The Accretion-Ejection connection

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Collaborators

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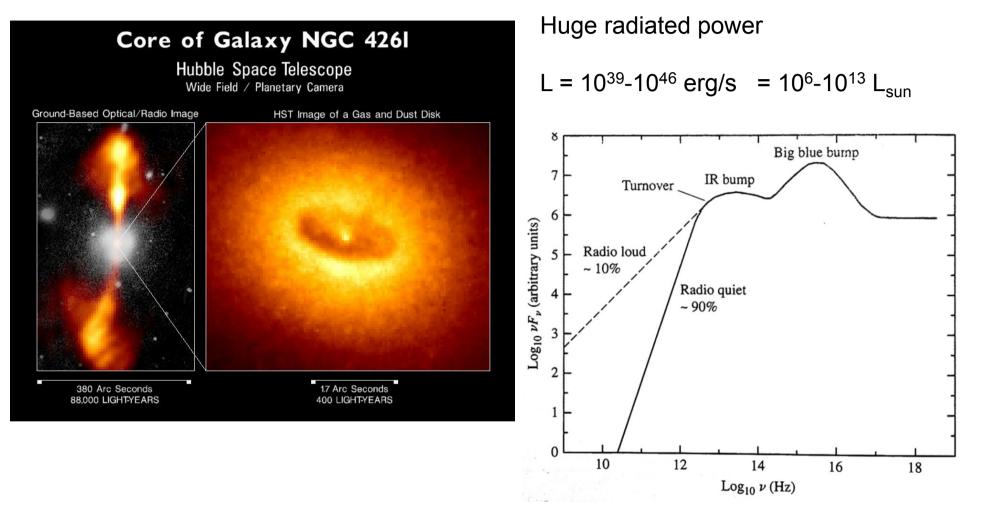
- Accretion ?
- Ejection ?
- Accretion-ejection correlations
- Jet-driven accretion (JED mode)
- Turbulence-driven accretion (SAD mode)
- Accretion-ejection cycles in X-ray binaries
- Role of the central object ?
- Conclusions



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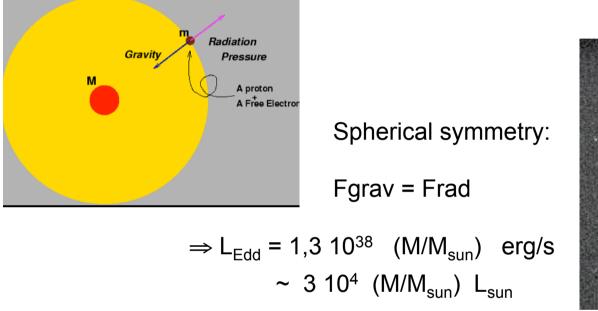
AGN (Active Galactic Nuclei) & Quasars



 \Rightarrow **Big Blue Blump**: spectrum cannot just be the sum of stars (« starburst » scenario)

 \Rightarrow Huge luminosity implies huge radiative pressure: how can this material remain there?

The Eddington luminosity limit





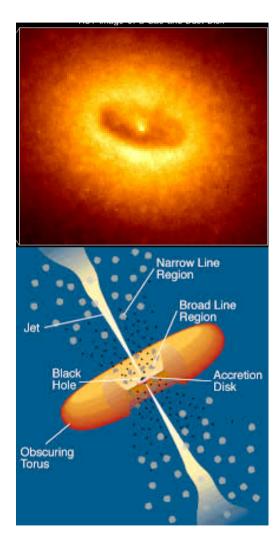
Ex: quasar 3C 273 has L= 3 10^{13} Lsun requires a **minimum** mass of 10^{9} M_{sun} !!

Variability requires luminosity emitted from region of size ~ few 100 au = few 10^{9} - 10^{10} km

$$\Delta t = r_g/c \qquad \qquad r_g = \frac{GM}{c^2} = 15km \left(\frac{M}{10M_{\odot}}\right)$$

⇒ Need of a central supermassive black hole M=10⁶-10¹⁰ M_{sun} for AGN & quasars ⇒ Where does this energy come from ?

The accretion disk paradigm: Lynden-Bell (1969)



Assume a rotating Keplerian disk around BH with
$$r_g = \frac{GM}{c^2}$$

 $E = \frac{u^2}{2} - \frac{GM}{r} = -\frac{GM}{2r}$
 $\Delta E = E(r_{in}) - E(r_{out}) = \frac{GM}{2r_{in}} \left(1 - \frac{r_{in}}{r_{out}}\right) \simeq \frac{GM}{2r_{in}}$

Assuming a mass flux through the disk $M_a\;$ leads to a released accretion luminosity

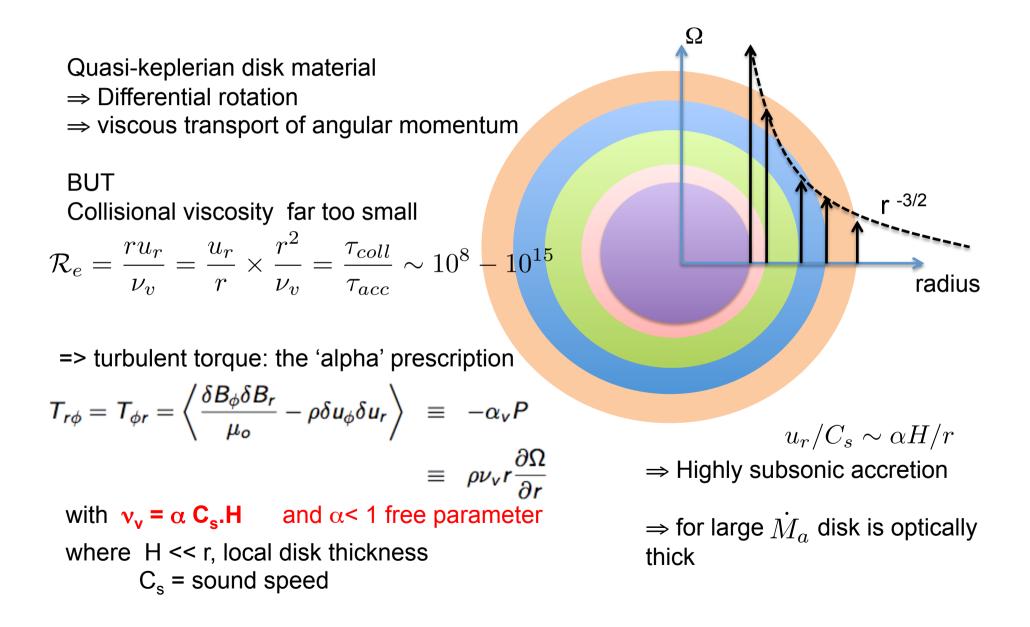
$$L = \dot{M}_a \Delta E = \frac{GM\dot{M}_a}{2r_{in}} = \dot{M}_a c^2 \left(\frac{r_g}{r_{in}}\right)$$
$$L = \eta \dot{M}_a c^2$$

with $\,\eta\,$ ~ 5 to 40% efficiency, depending on BH spin

Typical luminosities require BH fed with up to $10^{-2} - 1 M_{sun}/yr$

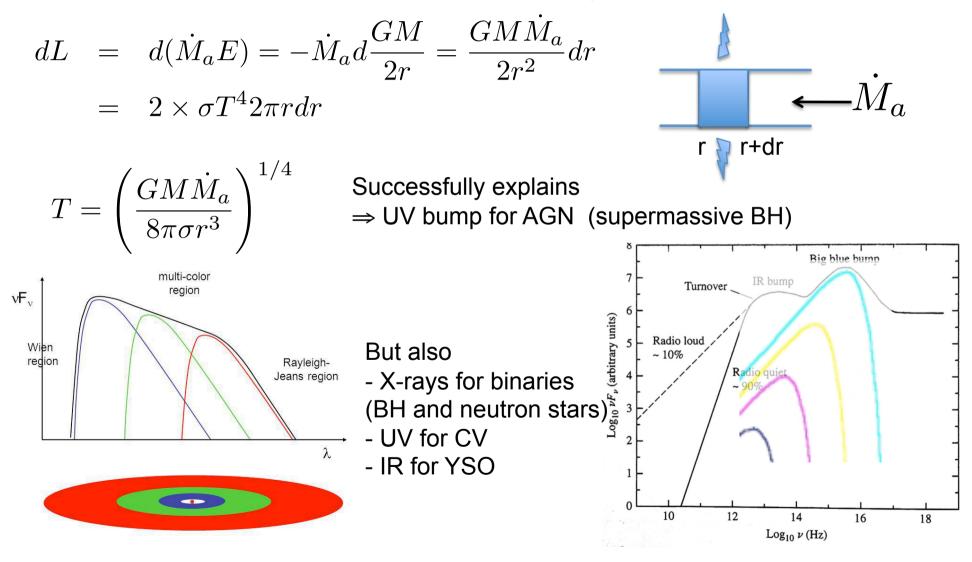
=> Need to find a way to brake down the rotating disk

The Standard Accretion Disk (SAD): Shakura & Sunyaev 1973

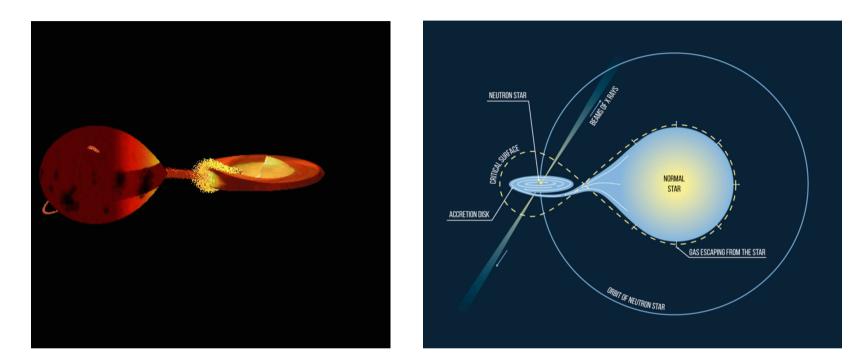


The Standard Accretion Disk (SAD): Shakura & Sunyaev 1973

Emitted broadband spectrum: sum of local blackbody of temperature T



Binary systems with mass transfer



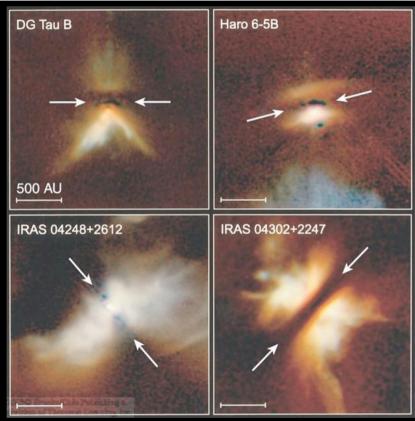
Compact object + normal star => accretion disk around compact object

- Compact object = White Dwarf => Cataclysmic Variable, seen in UV
- Compact object = BH or neutron star => X-ray Binary... seen in X-rays

Mass transfer via

- Roche-lobe overflow, for low-mass (M< 2 M_{sun}) star companion
- wind-fed, for high-mass (O/B M> 8 M_{sun}) companion

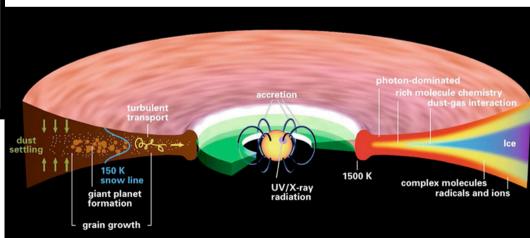
Young Stellar Objects (YSO) also



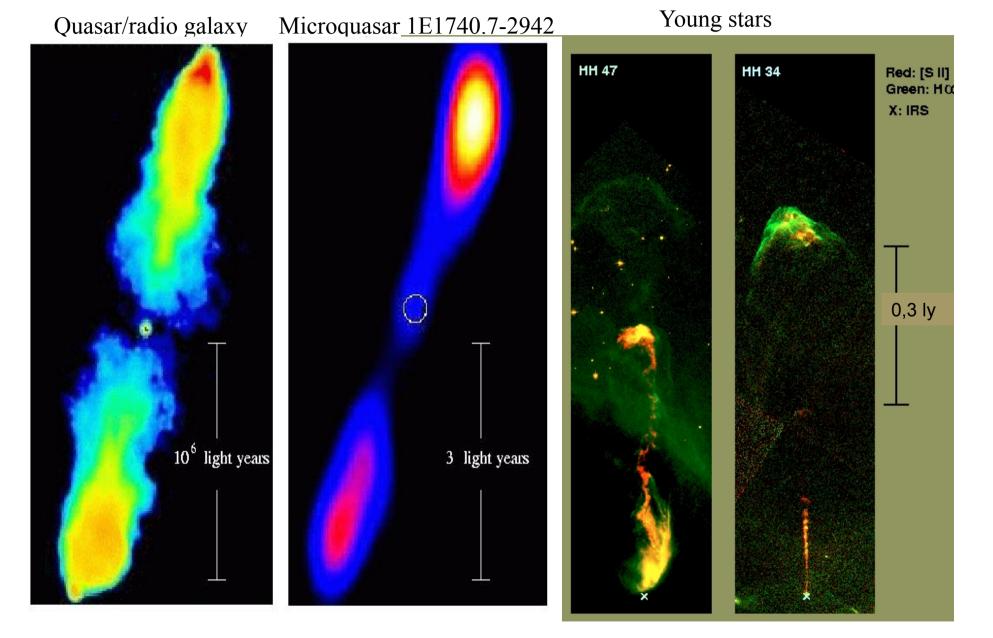
Gravitational collapse of a rotating cloud

=> Circumstellar disk= nursery of planets

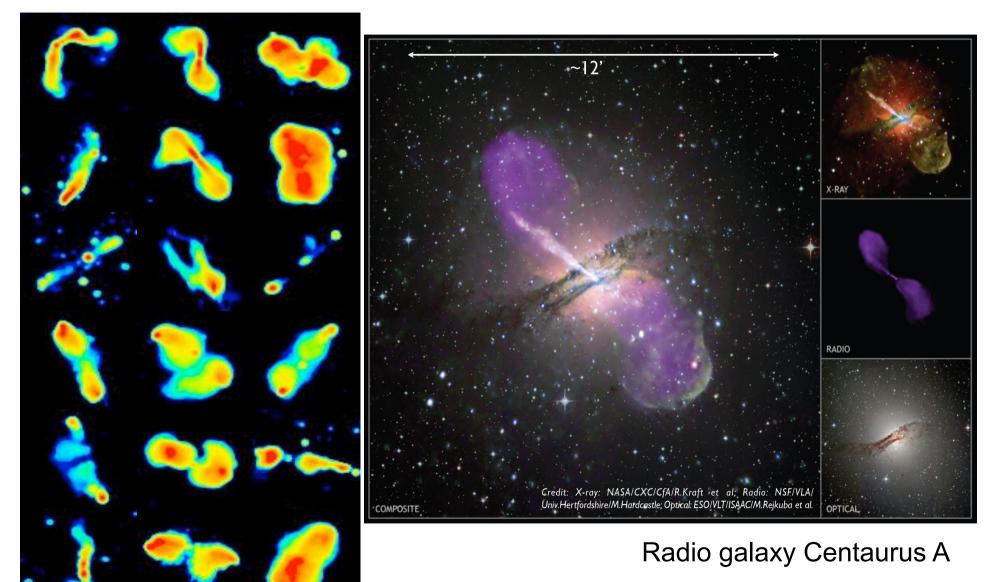
⇒ Disk formation around a protostar, seen as
 (i) absorbing (dust) layer in optical
 (ii) an infrared excess



Jets in all classes of accreting objects



Radio galaxies & Quasar gallery

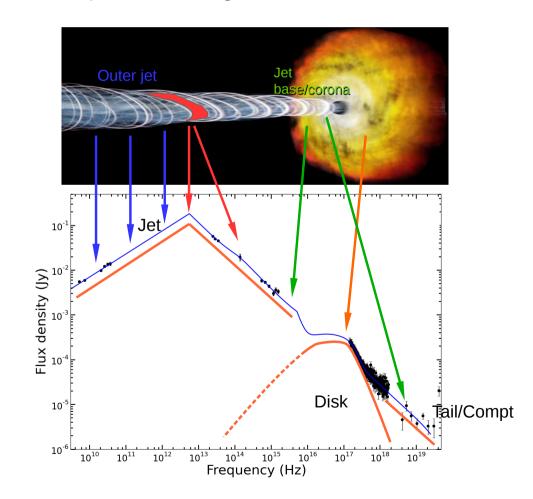


Leahy, JP

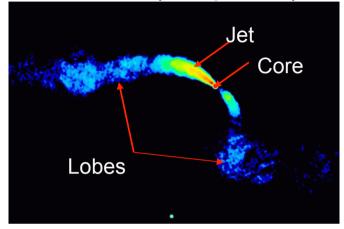
Jets from AGN & binary systems

Seen in Radio: synchrotron emission

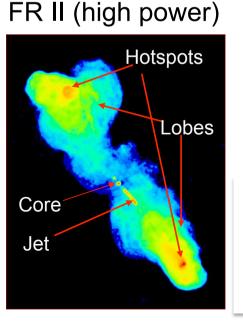
from non-thermal electron population ⇒ Magnetic fields present ⇒ Spectra + images : collimated flows



FR I (low power)



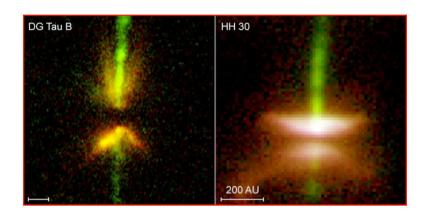
Flows slow to <0.3c on ~10 kpc scales

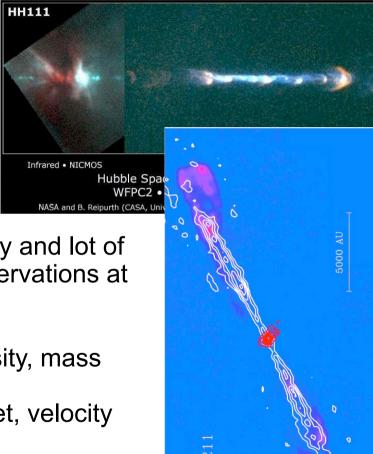


Flows likely still ≥0.7c on Mpc scales.

Brightest in the lobes

Jets from Young Stellar Objects





Large number of known YSOs, nearby and lot of information can be obtained from observations at different wavelengths

Optical & IR \rightarrow Temperature, density, mass

Radio \rightarrow ionized gas, base of the jet, velocity

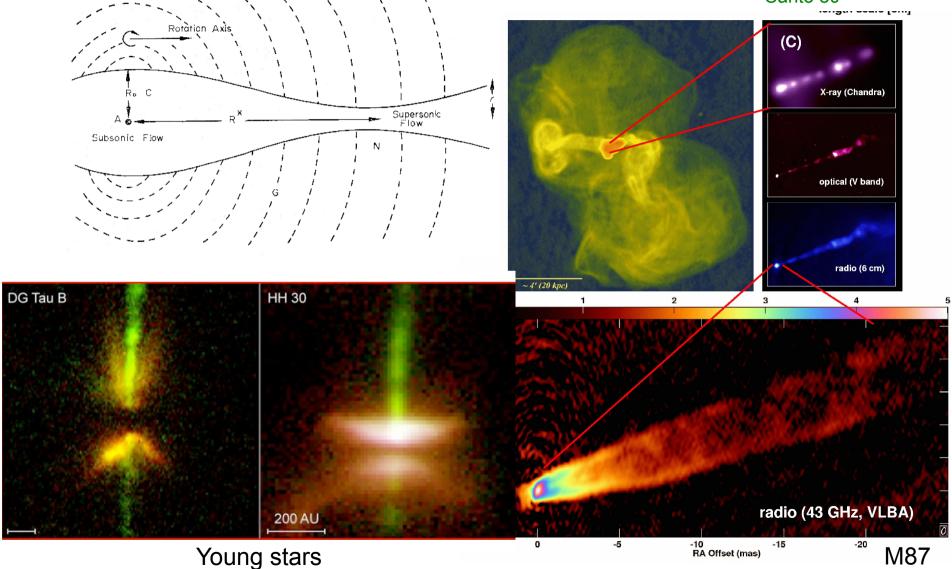
mm/submm \rightarrow Disk, molecular outflow

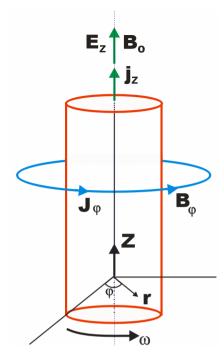
But **magnetic field**, very difficult to observe, specially in the jet, and we do not know very much about it

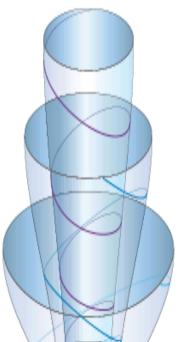
Same collimation issue

First tentative: a de Laval nozzle ?

Blandford & Rees 74 Canto 80







Magnetized jets

Blandford 76, Lovelace 76 Blandford & Payne 82

Jet = electron-proton **plasma** carrying a large scale helicoidal (Bz and Bphi) magnetic field

=> Magneto-hydrodynamics (MHD)

Axisymmetry => magnetic surfaces nested around each other, anchored onto a rotating object

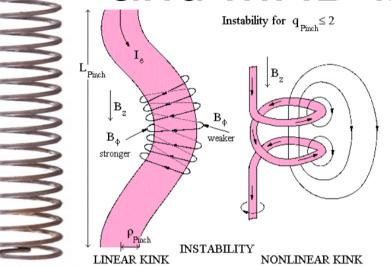
- central mass (BH, star)
- or surrounding accretion disk

Collimation = usual **hoop-stress** (Bphi) as in **Z-pinch** Controled by generalized Grad-Shafranov equation

Power = conversion of initial MHD Poynting flux into plasma kinetic energy (Bernoulli invariant)

Theory of steady-state jets is known... (it depends on 5 MHD invariants whose radial distribution must be given) but not solved yet :-/

and MHD instabilities !?



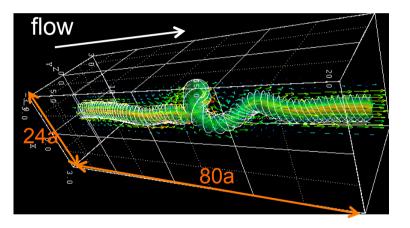
Since 60's, Z-pinch are known to be highly unstable to current-driven instabilities:

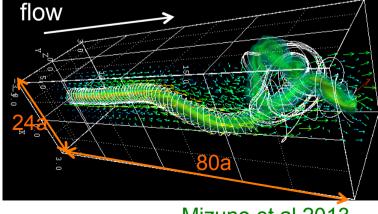
- sausage
- kink modes

May potentially destroy the jet, as in numerical simulations...

Why are real jets so stable ?

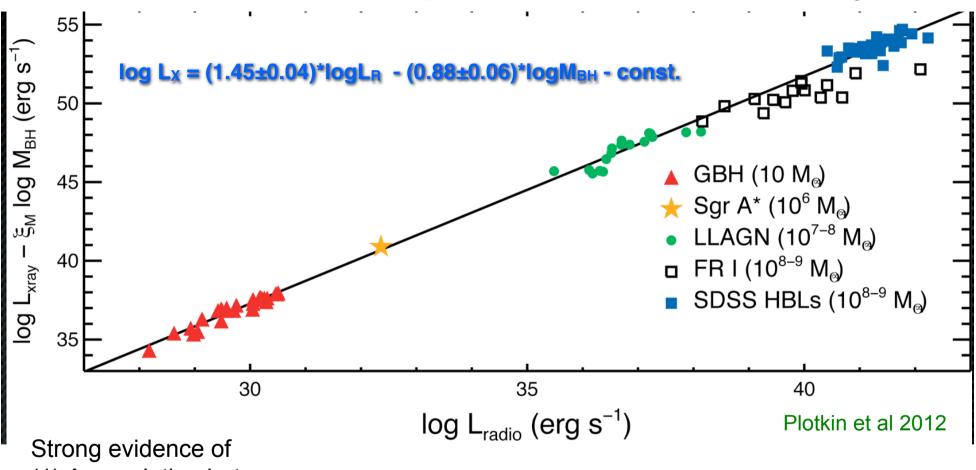
HINT: transport barrier due to differential rotation of magnetic surfaces => disk ?





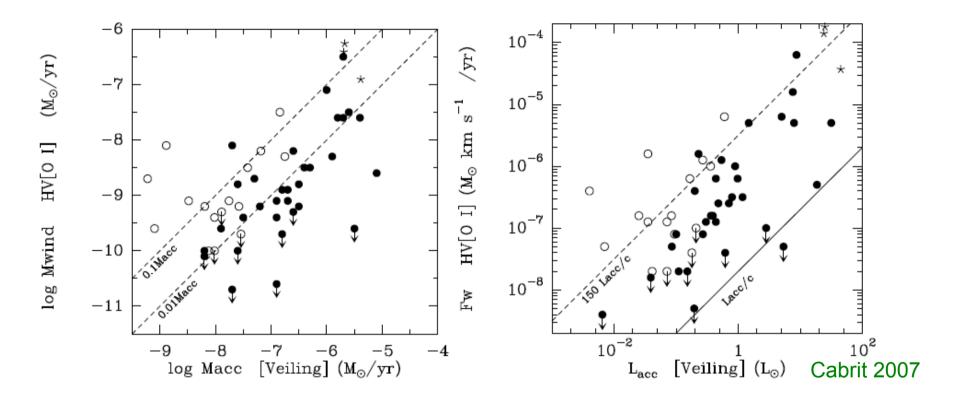
Mizuno et al 2013

Fundamental plane of BH activity



- (1) A correlation between
 - Accretion (using X rays as a proxy)
 - Ejection= steady jets, emitting self-absorbed synchrotron emission (radio)
- (2) Physics scaling with BH mass
 - => X-ray Binaries could be seen as micro- or even nano-quasars

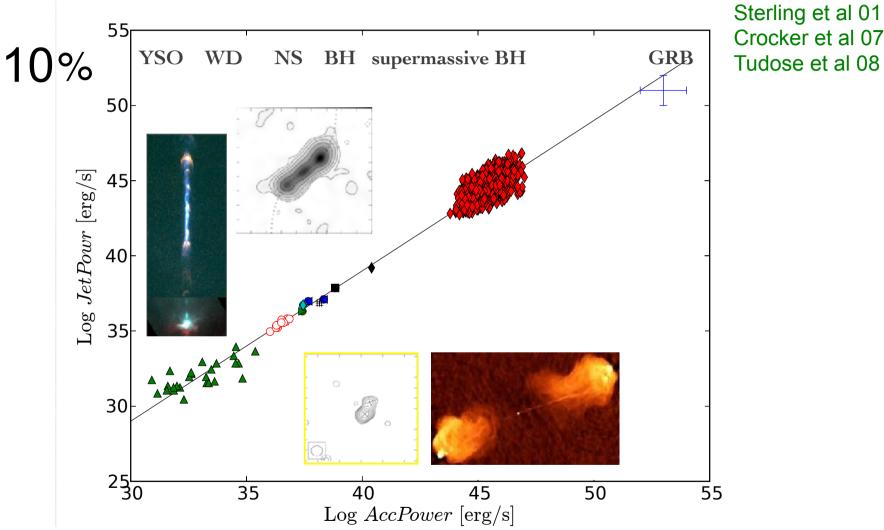
Accretion-Ejection correlation in YSO



(i) Mass loss in wind correlated with disk accretion rate

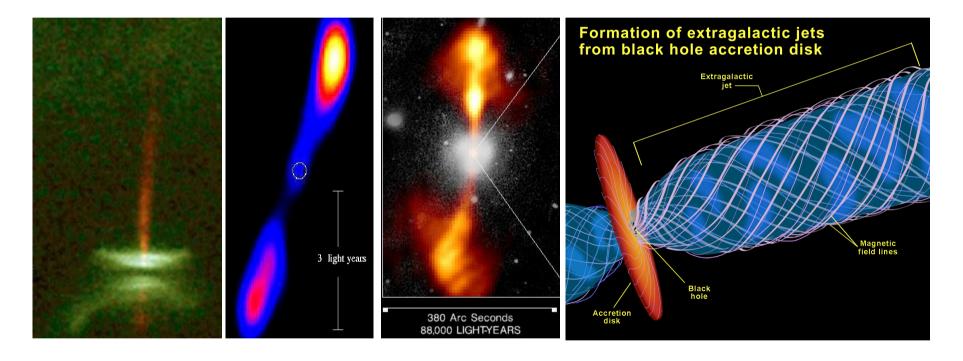
(ii) $F_w = M_{wind}$. V_{wind} jet momentum thrust >> radiation thrust: YSO jets cannot be radiatively driven

A universal correlation..?



Regardless of the nature of the central object ! => Look for an **interdependent accretion-ejection process**

Accretion-ejection in Astrophysics



Main assumption: a large scale magnetic field threads the disk



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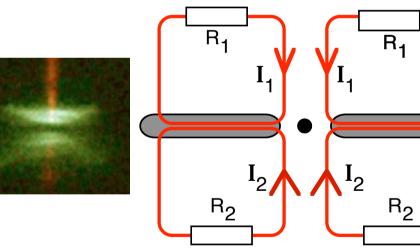
Disk as a unipolar inductor: 2 jets

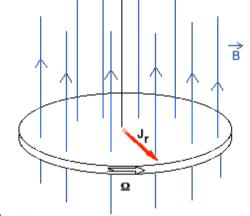
Barlow wheel (1822): unipolar induction effect

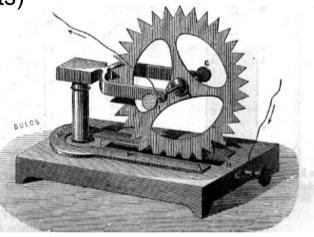
1) Gravitation + Magnetic Field => e.m.f

$$e = \int \Omega r \ B_z \ dr$$

2) e.m.f => electric current (2 independent circuits)

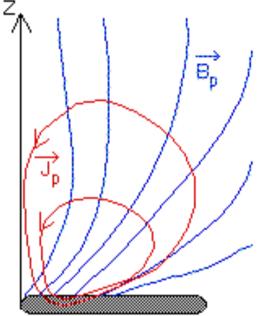


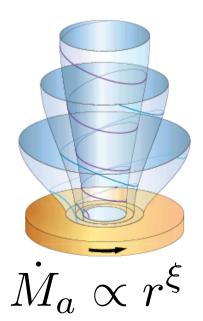




- 3) Conversion of mechanical energy into MHD Poynting flux
- 4) Existence of a torque braking down the disk => **accretion**
- 5) If $R_1 \neq R_2$, asymmetric jets are produced (mass flux, velocity)

The role of the poloidal electric current (Bphi)





Ideal MHD: Jet acceleration and confinement

Collimation due to magnetic hoop-stress (toroidal field) Heyvaerts & Norman 89, 03, Ferreira 97, Okamoto 01

$$F_{\phi} = \frac{B_{p}}{2\pi r} \nabla_{\parallel} I$$

$$F_{\parallel} = -\frac{B_{\phi}}{2\pi r} \nabla_{\parallel} I$$

$$F_{\perp} = B_{p} J_{\phi} - \frac{B_{\phi}}{2\pi r} \nabla_{\perp} I$$



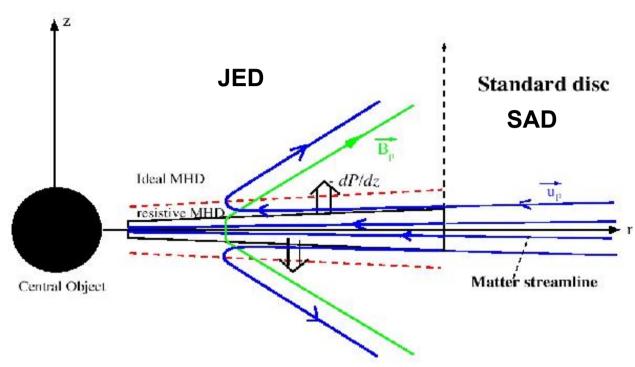
- Depends on *asymptotic* current distribution I(r)
- Not all field lines can be collimated: outer pressure required

Resistive MHD: Disc torque and mass loss

The disc ejection efficiency $\boldsymbol{\xi}$ must be computed as function of the disc parameters

=> NEW MHD flow model where parameter space is constrained by smoothly crossing critical points

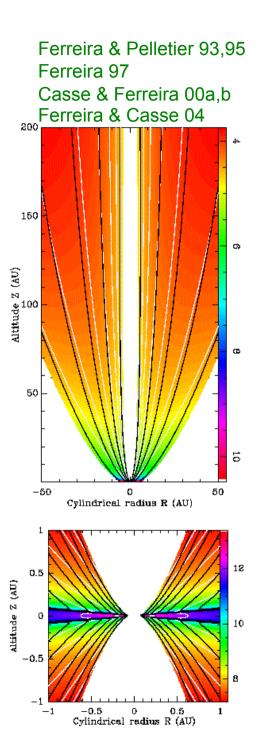
Jet Emitting Disks (JEDs)

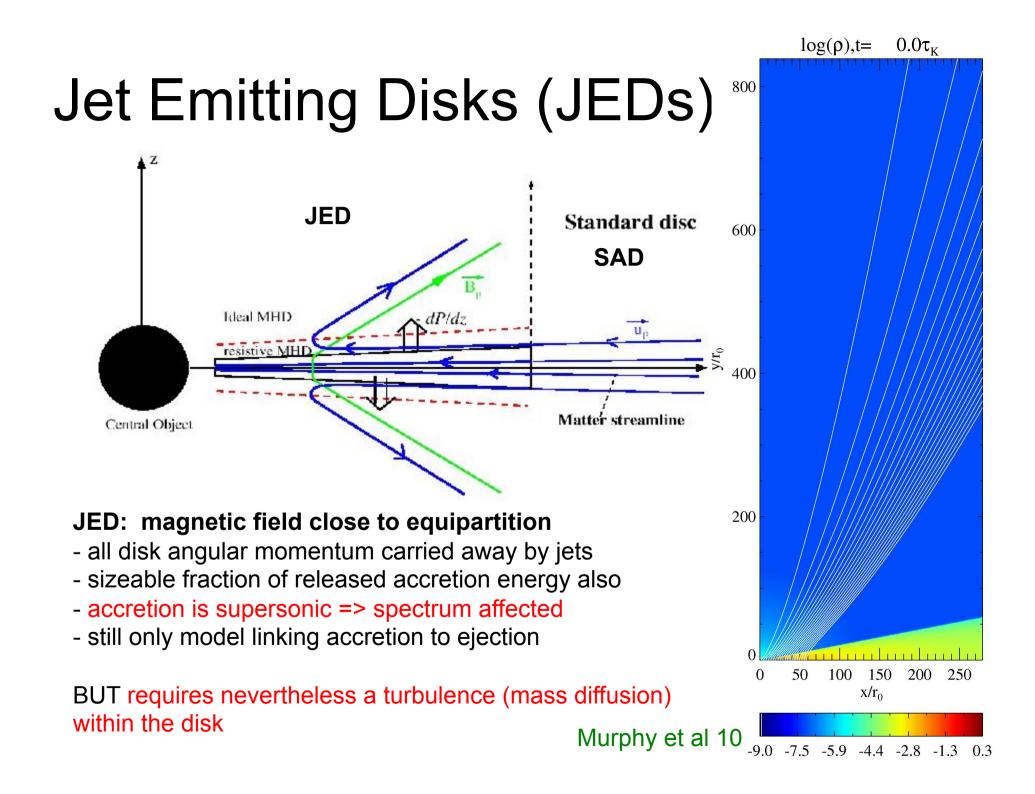


JED: magnetic field close to equipartition

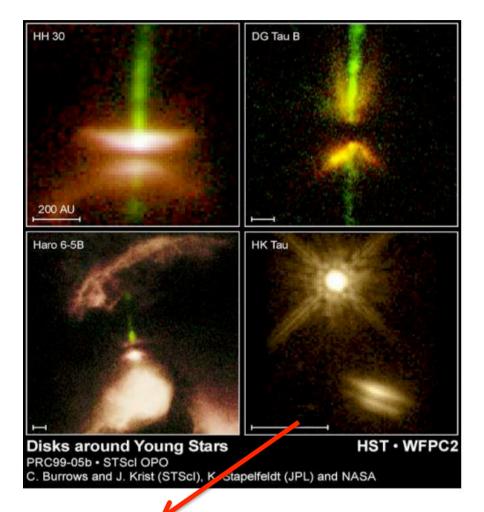
- all disk angular momentum carried away by jets
- sizeable fraction of released accretion energy also
- accretion is supersonic => spectrum affected
- still only model linking accretion to ejection

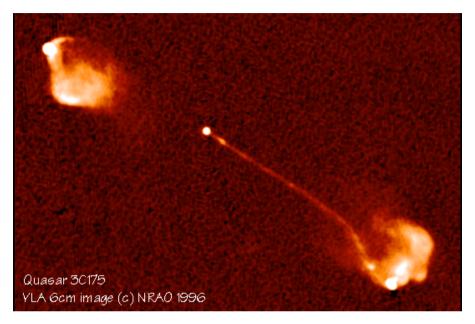
BUT requires nevertheless a turbulence (mass diffusion) within the disk





But JEDs are not the whole story





Only ~ 10% of AGN have jets

Not all YSO accretion disks have jets

=> Another mechanism of disk angular momentum removal must be at work

Back to the old idea of radial transport via turbulence (SAD)

Turbulence: ok, but which instability?

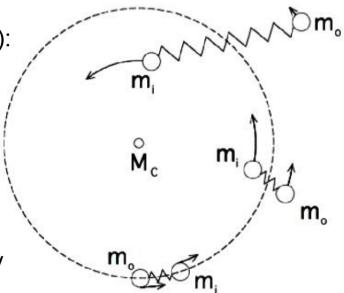
Shakura & Sunyaev 1973: the alpha prescription

BUT Keplerian disks are Rayleigh stable: 20 years of theoretical efforts within the context of hydro disks...

... Until Balbus & Hawley 1991: magnetic fields where introduced in disks

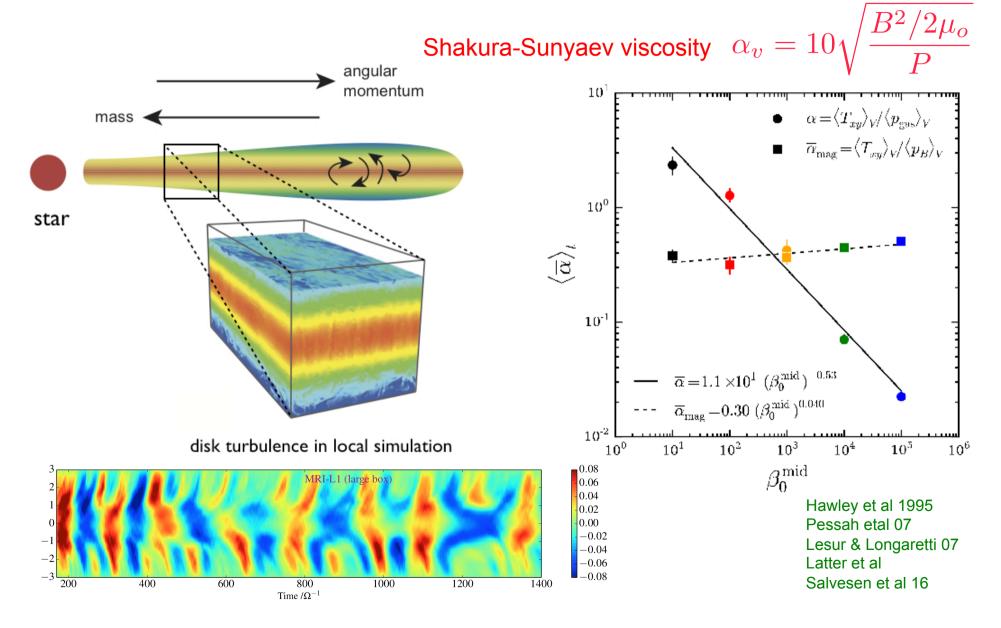
⇒ Existence of an ideal MHD instability (*): Magneto-Rotational Instability (MRI)

- Requires a sub-equipartition field
- Non-linear stage is a self-sustained TURBULENCE

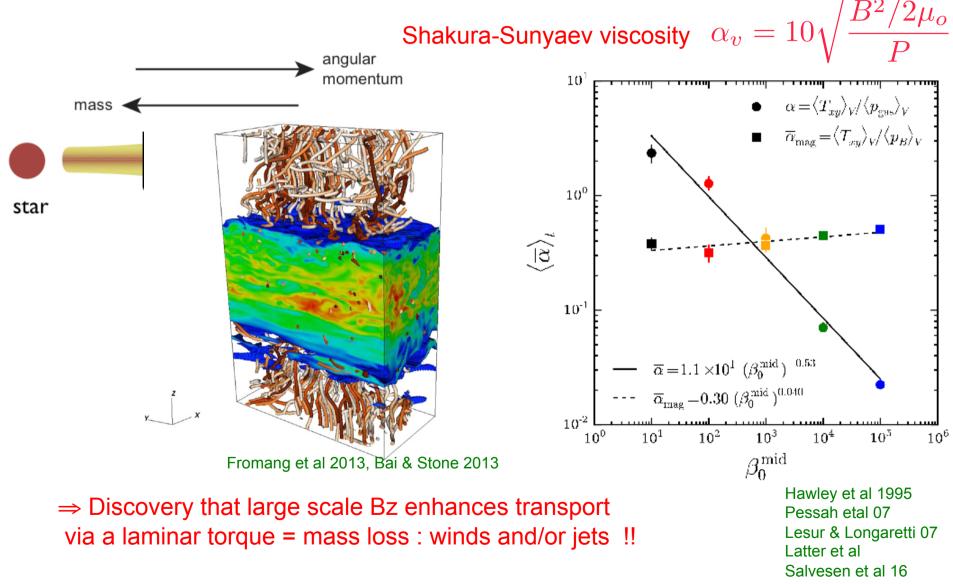


(*): requires a fully ionized plasma, partially quenched in non-ideal contexts (outer CV and YSO disks)

Shearing box (local) simulations



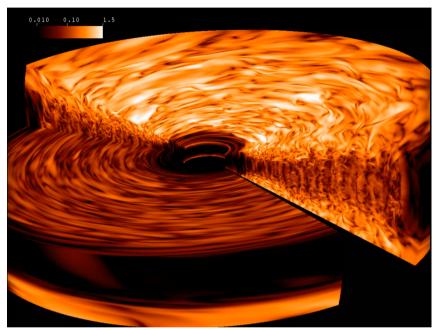
Shearing box (local) simulations



Need to go for global 3D simulations

« MRI-driven » winds: global simulations

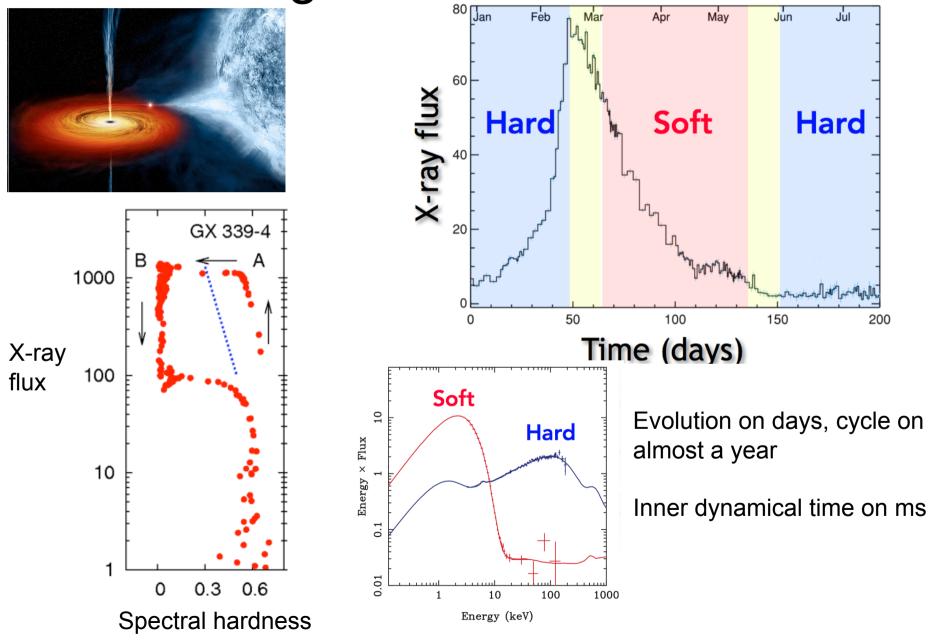
Flock et al 11



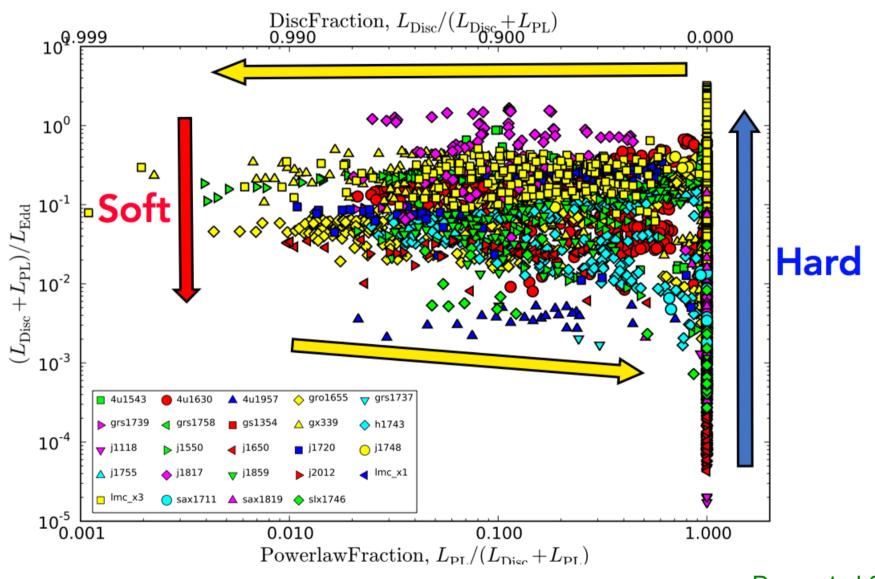
Without large scale Bz: accretion with no wind With large scale Bz: enhanced accretion speed and winds... or self-confined jets ??

Numerical challenge: following 3D turbulence and addressing large spatial scales for flow collimation Hint of flux accumulation: increasing magnetization? Suzuki & Inutsuka 14 Gressel et al 15 Zhu & Stone 17 Béthune et al 17

Outbursting cycles in XrB: GX339-4



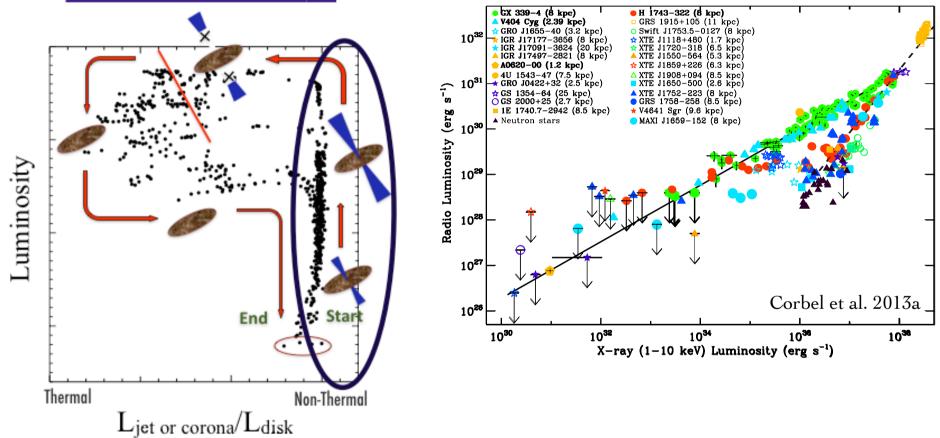
A quite generic behavior



Dunn et al 09

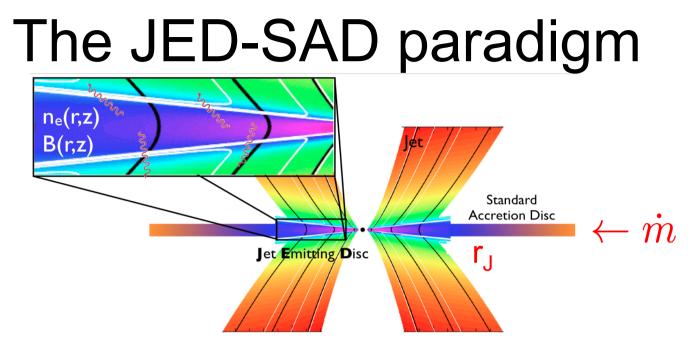
Accretion-ejection correlation

Stellar Mass: 10 Msun



- Jets always associated with HARD states, no-jet always in SOFT states

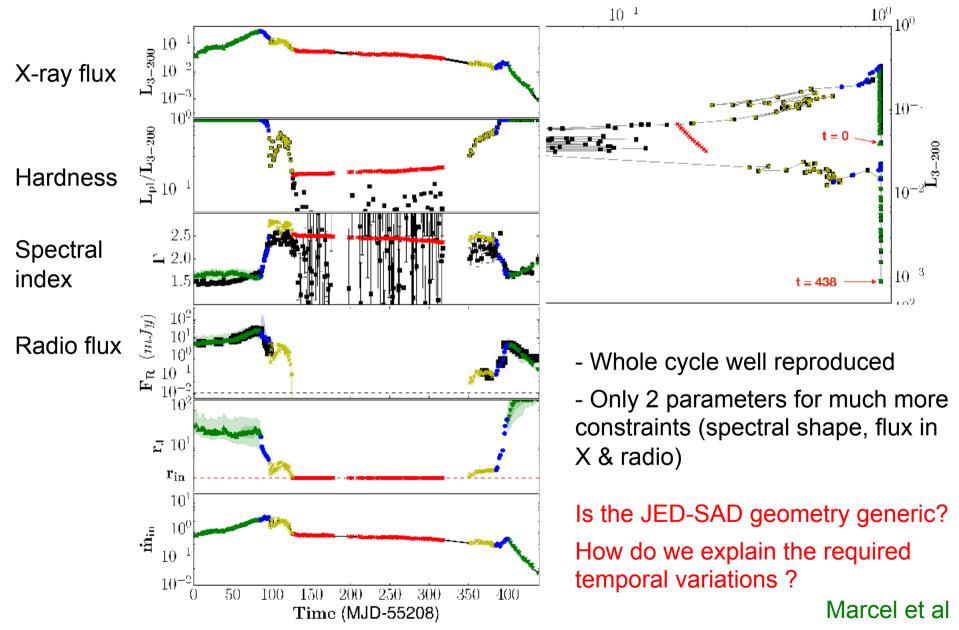
- Each « state » lasts for several days, object evolves on time scales >> local dynamical time scale



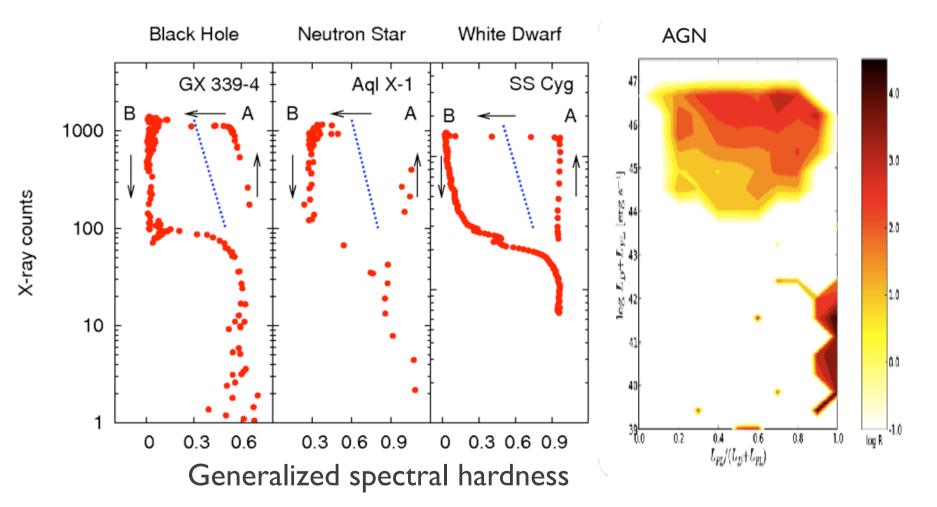
Assume that disk magnetization varies radially such that

- MRI-driven accretion from outer regions down to rJ (SAD)
- Jet-driven accretion from rJ down to BH (JED)
- => Use disk accretion rate \hat{m} and transition radius r_J as free parameters => Compute self-consistent energy equation + spectrum taking into account:
 - JED and SAD dynamical properties
 - optically thin emission (Synchrotron, Bremsstrahlung)
 - local and external comptonization of soft photons
 - collisional Coulomb coupling between ions and electrons
 - advection of energy

The 2010-2011 outburst of GX339-4

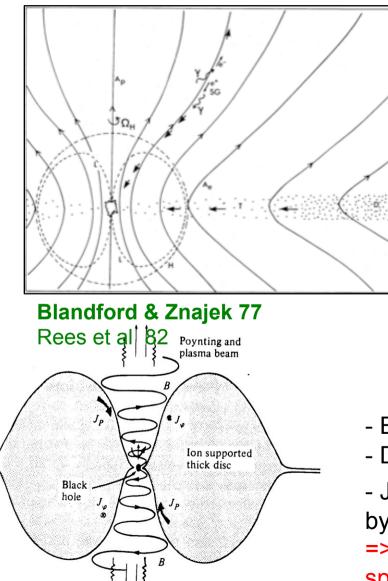


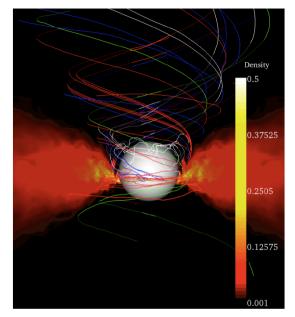
Accretion states of compact objets Körding et al, 2006, 2008



Does NOT seem to require a black hole, only the surrounding accretion disk. But what would be its influence ?

Large scale Bz field and rotating black holes: the Blandford-Znajek (1977) process



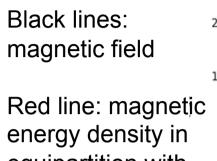


Punsly, Igumenshchev & Hirose 09 Tchekhovskoy et al 10,11 McKinney et al 12

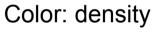
- Extract BH rotational energy
- Drive relativistic jet (spine)

- Jet power depends on magnetic flux brought in by outer accretion disk

=> Numerical challenge: density floor and huge spatial scales in 3D GRMHD



equipartition with rest mass energy density

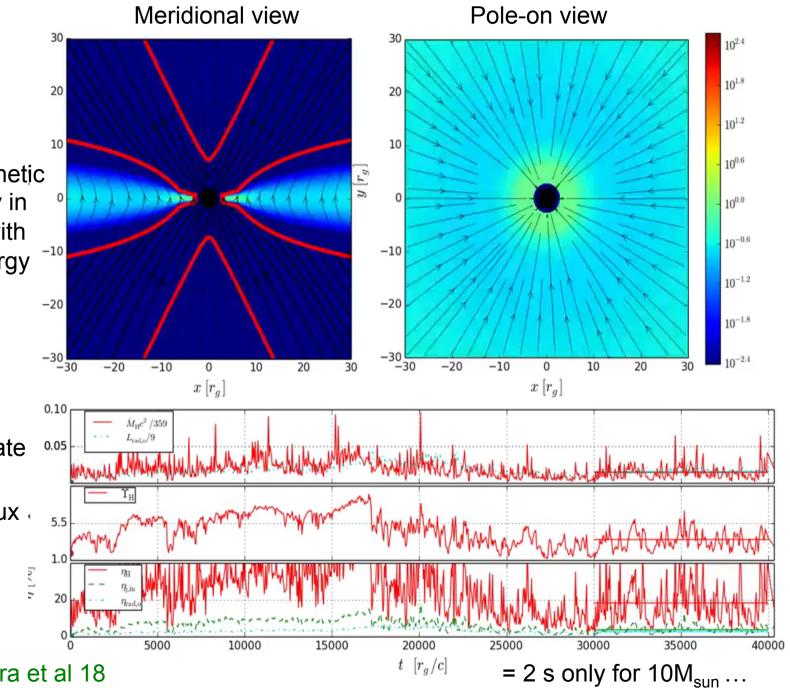




Magnetic flux

Efficiency

Morales Teixeira et al 18

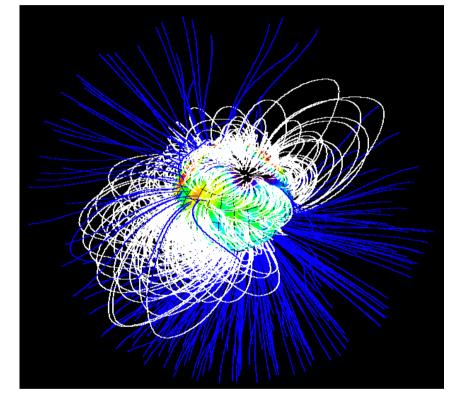


Magnetic star-disk interaction: YSO, neutron star, white dwarf

t = 0.0-1.25 -1.8 1 2

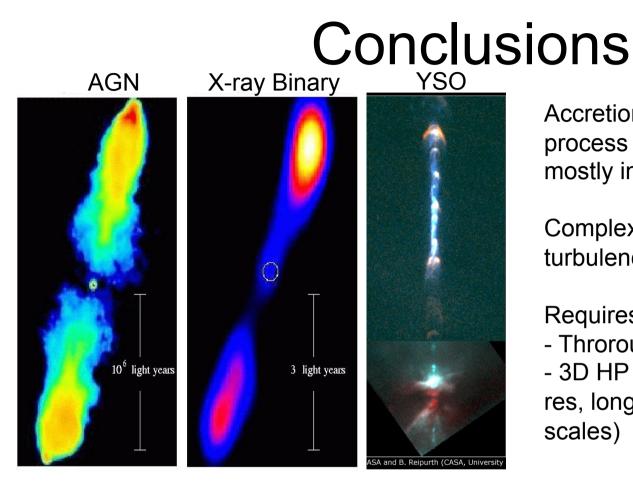
2D MHD simulations Zanni & Ferreira 09,13

3D YSO magnetic field maps: Donati et al



Unsteady ejecta @ interface:

- May provide efficient spin down of rotating object
- May affect large scale jet dynamics (collimation, jet emission via shocks)
- Numerical challenge: need to go 3D



Accretion-Ejection is a universal process (possibly also GRB, TDE), mostly independent of central object

Complex interplay between disk turbulence and large scale jets

Requires a feedback between - Throrough analytical models - 3D HP MHD computations (high res, long time scales, large spatial scales)

Process relies on the existence of a large scale magnetic field

- of unknown origin
- barely detectable

But this invisible agent is ultimately shaping the accretionejection process and its long term variability

