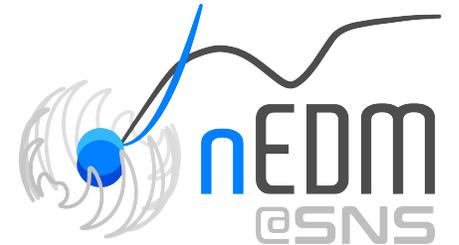


# Cryogenic UCN storage, super-thermal production, and live spin analysis for measurements of static nEDMs, oscillating EDMs, and the neutron magnetic moment

Kent Leung

“Les Houches” nEDM2021 workshop, 2021-02-16



# The nEDM@SNS experiment

Golub & Lamoreaux's technique: polarized UCN + polarized  $^3\text{He}$  atoms + superfluid  $^4\text{He}$

$$\text{Statistical "shot noise" limit: } \sigma(d_n) \sim \frac{\hbar}{2\alpha ET\sqrt{N}}$$

$E$  = electric field (75 kV/cm thanks to sf- $^4\text{He}$ . See Ito, Riley, Blatnik, Korsch talks.)

$\alpha$  = polarization contrast (UCN polarization  $\sim 98\%$ )

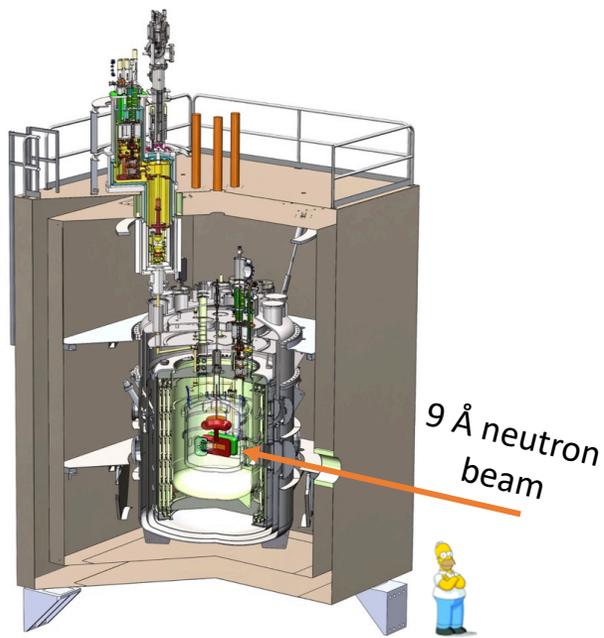
$T$  = precession time (goal to use 1000 sec if wall loss times low enough)

$N$  = no. detected neutrons (high density from in-situ super-thermal production & accumulation. Density increases with  $\tau_{\text{storage}}$ . No need to transport. But "filling" time is  $\sim 1000$  sec)

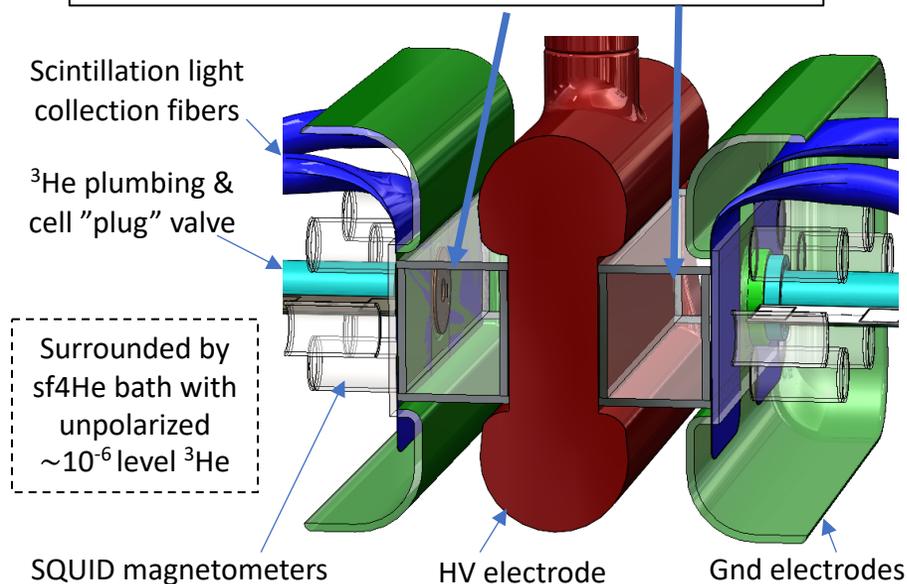
Polarized  $^3\text{He}$  ( $\sim 98\%$ ) serves as in-situ & live UCN spin analyzer (new type of signal!)

Low temperature ( $\sim 0.4$  K): SQUIDs and superconducting magnetic shielding

$^3\text{He}$  + SQUIDs as co-magnetometer. Small changes of sf $^4\text{He}$  temperature causes large changes ( $\sim T^{7.5}$ ) of  $^3\text{He}$  mean-free-path as dominated by  $^3\text{He}$ -phonon scattering. Great for co-magnetometer systematic checks .



2x measurement cells:  
inner 7.5 cm (W), 10 cm (H), 40 cm (D)  
filled with UCN,  $\sim 0.4$  K sf $^4\text{He}$ , and  $^3\text{He} \sim 10^{-10}$



# Overview

- Measurement cells: UCN storage and reduction of neutron beam-induced activation
- *in-situ* & live UCN spin analysis: a new type of precession signal (simulations)
- Advantages in searching for time-oscillating nEDMs & extraction of  $\gamma_n/\gamma_{3\text{He}}$

*Central themes: UCN storage, measurement cells, super-thermal production in  $4\text{He}$ , and signal analysis*

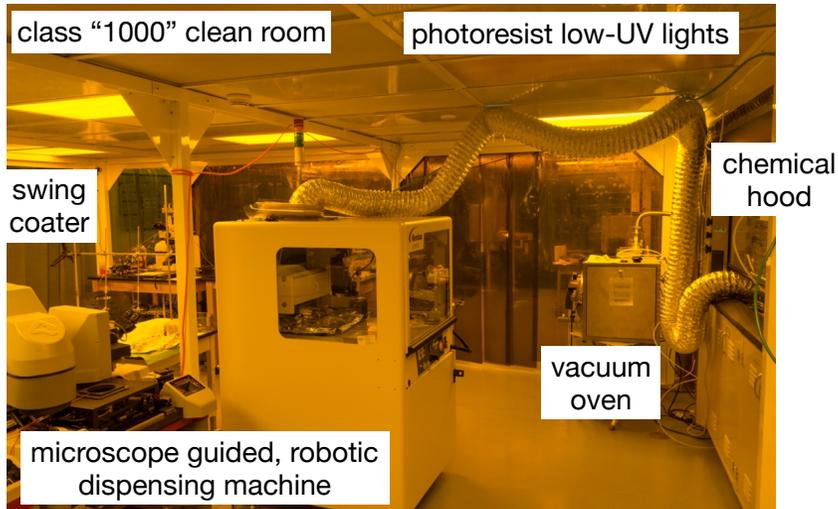


# Measurement cells

- Dimensions optimized for statistics and systematics. (e.g. E-field, diffusion-limit  $L^{-4}$  dependence of  $T_2$  time &  $L^2$  dependence of false EDM, magnetic field gradients, etc.).
- Cells separate from electrodes (cooling needed for heat flush)

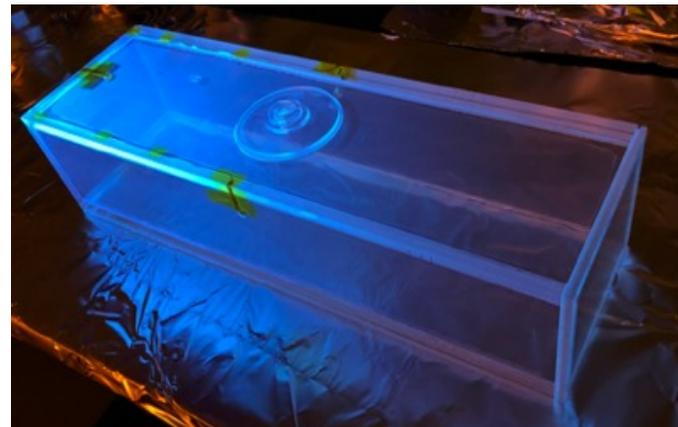
<b>Inner coating</b>	~ 1 $\mu\text{m}$ thick d-TPB embedded in d-PS matrix	Non-magnetic and not electrically conductive. Deuteration for high Fermi potential (165 neV from neutron reflectometry). TPB for detection of 80 nm $^4\text{He}$ scintillation light (see Ito, Cianciolo, Loomis talks)
<b>Side walls</b>	~ 1 cm thick p-PMMA plates	Optical photon transmission, mechanical strength, and purity. "Swing coating" on flat plates produce surface finish ~ 5 nm RMS roughness (AFM measured)
<b>End windows</b>	~ 5 mm thick d-PMMA	reduce activation caused by scattering of 9 $\text{\AA}$ beam
<b>Sealed cell</b>	deuterated solvent cemented	to separate $10^{-10}$ polarized $^3\text{He}$ inside cell from unpolarized natural-abundance $^3\text{He}$
<b>Cell hole</b>	~ 1 cm opening	initial $^4\text{He}$ filling then loading/unloading $^3\text{He}$ via heat flush. Low UCN loss.

Cell production facilities in clean room



Full-sized test measurement cell (illuminated by 360 nm UV lamp).

Temperature sensors embedded from outside for tests.



\*p- & d- refers to "protonated" or fully-deuterated versions of a material.

PMMA = poly(methyl methacrylate), more commonly-known by tradenames acrylic or plexiglass.

PS = polystyrene.

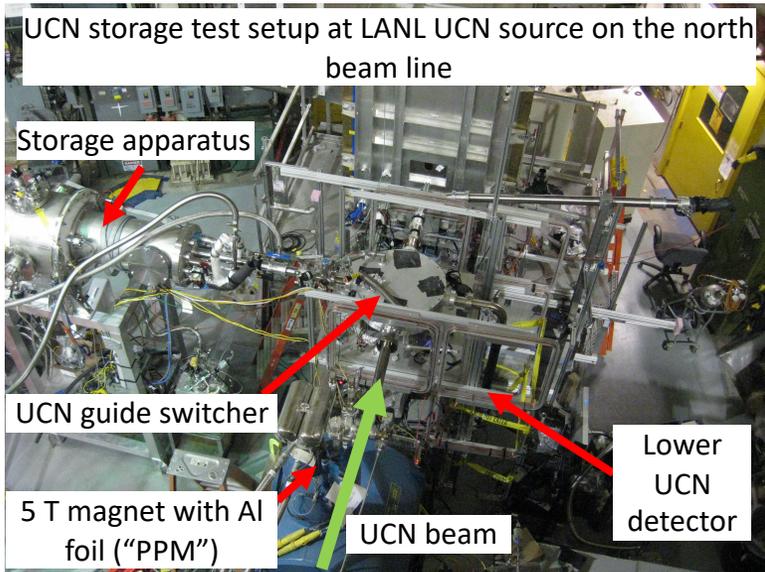
TPB = tetraphenyl butadiene, an electroluminescent dye

# Cryogenic UCN storage tests with external UCN source



Thanks to collaborators: Cianciolo, Clayton, Cooper, Curie, Golub, Griffith, Huffman, Ito, Korobkina, Makela, McDonald, O'Shaughnessy, Pentilla, Ramsey, Smith, Stanislaus, Tang, Weaver, Wei

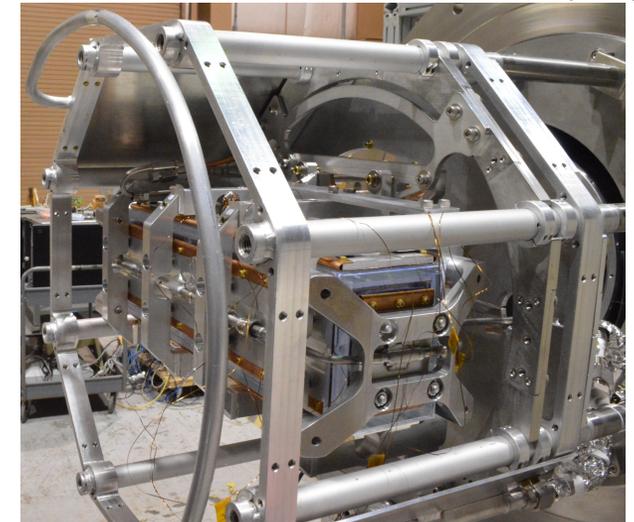
- A **cryogenic UCN storage apparatus** used with the **Los Alamos National Lab. solid-D<sub>2</sub> UCN source**.
- **Standard fill-and-empty setup** with cell valve & UCN guide switcher to connect (1) UCN source to cell or (2) cell to lower UCN detector
- **Cells cooled by spring-loaded cooling plates coupled to flowing L<sup>4</sup>He line**. Cells reach nominal 30 K in ~ 3 days (but can reach 15 K with high L<sup>4</sup>He rate).
- Dry & clean vacuum system: in vacuum outside cell,  $P \sim 5 \times 10^{-7}$  mbar before starting to cool cell. Outgassing dominated by plastic.
- Pumping inside cell restricted due to 1-cm cell hole & UCN guides. Put UCN guide switcher in “intermediate” position for increased cell pumping speed.
- **Cells are heated to 50 degC** (coating limit) and **pumped for 10-14 days** before cooling.
- The input L4He flow is split, **a region of stainless steel UCN guide ~20 cm away from the cell is maintained to be coldest point in the system** (nominally > 20 K colder than cell). This SS surface **acts as a cryopump** for pumping out cell and **traps condensable-contamination** from reaching inside cell.



Slide-openable vacuum vessel with flowing LN<sub>2</sub> radiation shield (80 K & ~ 3π sr coverage)

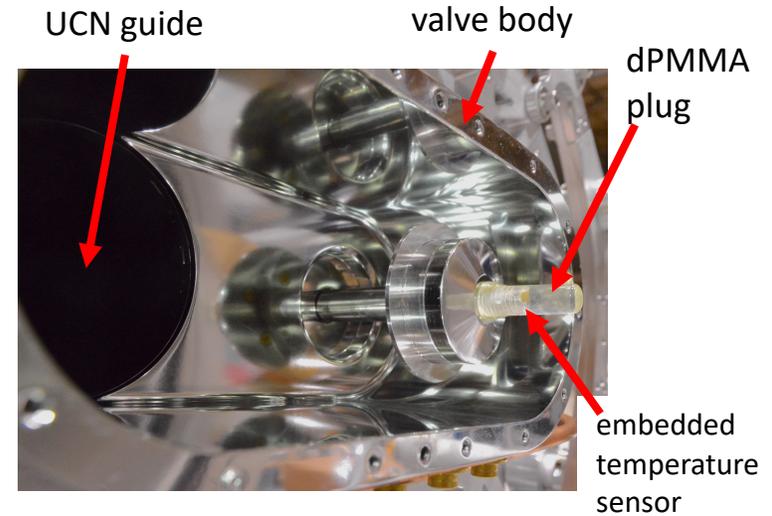


Spring loaded cooling plates used to maintain contact with cell  
Return L4He cools another radiation shield (50 K)

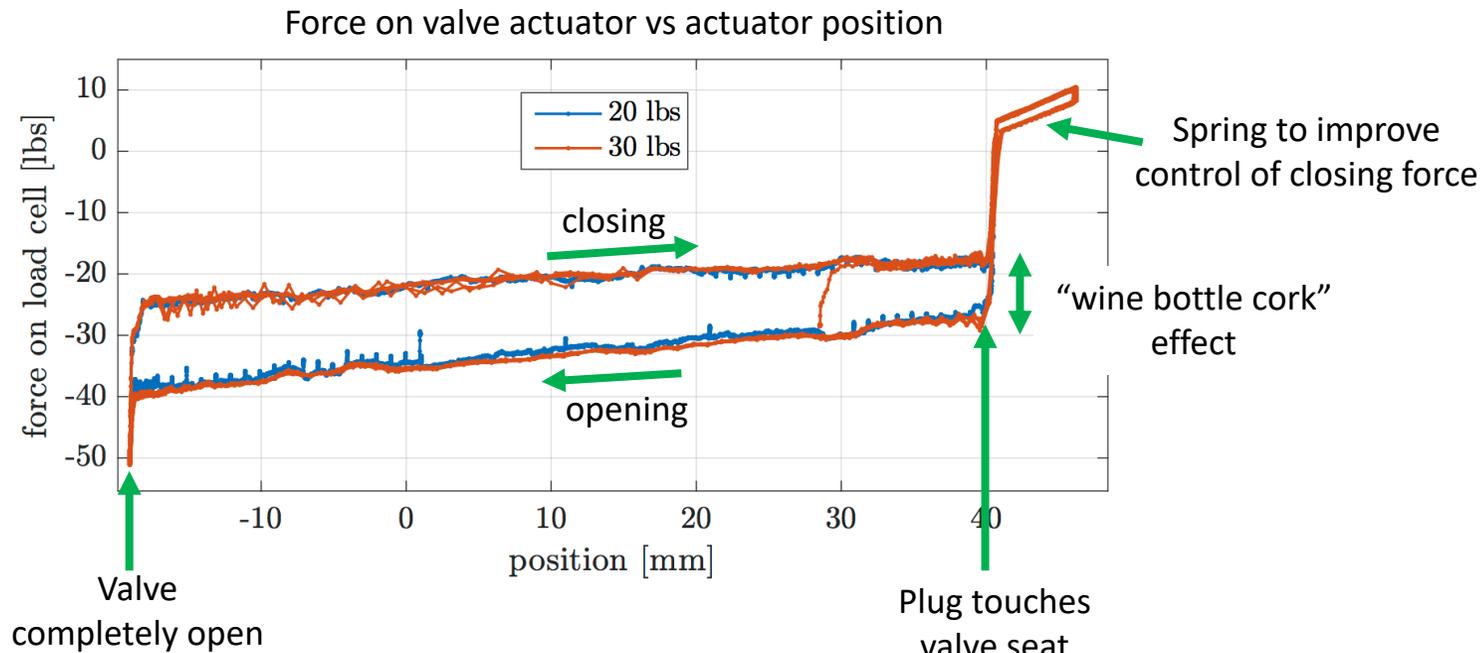
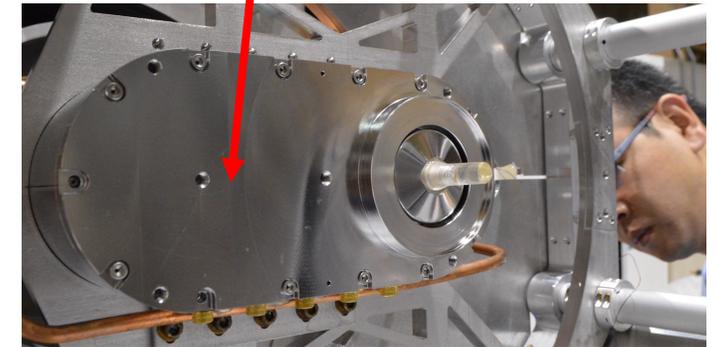


# Deuterated plastic cell valve

- The cell-valve body is stainless steel 180 neV (eventually it will be plastic)
- Valve **plug is bulk-dPMMA** & **valve seat is bulk-dPMMA ring** solvent-welded to cell.
- Linear motion controlled by servo-motor actuator with optical encoder (< 0.1 mm precision). Force is measured with a load cell. **Precise control of position, force, and speed of the valve.**



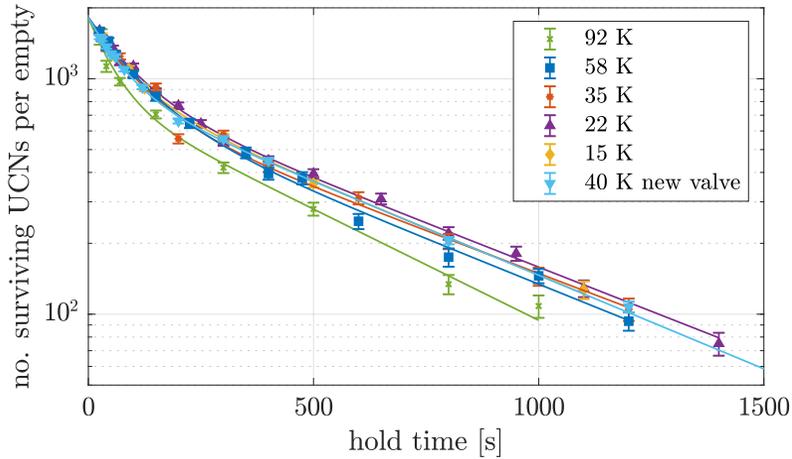
SS plate closing off valve body



# UCN storage results

(Data normalized so number of UCNs at t=0 extrapolated with double exponential fit are matched. The value chosen at t=0 is average of the different storage curves with same cell, measured over consecutive days)

2<sup>nd</sup> generation cell #2



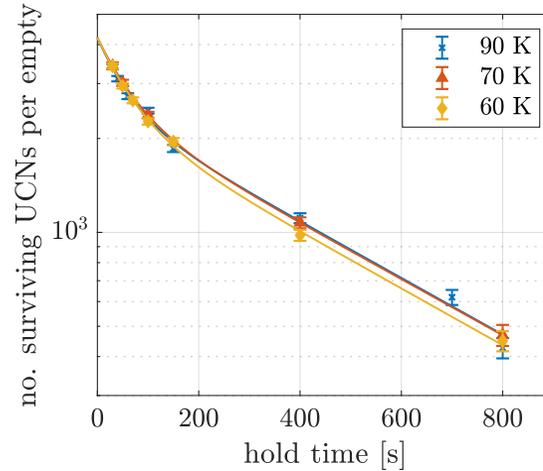
- Beam position: behind UCN $\tau$
- Measured with old and new valve systems

### Fitting 40 K data:

$$N_{\text{short}} = (50 \pm 2) \% , \tau_{\text{short}} = (73 \pm 7) \text{ s}$$

$$N_{\text{long}} = (50 \pm 2) \% , \tau_{\text{long}} = (546 \pm 16) \text{ s}$$

2<sup>nd</sup> generation cell #3



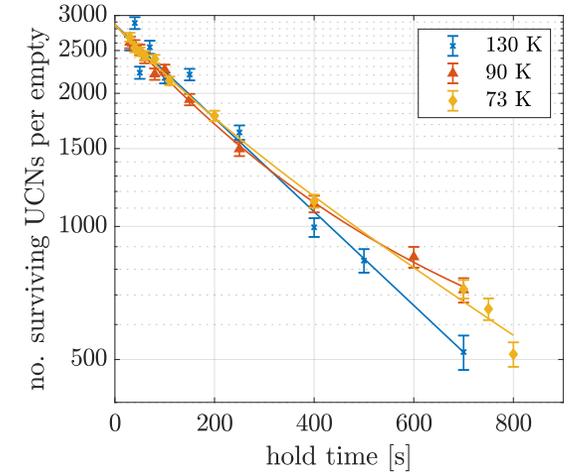
- Beam position: north beam line
- dPMMA valve system

### Fitting 60 K data:

$$N_{\text{short}} = (45 \pm 4) \% , \tau_{\text{short}} = (66 \pm 16) \text{ s}$$

$$N_{\text{long}} = (55 \pm 4) \% , \tau_{\text{long}} = (477 \pm 47) \text{ s}$$

2<sup>nd</sup> generation cell #4



- Beam position: north beam line
- dPMMA valve system
- Reduced dust contamination during cell production

### Fitting 73 K data:

$$N_{\text{short}} = (24 \pm 28) \% , \tau_{\text{short}} = (170 \pm 190) \text{ s}$$

$$N_{\text{long}} = (76 \pm 28) \% , \tau_{\text{long}} = (590 \pm 190) \text{ s}$$

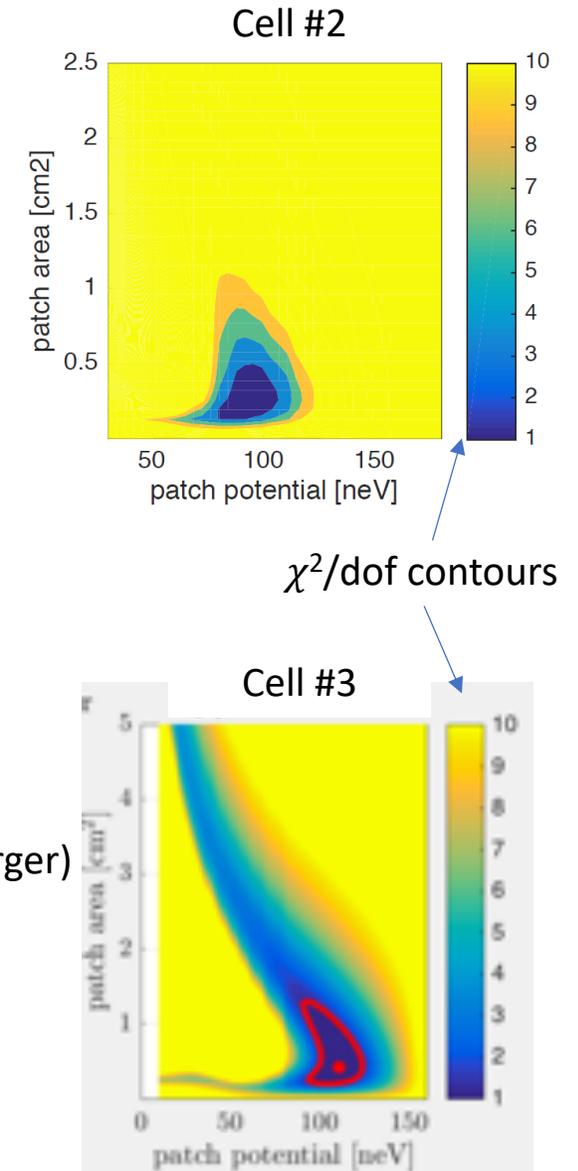
OR

$$\tau_{\text{single-exp}} = (490 \pm 11) \text{ s}$$

- Either sum of two exponential decays ("long" and "short") or single-exponential fit.
- When sum of two exp decays, strong correlations between fitted parameters
- When cell not cooled,  $\tau_{\text{single-exp}} = (130 \pm 5) \text{ s}$
- Storage times here are  $\beta$ -decay + cell wall losses

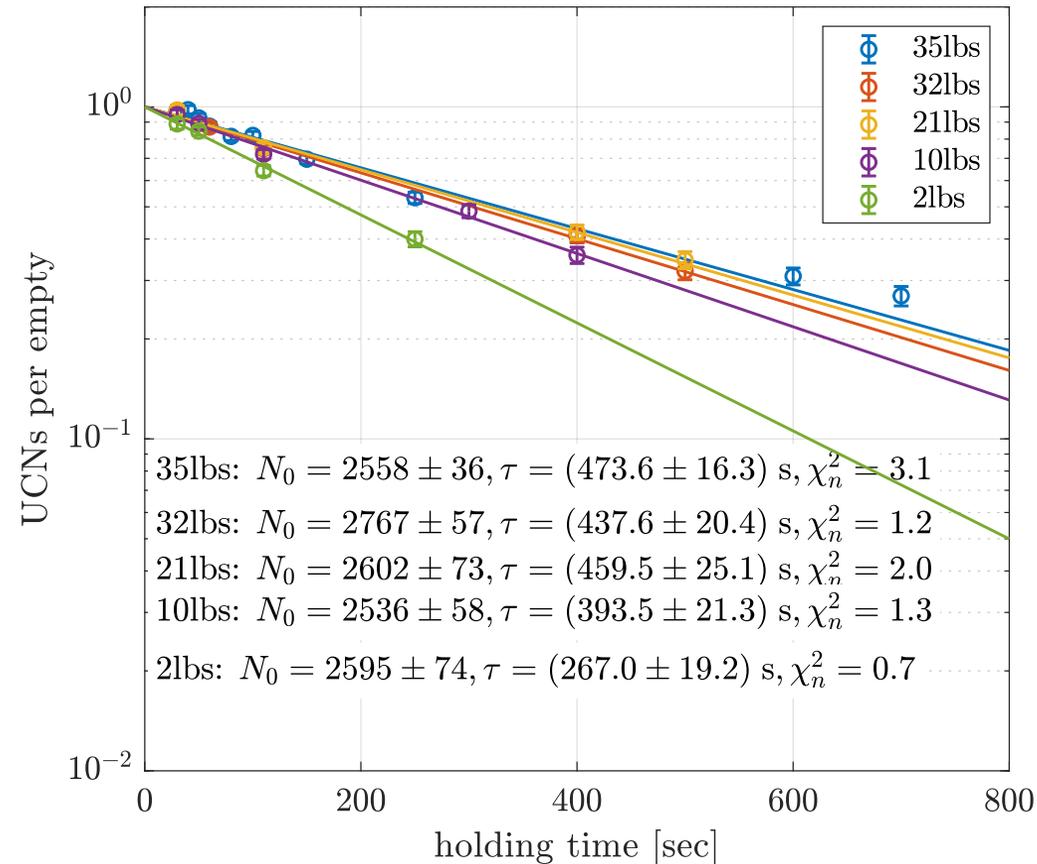
# Interpretation of results

- High quality set of cryogenic UCN storage data with a new type of deuterated polymer coated surface ( $V_F = 165$  neV, fluorescent,  $^3\text{He}$  polarization friendly) exhibiting excellent UCN storage properties. Cell #4 with  $\tau_{\text{single-exp}} = (490 \pm 11)$  s => **loss per reflection  $\sim 2 \times 10^{-5}$**
- **Standard UCN-energy-dependent wall loss theory** (i.e.  $f = W/V_F$  description) **insufficient** for describing data exhibiting **strong double-exponential decay**.
- From crude information on UCN spectrum, we deduce it is the higher-energy UCNs in the  $\tau_{\text{short}}$  component. Commonly explained away by above-trapping threshold UCNs but the time scales here are very long for our rectangular cells.
- Most promising model found so far has been fitting with a small area “patch” ( $\sim 0.3$  cm<sup>2</sup>) with low  $V_F$  ( $\sim 100$  neV) exposed to UCNs
- Patch parameters very similar between cells #2 & #3
- Possible theories for cell #4:
  - (1)  $\tau_{\text{short}}$  component very short, so not visible since shortest hold time is 25 sec (patch area > 3x larger)
  - (2)  $\tau_{\text{short}}$  component very long or non-existent (difficult to fit for. Would need patch 5x smaller)
  - (3) higher UCN energies not loaded into the cell in the first place.
- Currently developing UCN tracking simulations to increase understanding. Possibly use UCN spectrometer in future experiments.



# dPMMA UCN cell valve results

- These results show we have **learned to implement a UCN valve with negligible UCN loss**
- The leak tightness of the valve is more stringent for containing  $^3\text{He}$  inside the cell in final experiment ( $v \sim 30$  m/s)



# 9 Å neutron beam induced cell activation in final experiment

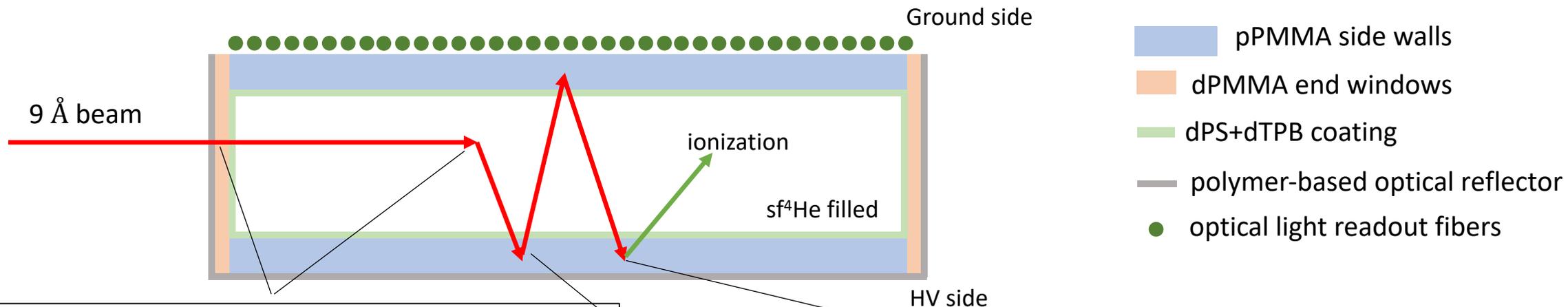
- In-situ UCN production via super-thermal production in  $\text{sf}^4\text{He}$  by 9 Å beam (via choppers) offers many advantages:

High UCN density  $\sim 150 \text{ UCN/cm}^3$   
after 1000 s of “filling time”

No UCN or depolarization loss from transport

Spatial spread of beam and filling time reduces phase-space evolution at start of precession measurements

- The beam intensity is  $5 \times 10^5 \text{ s}^{-1}$  in each cell (UCN production  $P = 0.31 \text{ UCN/cm}^3/\text{s}$ ), so  $T_{\text{fill}} = 1000 \text{ s} \rightarrow 5 \times 10^5$  neutrons incident



Scattering in  $\text{sf}^4\text{He}$  (coherent inelastic scattering).

Mean-free-path = 20 m, 40 cm cell, 2%.

OR

Scattering in dPMMA windows (mix coherent elastic scattering & incoherent elastic scattering).  $\sim 1\%$  depending on window thickness

OR

Direct activation of impurities in dPMMA with decay chains

Cold neutrons bounce around via incoherent elastic off  $^1\text{H}$

Eventually  $n+p \rightarrow D + \gamma$  (2.2 MeV) + other prompt-reactions

=> can produce free charge in cell

OR

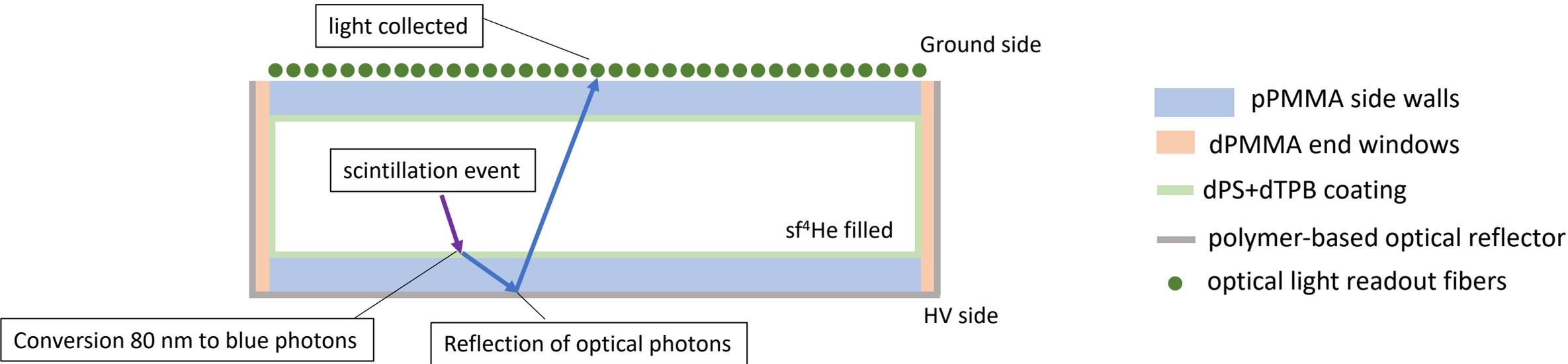
$n + \text{contaminant} \rightarrow$  short or long lived decay chain => delayed background scintillation events

(see Korsch & Loomis talks)

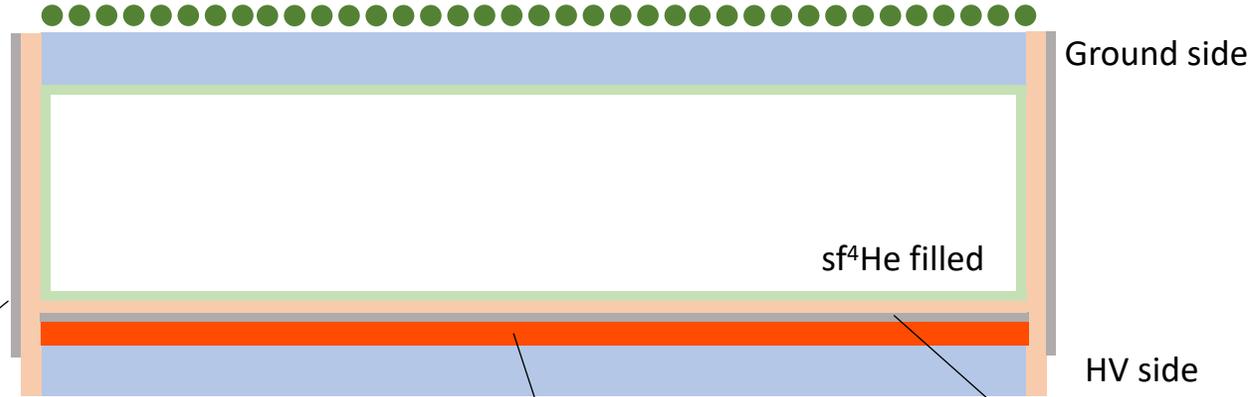
**Note:** we developed  $\text{L}^4\text{He}$  MCNP neutron scattering kernels to 4 K. (Thanks to C. Lavelle from Johns Hopkins!)

# How scintillation light is collected in measurement cell

(see Cianciolo talk)



# Solution to problem



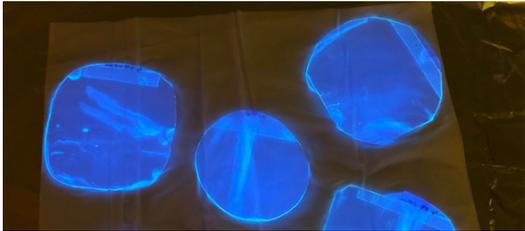
- pPMMA side walls
- dPMMA end windows
- dPS+dTPB coating
- polymer-based optical reflector
- optical light readout fibers

Make dPMMA end windows (~ 0.5 - 1 mm) & very pure

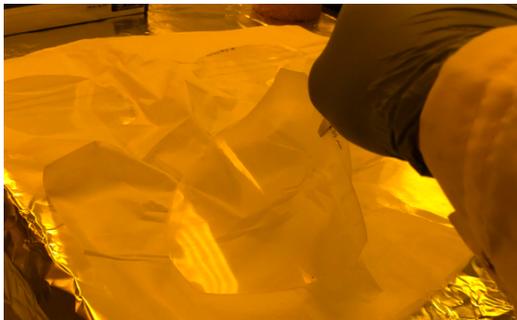
Embed  $^6\text{Li}$  cold neutron absorbing layer (1-2mm thick) close to inside of cell (but UCNs can't reach)

Place optical reflector between dPS+dTPB coating and (non-optically transparent) n-absorbing layer

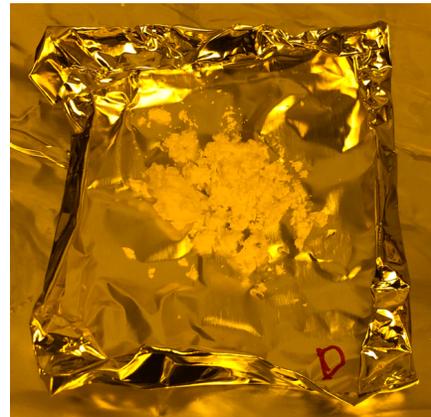
*Coated thin end PMMA windows*



*LN2 dunk and strength tests*



*In-house synthesis of high-purity dPMMA*



*1/3<sup>rd</sup> scale prototype side wall with embedded LiOH absorber and optical reflector*



# High UCN-current, “inverse spallation target geometry”, sf-<sup>4</sup>He source (small digression)

- Take **1 MW proton beam @ 800 MeV energy** (available SNS/LANL/Fermilab)
- 40 L sf-<sup>4</sup>He bath with **100 W @ 1.6 K using sub-cooled helium technology** (e.g. common at LHC, JLab, SNS, Fermilab, ESS)
- At 1.6 K, up-scattering  $\tau_{up} \approx 3$  s, **optimize for a maximum current UCN source (pessimistically assume single passage transport only)**
- **Raster proton beam** on “ring” shaped tungsten spallation target allows **edge water cooling**
- Total UCN production rate in sf-<sup>4</sup>He is **1.8 x 10<sup>9</sup> UCN/s**.
- UCN delivery, **5 m away in 18 cm diameter guide**, from UCN simulations is **5 x 10<sup>8</sup> UCN/s** (1-passage efficiency 28%. Foil ~ 68% transmission.)

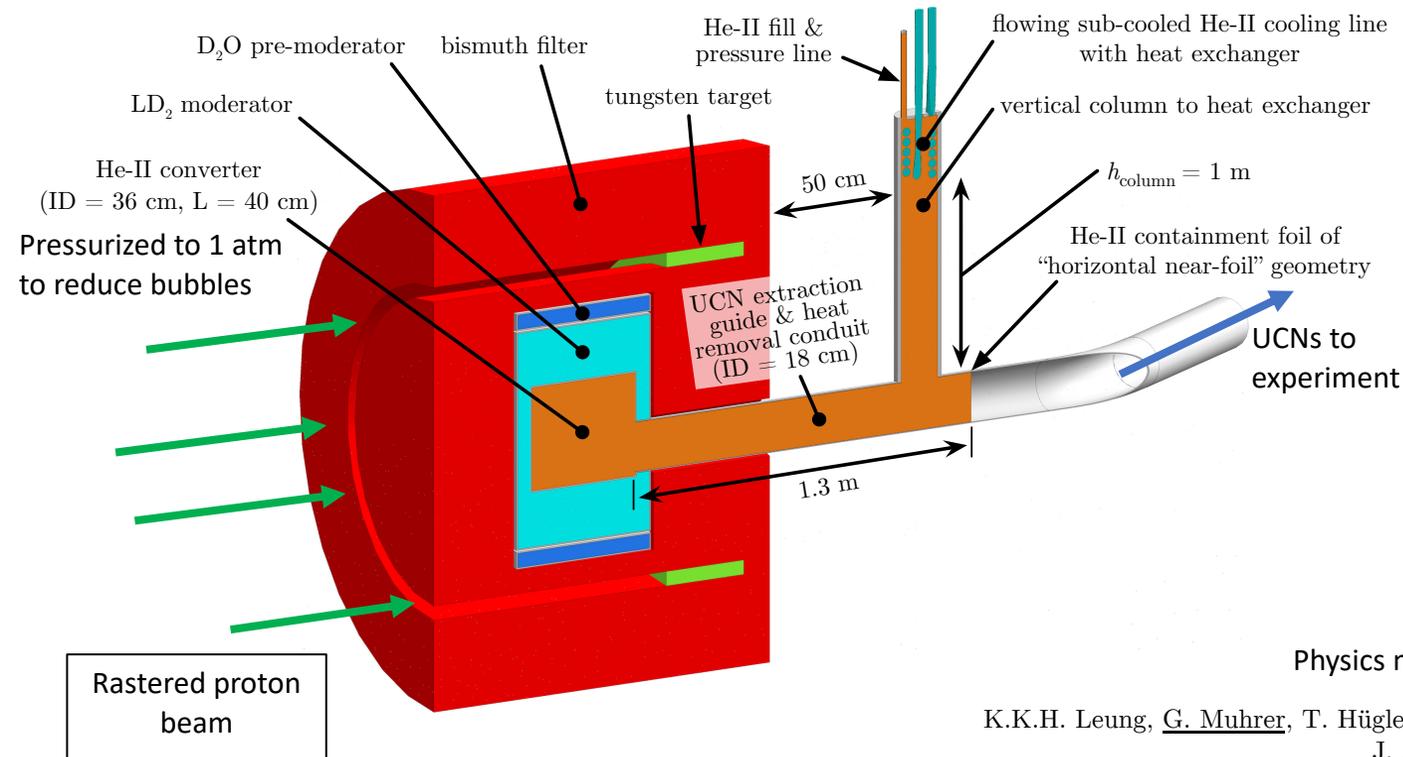
\*<sup>58</sup>Ni cut-off potential & unpolarized

## Delivered equilibrium UCN density to external bottle with 18 cm hole:

$V_{\text{bottle}}$ (l)	5	50	500	$5 \times 10^3$	$5 \times 10^4$
$\rho_{\text{bottle}}$ ( $\times 10^4$ UCN $\text{cm}^{-3}$ )	1.12	1.11	1.05	0.80	0.31
$\tau_{\text{bottle}}$ (s) (Filling time)	0.11	1.1	10	80	315

(pessimistically assume if UCN leaves volume then it is lost)

- Density doesn't drop until extremely large volumes
- If filling restricted by a small hole (e.g. 1 cm  $\varnothing \Rightarrow$  200 s fill time) then need **1 MW beam needed at ~ 10% duty**



Physics model (w/ minor engineering considerations) published in:

K.K.H. Leung, G. Muhrer, T. Hügle, T.M. Ito, E.M. Lutz, M. Makela, C.L. Morris, R.W. Pattie, A. Saunders, A.R. Young  
 J. Appl. Phys. 126, 224901 (2019). doi:10.1063/1.5109879

a new type of UCN precession signal:  
***in-situ* & live spin analysis**

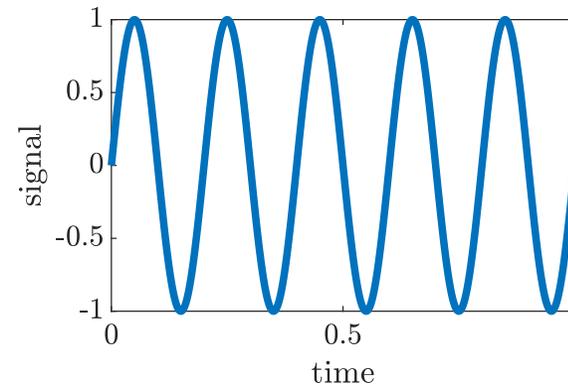
# Difference between measuring phase vs frequency

- **Ramsey technique's observable is final phase (via UCN count, which is an amplitude):**  $\phi_{nf} \equiv \phi_n(T_f) = \phi_{n0} + \int_{t'=0}^{T_f} \omega_n(t') dt'$

Can only deduce *average* frequency:  $\bar{\omega}_n = (\phi_{nf} - \phi_{n0})/T_f$

Can only tell you frequency with time if frequency is constant (i.e. B-field & other systematics are constant within precession time)

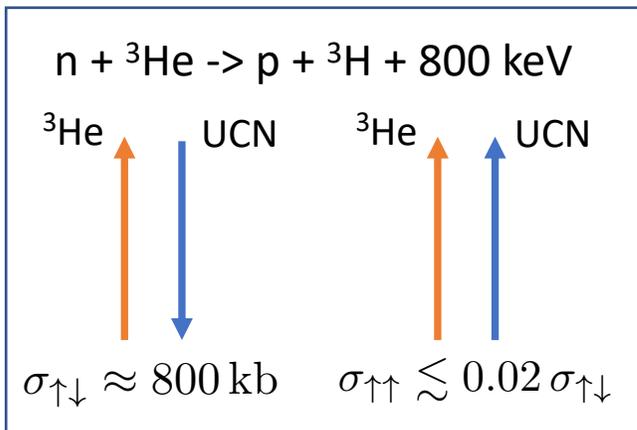
- **“Always measure frequency...”** (Rabi? Ramsey? Wieman?) Why?



Changes of signal over short times suppresses drifts in system.

Don't need absolute signal size.

- **Live and in-situ UCN spin analysis with polarized  $^3\text{He}$ :** measures throughout precession phase difference between  $^3\text{He}$  analyzer & UCNs:



$$\dot{N}_3(t) = \frac{N_n(t)}{\bar{\tau}_3} [1 - P_n(t) P_3(t) \cos \phi_{3n}(t)]$$

500 s is our optimum  $\sim 10^{-10}$   $^3\text{He}$  concentration

$$\phi_{3n}(t) \equiv \phi_3(t) - \phi_n(t)$$

$\phi_3(t)$  measured with SQUIDS  
( $^3\text{He}$  is a co-magnetometer also.)

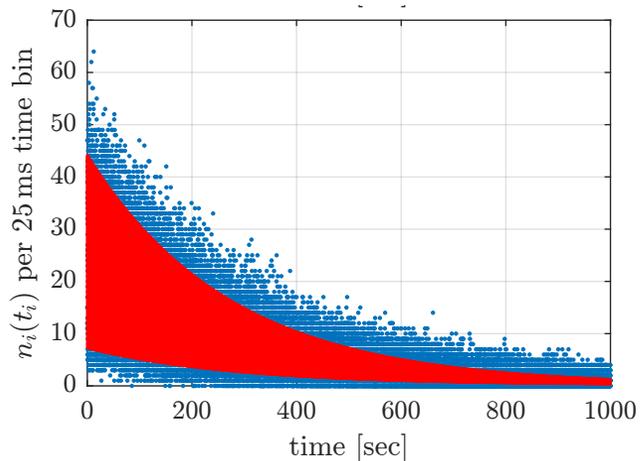
The spirit of frequency measurement lives!  
(via derivative)

# Free-precession mode scintillation light

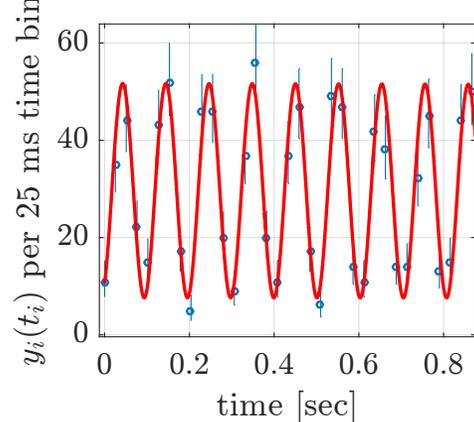
- After  $\pi/2$  pulse, UCN &  $^3\text{He}$  free precession:  $\phi_{3n}(t) = \left[ (|\gamma_3| - |\gamma_n|)B_0 \pm \frac{2d_n|E|}{\hbar} \right] t + \phi_{3n0}$  (ignoring shifts & drifts for now)  
 $= \omega_{3n}^\pm t + \phi_{3n0}$  ( $\omega_3/(2\pi) \approx 100 \text{ Hz}$  &  $\omega_{3n}/(2\pi) \approx 10 \text{ Hz}$  @ 30 mG)

- “accepted” rate of scintillation light events:  $\dot{N}_{ac-s}(t) = N_n(t) \left( \frac{\epsilon_\beta}{\tau_\beta} + \frac{\epsilon_3}{\bar{\tau}_3} \left\{ 1 - P_3(t)P_n(t) \cos[\theta_{3n}(t) + \phi_{3n0}] \right\} \right) + R_{BG}$ 
    - acceptance probabilities 0.33 & 0.93
    - no. UCNs in cell
    - recall: 500 s
    - polarizations
    - Background rate (could be time-dependent)
    - $P_3(t)P_n(t) = P_{30}P_{n0} \exp(-t/T_{2,tot}) \approx (0.98)(0.98) \exp(-t/[10,000 \text{ s}])$
- UCN spectrum
- $$N_n(t) = \int_0^{E_{\max}} dE n_{n0}(E) \exp \left[ -\frac{t}{\tau_{\text{cell}}(E)} - \frac{t}{\bar{\tau}_\beta} + \frac{P_n(t)P_3(t)}{\bar{\tau}_3} \int_0^t \cos \phi_{3n}(t') dt' \right]$$
- Oscillating term due to previous n- $^3\text{He}$  absorption

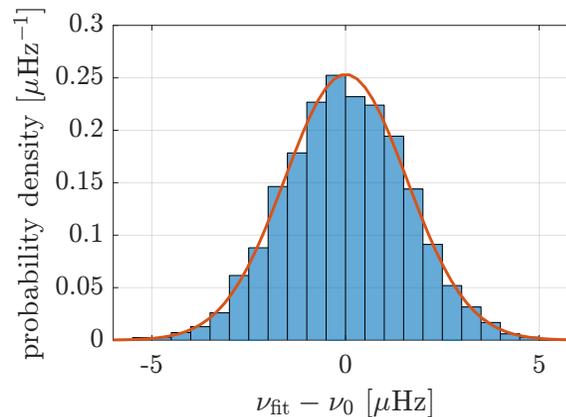
Simulated signal 1-cycle (accepted events)



Zoom in: red line = fitted



Repeat generation & fit:



Generated “full statistics” data set with constant  $B_0$

From distribution:  $1\sigma$  error bar =  $1.8 \mu\text{Hz}$   
 ➤ Free precession mode 300 live days data taking (e.g. 2-3 years)  $1\sigma$  nEDM error =  $3 \times 10^{-28} \text{ e.cm}$

Dressed-spin mode gives  $1.6 \times 10^{-28} \text{ e.cm}$

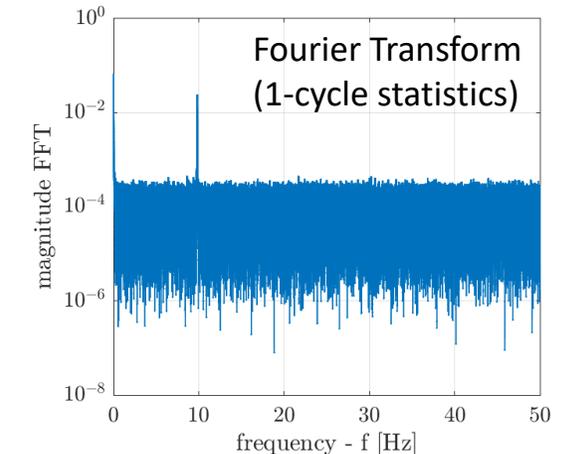
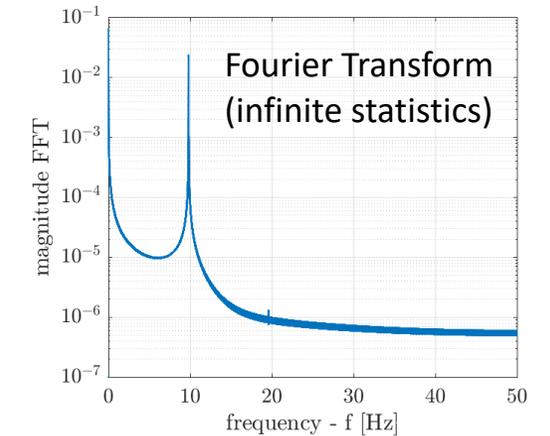
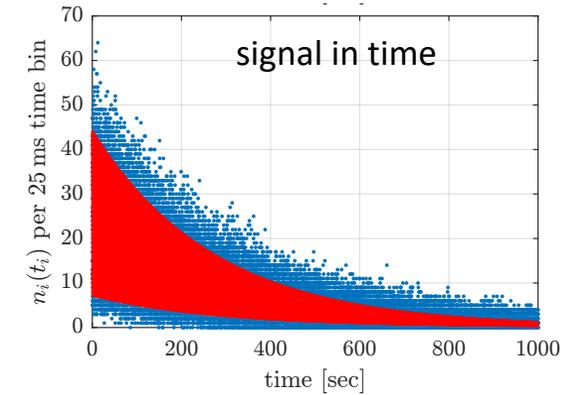
# Data analysis simulations

We (re-)launched a team to study statistics and systematic effects in our data analysis.  
Effects studied or under-study:

- Neutron decay  **$\beta$ -asymmetry** (Plaster)
- **Spatial-variation scintillation light** detection efficiency (see talks Cianciolo, [Loomis](#), [Leggett](#))
- **Oscillation in  $N_n(t)$**  due to history of n- $^3\text{He}$  absorption (moi)
- Difficult to analytically-fit **UCN energy-dependent wall loss** (moi)
- Generation of scintillation light data with **magnetic field drifts** (moi & Cianciolo)
- Generation of SQUID  $^3\text{He}$  **signal with noise and drifts** ([Behzadipour](#), Cianciolo, Plaster, Clayton)
- **Magnetic field noise** in spin-dressing mode (see talks Swank and [Tat](#))
- UCN spin-tracking on **Graphics Processing Units** (Broussard, [Kline](#), Matthews)
- UCN center-of-mass **gravitational offset time-evolution** (moi)

Data analysis techniques being developed:

- Maximum likelihood fitting time-binned counts and **un-binned counts**
- Markov Chain Monte Carlo random walk fitting techniques
- **Fourier Transform** analyses, digital filtering and **matched-filter** techniques



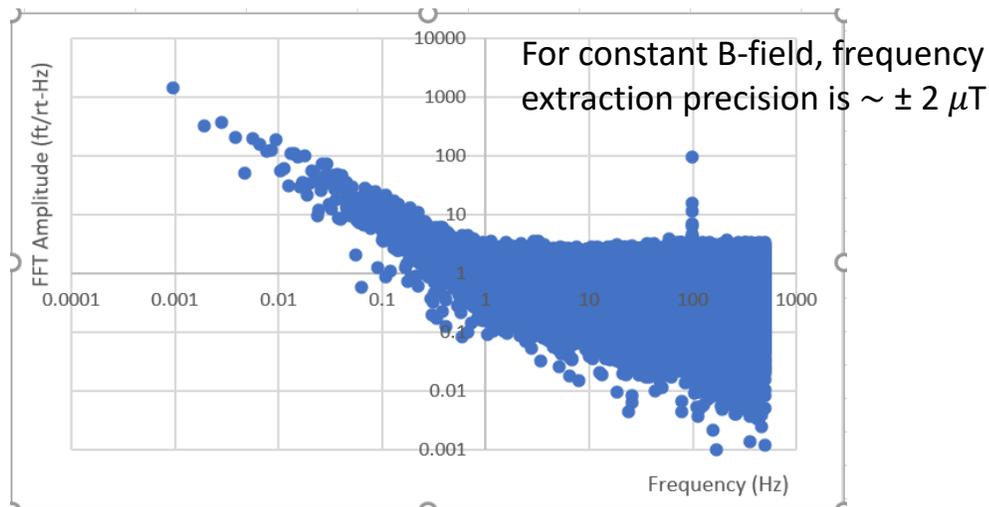
# Handle field drifts in free-precession mode

- Highlight work led by Mojtaba Behzadipour (U. Kentucky grad student) **simulation of SQUID  $^3\text{He}$  signal with white+pink electronic noise** (based on data) and **B-field slow drifts and jumps** (based on magnetic shielding specs)
- Advantage of live UCN spin analysis is do not necessarily need to throw away data if field jump
- With live & in-situ UCN spin analysis the **time-dependent  $B(t)$  from the (co-)magnetometers needed**. (In Ramsey technique need to think a bit about this too if there's a statistics versus time variations in B-field extraction from magnetometers due to relaxation)
- General procedure under development: approximate  $B(t)$  with 2<sup>nd</sup> order polynomial:

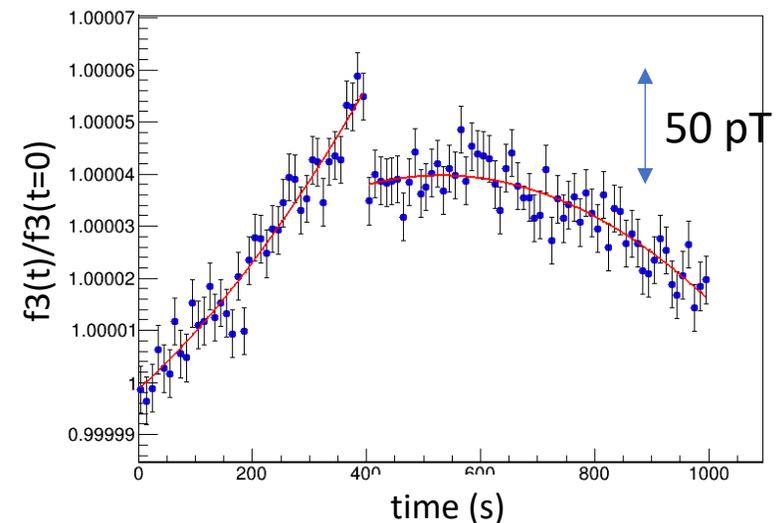
$$B(t) = B(0)(1 + \epsilon_1 t + \epsilon_2 t^2) \longrightarrow \text{SQUID signal from superconducting pick-up coils} \quad S(t) = \underline{A_0} \exp(-t/T_{2,3\text{He}}) \cos \left[ |\gamma_3| \underline{B(0)} \left( t + \frac{1}{2} \underline{\epsilon_1} t^2 + \frac{1}{3} \underline{\epsilon_2} t^3 \right) + \phi_{30} \right]$$

- (1) fit  $B(0)$ ,  $\epsilon_1$  and  $\epsilon_2$  from  $^3\text{He}$  comagnetometer and then used these values fixed in scintillation light fitting, OR
- (2) "Global" fit of  $^3\text{He}$  and scintillation light signal

Simulated FFT signal from SQUID for 1000 s, with SQUID noise based on data from LANL (1 fT/rt-Hz, 1 Hz cross-over frequency)



Split signal into 10 s windows and extract frequency in each



# Advantages in alternative physics extraction with live & in-situ UCN phase information

# Searching for time-oscillating nEDM signals

- Axion-like particles can induce an oscillating nEDM signal [Abel et al., Phys. Rev. X 7, 041034 (2017)]:

$$\omega_n(t) = |\gamma_n|B_0 \pm \frac{2d_n|E|}{\hbar} + \frac{2|E|\alpha_{ax}}{\hbar} \cos(\omega_{ax}t + \phi_{ax})$$

Amplitude of oscillation (units: e.cm)  
phase (free parameter in fit)

- Sensitivity scales with nEDM sensitivity. nEDM@SNS every measurement cycle (2400 s) get a  $1\sigma$  precision of  $\pm 4 \times 10^{-26}$  e.cm. Expected 1-2 orders-of-magnitude improvement in “standard” base-line search technique.

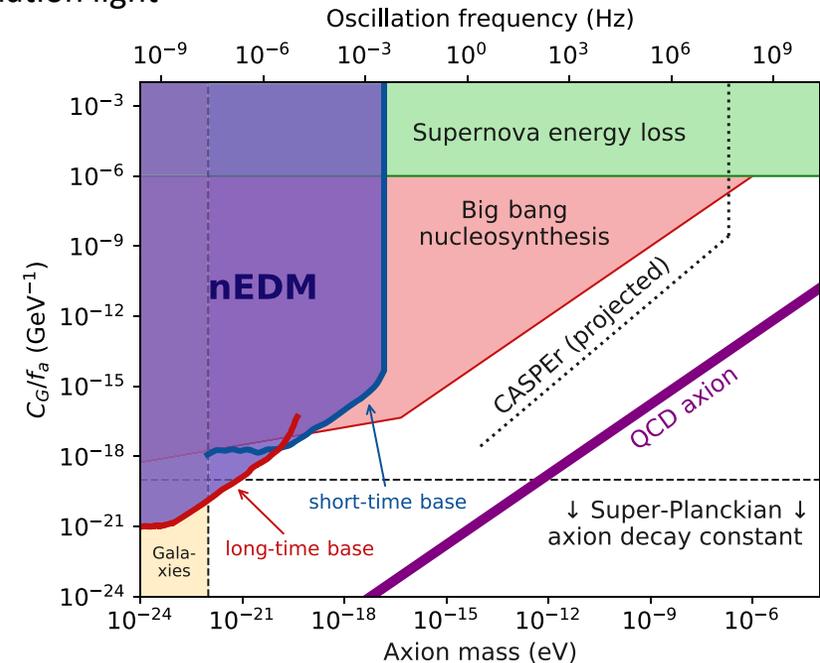
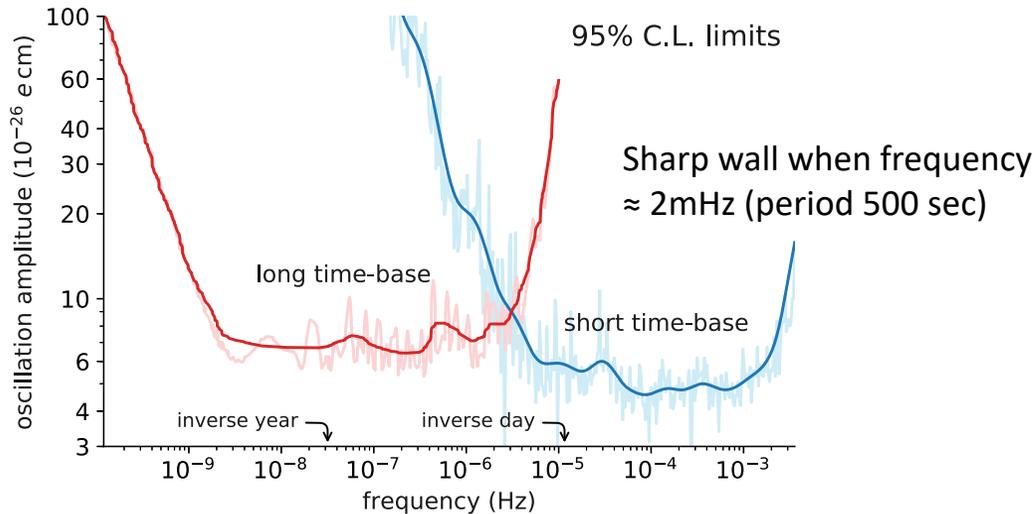
## Ultra-short time-base search:

- $\phi_n(t)$  information allows searches for oscillating nEDM signal at frequencies higher than the measurement cycling:

$$\phi_{3n}(t) = \left[ (|\gamma_3| - |\gamma_n|)B_0 \pm \frac{2d_n|E|}{\hbar} \right] t + \frac{2|E|\alpha_{ax}}{\hbar\omega_{ax}} [\sin(\omega_{ax}t + \phi_{ax}) - \sin\phi_{ax}] + \phi_{3n0}$$

↑ determines n-3He scintillation light

From Abel et al.:



# High-precision measurement of the ratio $\gamma_n/\gamma_{3\text{He}}$

- Only minor improvements in knowledge of the neutron magnetic moment over the past 5 decades

PHYSICAL REVIEW D      VOLUME 20, NUMBER 9      1 NOVEMBER 1979

**Measurement of the neutron magnetic moment**

G. L. Greene\* and N. F. Ramsey      **+ et al.**  
*Harvard University, Cambridge, Massachusetts 02138*

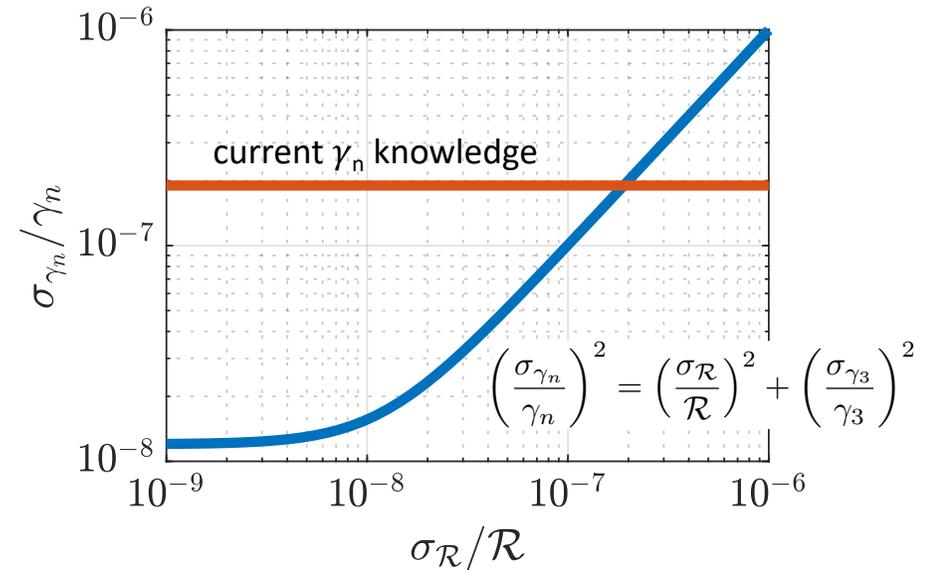
$$\frac{\sigma_{\gamma_n}}{\gamma_n} = 0.24 \text{ ppm via } \frac{\gamma_n}{\gamma_p}$$

Afach et al., Phys. Lett. B 739 (2014)  $\rightarrow \frac{\sigma_{\gamma_n}}{\gamma_n} = 0.19 \text{ ppm via } \frac{\gamma_n}{\gamma_{\text{Hg}}}$

- In nEDM@SNS, we measure:  $\mathcal{R} \equiv \frac{\gamma_n}{\gamma_{3\text{He}}}$       CODATA 2018:  $\frac{\sigma_{\gamma_{3\text{He}}}}{\gamma_{3\text{He}}} = 1.2 \times 10^{-8}$

## STATISTICAL ERROR:

- In one measurement cycle:  $\frac{\sigma_{\mathcal{R}}}{\mathcal{R}} \approx \frac{\sigma_{\omega_{3n}}}{\omega_{3n}} \approx \frac{1.8 \text{ uHz}}{10 \text{ Hz}} = 2 \times 10^{-7} (0.2 \text{ ppm})$
- To reach  $\pm 0.01 \text{ ppm}$  ( $\sim 20 \times$  improvement in  $\gamma_n$ ) need 11 days of statistics

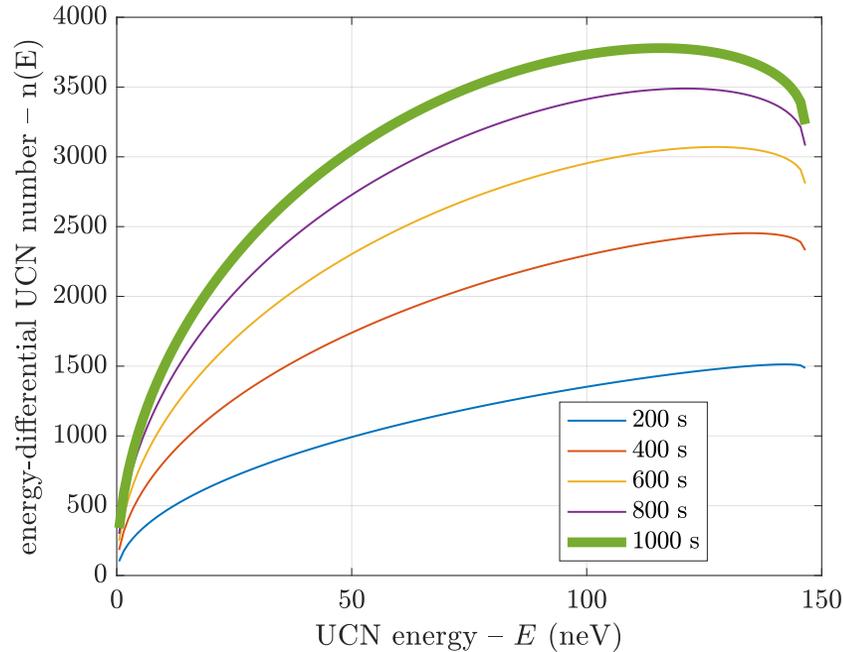


**KEY SYSTEMATIC:** UCN center-of-mass gravitational offset effect

# Evolution of UCN spectrum during filling and precession

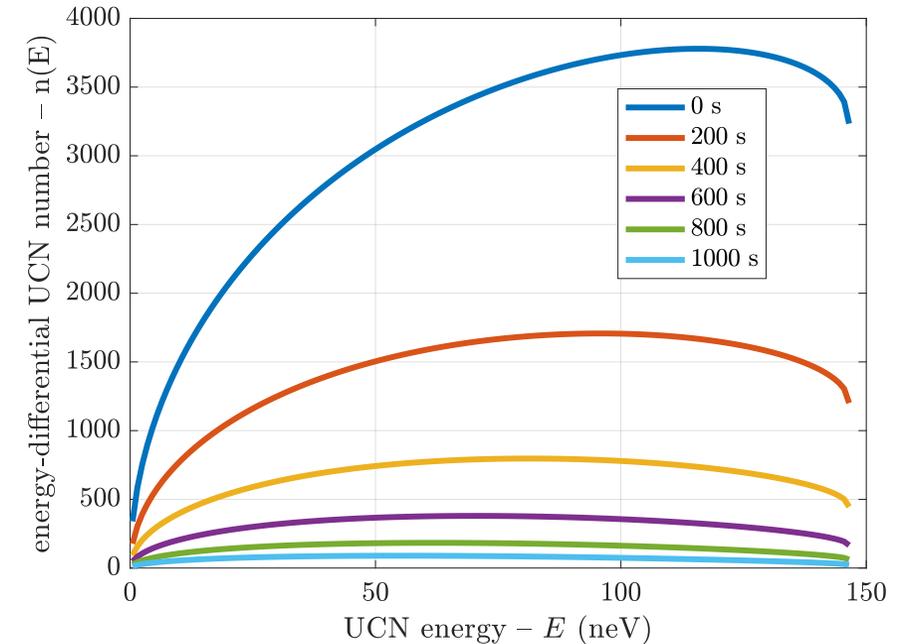
Example: used  $f = W/V = 0.8 \times 10^{-5}$

UCN spectrum during “filling” (low  $^3\text{He}$  absorption)



Statistically optimized  $T_{\text{fill}} = 1000$  s:

UCN spectrum evolution during free precession measurement time (includes  $\tau_3 = 500$  s)



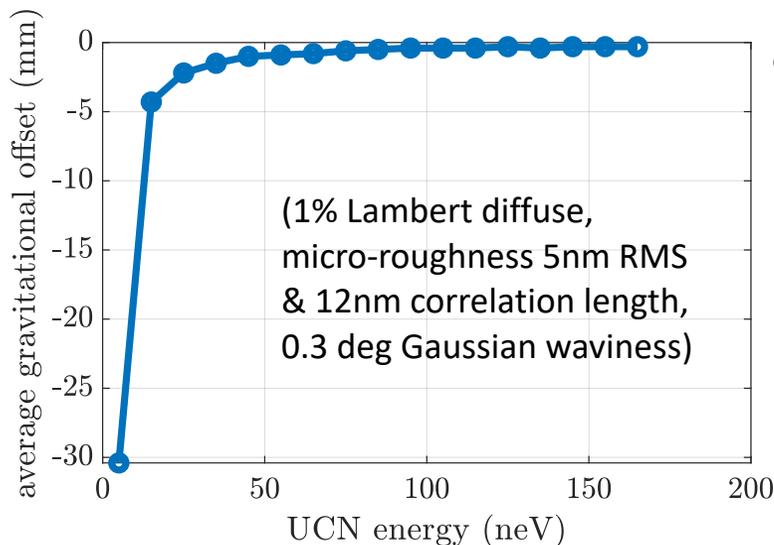
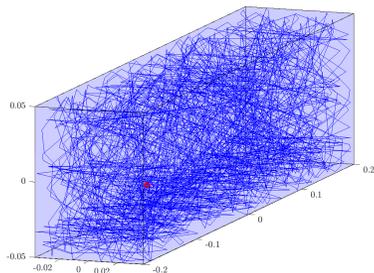
- Produced **UCN spectrum in  $\text{sf}^4\text{He}$  is well-described**. Transport in guides not so much.
- Above assumes mechanical equilibrium. **Phase-space evolution will be small in nEDM@SNS** (3 L cell, with UCNs produced with approximately isotropic momentum and close to uniformly throughout cell, filled over 1000 s). Next step is to confirm with simulations
- These are the UCN spectra inside the cell. **Since spin analysis is in-situ, no need to correct for UCN-energy dependent transport loss** (and depolarization loss) during transport to interpret any UCN spectral measurements
- Change filling time (with reduction on statistics) to change initial UCN spectrum slightly for systematics

# Use time-evolution of UCN gravitational offset to study UCN spectrum

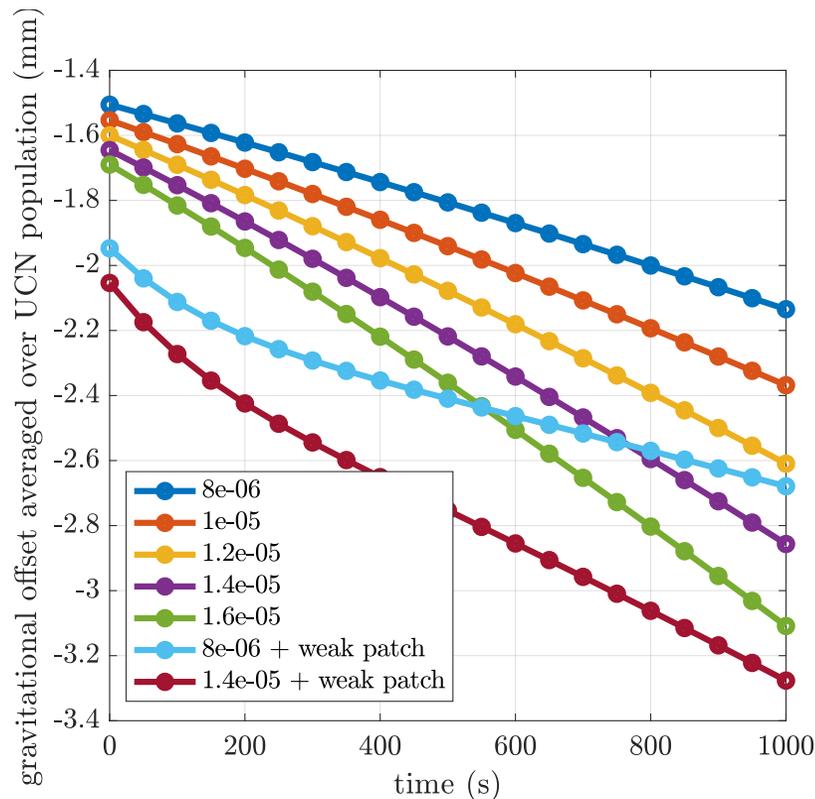
## UCN tracking simulations:

Shown E = 100 neV

Straight lines between wall collisions displayed.



Combine with UCN spectral evolution from previous slide



UCN center-of-mass offset when averaged over UCN population

$$f_n = \frac{\gamma_n}{2\pi} \left( B_0 + \frac{d|B|}{dz} h_{\text{off}}(t) \right)$$

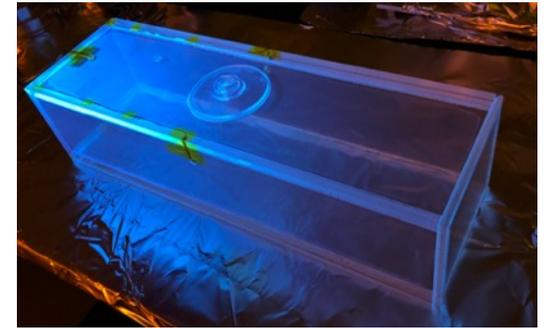
Vertical gradient

( $B_0 \rightarrow$  from  $^3\text{He}$ )

- Apply known gradient, measure  $f_n(t)$  to get information on the time-evolution of average UCN energy during precession.
- **Combine with total number UCN storage time measurement:** just watch  $\beta$ -decay without loading  $^3\text{He}$ .
- More direct than spin-echo since since measuring precession frequency
- Phase-space evolution (which will impact this systematic effect as well as others!) less in the nEDM@SNS experiment
- A good chance for controlling this important systematic effect! Currently in early days of development still.

# Summary

- Below a “warm” 90 K, UCN storage improves by > 4 times compared to room temperature
- Very long UCN storage times for multiple full-sized cells:  $\tau_{\text{single-exp}} = (490 \pm 11) \text{ s}$  or  $\tau_{\text{long}} \approx 550 \text{ s}$
- In-situ superthermal UCN production  $^4\text{He}$  offers many advantages (no need to transport, high density, phase-space evolution reduction)
- However, there is **activation caused by the 9 Å neutron beam**. Need to add neutron absorber, vikuiti reflector, and thin and pure dPMMA end windows. (*Should not affect UCN storage.*)
- Polarized  $^3\text{He}$  used as a co-magnetometer and live and in-situ UCN spin analyzer allows the **continuous measurement the neutron phase**. (In order to “measure frequency”: this is the way...)
- This is a new type of signals in nEDM experiments. **Detailed simulations are being performed to understand how to best analyze this data**
- $\phi_n(\mathbf{t})$  information allows a new way for studying UCN spectral evolution inside the cell by coupling to the UCN center-of-mass gravitational offset coupling to an applied B-gradient



Simulated signal 1-cycle (accepted events)

