

# Hamiltonian Limit of Lattice QED in 2+1 Dimensions

Christiane Groß

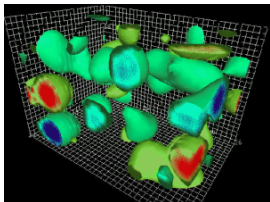
Helmholtz-Institut für Strahlen- und Kernphysik, Universität Bonn

4. July 2023

# Motivation

## Lagrangian

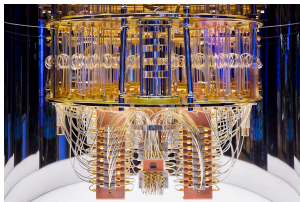
Wilson 1974 [1]  
discrete time  
classical computers



Deep Hybrid [2]

## Hamiltonian

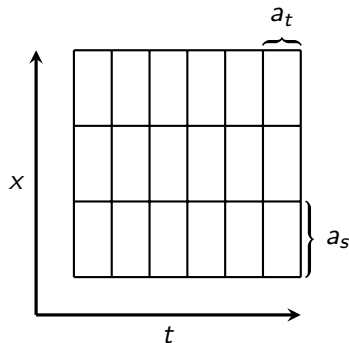
Kogut, Susskind 1975 [3]  
continuous time  
Exact Diagonalization  
Quantum Computers



Engadnet[4]

Match to test universality

# What is anisotropy?



adapted from [5], [6]

$$S = \sum_{r,i} \frac{\beta}{\xi_{\text{input}}} \text{Re Tr} (1 - P_{0i}(r)) \\ + \sum_{r,i>j>0} \beta \xi_{\text{input}} \text{Re Tr} (1 - P_{ij}(r))$$

$$\xi_{\text{ren}} = \frac{a_t}{a_s}$$

$$\xi_{\text{input}} \text{ set as } \frac{L}{T}$$

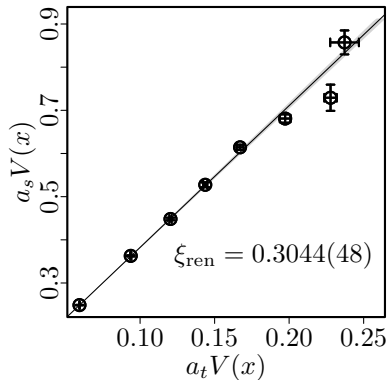
## How to determine anisotropy?

$$\lim_{x \rightarrow \infty} \frac{W_{ss}(x, y+1)}{W_{ss}(x, y)} = \exp(-a_s V_s(x[a_s]))$$

$$\lim_{t \rightarrow \infty} \frac{W_{st}(x, t+1)}{W_{st}(x, t)} = \exp(-a_t V_s(x[a_s]))$$

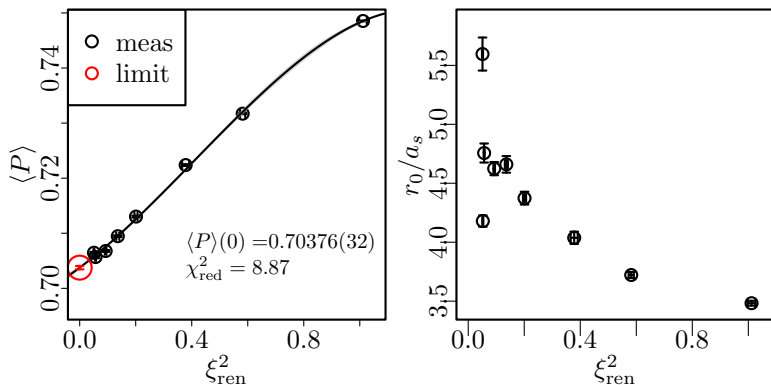
$$a_s V(x[a_s]) = \frac{1}{\xi} a_t V(x[a_s]) + se'$$

adapted from [5]



$$L = 16, T = 48, \beta = 1.7$$

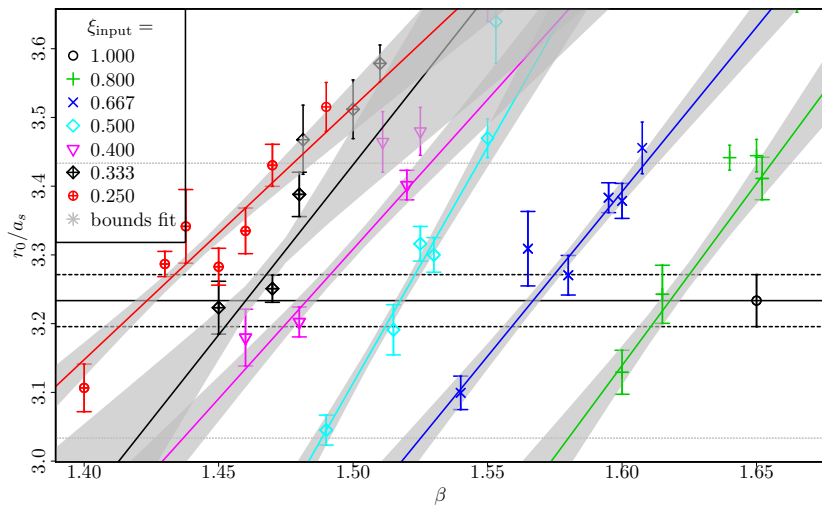
# Scale setting



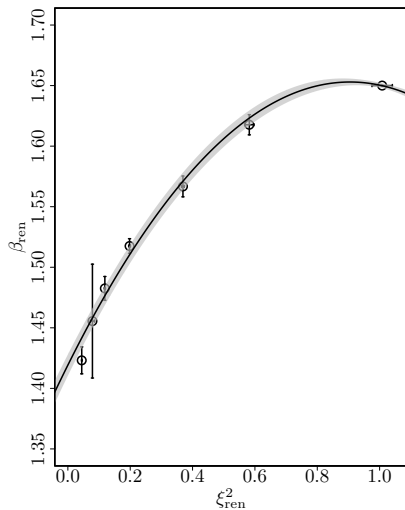
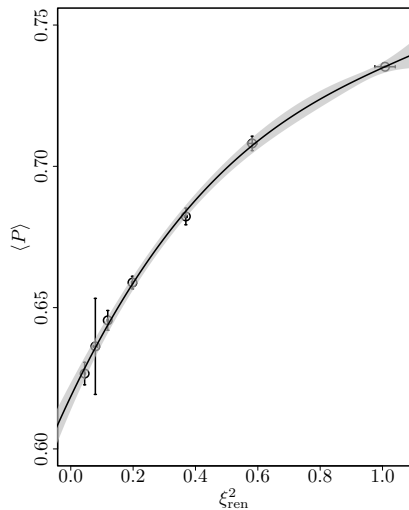
$$L = 16, \beta = 1.7$$

$$\langle P \rangle(\xi_{\text{ren}}^2) = a_0 + a_1 \xi_{\text{ren}}^2 + a_2 \xi_{\text{ren}}^4 + a_3 \xi_{\text{ren}}^6$$

# Scale setting



# Continuum limit



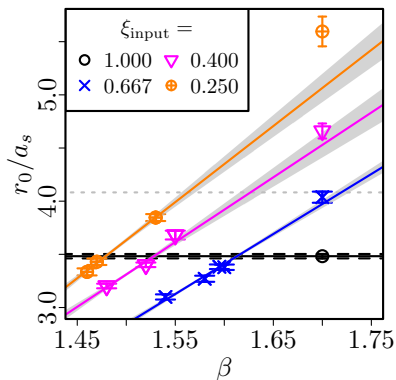
# Conclusion

## Done

Determined anisotropy  
nonperturbative  
renormalization of  $\xi$ ,  $\beta$   
2 continuum limits in  $a_t$   
Published proceeding[7]

## To Do

Finite Volume





# Bibliography I

- [1] Kenneth G. Wilson. 'Confinement of quarks'. In: *Phys. Rev. D* 10 (8 Oct. 1974), pp. 2445–2459. DOI: 10.1103/PhysRevD.10.2445. URL: <https://link.aps.org/doi/10.1103/PhysRevD.10.2445>.
- [2] *Deep hybrid use cases*. URL: <https://deep-hybrid-datacloud.eu/use-cases/>.
- [3] John Kogut and Leonard Susskind. 'Hamiltonian formulation of Wilson's lattice gauge theories'. In: *Phys. Rev. D* 11 (2 Jan. 1975), pp. 395–408. DOI: 10.1103/PhysRevD.11.395. URL: <https://link.aps.org/doi/10.1103/PhysRevD.11.395>.
- [4] *IBM Quantum computer*. URL: <https://www.engadget.com/ibm-quantum-computing-speedup-050134678.html>.

## Bibliography II

- [5] Mushtaq Loan and Chris Hamer. 'Hamiltonian study of improved  $U(1)$  lattice gauge theory in three dimensions'. In: *Phys. Rev. D* 70 (1 July 2004), p. 014504. DOI: 10.1103/PhysRevD.70.014504. URL: <https://link.aps.org/doi/10.1103/PhysRevD.70.014504>.
- [6] Mushtaq Loan et al. 'Path integral Monte Carlo approach to the  $U(1)$  lattice gauge theory in 2+1 dimensions'. In: *Physical Review D* 68.3 (Aug. 2003). ISSN: 1089-4918. DOI: 10.1103/physrevd.68.034504. URL: <http://dx.doi.org/10.1103/PhysRevD.68.034504>.
- [7] L. Funcke et al. *Hamiltonian limit of lattice QED in 2+1 dimensions*. 2022. arXiv: 2212.09627 [hep-lat].

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- [8] Christof Gattringer and Christian B Lang. *Quantum chromodynamics on the lattice: an introductory presentation*. Lecture notes in physics ; 788. Berlin u.a.: Springer, 2010. DOI: 10.1007/978-3-642-01850-3. URL: [http://digitale-objekte.hbz-nrw.de/storage2/2019/05/04/file\\_79/8446751.pdf](http://digitale-objekte.hbz-nrw.de/storage2/2019/05/04/file_79/8446751.pdf).
- [9] Nicholas Metropolis et al. 'Equation of State Calculations by Fast Computing Machines'. In: *The Journal of Chemical Physics* 21.6 (1953), pp. 1087–1092. DOI: 10.1063/1.1699114. eprint: <https://doi.org/10.1063/1.1699114>. URL: <https://doi.org/10.1063/1.1699114>.

## Bibliography IV

- [10] Roger Horsley and Ulli Wolff. 'Weak coupling expansion of Wilson loops in compact QED'. In: *Physics Letters B* 105.4 (1981), pp. 290–296. ISSN: 0370-2693. DOI: [https://doi.org/10.1016/0370-2693\(81\)90891-1](https://doi.org/10.1016/0370-2693(81)90891-1). URL: <https://www.sciencedirect.com/science/article/pii/0370269381908911>.
- [11] R. Balian, J. M. Drouffe and C. Itzykson. 'Erratum: Gauge fields on a lattice.III.Strong-coupling expansions and transition points'. In: *Phys. Rev. D* 19 (8 Apr. 1979), pp. 2514–2515. DOI: [10.1103/PhysRevD.19.2514](https://doi.org/10.1103/PhysRevD.19.2514). URL: <https://link.aps.org/doi/10.1103/PhysRevD.19.2514>.

## Bibliography V

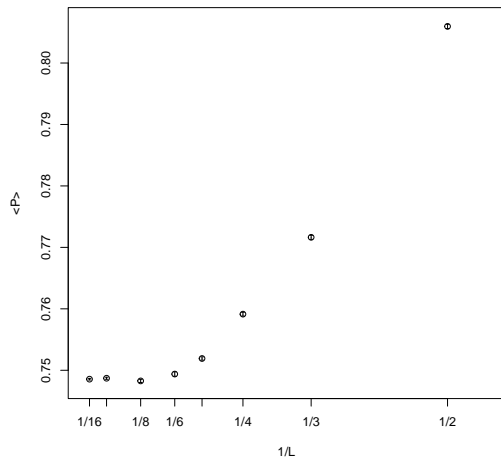
- [12] R. Sommer. 'A new way to set the energy scale in lattice gauge theories and its application to the static force and  $\sigma_s$  in SU (2) Yang-Mills theory'. In: *Nuclear Physics B* 411.2-3 (Jan. 1994), pp. 839–854. DOI: [10.1016/0550-3213\(94\)90473-1](https://doi.org/10.1016/0550-3213(94)90473-1). URL: <https://doi.org/10.1016%2F0550-3213%2894%2990473-1>.
- [13] Angus Kan et al. 'Investigating a (3+1)D topological  $\theta$ -term in the Hamiltonian formulation of lattice gauge theories for quantum and classical simulations'. In: *Phys. Rev. D* 104.3 (2021), p. 034504. DOI: [10.1103/PhysRevD.104.034504](https://doi.org/10.1103/PhysRevD.104.034504). arXiv: 2105.06019 [hep-lat].
- [14] T. M. R. Byrnes et al. 'Hamiltonian limit of (3+1)-dimensional SU(3) lattice gauge theory on anisotropic lattices'. In: *Physical Review D* 69.7 (Apr. 2004). ISSN: 1550-2368. DOI: [10.1103/physrevd.69.074509](https://doi.org/10.1103/physrevd.69.074509). URL: <http://dx.doi.org/10.1103/PhysRevD.69.074509>.

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- [15] M. Alford et al. 'Measuring the aspect ratio renormalization of anisotropic-lattice gluons'. In: *Phys. Rev. D* 63 (7 Feb. 2001), p. 074501. DOI: [10.1103/PhysRevD.63.074501](https://doi.org/10.1103/PhysRevD.63.074501). URL: <https://link.aps.org/doi/10.1103/PhysRevD.63.074501>.

# Backup slides - finite volume effects

$\xi_{in} = 1$ ,  $\beta = 1.7$



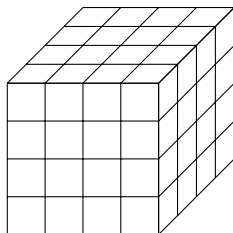
# Backup slides

Master thesis talk, October 2022



# Motivation

- Lattice: non-perturbatively
- Lagrangian formalism[1]: discrete time
- Hamiltonian formalism[2]: continuous time
- Test universality
- $U(1)$  in  $(2+1)$  dimensions



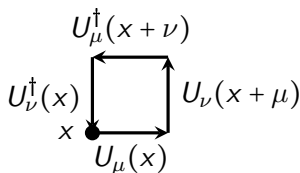
# Discretizing the Action

taken from [3]

$$S[A] = \frac{1}{2g^2} \int d^3x \operatorname{Tr} [F_{\mu\nu}(x) F_{\mu\nu}(x)]$$

$$U_\mu(n) = \exp(iaA_\mu(n))$$

$$P_{\mu\nu}(x) = U_\mu(x) U_\nu(x + \mu) U_\mu^\dagger(x + \nu) U_\nu^\dagger(x)$$



$$S = \beta \sum_r \sum_{\mu > \nu} \operatorname{Re} \operatorname{Tr} (1 - P_{\mu\nu}(r))$$

# Metropolis-Hastings

taken from [3]

$$\langle O \rangle = \frac{1}{Z} \int \mathcal{D}[U] O(U) e^{-S[U]}$$

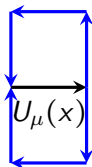
$$Z = \int \mathcal{D}[U] e^{-S[U]}$$

importance sampling:

$$\langle O \rangle = \frac{1}{N} \sum_{n=1}^N O[U_n]$$

transition probability[6]:

$$p = \min(1, \exp(-\Delta S))$$



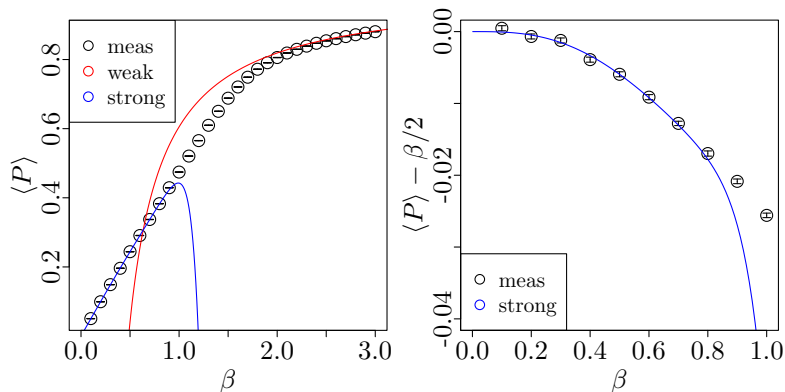
$$\Delta S / \beta$$

$$= \text{Re Tr} \left( (U_\mu(x) - U'_\mu(x)) \cdot \Gamma_\mu(x) \right)$$

$$U'_\mu(x) = U_\mu(x) \cdot R(\delta)$$

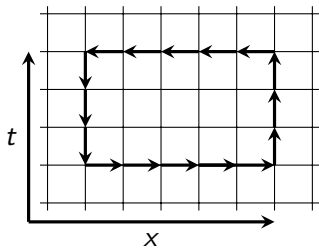
# Plaquette: Code check

Configurations with **SU(2)**-package, analysis with **hadron**-package



Predictions: [7], [8],  $L = 16$

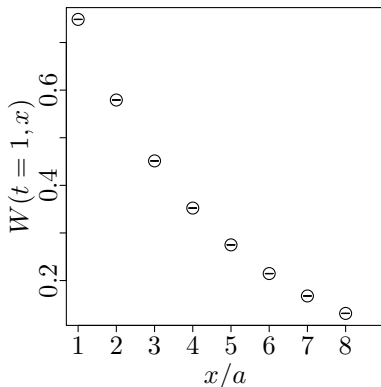
# Wilson Loops and Static Potential



$$W(t, r) = \sum_{k=0}^{\infty} C(r) \exp(-atE_k(r))$$

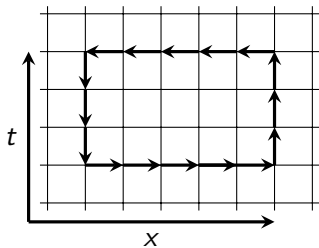
$$\lim_{t \rightarrow \infty} \frac{W(t+1, r)}{W(t, r)} = \exp(-aV(r[a]))$$

[3], [4]



$L = 16, \beta = 1.7$

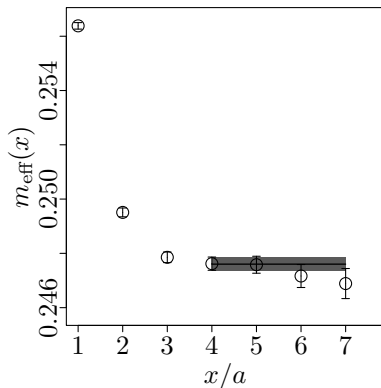
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[3], [4]



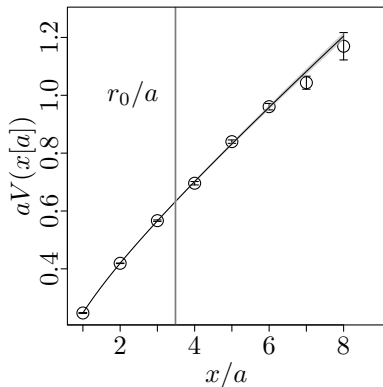
$L = 16, \beta = 1.7$

# Setting the scale

$$V(r) = d + \sigma \cdot r + b \cdot \ln(r)$$

$$-r^2 \frac{d}{dr} V(r)|_{r=r_0} = c = -1.65$$

[9], [5]



$$L = 16, \beta = 1.7$$

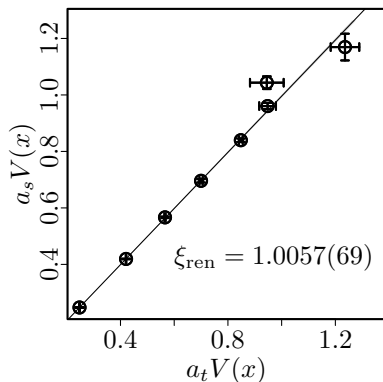
# Normal potentials

$$\lim_{x \rightarrow \infty} \frac{W_{ss}(x, y+1)}{W_{ss}(x, y)} = \exp(-a_s V_s(x[a_s]))$$

$$\lim_{t \rightarrow \infty} \frac{W_{st}(x, t+1)}{W_{st}(x, t)} = \exp(-a_t V_s(x[a_s]))$$

$$a_s V(x[a_s]) = \frac{1}{\xi} a_t V(x[a_s]) + se'$$

adapted from [4]



$$L = 16, T = 16, \beta = 1.7$$



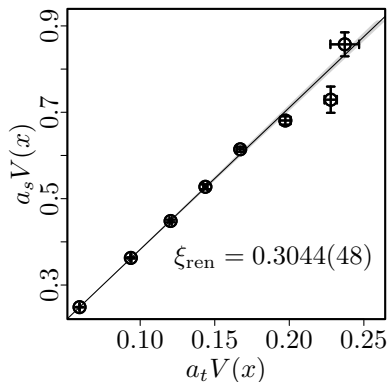
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$$a_s V(x[a_s]) = \frac{1}{\xi} a_t V(x[a_s]) + se'$$

adapted from [4]



$$L = 16, T = 48, \beta = 1.7$$

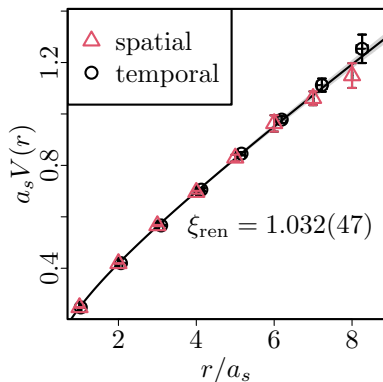
## Sideways potentials

$$\lim_{x \rightarrow \infty} \frac{W_{ss}(x+1, y)}{W_{ss}(x, y)} = \exp(-a_s V_s(y[a_s]))$$

$$\lim_{x \rightarrow \infty} \frac{W_{st}(x+1, t)}{W_{st}(x, t)} = \exp(-a_s V_s(t[a_t]))$$

$$V(y[a_s]) \stackrel{!}{=} V\left(\frac{1}{\xi_{\text{ren}}} t[a_t]\right)$$

adapted from [10]



$$L = 16, T = 16, \beta = 1.7$$

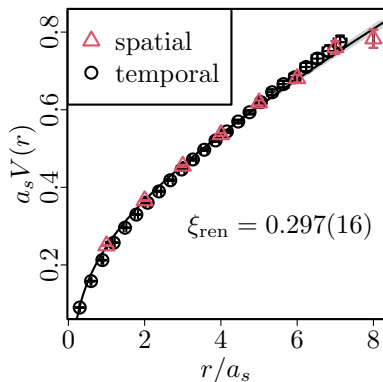
# Sideways potentials

$$\lim_{x \rightarrow \infty} \frac{W_{ss}(x+1, y)}{W_{ss}(x, y)} = \exp(-a_s V_s(y[a_s]))$$

$$\lim_{x \rightarrow \infty} \frac{W_{st}(x+1, t)}{W_{st}(x, t)} = \exp(-a_s V_s(t[a_t]))$$

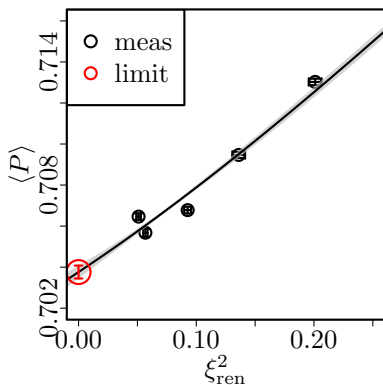
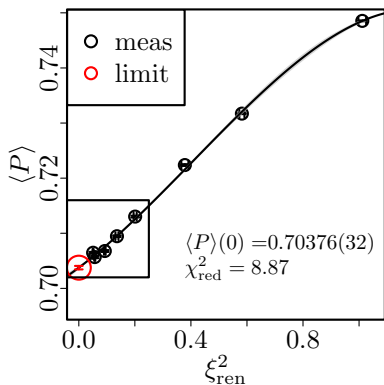
$$V(y[a_s]) \stackrel{!}{=} V\left(\frac{1}{\xi_{\text{ren}}} t[a_t]\right)$$

adapted from [10]



$$L = 16, T = 48, \beta = 1.7$$

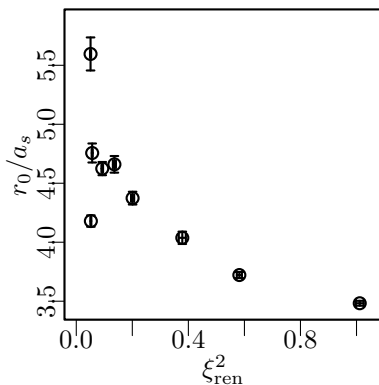
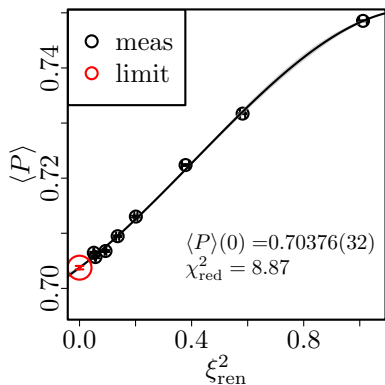
## Naive continuum limit



$$L = 16, \beta = 1.7$$

$$\langle P \rangle(\xi_{\text{ren}}^2) = a_0 + a_1 \xi_{\text{ren}}^2 + a_2 \xi_{\text{ren}}^4 + a_3 \xi_{\text{ren}}^6$$

## Naive continuum limit

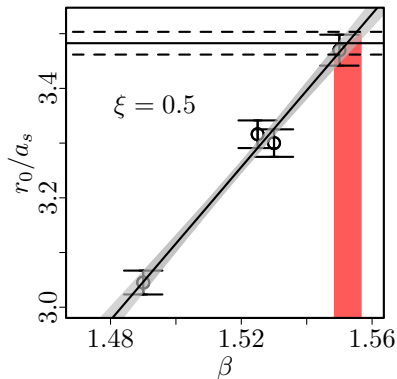


$$L = 16, \beta = 1.7$$

$$\langle P \rangle(\xi_{\text{ren}}^2) = a_0 + a_1 \xi_{\text{ren}}^2 + a_2 \xi_{\text{ren}}^4 + a_3 \xi_{\text{ren}}^6$$

# Renormalization of $\beta$

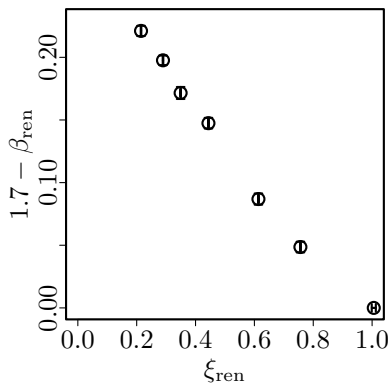
- $r_0/a_s$  for constant lattice spacing
- Linear interpolations of  $r_0/a_s(\beta)$ ,  $\langle P \rangle(\beta)$ ,  $\xi_{\text{ren}}(\beta)$
- $\beta_{\text{ren}} : r_0(\beta_{\text{ren}}) = r_0(\xi_{\text{input}} = 1)$



$$L = 16, T = 32$$

# Renormalization of $\beta$

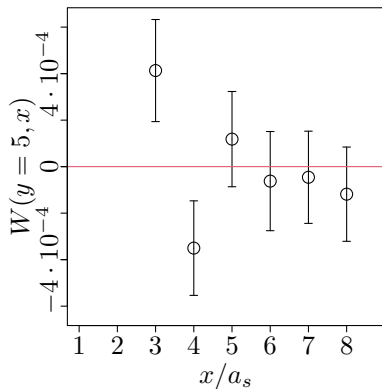
- $r_0/a_s$  for constant lattice spacing
- Linear interpolations of  $r_0/a_s(\beta)$ ,  $\langle P \rangle(\beta)$ ,  $\xi_{\text{ren}}(\beta)$
- $\beta_{\text{ren}} : r_0(\beta_{\text{ren}}) = r_0(\xi_{\text{input}} = 1)$
- Pairs  $(\xi_{\text{ren}}, \langle P \rangle)(\beta_{\text{ren}})$



$$L = 16$$

# Difficulties with this method

**Fluctuations at  
small  $\beta$**



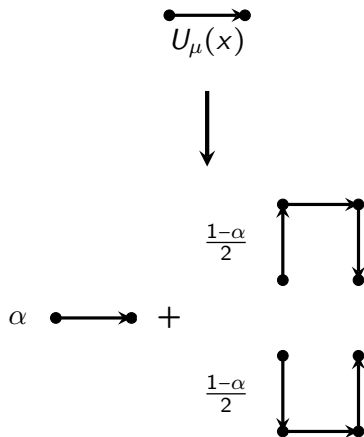
$$L = 16, T = 16, \beta = 1.0$$



# Difficulties with this method

## Fluctuations at small $\beta$

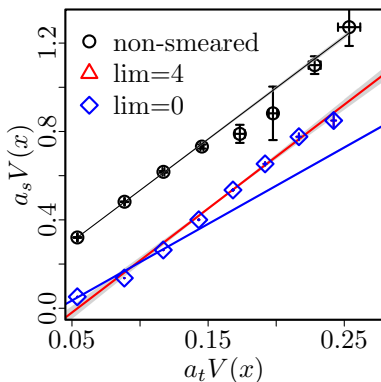
- smearing[5]



# Difficulties with this method

## Fluctuations at small $\beta$

- smearing[5]

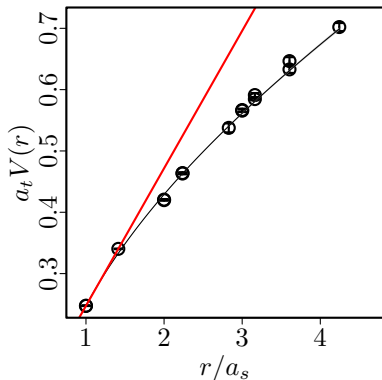


$$\beta = 1.48, n = 3, \alpha = 0.7$$

# Difficulties with this method

## Fluctuations at small $\beta$

- smearing[5]
- use non-integer distances



$$L = 16, T = 16, \beta = 1.7$$