

**Virtual states, halos and resonances in three-body
atomic and nuclear systems**

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OUTLINE

Physics of weakly bound three-body systems

Two neutron weakly bound three-body halo nuclei/
Atomic weakly bound three-body systems

Scales of n-n-c system: contact interaction

Classification scheme

Threshold conditions for excited Efimov states

n-¹⁹C scattering and Efimov physics

Summary and outlook

Physics of weakly bound three-body systems

Charateristic phenomena three-body systems:

Thomas collapse (1935) and Efimov effect (1970)

$$r_0 \rightarrow 0$$

$$|a| \rightarrow \infty$$

??? infinitely many weakly bound states

$$|a|/r_0 \rightarrow \infty$$

Thomas-Efimov effect!

S.K. Adhikari, A. Delfino, T. Frederico, I.D. Goldman, and L. Tomio, Phys. Rev. A **37**, 3666 (1988).

One three-body scale is necessary to represent short-range physics !!!!

A. S. Jensen, K. Riisager, D. V. Fedorov, and E. Garrido, Rev. Mod. Phys. **76**, 215 (2004).

E. Braaten, H.-W. Hammer, Phys. Rep. **428**, 259 (2006)

Zero-range 3-boson equation: Thomas-Efimov effect

Skorniakov and Ter-Martirosian equations (1956)

$$\chi(\vec{y}) = \frac{-\pi^{-2}}{\pm \sqrt{\epsilon_2} - \sqrt{\epsilon_3 + \frac{3}{4}y^2}} \int d^3x \left(\frac{1}{\epsilon_3 + y^2 + x^2 + \vec{y} \cdot \vec{x}} - \frac{1}{1 + y^2 + x^2 + \vec{y} \cdot \vec{x}} \right) \chi(\vec{x})$$

$$\epsilon_3 = E_3 / \mu_{(3)}^2 \quad \epsilon_2 = E_2 / \mu_{(3)}^2 \quad \mu_{(3)}^2 = 1$$

Adhikari,TF,Goldman, PRL74 (1995) 487

Thomas collapse: $\mu_{(3)}^2 \rightarrow \infty$

$$\epsilon_2 = E_2 / \mu_{(3)}^2$$

Efimov effect: $E_2 \rightarrow 0$

$$\epsilon_3^{(N)} \equiv \epsilon_3^{(N)} (\pm \sqrt{\epsilon_2})$$

Scaling limit & limit cycle

$$\epsilon_3^{(N)} \equiv \epsilon_3^{(N)} (\pm \sqrt{\epsilon_2})$$

$$\xi \equiv \pm \sqrt{\epsilon_2} = \pm (E_2 \epsilon_3^{(N)} / E_3^{(N)})^{1/2}$$

$$\frac{E_3^{(N+1)}}{E_3^{(N)}} = \lim_{N \rightarrow \infty} \frac{\epsilon_3^{(N+1)}(\xi)}{\epsilon_3^{(N)}} = \mathcal{F} \left(\pm \sqrt{\frac{E_2}{E_3^{(N)}}} \right)$$

$$\mathcal{F}(0) = 1/515$$

Scaling function

Efimov 1970

Scaling limit:

Frederico et al PRA60 (1999)R9

Yamashita et al PRA66(2003)052702

Limit cycle:

Mohr et al Ann.Phys. 321 (2006)225

Scaling functions: Correlation between observables

$$O(E, E_3, E_2) = (E_3)^\eta \mathcal{A}(\sqrt{E/E_3}, \sqrt{E_2/E_3})$$

Scaling function

Scaling limit:

Frederico et al PRA60 (1999)R9

Yamashita et al PRA66(2003)052702

Limit cycle:

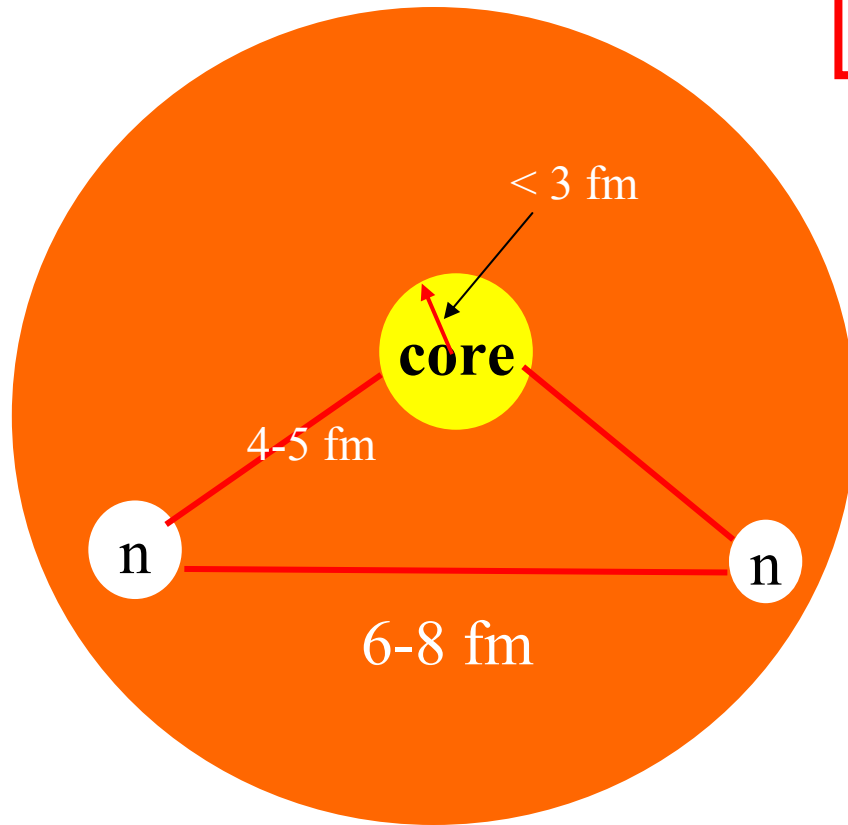
Mohr et al Ann.Phys. 321 (2006)225

Correlation between S-wave observables

- Phillips plot: triton B.E. versus doublet scattering length
- 2nd order n-d polarization observables versus triton B.E.
- Trapped atomic trimer B.E. versus recombination rate

Two neutron weakly bound three-body halo nuclei

core-neutron-neutron halo nuclei



^{11}Li ^{14}Be ^{20}C

Binding energy \sim MeV or $<$ MeV

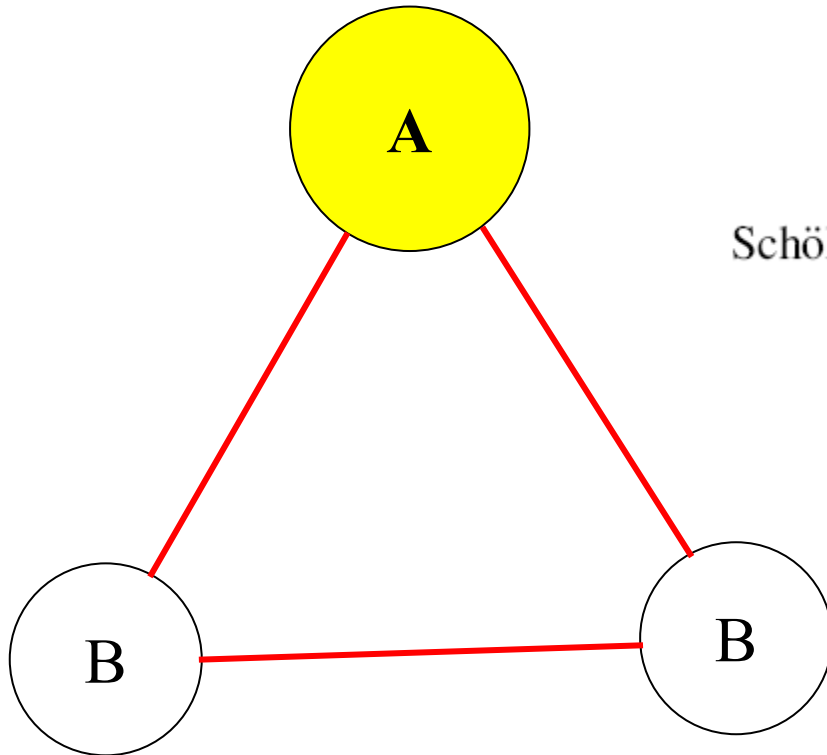
$R_{nn}(\text{Exp}) \sim 6 - 8 \text{ fm}$ (^{11}Li)

F. M. Marqués et al. Phys. Rev. C **64**, 061301 (2001)

M. Petrascu et al. Nucl. Phys. A **738**, 503 (2004)

Tanihata et al., Phys. Rev. Lett. **55**, 2676 (1985)

Atomic weakly bound three-body systems



Schöllkopf, W., Toennies, J. P.: Science **266**, 1345 (1994)

dimer $R_{\text{He-He}} \sim 50 \text{ \AA}$

A-B-B weakly bound molecules

ultra-low binding $\sim \text{mK}$ or $< \text{mK}$

$^{133}\text{Cs}_3$ (trapped ultracold gas near a Feshbach resonance)

$^4\text{He}_3$ $^4\text{He}_2 - ^7\text{Li}$ $^4\text{He}_2 - ^6\text{Li}$ $^4\text{He}_2 - ^{23}\text{Na}$

Scales of n-n-c system: contact interaction

E_{nn} Energy of the bound/virtual nn system

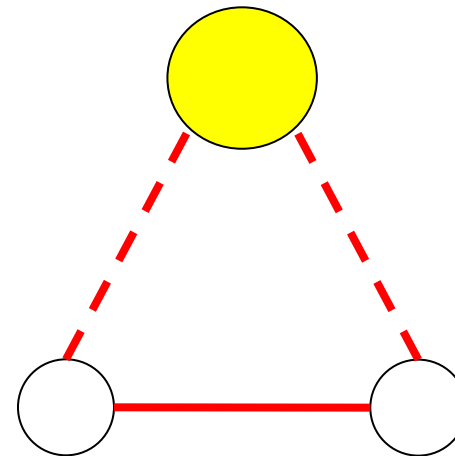
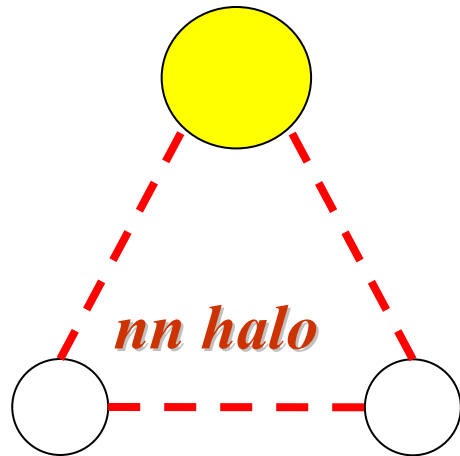
E_{nc} Energy of the bound/virtual nc system

$B_N = |E_3^{(N)}|$ Energy of the Nth state of the nnc system

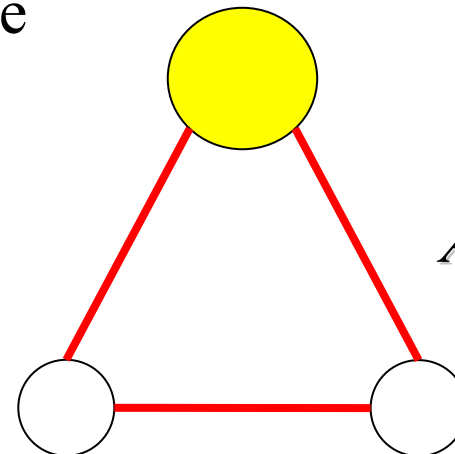
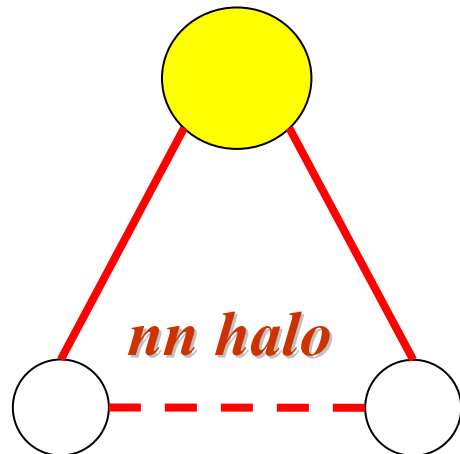
A = mass of the core

Classification scheme

Yamashita, Tomio and T. F. Nucl. Phys. A 735, 40 (2004)



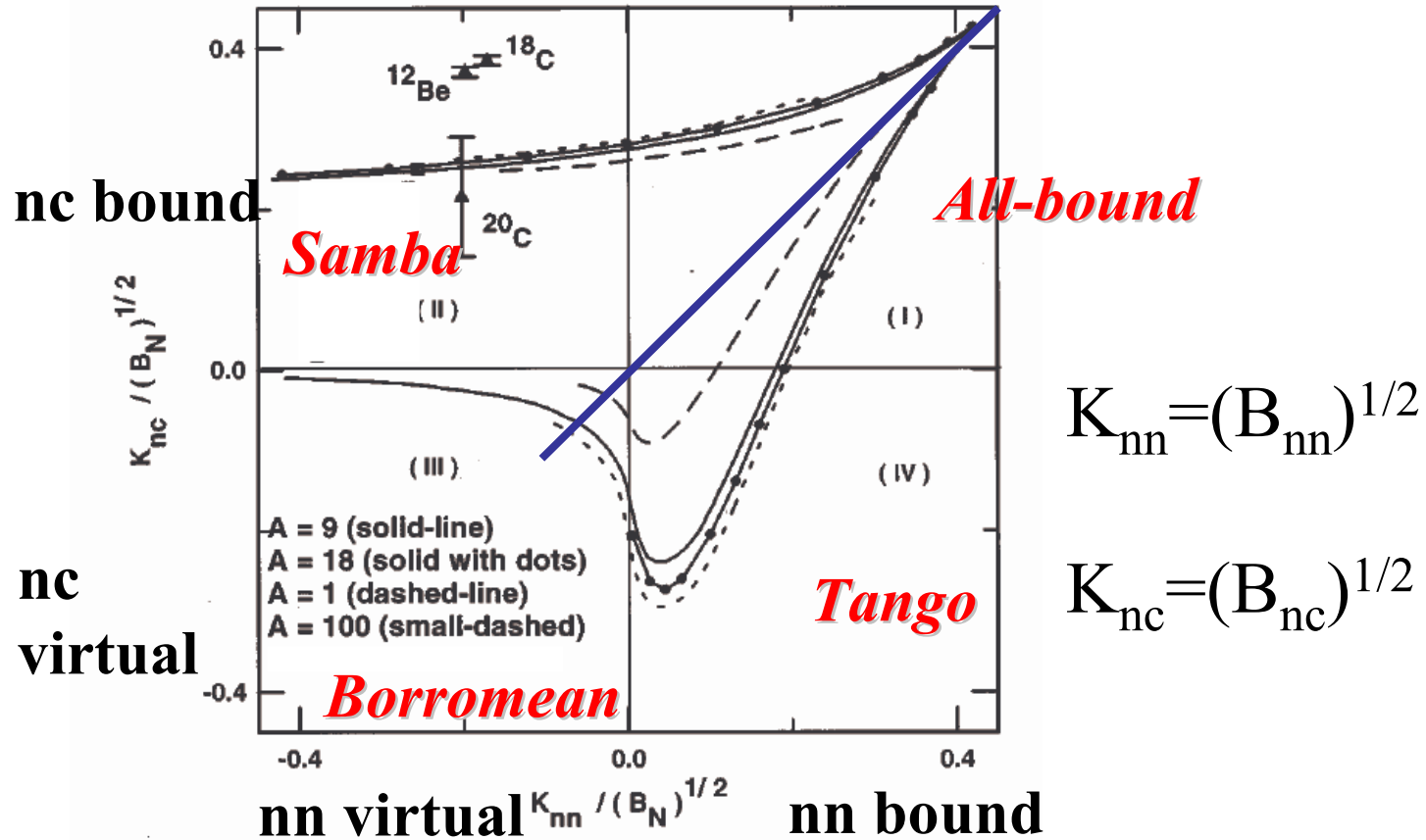
— bound state
- - - virtual state



A.S. Jensen, K. Riisager, D.V. Fedorov, E. Garrido, Europhys. Lett. 61 (2003) 320.
F. Robicheaux, Phys. Rev. A 60 (1999) 1706.

Threshold for an excited Efimov state: Halo-nuclei

Critical condition for an excited (N+1)-th Efimov state above the N-th one:



Amorim,TF,Tomio PRC56(1997)2378

D. L. Canham and H.-W. Hammer, *Universal properties and structure of halo nuclei*, arXiv:0807.3258.

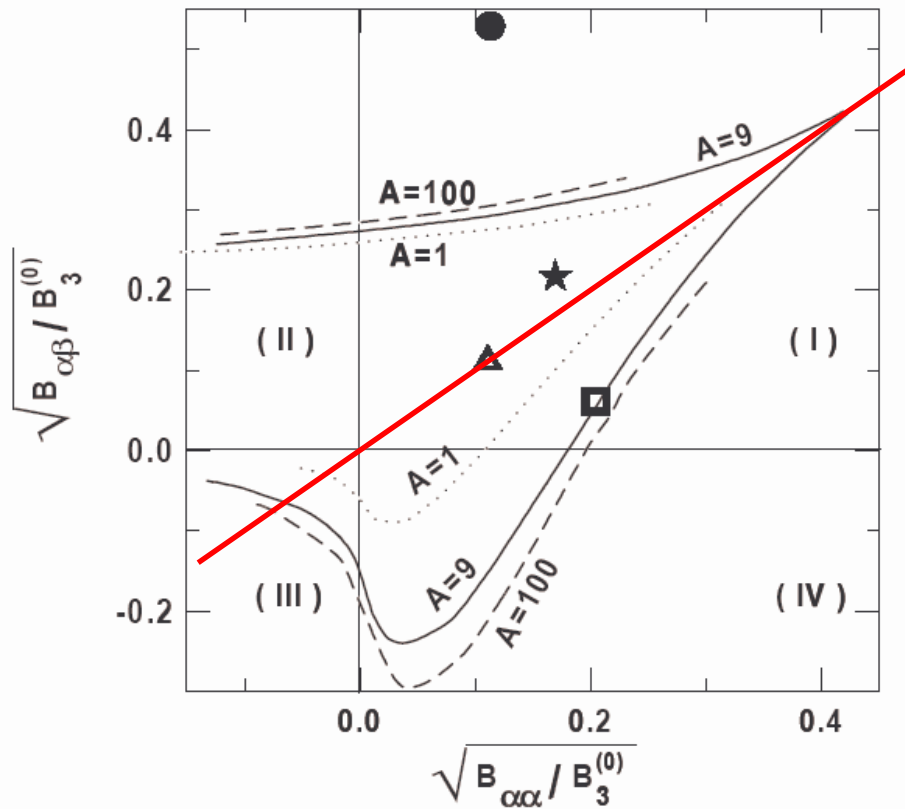
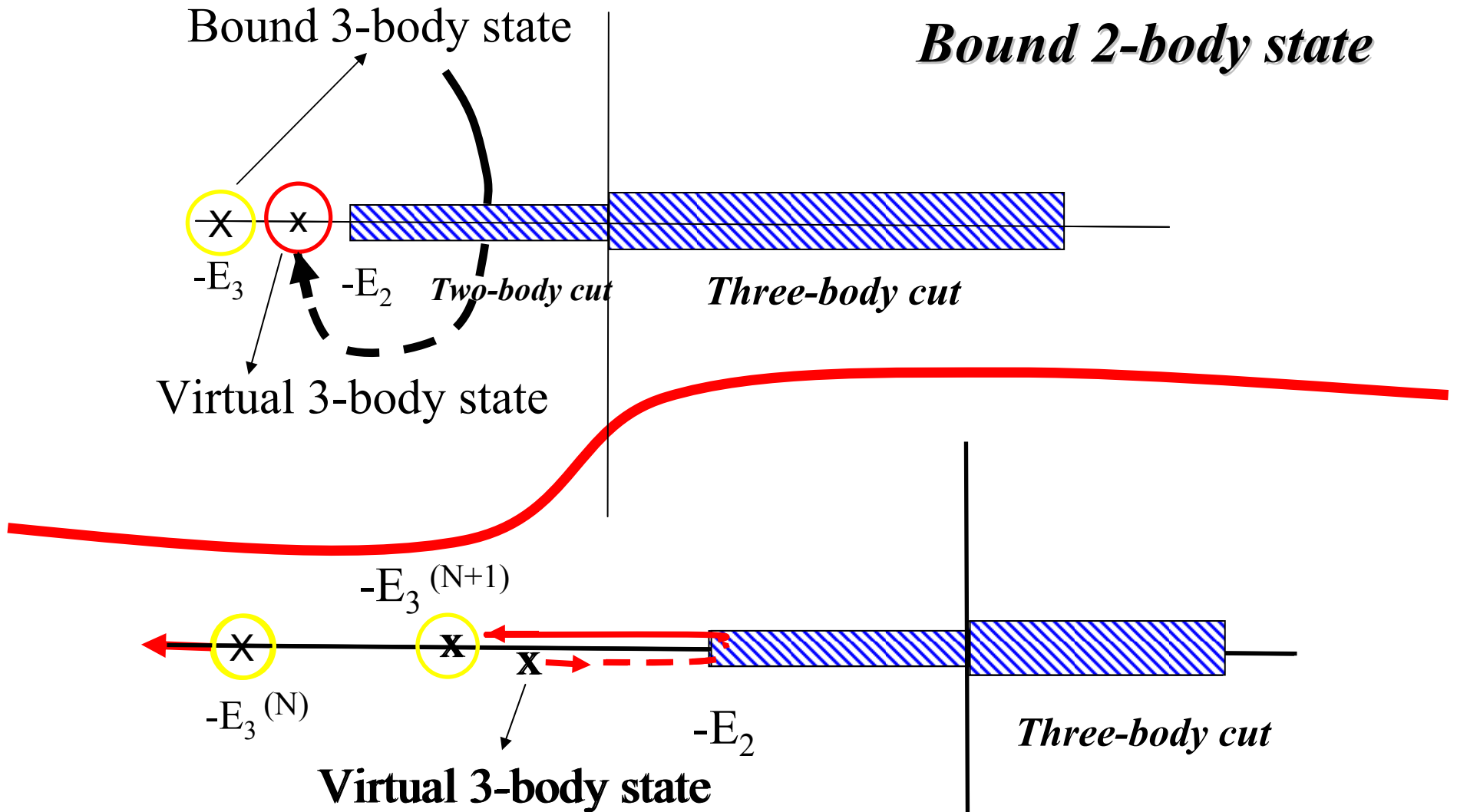
Threshold for an excited Efimov state: Weakly bound molecules

FIG. 1. The scaling approach for the three-body system $\alpha - \alpha - \beta$. The coordinates of the systems with $\alpha \equiv {}^4\text{He}$ and $\beta = {}^4\text{He}, {}^7\text{Li}, {}^6\text{Li}$ and ${}^{23}\text{Na}$ are respectively represented by a triangle, a star, a square and a full circle.

$$K_{aa} = (B_{aa})^{1/2}$$

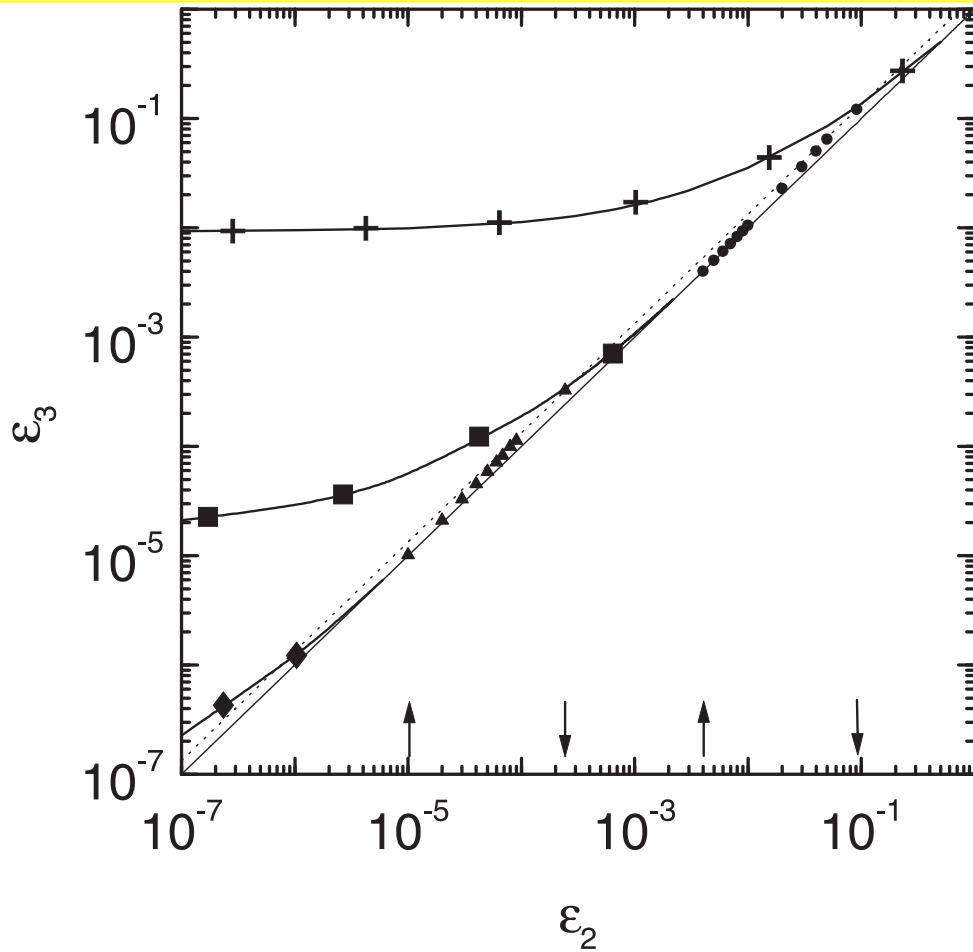
$$K_{ab} = (B_{ab})^{1/2}$$

Three-bosons: analytic structure & Efimov state trajectory



S.K. Adhikari and L. Tomio, Phys. Rev. C **26**, 83 (1982); S.K. Adhikari, A.C. Fonseca, and L. Tomio, *ibid.* **26**, 77 (1982).

Efimov States – Bound and virtual states (3 identical bosons)



Bound-states (lines with simbols)

With Plus - Fundamental state

With squares – 1st. excited

With diamonds – 2nd. excited

Virtual-states (just simbols)

circles – 1st. Excited state

triangles – 2nd. Excited state



the virtual state starts (dotted line)

Virtual-state turns to an excited state (continuous line)

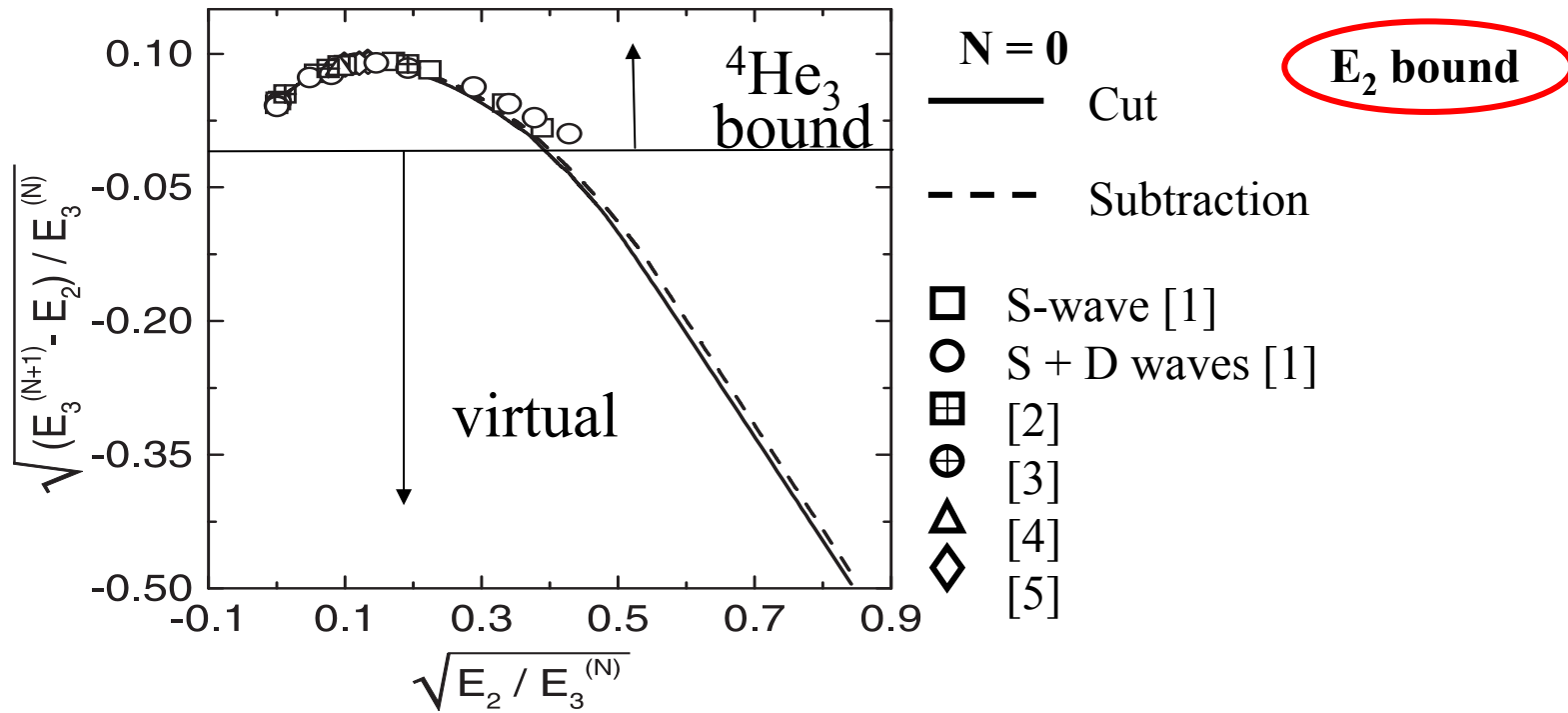
$$\epsilon_3 = \frac{4}{3} \epsilon_2$$

ϵ_2 bound

$$\epsilon_3 = \epsilon_2$$

Efimov States – Bound and virtual states (3 identical bosons)

T. Frederico, LT, A. Delfino and E. A. Amorim. *Phys. Rev. A* **60**, R9 (1999).



[1] Th. Cornelius e W. Glöckle. *J. Chem. Phys.* **85**, 1 (1996).
 [2] S. Huber. *Phys. Rev.* **A31**, 3981 (1985).
 [3] P. Barletta e A. Kievsky. *Phys. Rev.* **A64**, 042514 (2001).
 [4] D. V. Fedorov e A. S. Jensen. *J. Phys.* **A34**, 6003 (2001).
 [5] E. A. Kolganova, A. K. Motovilov e S. A. Sofianos. *Phys. Rev.* **A56**, R1686 (1997).

Range correction: Thogersen, Fedorov, Jensen PRA78(2008)020501(R)

Range correction: qualitative discussion

$$f(k) = (-a^{-1} + r_0 k^2 / 2 - ik)^{-1}$$

Bound state: $k = i (E_2)^{1/2} \rightarrow -a^{-1} - r_0 E_2 / 2 + (E_2)^{1/2} = 0$

$$f(k) \sim (-E_2^{1/2} - ik)^{-1} (1 + r_0 E_2^{1/2})$$

Effective interaction in 3-boson eq. **increases!**

\rightarrow Critical value of a to allow an Efimov state **decreases!**

Virtual state: $k = -i (E_2)^{1/2} \rightarrow -a^{-1} - r_0 E_2 / 2 - (E_2)^{1/2} = 0$

$$f(k) \sim (E_2^{1/2} - ik)^{-1} (1 - r_0 E_2^{1/2} / 2)$$

Effective interaction in 3-boson eq. **decreases!**

\rightarrow Critical value of a to allow an Efimov state **increases!**

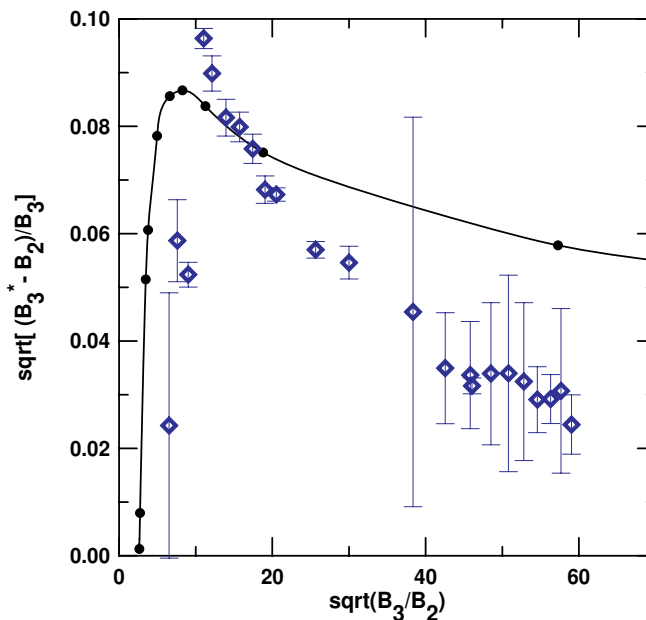
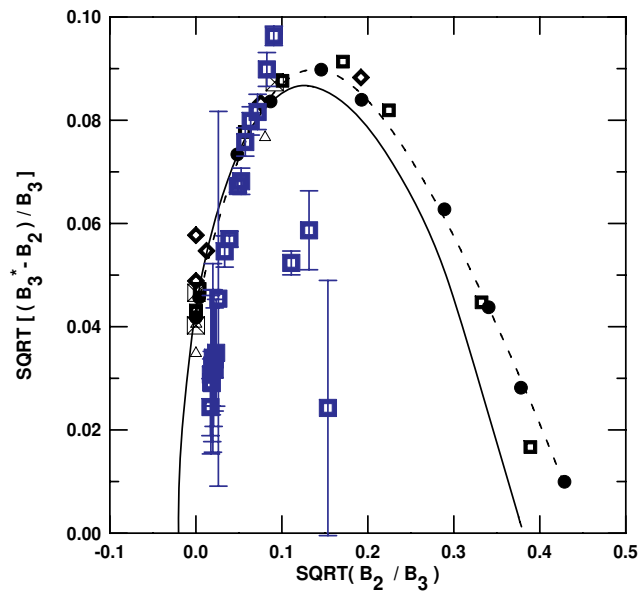
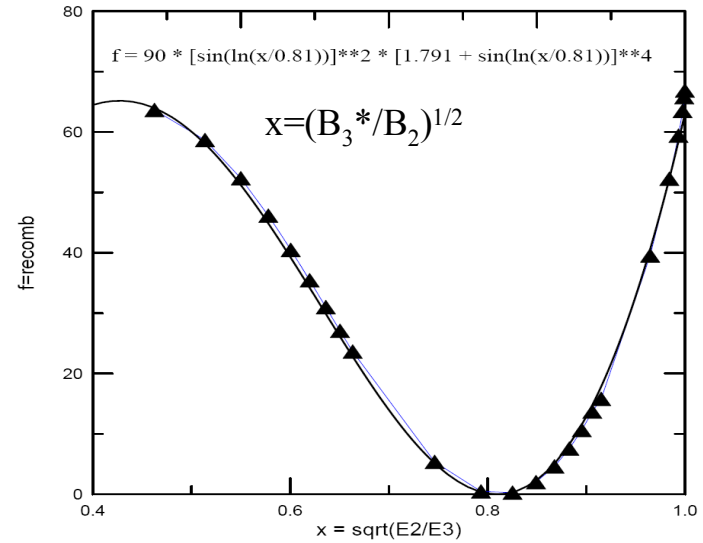
Three-atom recombination & trimer energy



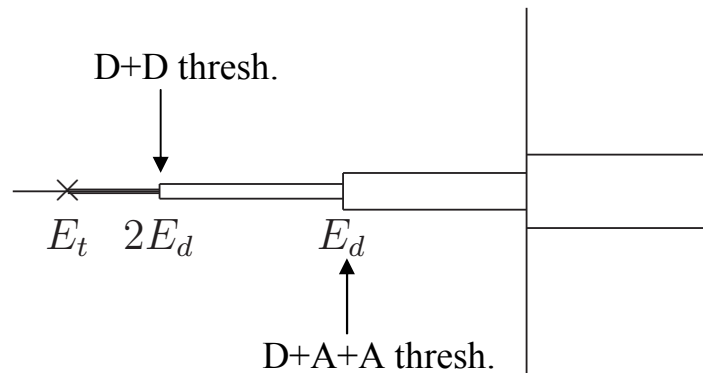
3-body recombination of Cesium ultracold gas data near Feshbach resonance v.s. scattering length

T. Kraemer et al, Nature 440, 315 (2006)

Resonance @ $T \sim 0$ and $a = -898a_0 \rightarrow B_3 = 1.31 \text{ mK}$



PRA68, 033406 (2003)
Trimer formation in D+D

Efimov states & D+D recombination

How many Efimov states between D+D and D+A+A thresholds ($E_d \rightarrow 0$)?

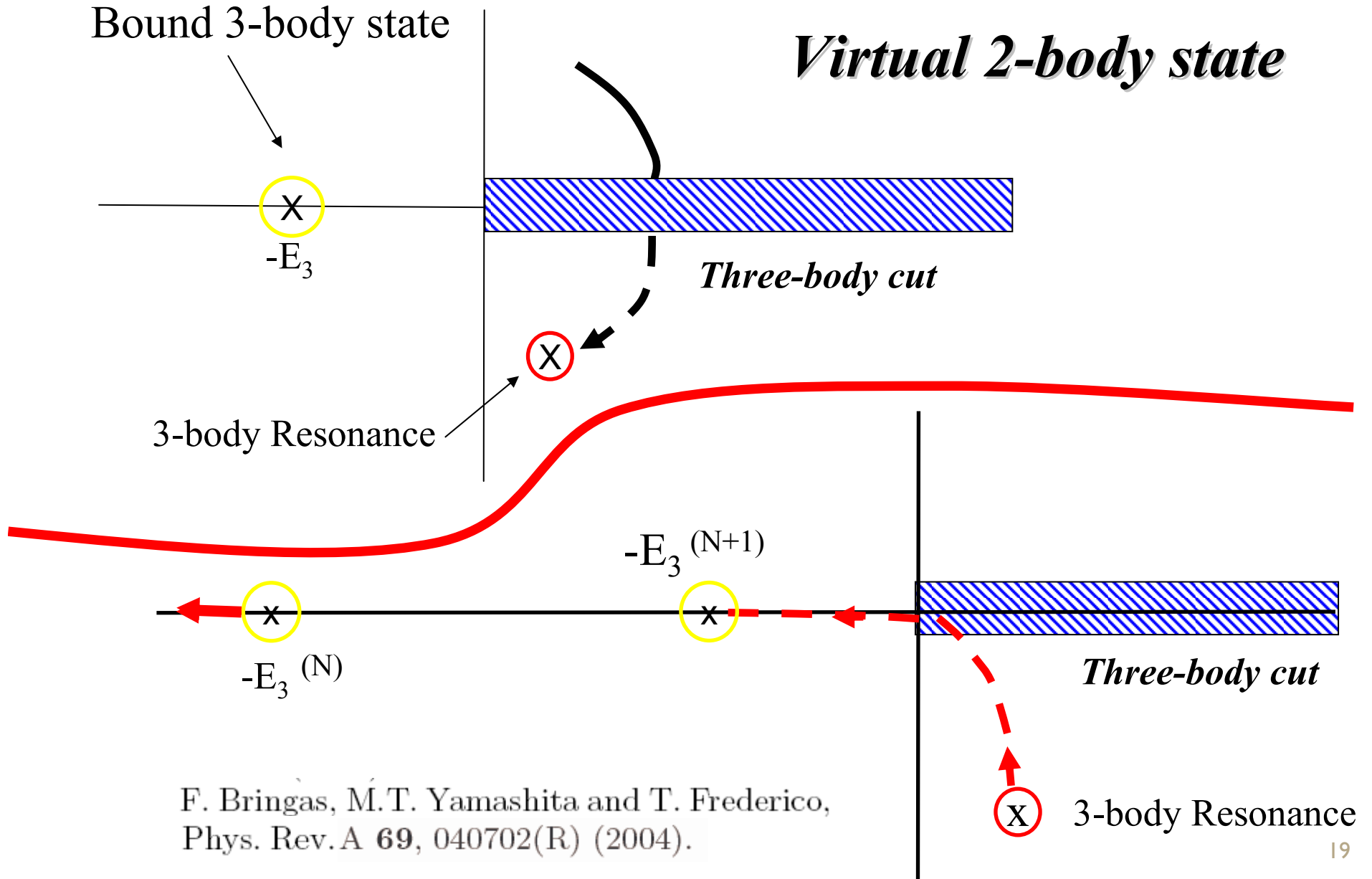
For $E_t = 6.9E_d \rightarrow E_t^* = E_d$ & $E_t > 6.9E_d \rightarrow |E_t^*| > |E_d|$

→ **At most one Efimov state appears in the range: $2|E_d| < |E_t| < |E_d|$!!!!**
(Range corrections hardly destroy that $\sim 30\%$!)

→ dimer+dimer recombination with trapped ultracold cesium?

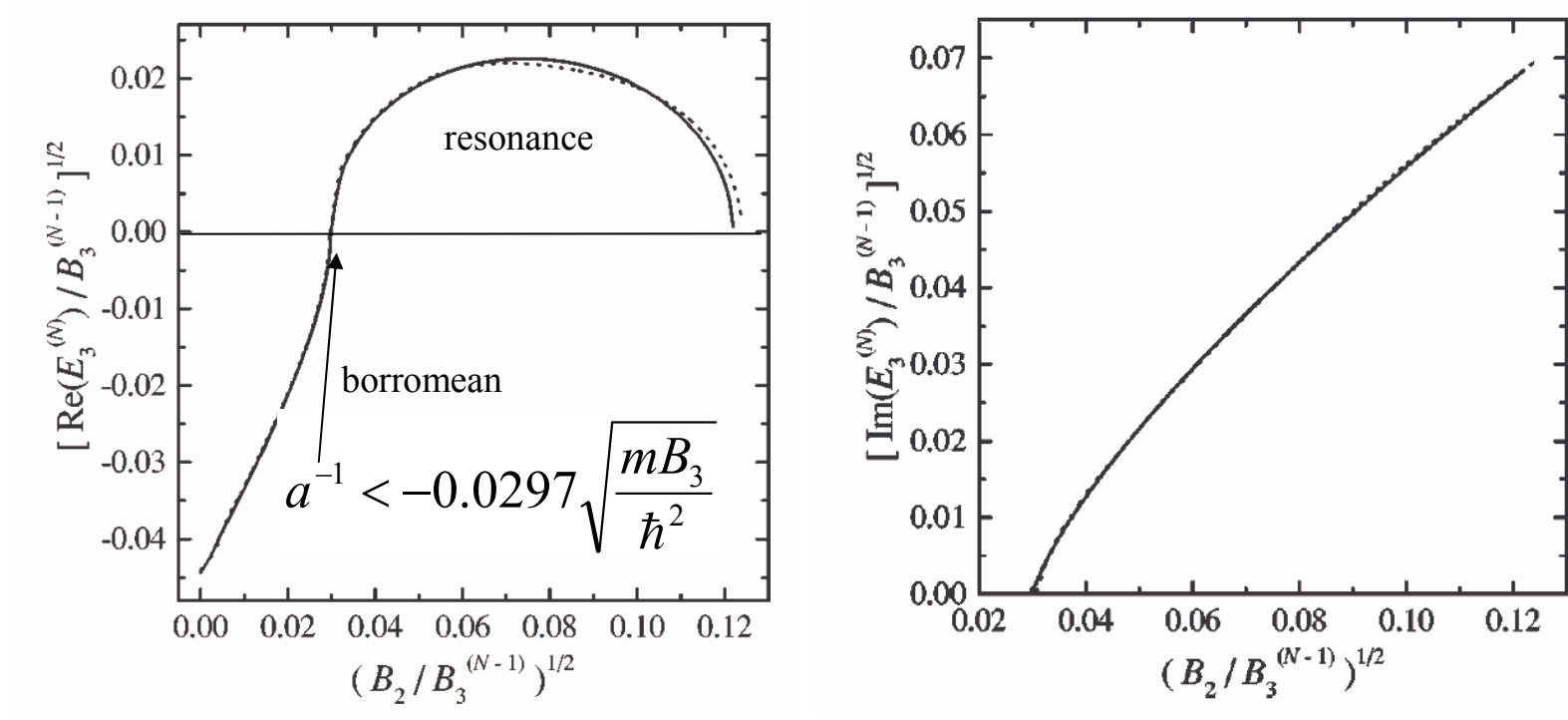
F. Ferlaino, S. Knoop, M. Mark, M. Berninger, H. Schobel,
H.-C. Ngerl, R. Grimm Phys. Rev. Lett. **101**, 023201
(2008).

Borromean configuration: analytic structure & Efimov state trajectory



Efimov state trajectory: borromean case

S-wave three-boson resonance



F. Bringas, M.T. Yamashita and T. Frederico, Phys. Rev. A **69**, 040702(R) (2004).

Evidence of continuum resonances in recombination of ultracold Cs atoms

T. Kraemer et al, Nature **440**, 315 (2006)

Evidence of continuum resonances in ultracold cesium gas

T. Kraemer et al, Nature 440, 315 (2006)

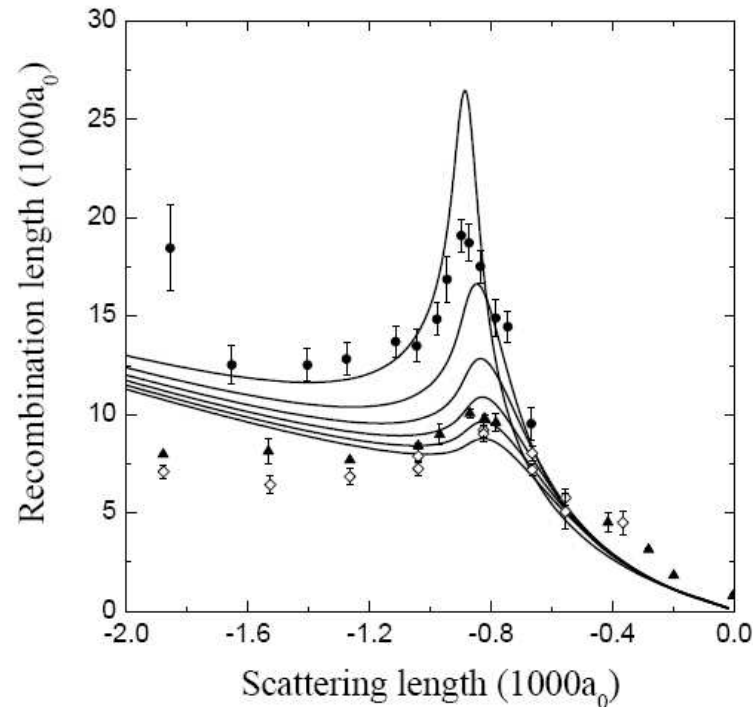
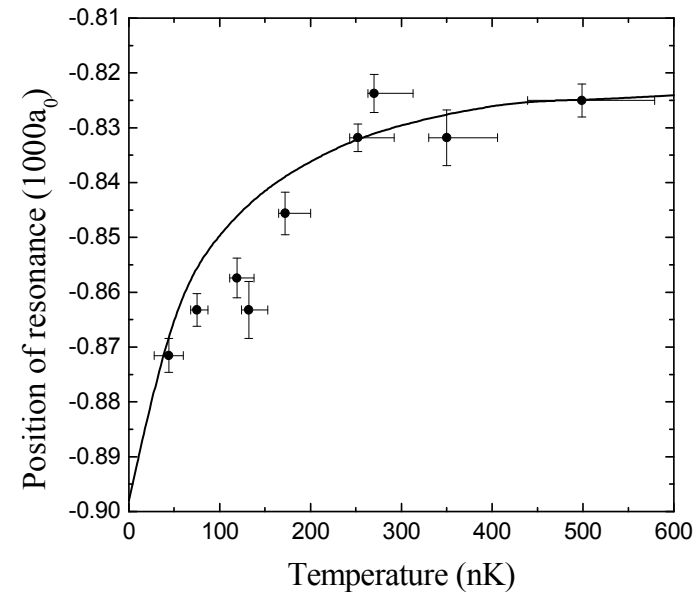


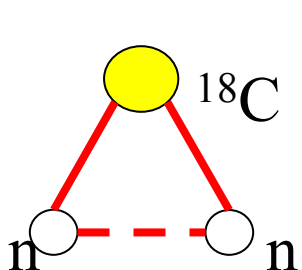
FIG. 2: Recombination length ($\rho_3 = [2m/(\sqrt{3}\hbar)\langle L_3 \rangle_T]$) in the cesium trapped gas as a function of the scattering length and temperature. The solid curves from up to bottom are the theoretical results for $T = 10$ nK, 100 nK, 200 nK, 300 nK, 400 nK and 500 nK. The symbols are the experimental results for $T = 10$ nK (full circles), 200 nK (full triangles) and 250 nK (open diamonds) from Ref. [4].

Position of the maximum of the recombination length as a function of the temperature. Experimental data from B. Engeser et al., *in preparation*.

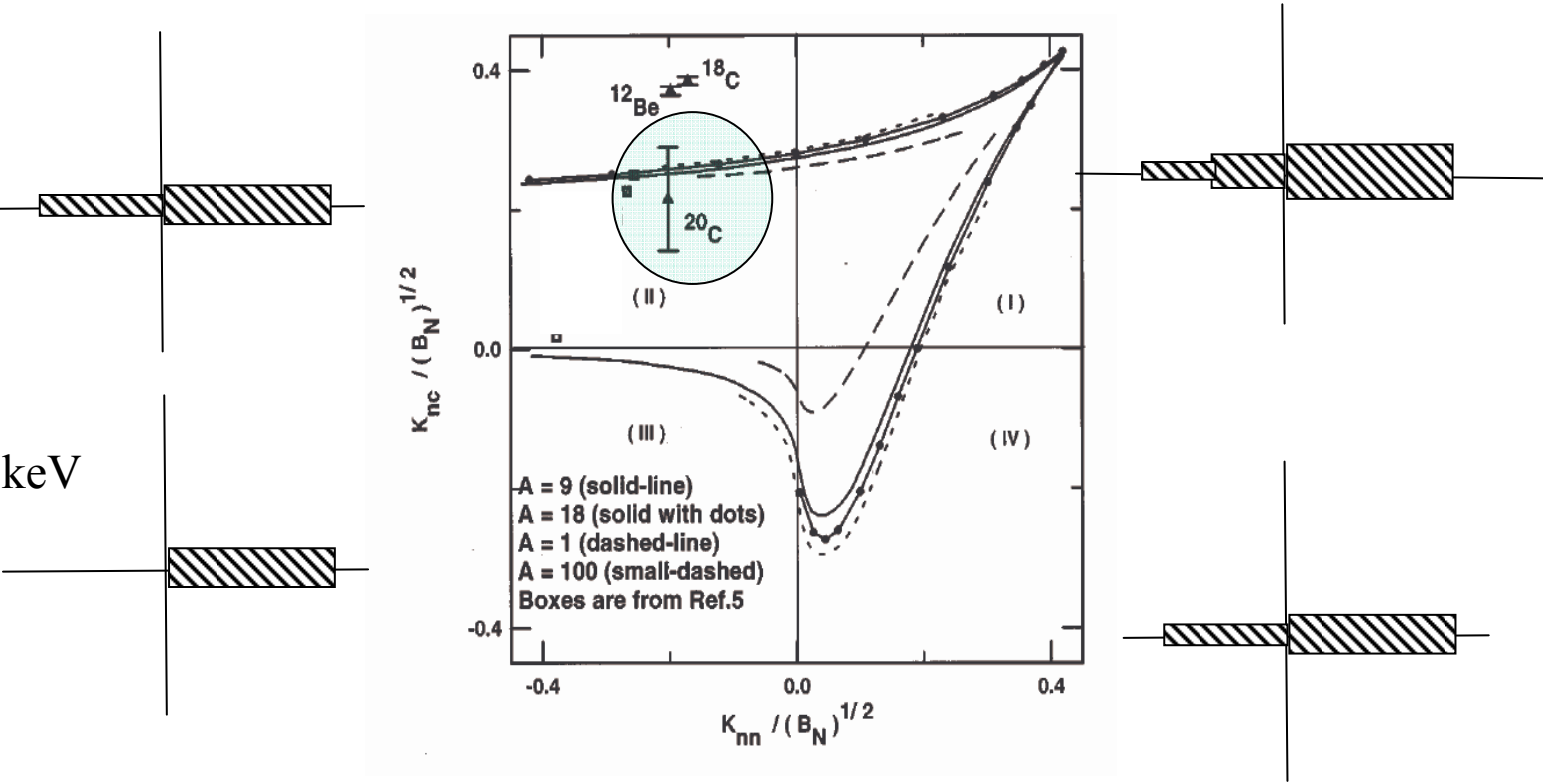


M.T. Yamashita et al. / Physics Letters A 363 (2007) 468.

Threshold for an excited Efimov state and trajectory: ^{20}C



$E_3 = 3.5 \text{ MeV}$
 $E_{nc} = 160 \pm 110 \text{ keV}$

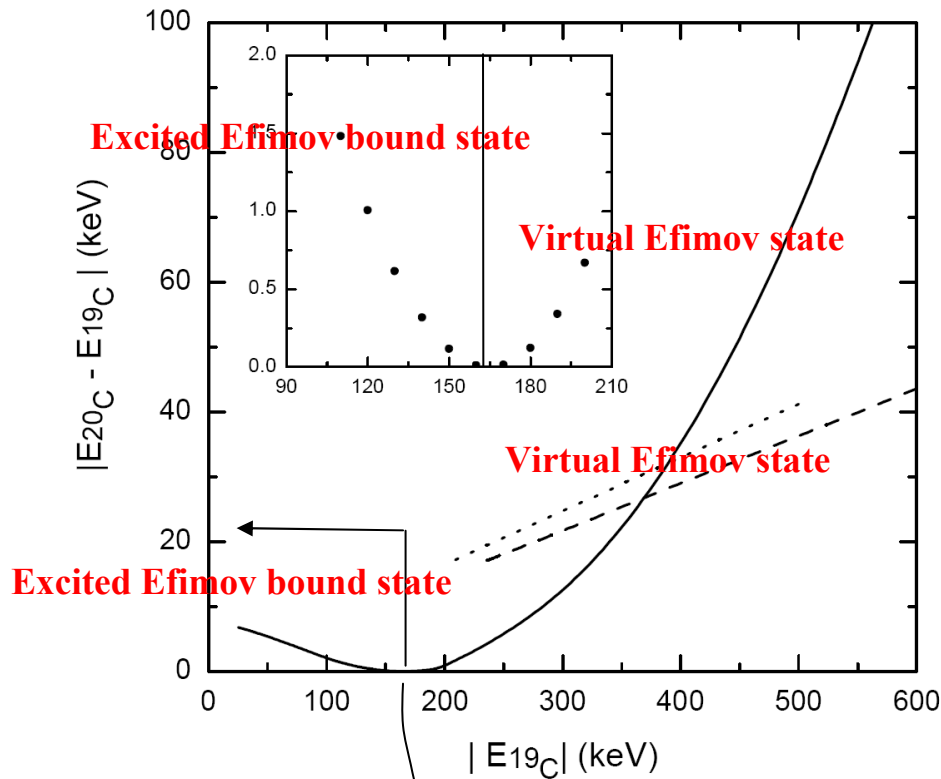


^{20}C can have a continuum resonance or virtual Efimov state?

Arora, Mazumdar, Bhasin PRC69 (2004)061301(R) Mazumdar, Rau, Bhasin PRL97(2006)062503
 Efimov state \rightarrow resonance of $n+^{19}\text{C}$ by changing K_{nc}

Threshold for an excited Efimov state and trajectory: ^{20}C

Yamashita, Frederico, Tomio,
PRL99 (2007)269201 Comment on "Efimov States and their Fano resonances..."
& PLB660(2008)339



Efimov state in ^{20}C goes to a virtual state for $|E_{19\text{C}}| > 165\text{keV}$!

Critical value: $|E_{19\text{C}}|=165\text{keV}$

n - ^{19}C scattering and Efimov physics

➔ Pole in s-wave $k\cot(\delta)$ for n-d system ! Well known ~ 50 years

Delves' 60, Van oers & Seagrave' 67, Girard & Fuda' 78

$$k\cot\delta_0 = -A + Bk^2 - \frac{C}{1 + Dk^2},$$

The existence of the triton virtual state was found on the basis of the effective range expansion.

➔ Universal property!

The atom-dimer (three-boson) scattering length is approximately given in Bratten and Hammer (Phys. Rep. 428 (2006) 259):

$$a_{AD} = (1.46 - 2.15 \tan[s_0 \ln(a\Lambda_*) + 0.09])a ,$$

where $s_0 = 1.00624$.

2-body scatt. lengths: }
 a_B for $a_{AD}=\text{infinity}$
 a_0 for $a_{AD}=0$

$$\frac{a_B}{a_0} = \exp\left(\frac{\pi/2 - 0.59654}{s_0}\right)$$

n - ^{19}C scattering and Efimov physics

Brief view...

$$\chi_n(\vec{q}) \equiv (2\pi)^3 \delta(\vec{q} - \vec{k}_i) + 4\pi \frac{h_n(\vec{q}; \mathcal{E}(k_i))}{q^2 - k_i^2 - i\epsilon},$$

$$h_n^\ell(q; \mathcal{E}_i) = \mathcal{V}^\ell(q, k_i; \mathcal{E}_i) + \frac{2}{\pi} \int_0^\infty dk k^2 \frac{\mathcal{V}^\ell(q, k; \mathcal{E}_i) h_n^\ell(k; \mathcal{E}_i)}{k^2 - k_i^2 - i\epsilon}.$$

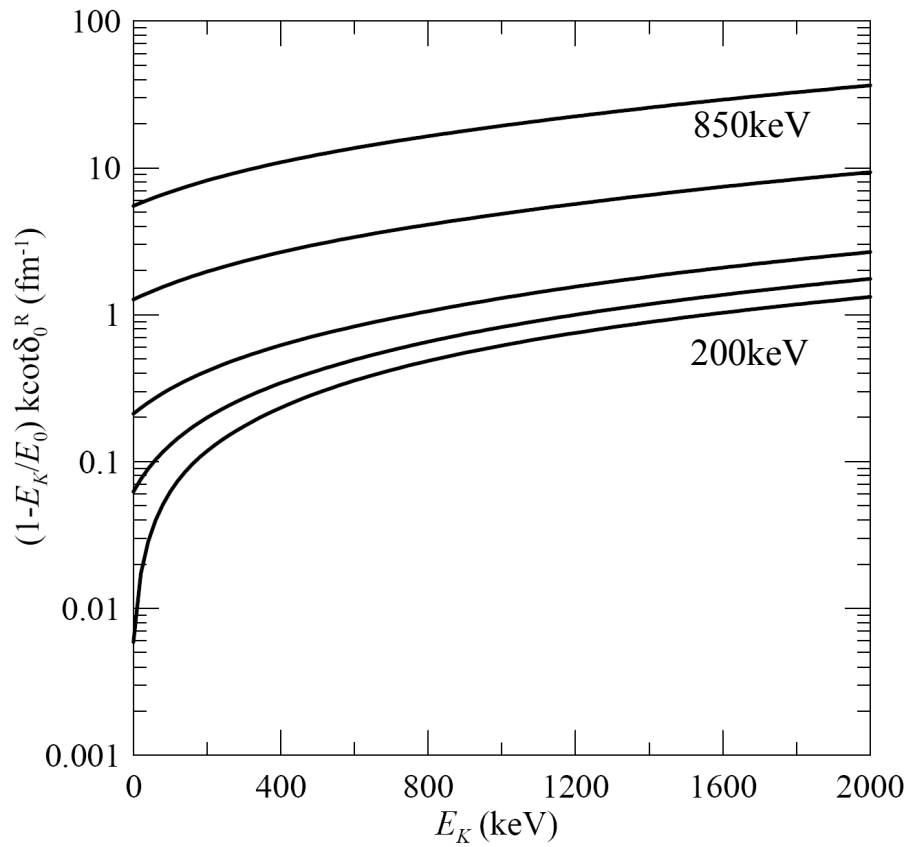
$$\mathcal{V}^\ell(q, k; \mathcal{E}) \equiv \pi \frac{(A+1)}{A+2} \left[K_2^\ell(q, k; \mathcal{E}) + \int_0^\infty dk' k'^2 K_1^\ell(q, k'; \mathcal{E}) \tau_{nn}(k'; \mathcal{E}) K_1^\ell(k, k'; \mathcal{E}) \right]$$

$$\tau_{nn}(q; \mathcal{E}) \equiv \frac{-2}{\pi} \left[\sqrt{|\epsilon_{nn}|} + \sqrt{\frac{A+2}{4A} q^2 - \mathcal{E}} \right]^{-1}, \quad \bar{\tau}_{nc}(q; \mathcal{E}) \equiv \frac{-1}{\pi} \left(\frac{A+1}{2A} \right)^{\frac{3}{2}} \left(\sqrt{|\epsilon_{nc}|} + \sqrt{\frac{(A+2)q^2}{2(A+1)} - \mathcal{E}} \right),$$

$$K_{i=1,2}^\ell(q, k; \mathcal{E}) \equiv G_i^\ell(q, k; \mathcal{E}) - \delta_{\ell 0} G_i^\ell(q, k; -\mu^2),$$

$$G_i^\ell(q, k; \mathcal{E}) = \int_{-1}^1 dy \frac{P_\ell(y)}{\mathcal{E} - \frac{A+1}{A+A^{i-1}} q^2 - \frac{A+1}{2A} k^2 - \frac{kqy}{A^{i-1}} + i\epsilon}.$$

n - ^{19}C scattering and Efimov physics



$$k \cot \delta_0^R = \frac{-a_{n-^{19}\text{C}}^{-1} + \beta E + \gamma E^2}{1 - E/E_0},$$

$ E_{19\text{C}} (\text{keV})$	$(a_{n-^{19}\text{C}})^{-1} (\text{fm}^{-1})$	$\beta (\text{fm} \cdot \text{keV})^{-1}$	$\gamma (\text{fm} \cdot \text{keV}^2)^{-1}$	$E_0 (\text{keV})$
200	$-0.591 \cdot 10^{-2}$	$5.685 \cdot 10^{-4}$	$4.673 \cdot 10^{-8}$	1442.745
400	$-0.624 \cdot 10^{-1}$	$6.743 \cdot 10^{-4}$	$8.821 \cdot 10^{-8}$	823.887
600	$-2.118 \cdot 10^{-1}$	$9.337 \cdot 10^{-4}$	$1.464 \cdot 10^{-7}$	451.398
800	-1.268	$3.11 \cdot 10^{-3}$	$4.424 \cdot 10^{-7}$	114.976
850	-5.510	$1.201 \cdot 10^{-2}$	$1.641 \cdot 10^{-6}$	28.845

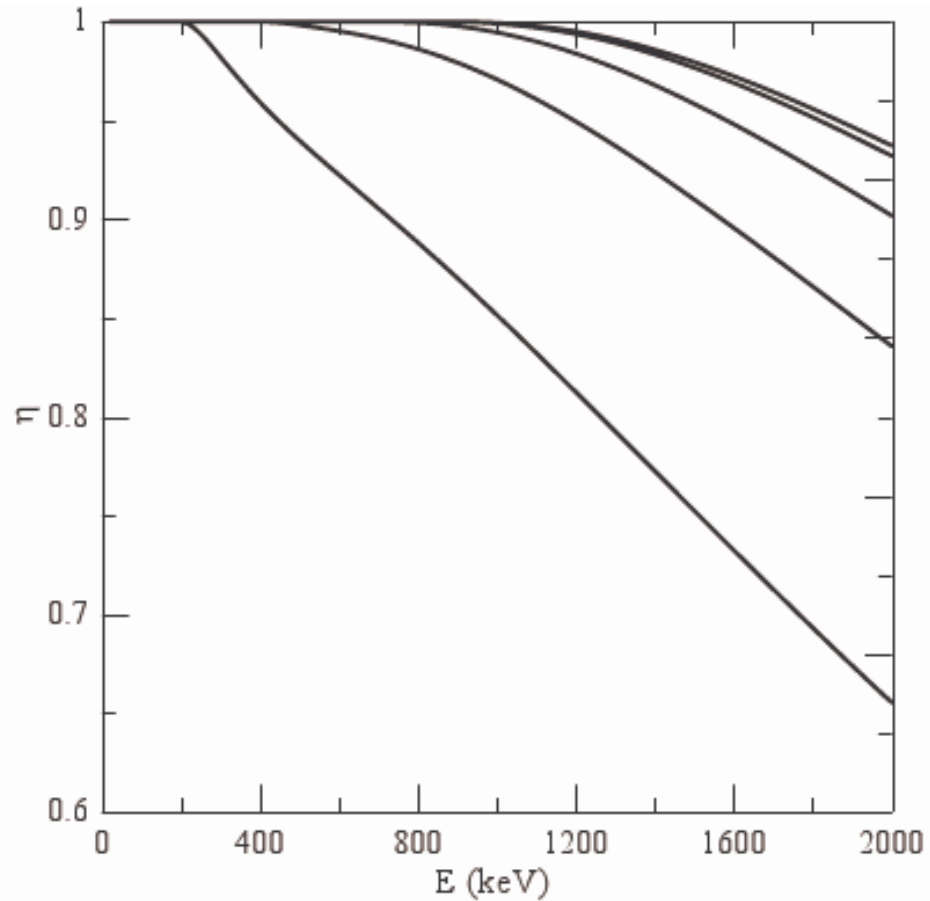
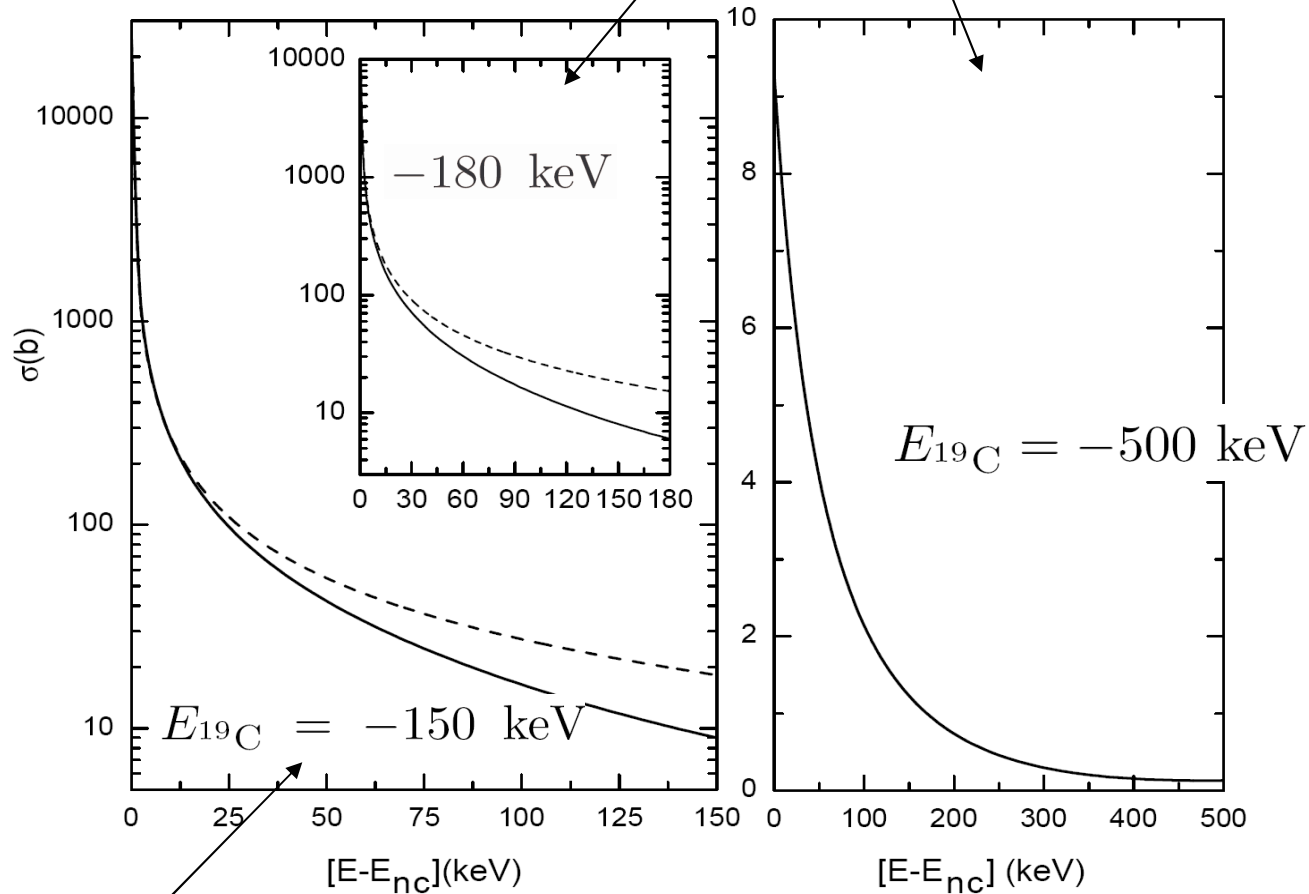
n - ^{19}C scattering and Efimov physics

Fig. 4. s -wave absorption parameter as a function of the CM kinetic energy. From left to right the curves corresponds to the following ^{19}C energies: 200, 400, 600, 800 and 850 keV.

n-¹⁹C scattering and Efimov physics

²⁰C has a virtual Efimov state



²⁰C has an excited bound Efimov state

n - ^{19}C scattering and Efimov physics

$$s_0 \sim 1.12 \text{ for } A = 18$$

$$a_0/a_B \simeq 0.42$$

Numerical solution of the scattering eq.s ~ 0.44

Summary and outlook

- ➔ Weakly bound & large systems: **few scales regime** in halo nuclei, molecules, trapped atoms
- ➔ Zero-range model: classification of weakly-bound systems
threshold conditions for excited states and resonances
(evidence of the trajectory of resonance in ultra-cold atoms)
- ➔ n-n-c systems:
Borromean configuration: **Efimov state → resonance**
All-bound, Samba, Tango (at least one subsystem is bound): **Efimov state → virtual state**
- ➔ ^{20}C **Efimov state → virtual state** $E_{19\text{C}} > 165 \text{ keV}$
- ➔ $n+^{19}\text{C}$ scattering: **pole in the s-wave phase-shift**
- ➔ Exploration of universality in scattering, breakup of halo nuclei and large molecules
- ➔ Resonances with different mass ratios in Borromean configurations?
- ➔ Trapped atoms with $a_{AD}=0$ decoupling of atom-dimer condensates? if $a_{AD}=\text{infinity}$?
- ➔ Four-boson excited states, resonances & scattering, formation of trimers in traps?
- ➔ Evidence for a four-boson scale?