coupled!
 Top Yukawa coupling O(1) !
 Top mass important SM Parameter
- New physics by compositeness? Higgs <u>and</u> top composite objects?

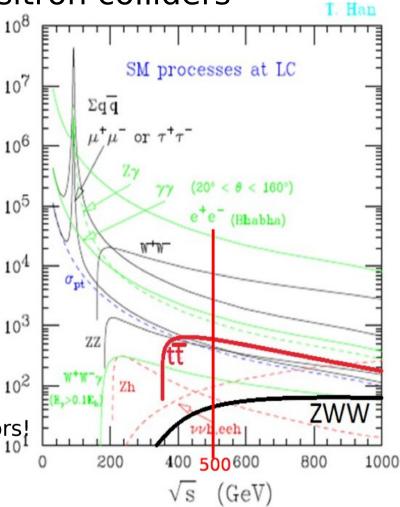
- LC perfectly suited to decipher both particles



Top quark physics at electron-positron colliders

σ





Top quark production through electroweak processes,

no competing QCD production => Small theoretical errors!

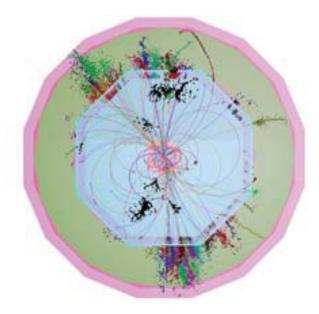
- High precision measurements

Top quark mass at ~ 350 GeV through threshold scan Polarised beams allow testing chiral structure at ttX vertex => Precision on form factors F

- Studies presented here deal with no or only mildly boosted tops, beta~0.7
 - A major difference between LC and LHC is that an LC will run triggerless
 - -> Unbiased event samples, all event selection happens off-line! JCL 2014 - LPSC Grenoble

Detector parameters

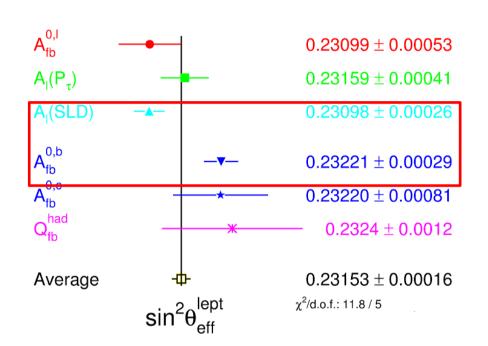
Track momentum: $\sigma_{1/p} < 5 \times 10^{-5}$ /GeV (1/10 × LEP) (e.g. Measurement of Z boson mass in Higgs Recoil) Impact parameter: $\sigma_{d0} < [5 \oplus 10/(p[GeV]sin^{3/2}\theta)] \mu m(1/3 \times SLD)$ (Quark tagging c/b) Jet energy resolution : dE/E = 0.3/(E(GeV))^{1/2} (1/2 × LEP) (W/Z masses with jets) Hermeticity : $\theta_{min} = 5 mrad$ (for events with missing energy e.g. SUSY)



Final state will comprise events with a large number of charged tracks and jets(6+).

- High granularity
- Excellent momentum measurement
- High separation power for particles

2> Two different approaches Detector concepts SiD et ILD The top quark and flavor hierarchy



Flavor hierarchy? Role of 3rd generation?

- A_{FB} anomaly at LEP for b quark Tensions at Tevatron?
- Heavy fermion effect

Strong motivation to study chiral structure of top vertex in high energy e+e- collisions

Why is it sooo heavy?

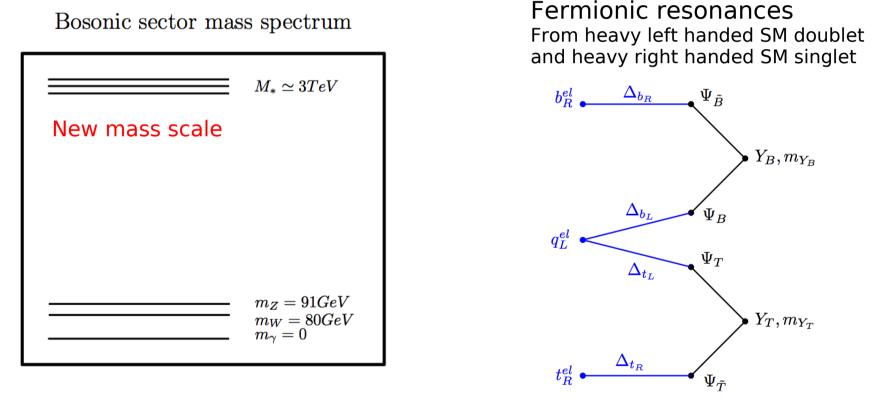
	10 ⁴	-			
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mass,	10 ⁰	-		c s µ	τ
8	10 ⁻¹ 10 ⁻² 10 ⁻³			μ	
	10 ⁻²	-	d u e		
	10 ⁻³	-	e		
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		100			

Towards New Physics

à la G.M. Pruna, LC 13, Trento

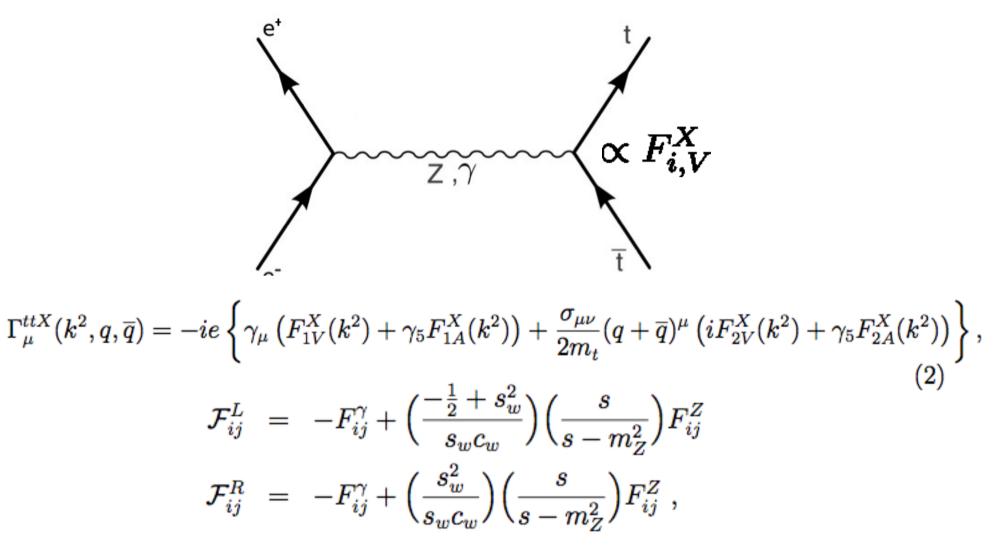
Compositeness:

- ... provides elegant solution for naturalness
- ... few tensions with SM predictions
- ... composite Higgs hypothesis has only been marginally studied in comparison with other "fundamental" scenarios
- ... all scalar objects observed in nature turned out to be bound states of fermions



Physics modify Yukawa couplings and Ztt, Zbb Heavy fermion effect!

Testing the chiral structure of the Standard Model



Pure γ or pure $Z^0: \sigma \sim (F_i)^2 \Rightarrow$ No sensitivity to sign of Form Factors Z^0/γ interference $: \sigma \sim (F_i) \Rightarrow$ Sensitivity to sign of Form Factors

Disentangling

At ILC **no** separate access to ttZ or ttγ vertex, but ...

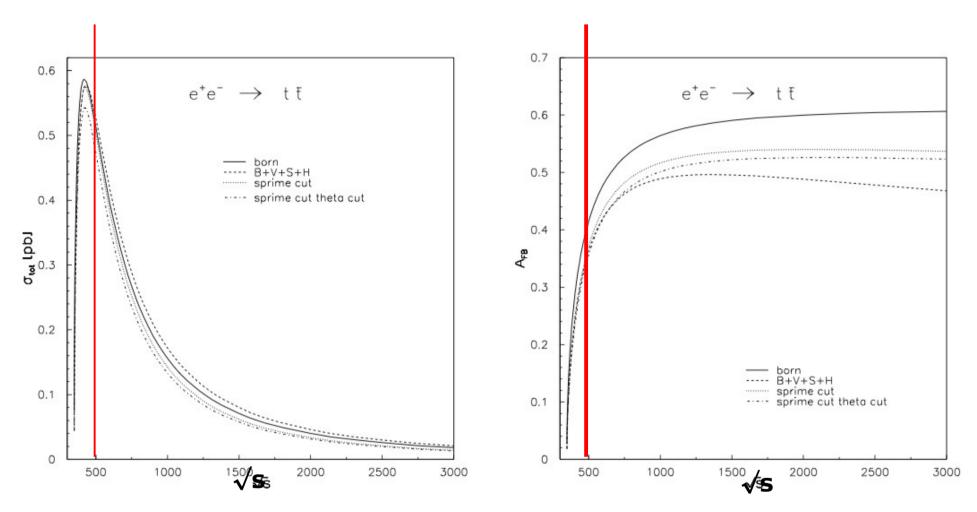
ILC 'provides' two beam polarisations

$$P(e^{-}) = \pm 80\%$$
 $P(e^{+}) = \mp 30\%$

There exist a number of observables sensitive to chiral structure, e.g.

$$\sigma_{I} \qquad A_{FB,I}^{t} = \frac{N(\cos\theta > 0) - N(\cos\theta < 0)}{N(\cos\theta > 0) + N(\cos\theta < 0)} \qquad (F_{R})_{I} = \frac{(\sigma_{t_{R}})_{I}}{\sigma_{I}}$$
x-section
Forward backward asymmetry
Fraction of right handed top quarks
Extraction of up to six (five) unknowns
 $F_{1V}^{\gamma}, F_{1V}^{Z}, F_{1A}^{\gamma} = 0, F_{1A}^{Z}$
 $F_{2V}^{\gamma}, F_{2V}^{Z}$

Measuring at 500 GeV

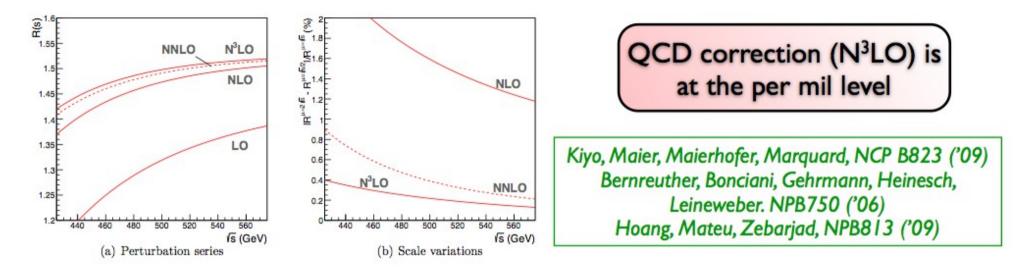


- Cross section close to maximum, $A_{_{FR}}$ well developed

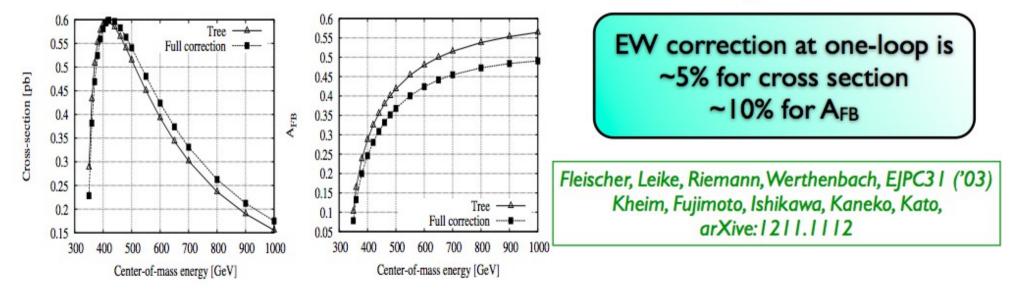
 Other remarks: Need some velocity to get sensitive to chiral obervables (see backup slides)

Theoretical uncertainties

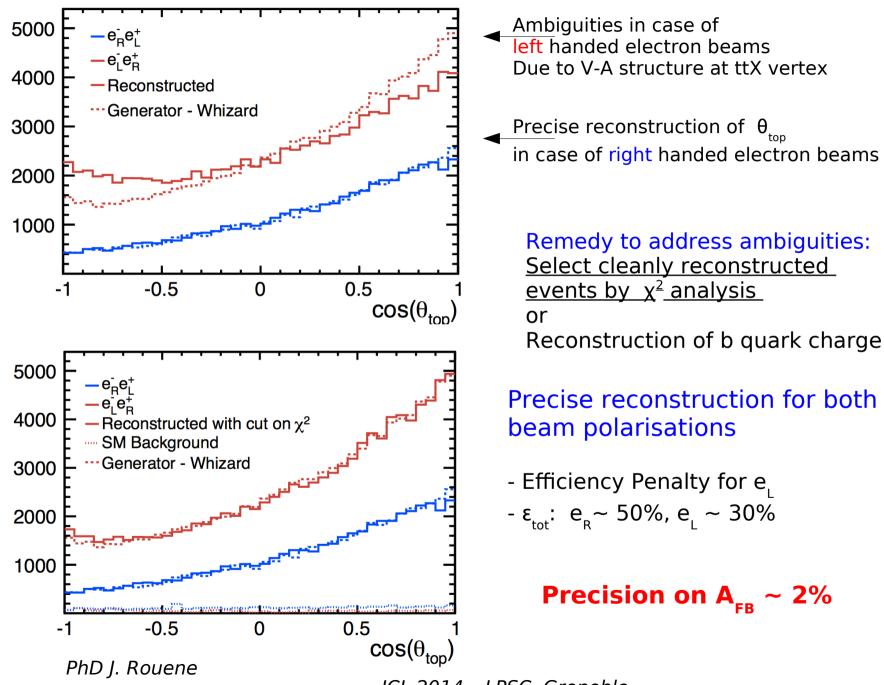
*QCD corrections are known up to N³LO



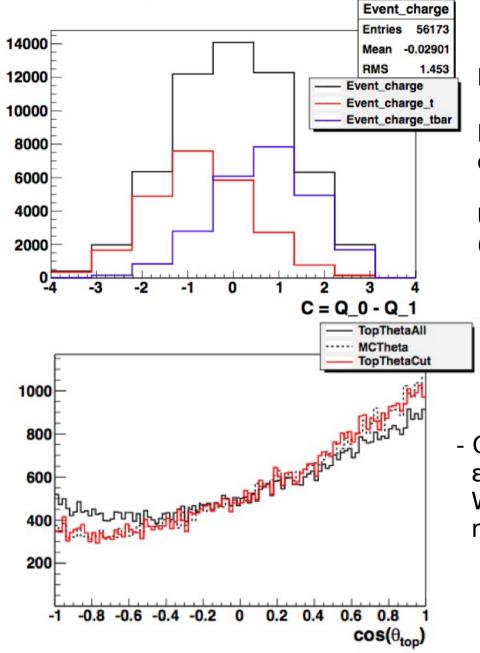
*Electroweak corrections are known at one-loop level



Semi Leptonic Analysis - Reconstruction of top quark production angle



Top polar angle using b charge (SL Analysis)



Event charge C = b1 - b2

In SL can compare charge C with lepton charge to select clean sample

Use only events with correct C or C=0 (plus another cut on the Lorentz Factor)

- Clean reconstruction of top quark direction $\epsilon \sim 30\%$ Will improve with improving charge reconstruction

JCL 2014 – LPSC Grenoble

B charge measurement - Potential

- b quark hadronises to about
- ${\sim}40\%$ to charged B mesons
- ${\sim}50\%$ to neutral B mesons
- ${\sim}10\%$ to Baryons

= 64% cases where there is at least one charged b = Should be recognisable

- neutral B mesons decay to about
 - \sim 50% into charged D Mesons => measurable
 - ~ 50% into neutral D mesons

 \sim 64% of these D neutral undergo prong decays => charged particles => measurable

=> Out of 36% cases remaining above ~75% can (in principle) be retrieved

=> 91% of the charges from top quark decays lead to signatures that are in principle measurable

Two tasks:

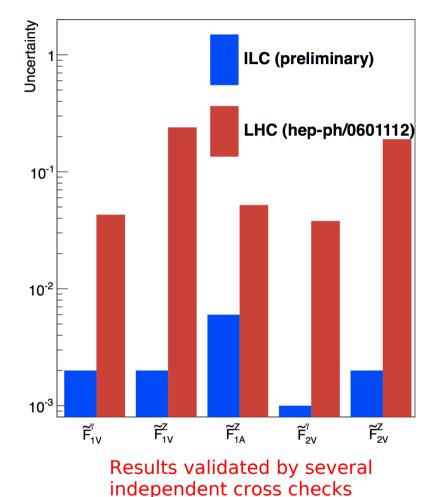
1) Understand why final state with charged B Meson are wrongly reconstructed Exact fraction depends on final state, looks as if SL is somewhat easier than fully hadronic

2) Tertiary vertices for neutral B Mesons

Results of full simulation study for DBD at $\sqrt{s} = 500$ GeV

ArXiv: 1307.8102

Precision: cross section ~ 0.5%, Precision A_{FB} ~ 2%, Precision λ_t ~ 3-4% Accuracy on CP conserving couplings



- ILC might be up to two orders of magnitude more precise than LHC ($\sqrt{s} = 14$ TeV, 300 fb⁻¹) Disentangling of vecto/axial vectol couplings for ILC One variable at a time For LHC However LHC projections from 8 years old study
- Need to control experimental (e.g. Top angle) and theoretical uncertainties (e.g. Electroweak corrections)
 > Dedicated work has started
- Journal paper of results in preparation

ILC will be indeed high precision machine for electroweak top couplings

Discussion of potential systematic uncertainties

Experimental

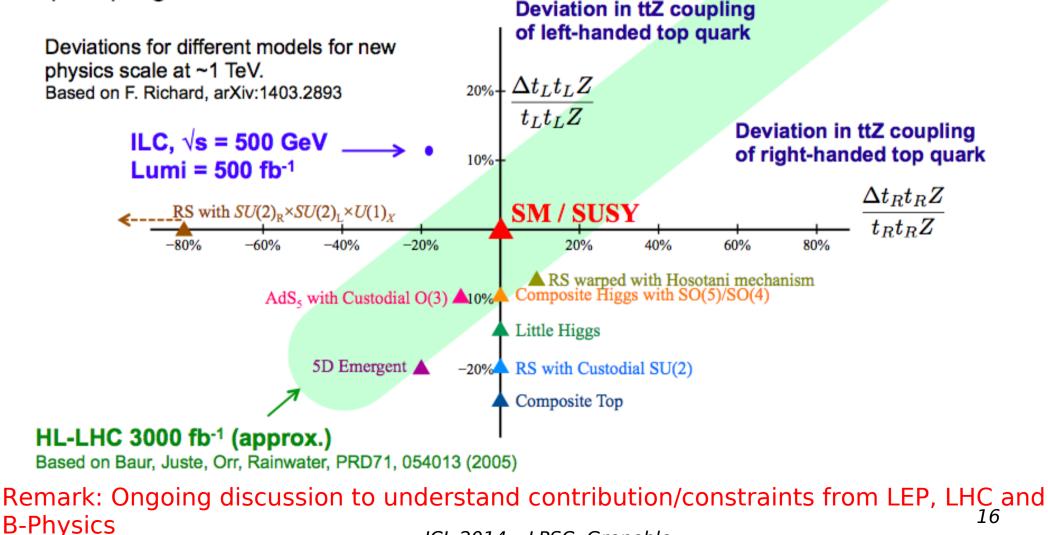
- Luminosity: Critical for cross section measurements Expected precision 0.1% @ 500 GeV
- Beam polarisation: Critical for asymmetry measurements Expected to be known to 0.1% for e- beam and 0.35% for e+ beam
- Migrations/Ambiguities: Critical for AFB: Need further studies but expect to control them better than the theoretical error
- Jet energy scale: Critical for top mass determination Systematic study CLIC states systematic error ~ statistical error
- Other effects: B-tagging, passive material etc. LEP claims 0.2% error on R_{h} -> guiding line for LC

Theory:

- See above
- Issue of single top under study

Sensitivity to New Physics

Composite Higgs theories have an impact on the top sector. Composite Higgs models can be tested at the ILC through precise measurements of the top couplings. Beam polarization (both e- and e+) is essential to distinguish the ttZ and tty couplings.



Sensitivities and constraints

q t	q' q
s-channel	$g \rightarrow \overline{b}$ <i>t</i> -channel

Model	dtR/tR %	dtL/tL %	dtLbL tLbL %	dɛb/ɛb	dɛ1/ɛ1	dOZtt/OZtt %
Carena	0	-20	-14	0.8	1.1	-30
Djouadi	-330	0	0	-1.4	1.1	70
Gherghetta	-20	-20	-14	0.7	2.1	-36
Grojean	0	10	7	-0.4	-1.0	17
Hosotani	18	-7	-5	-0.4	-0.8	-5
Little Higgs	0	-15	-10	0.6	1.0	-23
Pomarol	0	-25	-17	1.0	1.2	-37
Wulzer 1	25	25	17	-1.1	5.8	56
Wulzer 2	-10	-10	-7	0.4	1.3	-20

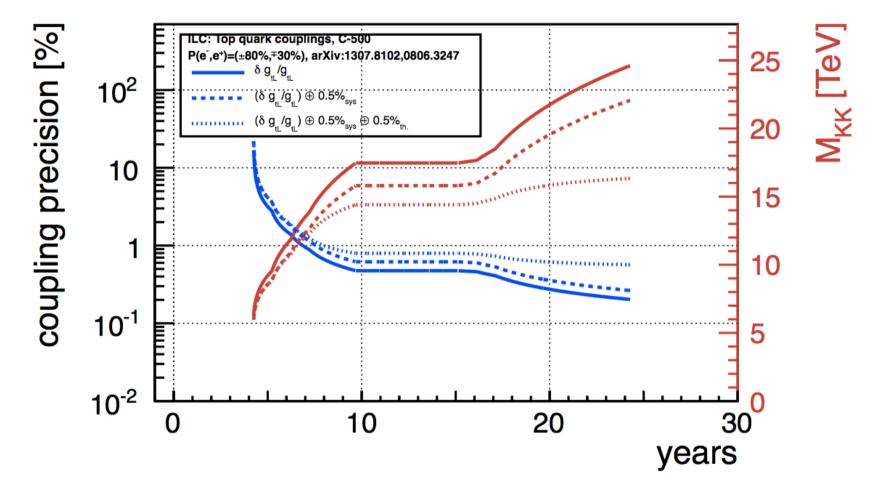
LEP constraints: $|\delta F^Z_{1A}| < 0.2, Q_{t_L} \rightarrow Q^{SM}_{t_L}$

=> LHC may see deviations but cannot distinguish Models
 => ILC will be able to distinguish at several sigma level

Example for physics reach

New physics reach for typical BSM scenarios with composite Higgs/Top And or extra dimenssions

Based on phenomenology described in Pomerol et al. arXiv:0806.3247



Can probe scales of ~25 TeV in typical scenarios (... and up tp 80 GeV for extreme scenarios)

Summary and outlook

- A LC is **the** machine for precision top physics

First machine to produce top pairs in electroweak production!!! Essential pillar of LC physics program

- Rich program of top quark physics with 'exciting' prospects
 - -Precision on top mass \sim 50 MeV => 'Final word' on vacuum stability of the universe
 - Test of models with extra dimensions and/or compositeness
 - Top elw. Measurements are complementary to Higgs coupling Measurements
- Exploitation of potential requires huge experimental and theoretical efforts
 - Theoretical uncertainty on top mass >> Experimental uncertainty
 - Uncertainty of theoretical prediction of AFB NNLO would be 10 years of work !!!
 - Measurement of b quark charge still in infancy, may need revision of algorithms and detector
 - In general experimentalists will have to make sure that systematic errors can be kept small

Towards a coherent approach (Theory and experiment)

1st workshop on top physics at LC - March 2013 at LPNHE/Paris

Group Topic		Midterm goals
University of Vienna Top mass theory		Elw./unstable particle for σ_{tot} .
MPI Munich	Top mass experiment	
University Tohoku	$tar{t}$ threshold	A_{FB} at threshold
WHIZARD [1]	$tar{t}$ threshold	Correct NLL/NLO matching
	Anomalous couplings	
GRACE	Elw. corrections	Elw. NLO for polarised beams
KEK	Japanese contact for top studies	
	within TYL^1	
LAL	Top couplings experiment	b charge determination
	Elw. corrections	Collab. with GRACE/New observables
	Phenomenology	Interpretation of results
	French contact for top studies	
	within TYL	
IFIC	Top couplings experiment	Role of single top
	Elw. corrections	Collab. with Spanish theory groups
	Phenomenology	Interpretation of results
DESY Zeuthen	Top couplings theory	"Resurrection" of NLO calculations

¹ French-Japanese virtual laboratory

Table 1: Non exhaustive list of groups working on LC top quark physics. The table reflects the status of Spring 2014.

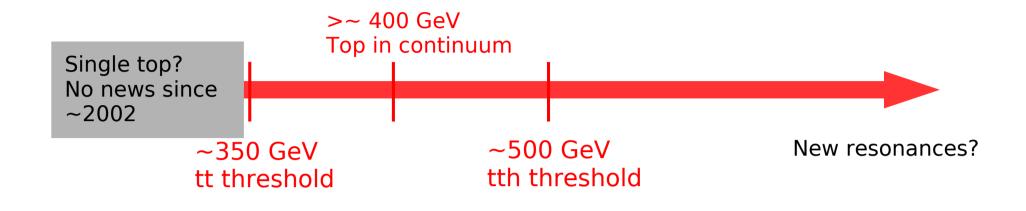
- Mailing list: topatlc-l@in2p3.fr

No new structure but lightweight forum on issues of top physics

- Stay tuned for 2nd workshop in Spring 2015

Backup

Relevant scales for Top physics and LC Physics programme



 After TDR and Japanese initiative, programme for ILC under discussion
 ILC in staged approach but which is first stage?

- Arguments to start at 350 GeV include Top physics programme

Relevant cross sections

type	final state	σ 500 GeV	σ 352 GeV
Signal ($m_{top} = 174 \text{ GeV}$)	tī	530 fb	450 fb
Background	WW	7.1 pb	11.5 pb
Background	ZZ	410 fb	865 fb
Background	$q\bar{q}$	2.6 pb	25.2 pb
Background	WWZ	40 fb	10 fb

Remarks:

- LC will have polarised beams => $(\sigma_{tt})_{L} \sim 1565 \text{fb}^{-1}$, $(\sigma_{tt})_{R} \sim 724 \text{fb}^{-1}$ at 500 GeV - Background varies differently with polarisations
- e.g. WW-Background $\rightarrow 26000 \text{ fb}^{-1}$ for e, and 150 fb⁻¹ for e_R

Form Factors and observables I

Form factors from previous pages re-written:

Cross section more explicitly (Stay in SM for the moment):

$$\begin{split} \sigma_{I} &= 2\mathcal{A}N_{c}\beta\left[(1+0.5\gamma^{-2})(\mathcal{F}_{1V}^{I})^{2} + (\beta\mathcal{F}_{1A}^{I})^{2} + 3F_{1V}^{I}F_{2V}^{I}\right] \\ &= \\ \sigma_{I} &= 2\mathcal{A}N_{c}\beta\left[(1+0.5\gamma^{-2})((Q_{e}Q_{t})^{2} + 2Q_{e}Q_{t}c_{L}^{e}f(s)F_{1V}^{Z} + 2Q_{e}Q_{t}c_{L}^{e}f(s)F_{1V}^{Z} + (\beta\mathcal{F}_{1A}^{I})^{2} + ...\right] \end{split}$$

Form Factors and observables II

Differential cross section:

 $\frac{d\sigma}{d\Omega} = \frac{\alpha^2}{4s} \left[A_0 (1 + \cos^2 \theta) + A_1 \cos \theta \right] \left\{ \begin{array}{c} \sim (1 + \cos^2 \theta) & \text{'Usual' Vector current, symmetric in } \cos \theta \\ & \sim \cos \theta & \text{Axial Vector current, asymmetric in } \cos \theta \end{array} \right.$

Forward Backward Asymmetry

$$(A_{FB}^{t})_{I} = \frac{-3\mathcal{F}_{1A}^{I'}(\mathcal{F}_{1V}^{I} + \mathcal{F}_{2V}^{I})}{2\left[(1+0.5\gamma^{-2})(\mathcal{F}_{1V}^{I})^{2} + (\mathcal{F}_{1A}^{I'})^{2} + 3\mathcal{F}_{1V}^{I}\mathcal{F}_{2V}^{I}\right]},$$

- Key observable to test chiral structure of Ztt (Zff) vertex
- Sensitive to amount of left-right asymmetry in interaction
 New physics may reduce asymmetry (→ left-right symmetric)

25

More on Form Factors

Vector/axial vector Form factors SM values (Better known as c_v, c_a):

$$F_{1V}^{Z} = \frac{2}{3} = Q_{t}, \ F_{1V}^{Z} = \frac{1}{4} - \frac{2}{3} \sin^{2}\theta_{W} / (\sin\theta_{W}\cos\theta_{W}), \ F_{1A}^{Z} = -\frac{1}{4} / (\sin\theta_{W}\cos\theta_{W})$$

 $F_{1A}^{\gamma} = 0$... and <u>always</u> 0 in SM No axial coupling to photon, QED gauge invariance

Tensorial couplings (all 0 at tree level):

Magnetic Dipole Moment:

$$F_{2V}^{\gamma}(0) = Q_t \frac{(g_t - 2)}{2}$$

(Anomalous) magnetic moment of top quark Scattering of particle in magnetic field, $(g_t - 2) \neq 0$ due to higher order corrections, 'not pointlike' anymore) Similar interpretation holds for F_{2V}^Z

Electrical Dipole Moment:

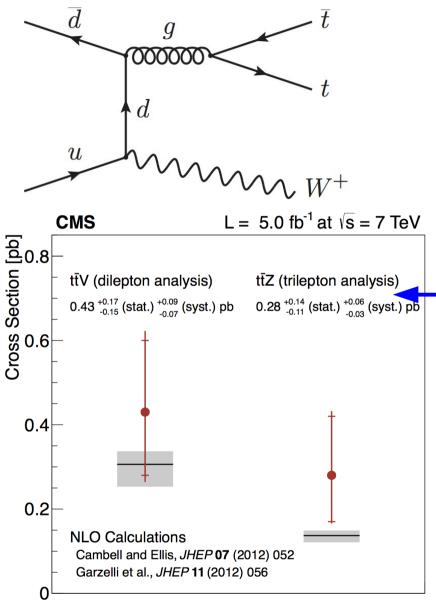
Bound state effects (e.g. vertex corrections) may create electrical dipole d extremely small in SM, O(10⁻¹⁴)

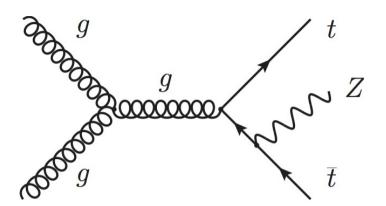
$$F_{2A}^{\gamma}(0) = ed_t^{\gamma}$$

CP Violating Any non-zero value measured 'today' is sign of BSM May receive contributions from CP Violating Higgs Similar interpretation holds for F_{2A}^Z

The race is open !

Recent result on ttV by CMS

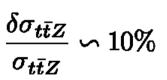




 $\sigma(t\bar{t}Z) = 0.28^{+0.14}_{-0.11} \,(\text{stat.}) \,{}^{+0.06}_{-0.03} \,(\text{syst.}) \,\text{pb}$

- Clearly, promising result
- How will it evolve with higher Luminosity?
- Revision of 'old' estimations of precisions are needed!

May expect:



Equations for cross section, A_{FB} and F_{R}

$$\sigma_I = 2\mathcal{A}N_c\beta \left[(1+0.5\gamma^{-2})(\mathcal{F}_{1V}^I)^2 + (\mathcal{F}_{1A}^{I'})^2 + 3\mathcal{F}_{1V}^I\mathcal{F}_{2V}^I \right],$$

$$(A_{FB}^{t})_{I} = \frac{-3\mathcal{F}_{1A}^{I'}(\mathcal{F}_{1V}^{I} + \mathcal{F}_{2V}^{I})}{2\left[(1+0.5\gamma^{-2})(\mathcal{F}_{1V}^{I})^{2} + (\mathcal{F}_{1A}^{I'})^{2} + 3\mathcal{F}_{1V}^{I}\mathcal{F}_{2V}^{I}\right]},$$

$$(F_R)_I = \frac{(\mathcal{F}_{1V}^I)^2 (1 + 0.5\gamma^{-2}) + (\mathcal{F}_{1A}^{I'})^2 + 2\mathcal{F}_{1V}^I \mathcal{F}_{1A}^{I'} + \mathcal{F}_{2V}^I (3\mathcal{F}_{1V}^I + 2\mathcal{F}_{1A}^{I'}) - \beta \mathcal{F}_{1V}^I \Re \mathfrak{e}(\mathcal{F}_{2A}^I)}{2 \left[(1 + 0.5\gamma^{-2}) (\mathcal{F}_{1V}^I)^2 + (\mathcal{F}_{1A}^{I'})^2 + 3\mathcal{F}_{1V}^I \mathcal{F}_{2V}^I \right]}$$

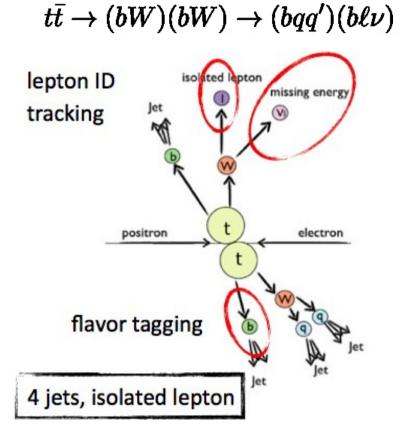
Elements of top quark reconstruction

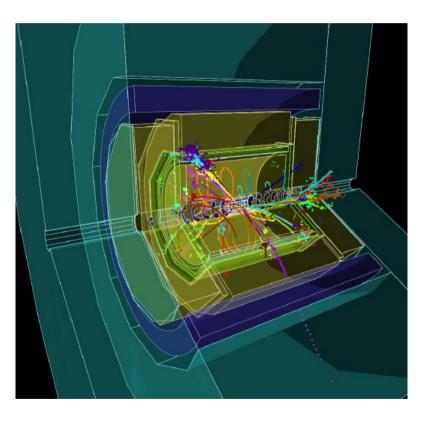
Three different final states:

1) Fully hadronic (46.2%) \rightarrow 6 jets

2) Semi leptonic (43.5%) → 4 jets + 1 charged lepton and a neutrino

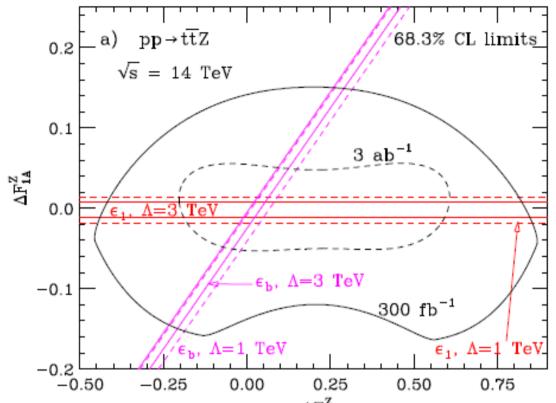
3) Fully leptonic (10.3%) \rightarrow 2 jets + 4 leptons



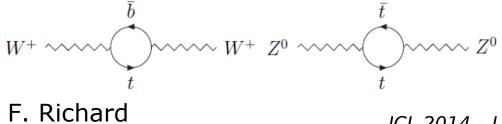


Results in the following mainly based on semi-leptonic decay Do however integrate results from fully hadronic study

Recap: LEP/SLD Constraints



- Recall that if one modifies the fermion \mathbb{E} couplings the SM loops becomes UV divergent and this requires introducing a **cutoff** $\pi \sim \text{TeV}$ to compute these contributions
- Given this cutoff the top EW couplings anomalies are limited by LEP/SLD measurements



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Constraints due to Gauge Invariance

Gauge invariance relates ZtLtL to WtLbL and ZbLbL

$$\kappa_{bL}^{NC} + \kappa_{tL}^{NC} \sim \kappa_{tL}^{NC} = 2\kappa_{tLbL}^{CC}$$

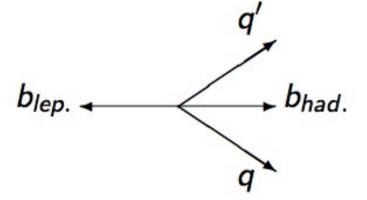
 From LEP1 we know that ZbLbL has no anomaly meaning that

$\frac{\delta WtLbL}{=}$	$0.72 \frac{\delta ZtLtL}{\delta ZtLtL}$	
WtLbL	$\frac{0.72}{ZtLtL}$	

- Loop contributions therefore fully constrain ZtLtL and ZtRtR and the only freedom left comes from BSM compensating contributions to ε1 and εb

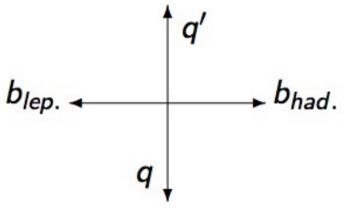
Experimental challenge b-charge reconstruction - Motivation

- To measure $A_{_{FB}}$ in fully hadronic decays there is no choice
- In semi-leptonic decays there is the charged lepton but



Right handed electron beam:

- mainly right handed tops In final state (V-A)
- <u>Hard W</u> in flight direction of Top and soft b's
- Flight direction of t from flight direction of W

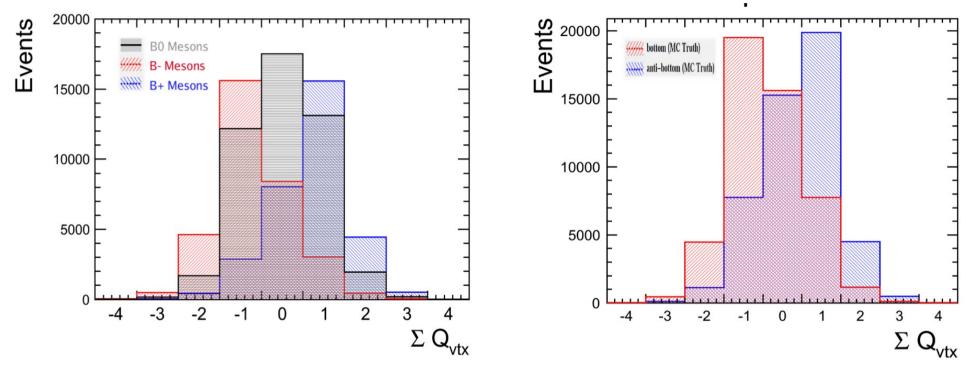


Left handed electron beam:

- mainly left handed tops
- <u>Hard b</u> in flight direction of Top and soft W's
- Flight direction of t from flight direction of b
- => Wrong association ↔ top flip

Measurement of b-charge to resolve ambiguities

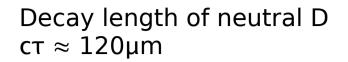
Measurement of b quark charge (N.B. At example of fully hadronic analysis, PhD M.S. Amjad)



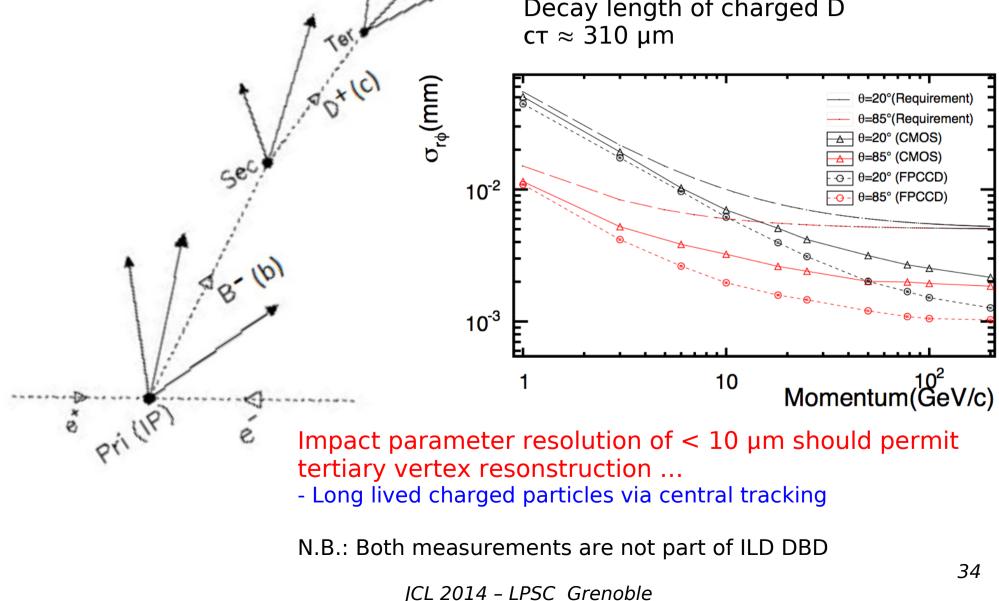
- Vertex charge measurement mandatory for fully hadronic top decays
- LC vertex and tracking system allows for determination of b-meson (b-quark) charge
 B-quark charge measured correctly in about 60% of the cases
 Can be increased to 'arbitrary' purity on the expense of smaller statistics
- LCFIPlus package not yet optimised for vertex charge measurement

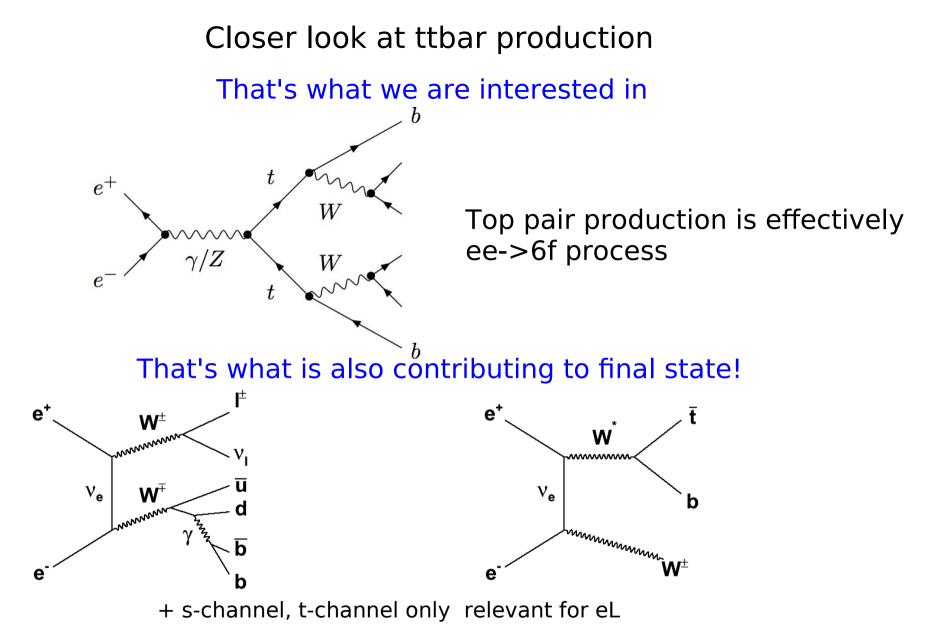
Optimisation of b-quark charge is major topic for future studies

Tertiary vertices – Principal considerations



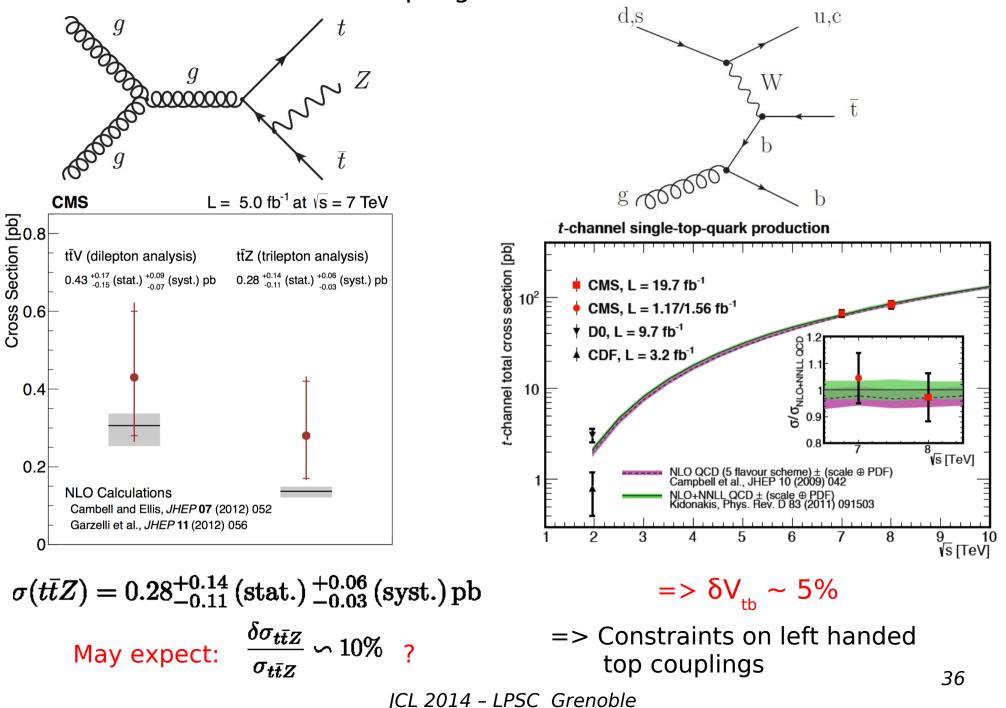
Decay length of charged D





- Can one really speak about a ttbar cross section?
- If only 6f is relevant: What are relations to ttX couplings?
- What selection cuts are (theoretically) save?

Electroweak couplings – LHC contributions



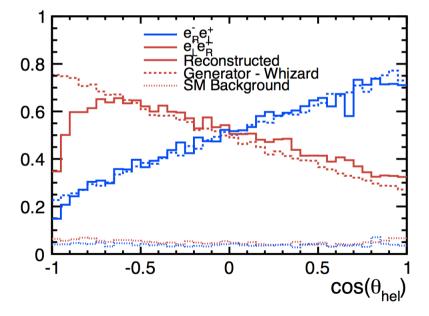
Measurement of top quark polarisation

Measure angle of decay lepton in <u>top quark rest frame</u> Lorentz transformation benefits from well known initial state (N.B. : Proposal for hadron colliders applied to lepton colliders)

Differential decay rate

 $\frac{1}{\Gamma}\frac{d\Gamma}{dcos\theta_{\ell}} = \frac{1+\lambda_t cos\theta_{\ell}}{2} \ \, {\rm with} \ \, \lambda_t = 1 \, {\rm for} \, t_{\rm R} \, {\rm and} \, \lambda_t = 1 \, {\rm for} \, t_{\rm L}$

Slope measures fraction of $t_{R,I}$ in sample



Measurement of decay lepton
almost 'trivial' at LC
High reconstruction efficiency for leptons
Reconstructed slope coincides
with generated slope

Slope λ_{L} can be measured to an accuracy of about 3-4%

$$F_{2A}^{\gamma}(0) = Q_t d_t^{\gamma}$$