3D (Y, M, Δ) scaling laws and projection effects in **The300-NIKA2 Sunyaev-Zeldovich Large Program Twin Samples**

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Motivation

- Galaxy clusters are the most massive, gravitationally collapsed structures in the Universe, serve as good as growth of structure tools to study cosmology (see e.g. Planck Collaboration 2016a; Reiprich et al. 2021)
- Galaxy cluster number density and distribution carries the imprint of structure formation and the underlying cosmology
- Requires accurate mass estimates with high precision but, mass is not directly observable:

=>Need for good calibration of observable-mass scaling relations =>Need to acquire detailed astrophysical information

 This requires a combination of high sensitivity and high angular resolution observations, with detailed hydrodynamical simulations



Theoretical galaxy cluster mass functions for different sets of cosmological parameters.



The NIKA2 Sunyaev–Zeldovich Large Program (LPSZ)* an The300**

LPSZ (talk by L. Perotto): Guaranteed time program of NIKA2, dedicated to SZ observations of clusters

- Sample of 38 clusters with 0.5 < z < 0.9, covering about one order of magnitude in mass
- Follow-up of SZ-selected clusters from the Planck and Atacama Cosmology Telescope (ACT) catalogs

The300 (talk by W. Cui): large sample of simulated galaxy clusters and their environment modelled using a range of simulation packages and physics modules that provides mock maps of the clusters in X-ray, optical, gravitational lensing, radio, and Sunyaev-Zeldovich (SZ) wavelengths

- 324 regions of radius 15 Mpc/h having a cluster identified at z = 0 with mass M_200 > 6.42 × 10¹⁴ M_o/h, at its centre
- Hydrodynamical simulations run for each region using standard SPH (GADGET-MUSIC) and modern SPH (GADGET-X, GIZMO-Simba)
- Semi-analytical model (SAM) counterparts produced using Galacticus, SAG, and SAGE

*L. Perotto et al. 2021, **<u>the300-project.org</u>



The IRAM 30-Meter telescope at Pico Veleta, Spain: site of the NIKA2 camera



An example of the clusters from The300. The coloured maps show the projected dark matter density, galaxy colours are calculated from the SDSS band magnitudes, and the sizes represent their stellar mass. W. Cui et al. 2018.

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The300-NIKA2 twin sample

The samples:

- Select snapshots in The300 with median z closest to z of each NIKA2 bin
- Create three "twin simples" (TS) based on following criteria
 closest total M_500: TS_TM
 - closest hydrodynamical mass M_500: **TS_HEM**
 - closest integrated Compton parameter Y_500: TS_Y

Available data:

• y-maps (multiple projections)

$$y \equiv \int \frac{k_{\rm B} T_{\rm e}}{m_{\rm e} c^2} \,\mathrm{d}\tau_{\rm e} = \int \frac{k_{\rm B} T_{\rm e}}{m_{\rm e} c^2} \,n_{\rm e} \sigma_{\rm T} \mathrm{d}l = \frac{\sigma_{\rm T}}{m_{\rm e} c^2} \int P_{\rm e} \,\mathrm{d}\tau_{\rm e}$$

- halo catalogues
- projected total mass density maps (multiple projections)
- ICM profiles (density, pressure, mass..)



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Integrated Y parameter estimates from y maps



To quantify: the spread in Y along different LoS and find the overdensity at which the spread is minimum Aishwarya Paliwal, mm Universe conference 2023 6/18

*All projected maps (y and M_total) from which the integrated (cylindrical) quantities are obtained, have an integration depth of 4 x r_200

For each twin sample we have y maps* along 29 different lines-of-sight, we estimate Y:



at different overdensities for all clusters in a TS

The shaded region indicates the spread of Y along the 29 LoS





Spread of Y along different LoS



We do not observe a special overdensity with minimum scatter in Y along different LoS: scatter decreases with decreasing overdensity

- The scatter induced on Y due to projections is quantified using the standard deviation in the value of Y calculated from
 - different LoS maps at a given overdensity.
 - We calculated this scatter for each cluster in the TSs.
 - Here we show the **mean scatter** of the sample. The shaded area is the 1σ region.

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*De Luca et al. 2021 NB: we do find correlation between the spread in y and χ



- NFW profiles (using Colossus, Diemer 2018): Projection 12

$$\Sigma(x) = \begin{cases} \frac{2\rho_{s}r_{s}}{x^{2}-1} \left(1 - \frac{2}{\sqrt{x^{2}-1}} \arctan \sqrt{\frac{x-1}{x+1}}\right) & (x > 1) \\ \frac{2\rho_{s}r_{s}}{x^{2}-1} \left(1 - \frac{2}{\sqrt{1-x^{2}}} \operatorname{arctanh} \sqrt{\frac{1-x}{1+x}}\right) & (x < 1) \\ \frac{2\rho_{s}r_{s}}{3} & (x = 1) \end{cases} \xrightarrow{r} \left(x < 1 \right) \\ (x = 1) \xrightarrow{r} \left(R_{500} \right) \\ Fxamples of NEW/ fite$$



Examples of 3D mass profiles obtained for some projections of cluster 27, snap 101, R_500 = 824.45 kpc, M_500 = 7.1x10^14 M_sol 3. We have the **spherical mass** estimates from the simulation mass profile: **M_sph**; used to compare our mass estimates

(1 and 2) with simulation

Projected mass estimates by directly integrating, within an aperture, the projected total mass density maps: **M_cyl** We construct total projected mass density profiles from the corresponding simulated maps and fit them to projected



And **de-project** the surface density profiles to 3D density profiles and consequently the **spherical mass**: **M_sph,NFW**

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 $M_{sph} - M_{sph_{NFW}}$ Comparison between spherical mass values obtained from simulation and NFW fits: $b_M = M_{sph}$

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Comparison between cylindrical and spherical masses obtained from simulation and NFW fits: this comparison becomes relevant when comparing scaling laws built with M_cyl vs M_sph



As expected, the cylindrical masses are higher than the spherical ones, with the ratio decreasing with decreasing overdensity

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Aishwarya Paliwal, mm Universe conference 2023

density e 2023

The scatter induced on mass due to the projection effect: higher in de-projected, spherical mass



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3D scaling law: Y, M, Δ

We build a scaling laws between Y, M, Δ :

where M Δ can either be M cy



The Y-M distribution, at different overdensities, of the TSs. The errors are the scatter induced due to projection. We use pylira (https://github.com/fkeruzore/pylira), the python wrapper of the LIRA code (M. Sereno 2016) to fit the scaling laws. We assume no redshift evolution in the slope, normalisation, or the intrinsic scatter.

*Planck Collaboration et al. 2014



$$: E(z)^{-\frac{2}{3}} \left[\frac{D_A^2 Y_\Delta}{10^{-4} \text{ Mpc}^2} \right] = A \left[\frac{M_\Delta}{6.10^{14} \text{ M}_\odot} \right]^B *$$

1 cyl or M sph,NFW and Y Δ is Y cyl

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3D scaling law: best-fit parameters



- \bullet
- No significant trend in the slope other than its dispersion decreasing with decreasing overdensity \bullet
- No departure in the value of slope from the self-similar value

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3D scaling law: best-fit parameters



- The scatter (and its dispersion) is **lower for M_cyl**
- derived with no observational/instrumental error and hence maybe underestimated

• The intrinsic scatter in the SL (at z=0) decreases with decreasing overdensity; e.g. the scatter at $\Delta(500)=1.2 \times \Delta(200)$

• The scatter due to projection is comparable to intrinsic scatter, hence cannot be ignored. Albeit, the intrinsic scatter is

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3D scaling law: comparison with existing scaling laws



The M_sph,NFW-Y_cyl scaling laws compared to literature. The errors on individual data points are omitted for visualisation purpose.

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Ongoing/future work

To include more realistic observational errors we extend this analysis in the following way

=> Create NIKA2 like maps for all the simulated y projections (example of one projection is shown on the right side for, cluster 29, snap 101)

=> Combine these maps with ICM electron density profiles from simulations to calculate HSE mass at different over-densities using:

$$M_{\text{HSE}}(< r) = -\frac{1}{\mu m_p G} \frac{r^2}{n_e(r)} \frac{\mathrm{d}P_e(r)}{\mathrm{d}r}$$

with the standard procedure followed for the LPSZ clusters (talk by C. Hanser, A. Moyer)

=> Fit the SL law using this mass estimate and compare the results to the present ones



z=0.82, M_500=6.76e14 M_sun, R_500=998.83 kpc, χ=0.78

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Takeaway

- is no special overdensity at which the scatter is minimum
- mass values M_sph,NFW as compared to M_cvl
- of the Y-M scaling law hence can not be ignored and the intrinsic scatter may be underestimated
- dynamical state

 Optimal overdensity: the scatter induced on Y due to projection and the Y-M scaling law intrinsic scatter decrease with decreasing overdensity i.e. there

• The projection induced scatter on both, Y and M is higher for de-projected

The scatter induced due to projection is comparable to the intrinsic scatter

Caveat: the scaling laws are calibrated without realistic observational errors

• The scatter induced on Y due to projection is not correlated with the cluster



Backup: The Sunyaev-Zeldovich (SZ) effect in galaxy clusters

$$y \equiv \int \frac{k_{\rm B} T_{\rm e}}{m_{\rm e} c^2} \,\mathrm{d}\tau_{\rm e} = \int \frac{k_{\rm B} T_{\rm e}}{m_{\rm e} c^2} \,n_{\rm e} \sigma_{\rm T} \mathrm{d}l = \frac{\sigma_{\rm T}}{m_{\rm e} c^2} \int P_{\rm e} \,\mathrm{d}\tau_{\rm e}$$

The derived pressure (profile) can be used to estimate mass of the cluster using the hydrostatic equilibrium equation

Further, the integrated Compton parameter is tightly related to the mass of the cluster.

$$Y\equiv\int_{\Omega} y\,d\Omega$$







Correlation between chi and spread in y

