

# Measurement of isolated photon-hadron correlations in Pb-Pb collisions at $\sqrt{s_{NN}}$ = 5.02 TeV with the ALICE experiment at LHC

Run 2 data: 2015-2018

Séminaires doctorants - Carolina Arata 27/03/23

#### Thesis supervisors

Gustavo Conesa Balbastre (CNRS - LPSC) Julien Faivre (UGA - LPSC)





## Dense QCD matter: quark-gluon plasma

- The QCD phase diagram predicts a phase transition at high energy density from <u>hadronic matter</u> to a colour-deconfined medium, the <u>quark-gluon plasma</u> (QGP)
- WHY? Study the strong interaction
- HOW? Via ultra-relativistic heavy-ion (nuclei) collisions
- WHERE? LHC at CERN, RHIC USA

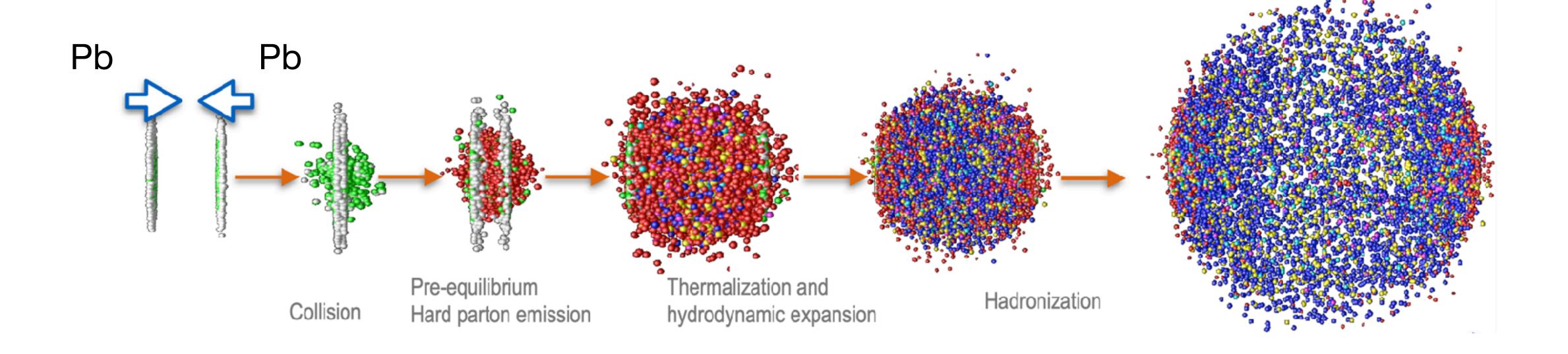
gh ined

150 MeV

RHIC

Atomic nuclei

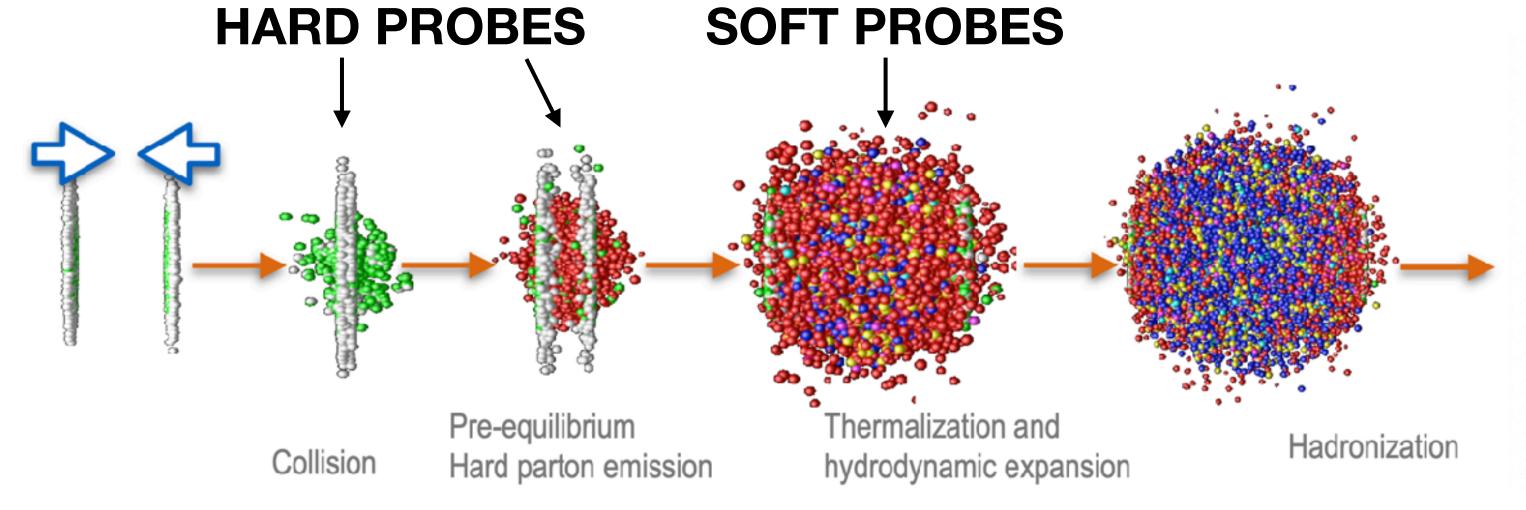
Baryon density

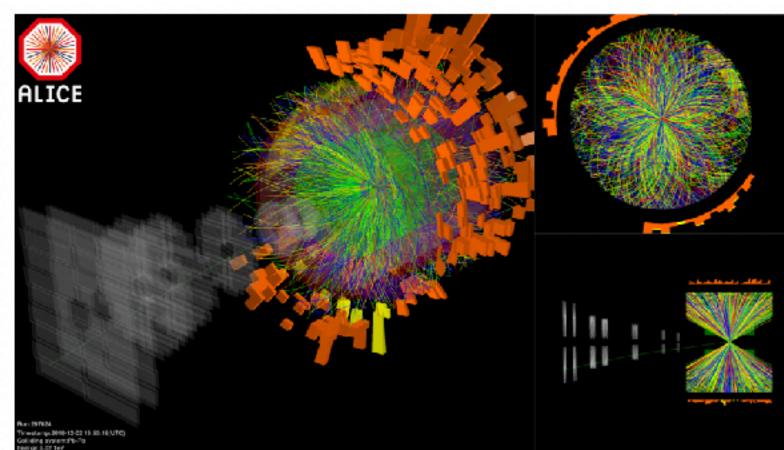


#### Probes for the QGP

- **SOFT PROBES**: low energy particles (*E* < 5 GeV)
  - to study thermodynamic and hydrodynamic properties and collectivity effects of the QGP
- **HARD PROBES**: High-energy partons (quarks and gluons) produced in the early stage of the collision to study how particles interact with the QGP medium and lose energy.

All particles not coming from the hard process are underlying event (UE) for the HARD PROBES

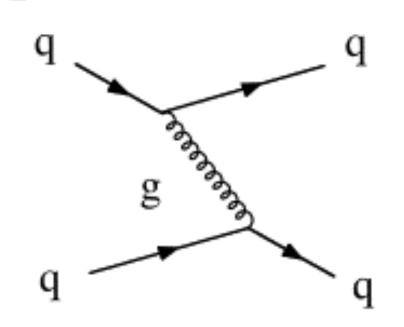




#### C. Arata

#### Hard Probes

In vacuum (collisions pp):
 partons fragment into a jet of particles



• In Pb-Pb collisions:

partons traverse the QGP,

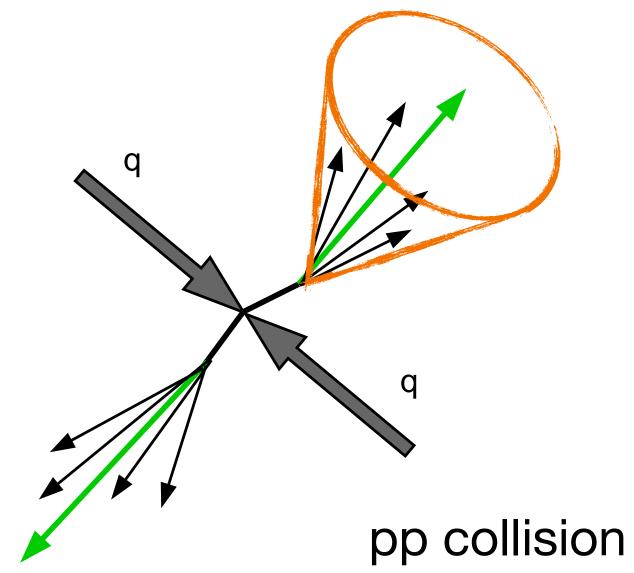
lose energy via *radiational* (gluon emission) + collisional processes, then fragment into a **jet of particles** 

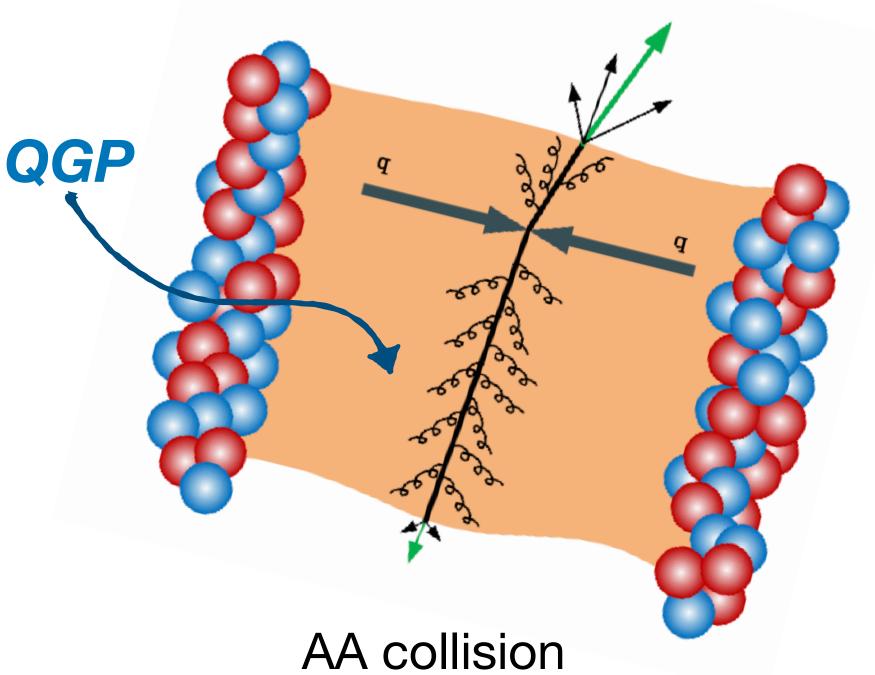
#### Loss of energy in medium = Jet Quenching

• Observable: fragmentation function  $D(z_{\mathrm{T}}) = \frac{1}{N^{trig}} \frac{dN^{ch}}{dz_{\mathrm{T}}}$  ,

where  $z_{\rm T}=p_{\rm T}^{hadrons}/p_{\rm T}^{parton}$  (  $p_{\rm T}=$  transverse momentum)

$$D(z_{\rm T})_{\rm AA} \neq D(z_{\rm T})_{\rm pp}$$



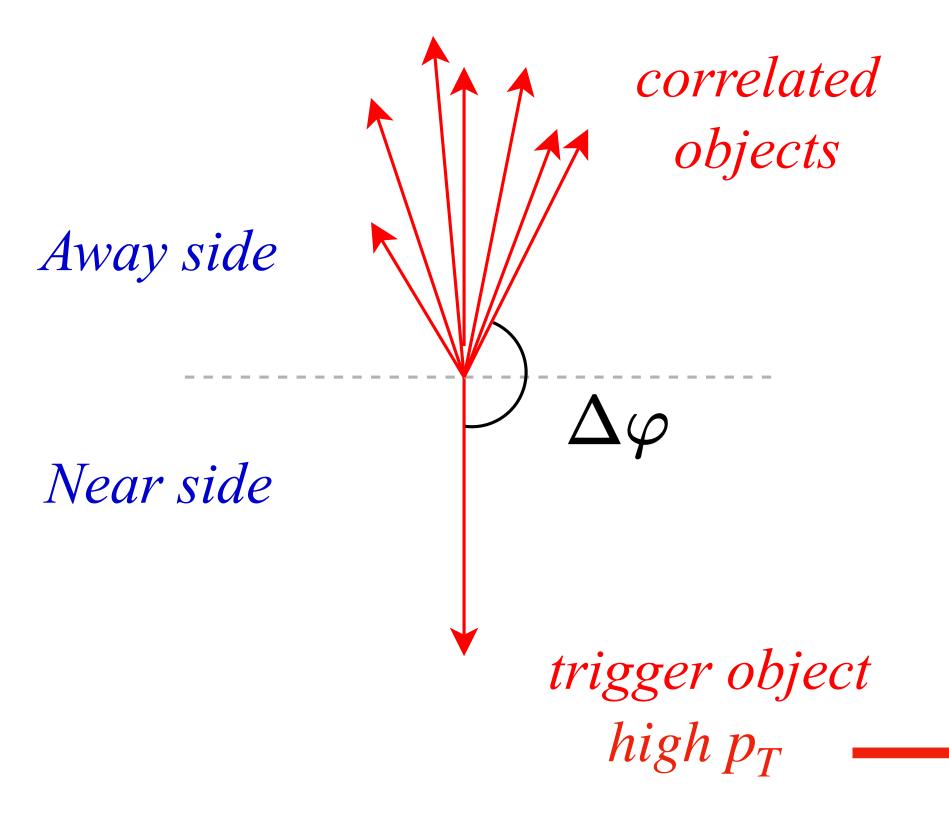


## How can we quantify jet quenching?

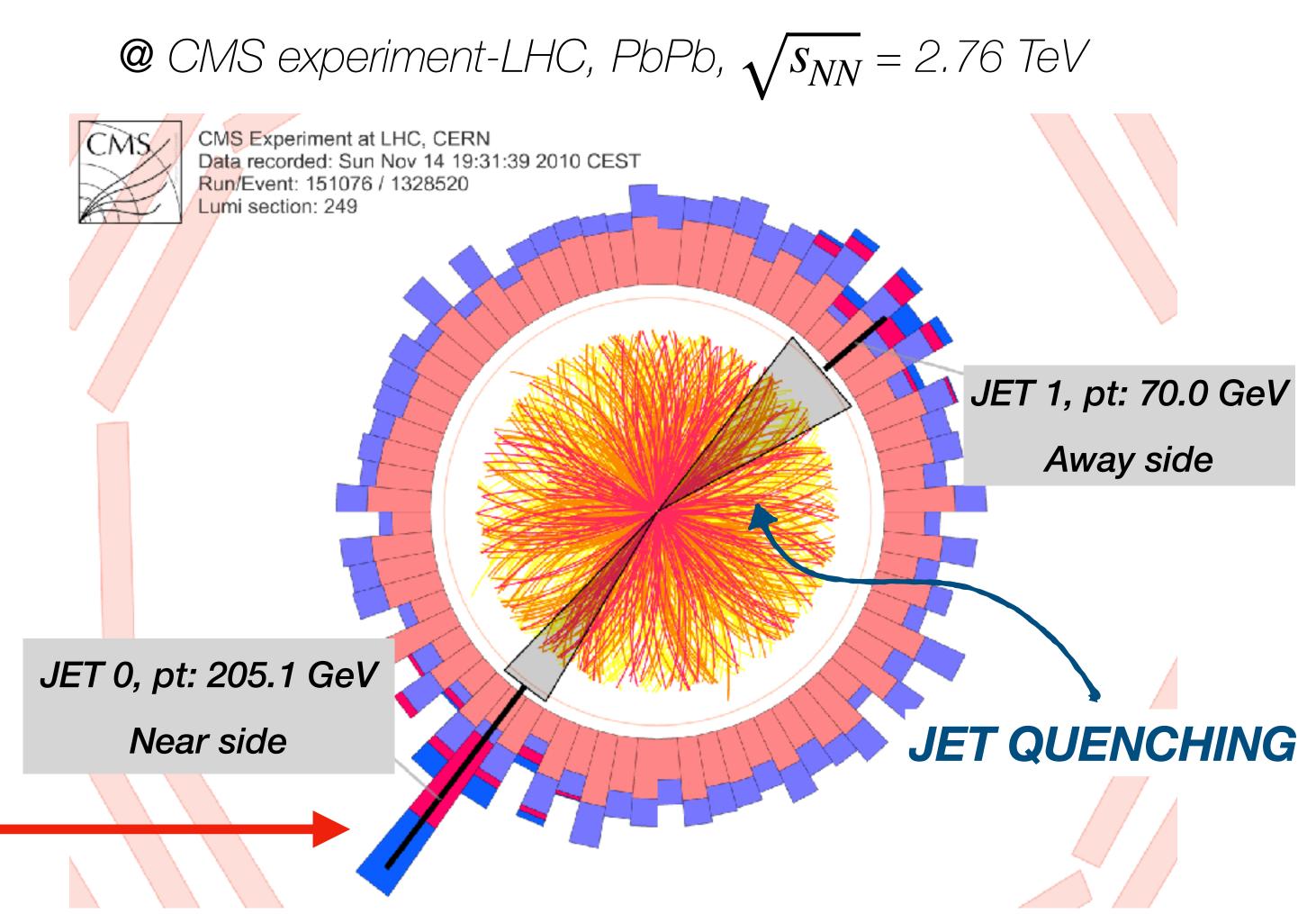
#### Azimuthal correlations distribution

between the trigger and associated particles

$$\Delta \varphi = (\varphi^{trig} - \varphi^{assoc})$$



#### Jet-jet correlations:

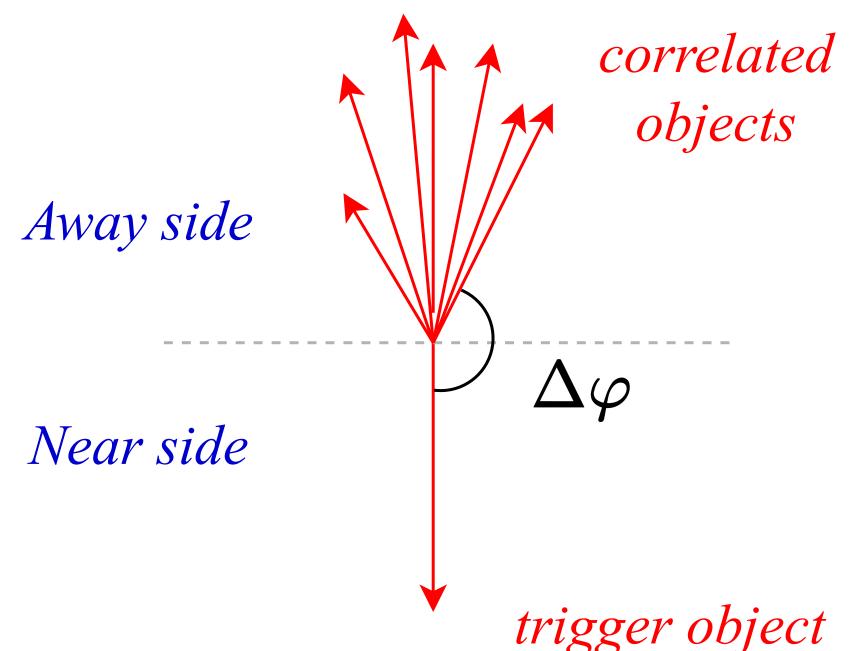


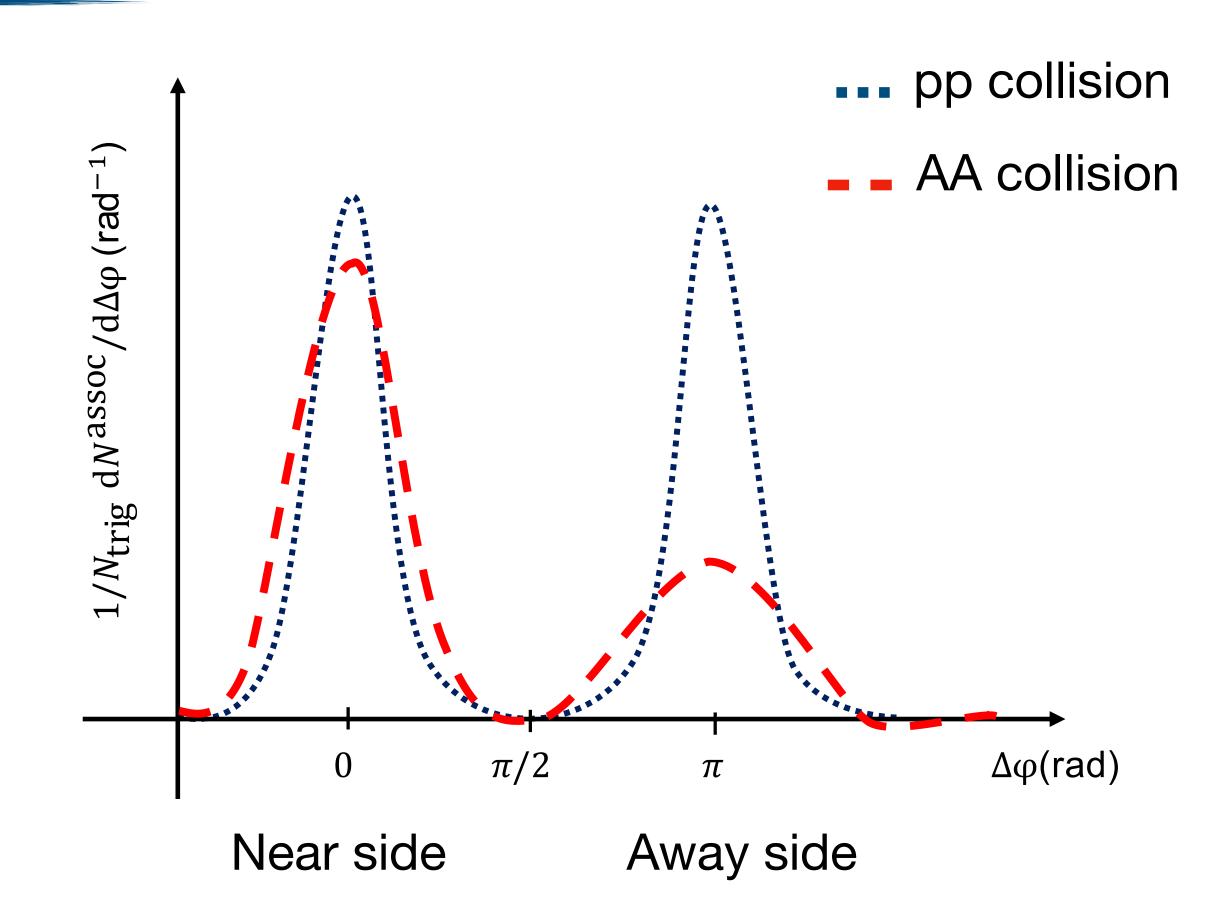
### How can we quantify jet quenching?

#### Azimuthal correlations distribution

between the trigger and associated particles

$$\Delta \varphi = (\varphi^{trig} - \varphi^{assoc})$$



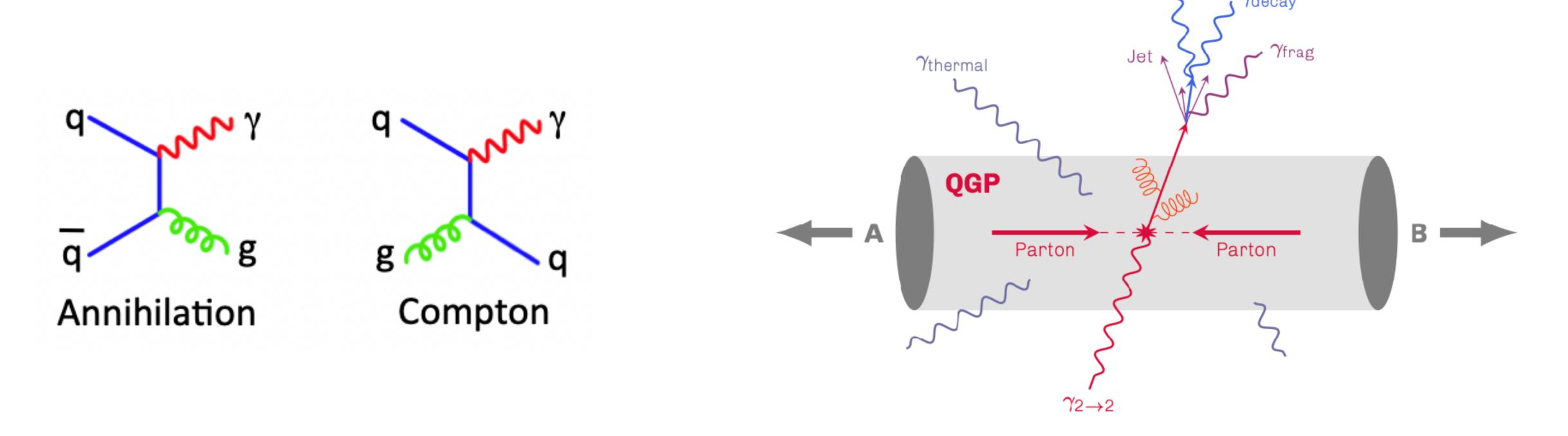


Trigger objects like jets or hadrons not ideal:

their parent partons can lose energy in the medium

$$p_T^{trigger} \neq p_T^{parton}$$
 BIASED REFERENCE

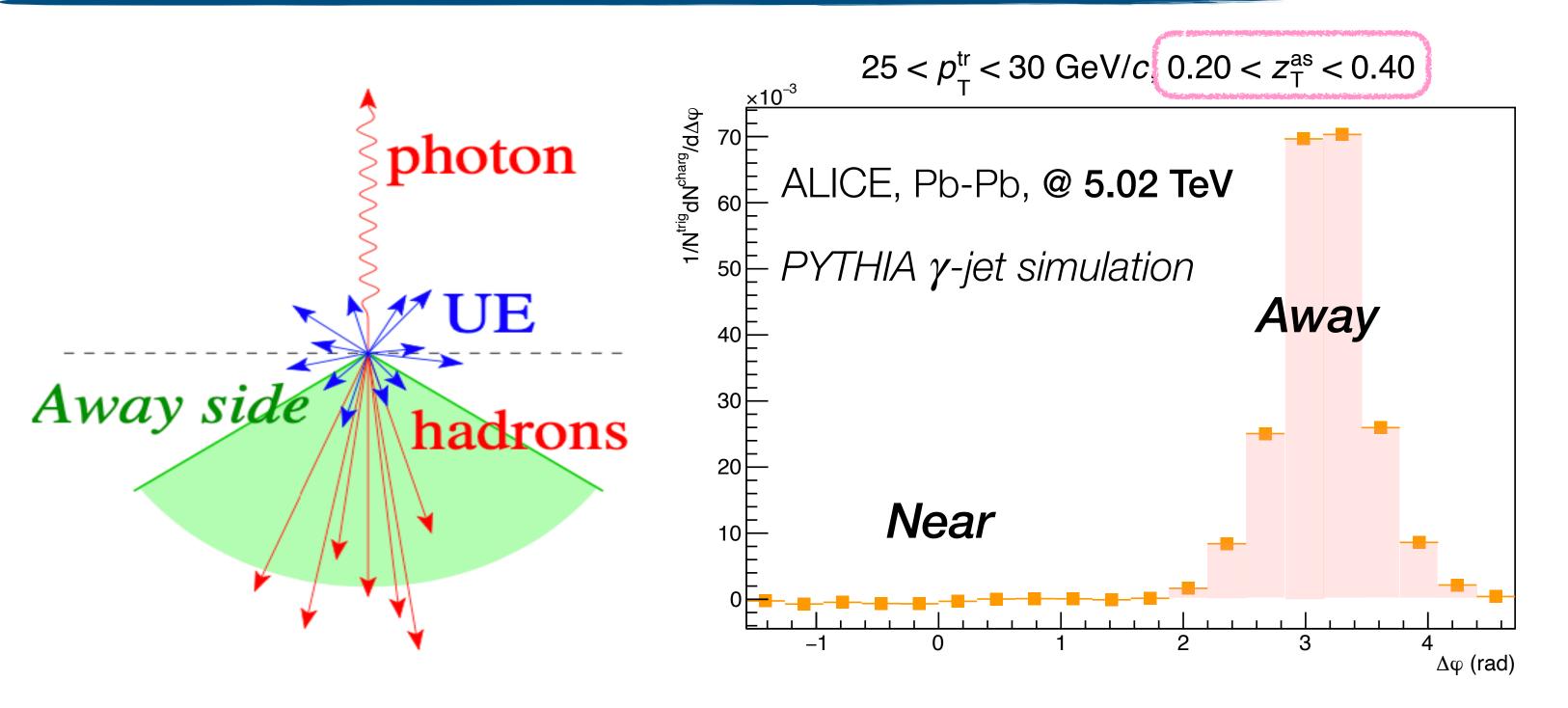
## Photons to study jet quenching



- Photons: neutral to colour, *not affected* by QCD medium → valuable probe
- $\gamma_{prompt}$  from 2o2 processes, *no hadrons* emitted in the direction of the photon o *ISOLATED*

 $\gamma_{2 \to 2}$  allow to tag the initial energy of the parton  $p_T^{\gamma} \approx p_T^{parton} = REFERENCE$ 

## γ<sub>2→2</sub> - hadron correlations

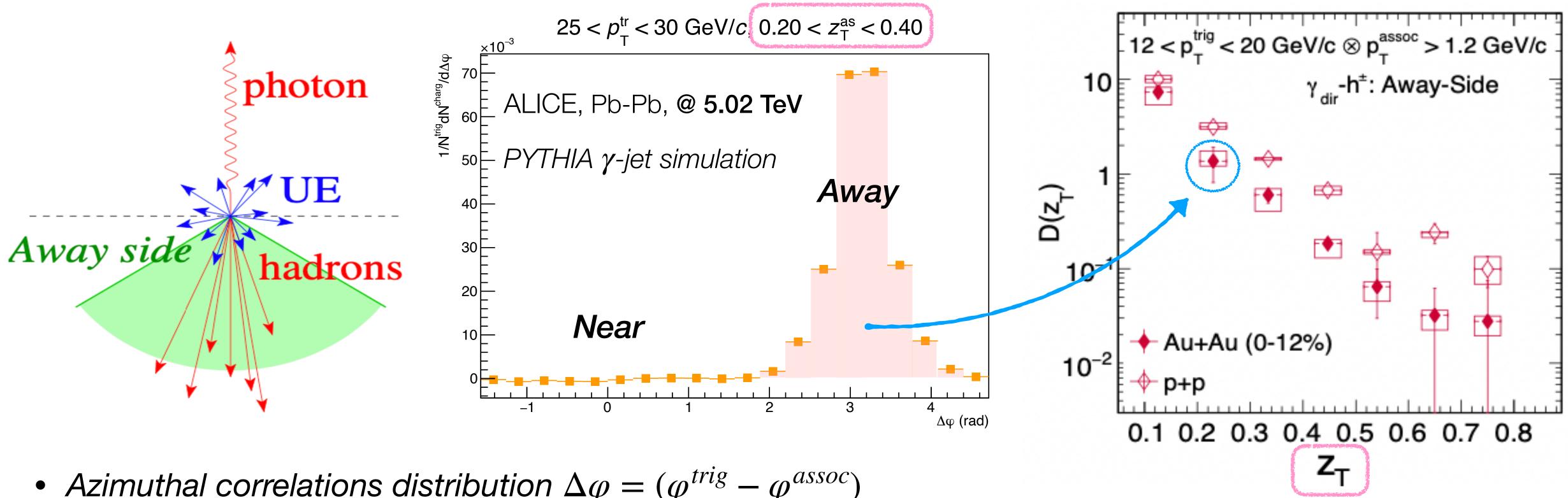


• Azimuthal correlations distribution  $\Delta \varphi = (\varphi^{trig} - \varphi^{assoc})$ 

• 
$$z_{\rm T}=p_{\rm T}^{hadr}/p_{\rm T}^{\gamma}$$
  $ightarrow$  Observable: the hadrons  $p_{\rm T}$  distribution  $D(z_{\rm T})=\frac{1}{N^{trig}}\frac{dN^{ch}}{dz_{\rm T}}$ 

STAR, @ 200 GeV, arXiv:1610.09568

## - hadron correlations

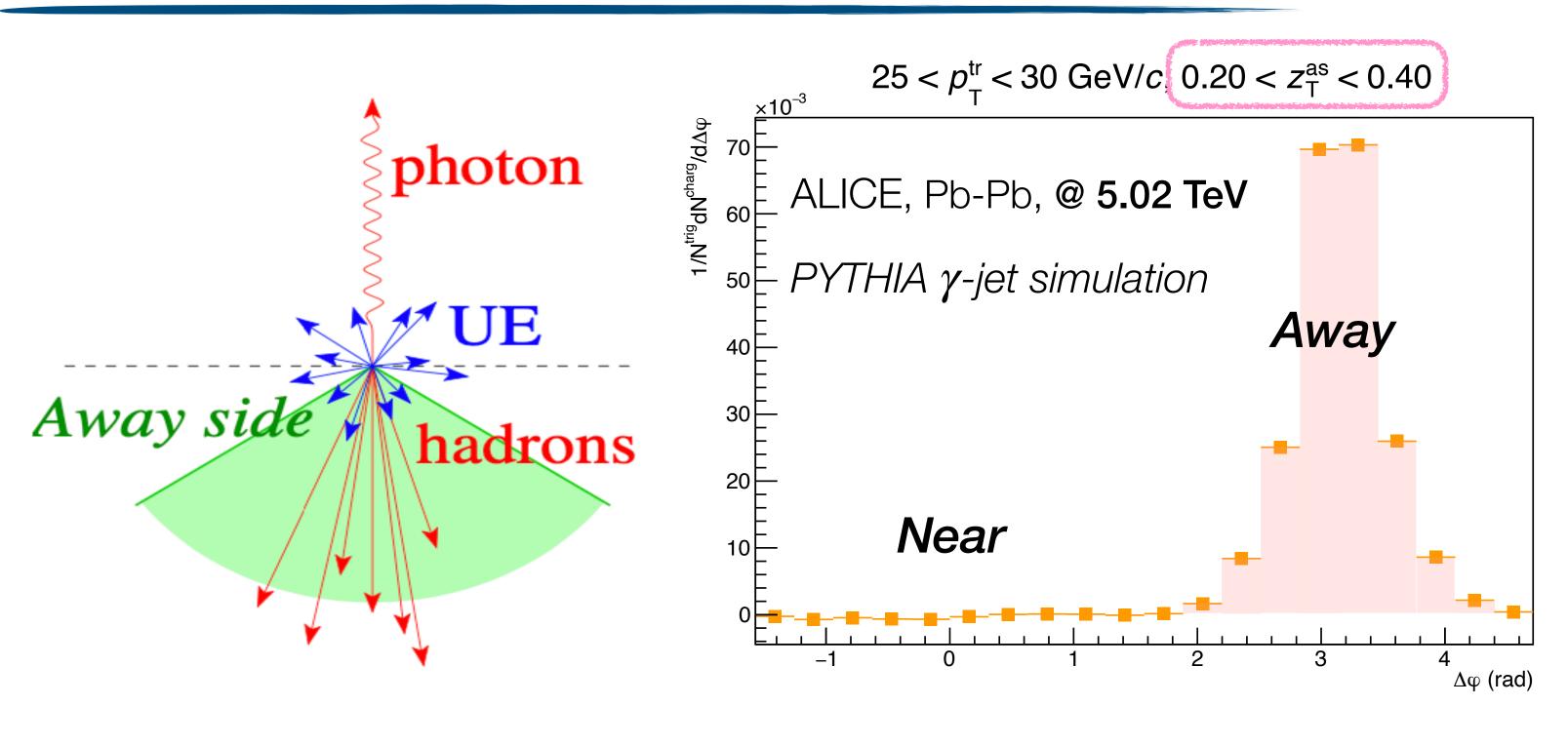


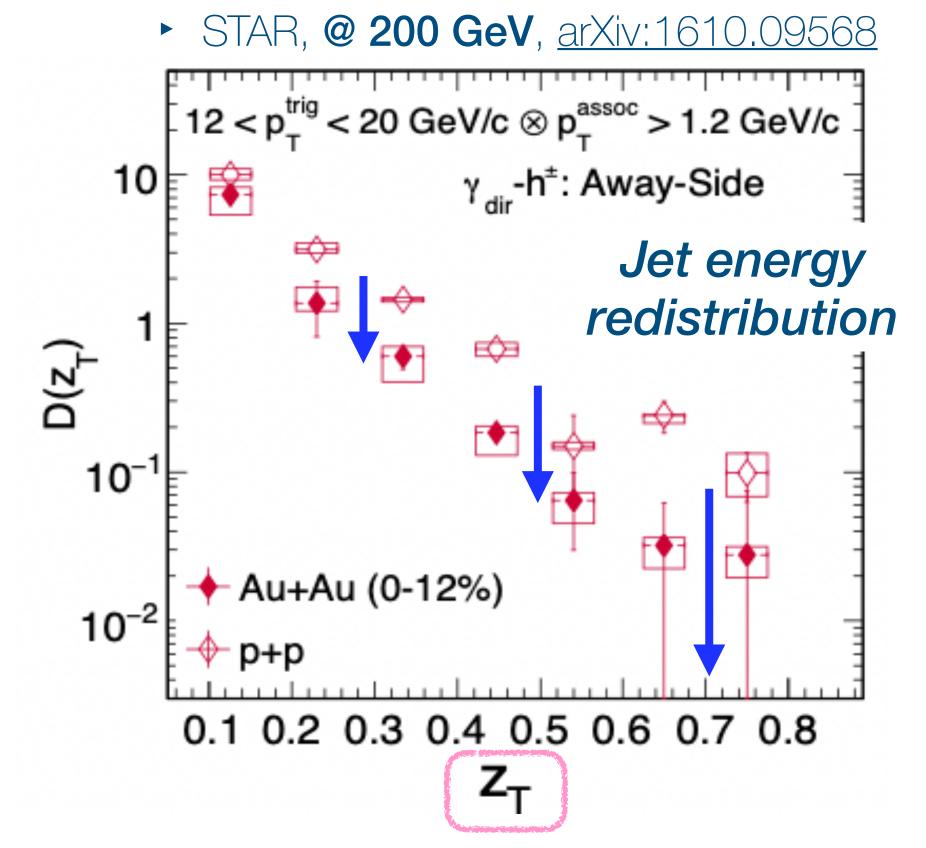
• Azimuthal correlations distribution 
$$\Delta \varphi = (\varphi^{trig} - \varphi^{assoc})$$

• 
$$z_{
m T}=p_{
m T}^{hadr}/p_{
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m T}$  distribution  $D(z_{
m T})=rac{1}{N^{trig}}rac{dN^{ch}}{dz_{
m T}}$ 

•  $D(z_{\rm T})$  is a proxy for the jet fragmentation function

## γ<sub>2→2</sub> - hadron correlations

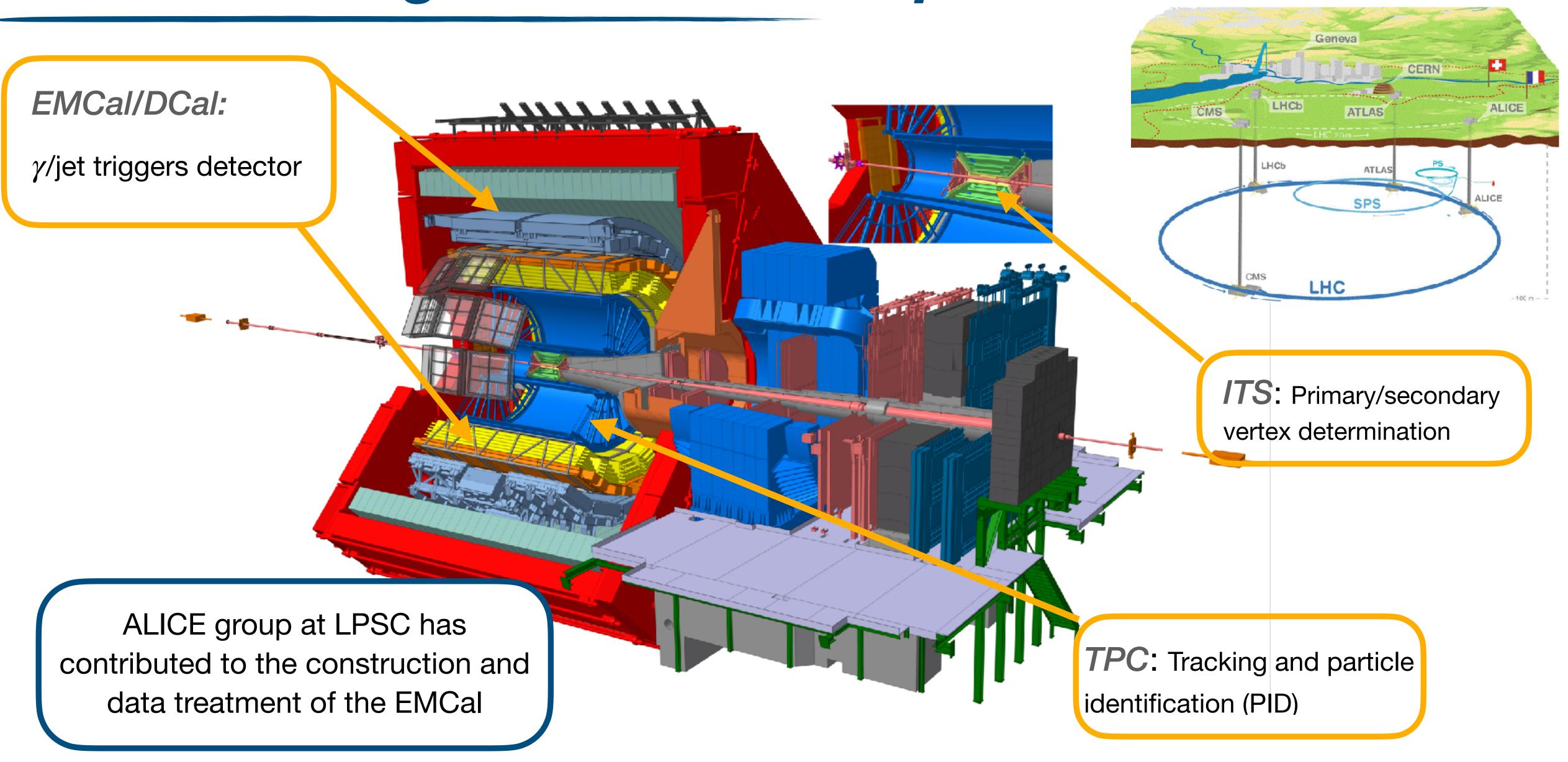




- Azimuthal correlations distribution  $\Delta \varphi = (\varphi^{trig} \varphi^{assoc})$
- $z_{
  m T} = p_{
  m T}^{hadr}/p_{
  m T}^{\gamma}$  ightarrow Observable: hadrons  $p_{
  m T}$  distribution  $D(z_{
  m T}) = \frac{1}{N^{trig}} \frac{dN^{ch}}{dz_{
  m T}}$
- $D(z_{\rm T})$  is a proxy for the jet fragmentation function  $\to D(z_{\rm T})_{\rm AA} \neq D(z_{\rm T})_{\rm pp}$

Distribution modified: information on parton energy loss (high  $p_{
m T}$  particle suppression)

#### ALICE - A Large Ion Collider Experiment



Calorimeter identification



Isolation

Y<sub>2→2</sub>
trigger
selection

Photon-hadron raw azimuthal distribution



Subtraction of the UE



Purity Correction



Systematic errors

Calorimeter identification



Isolation

γ<sub>2→2</sub>
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Photon-hadron raw azimuthal distribution



Subtraction of the UE



Purity Correction



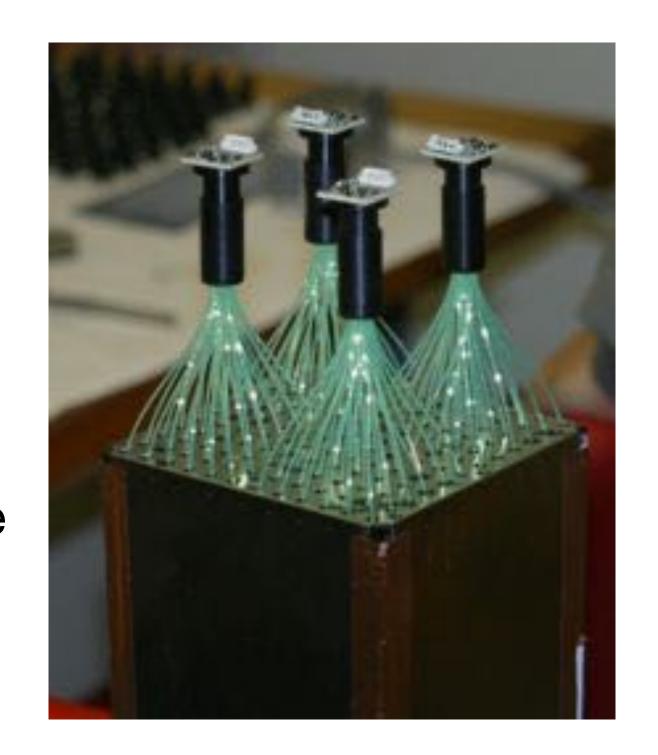
Systematic errors

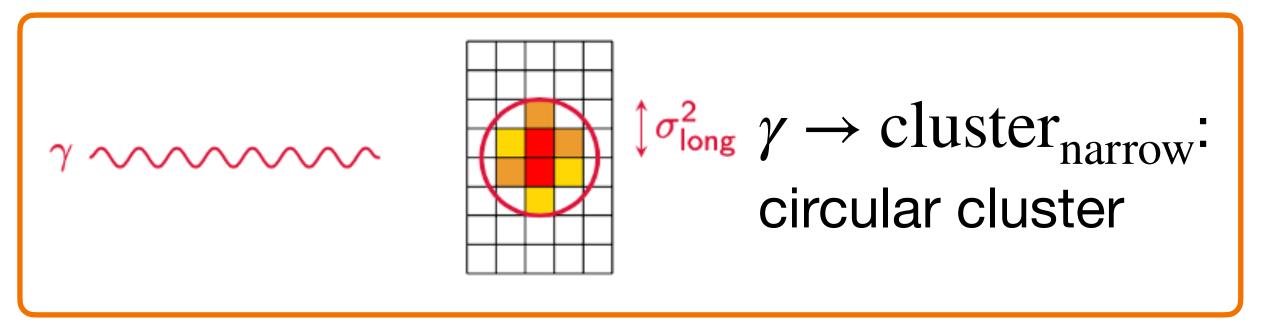
#### Photon identification with EMCal

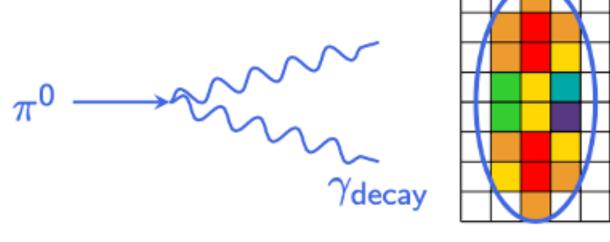
A particle interacting with the *cell material* produces a shower spreading its energy over *neighbouring cells*.

• Cluster: aggregate of cells

The **distribution of energy** within a cluster allows to discriminate between single photons  $\gamma$  shower and overlapping  $\gamma$  showers ( $\gamma_{decay}$ ) from high energy  $\pi^0 \rightarrow \gamma\gamma$ 



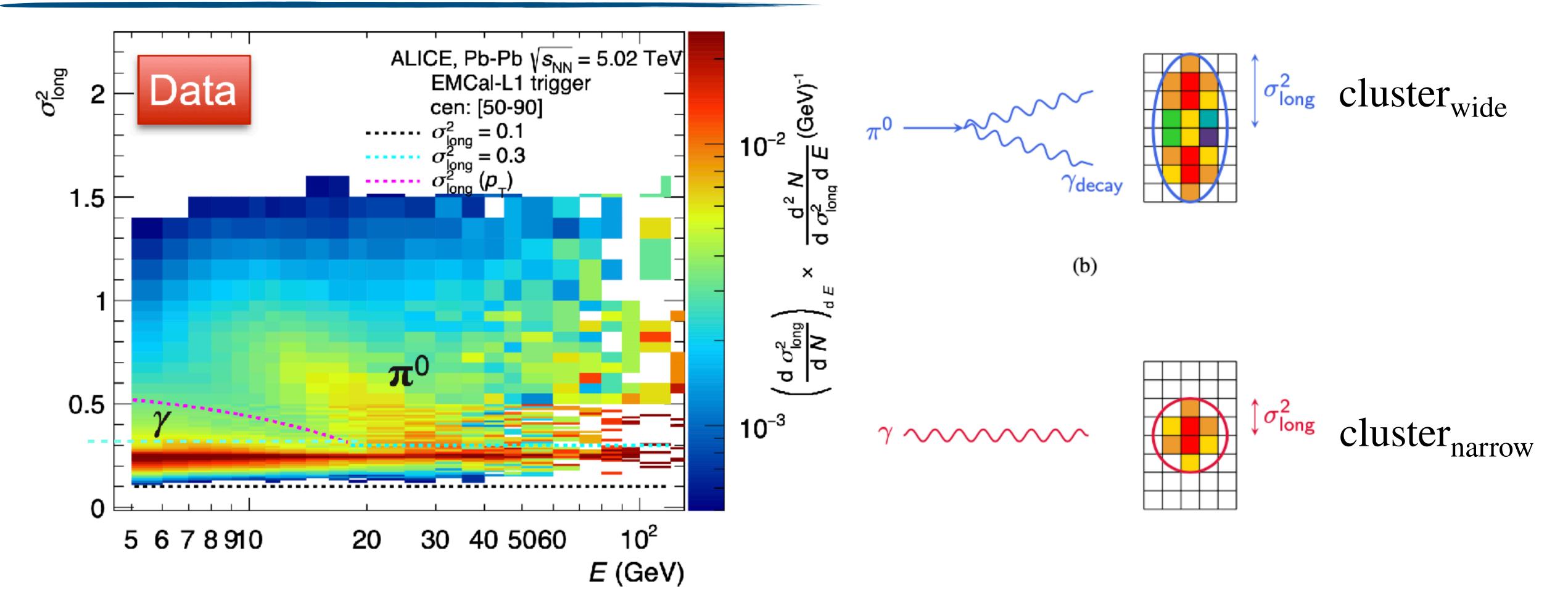




 $\sigma_{\text{long}}^2 \pi^0 \rightarrow \text{cluster}_{\text{wide}}$ :

elliptical cluster

## Photon identification: Cluster elongation

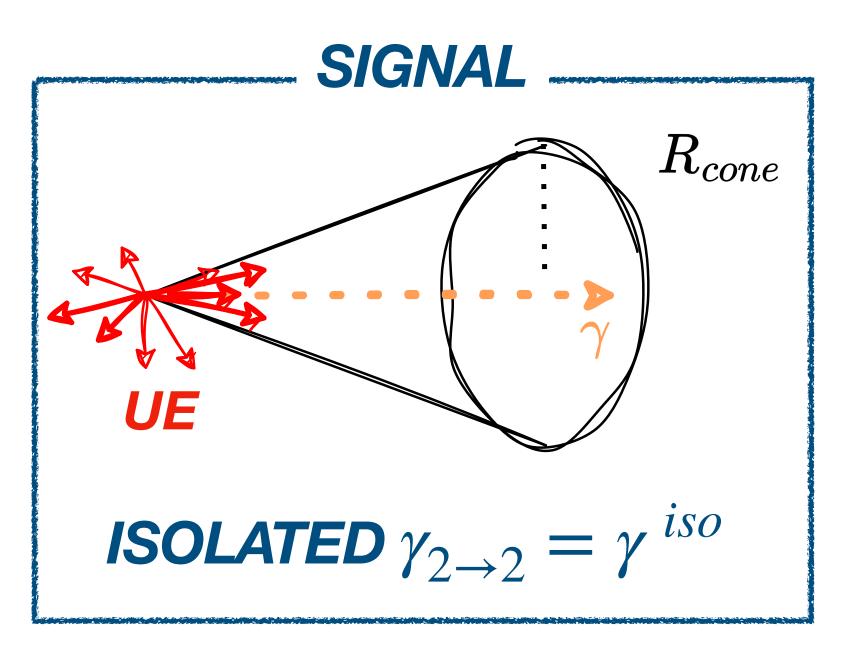


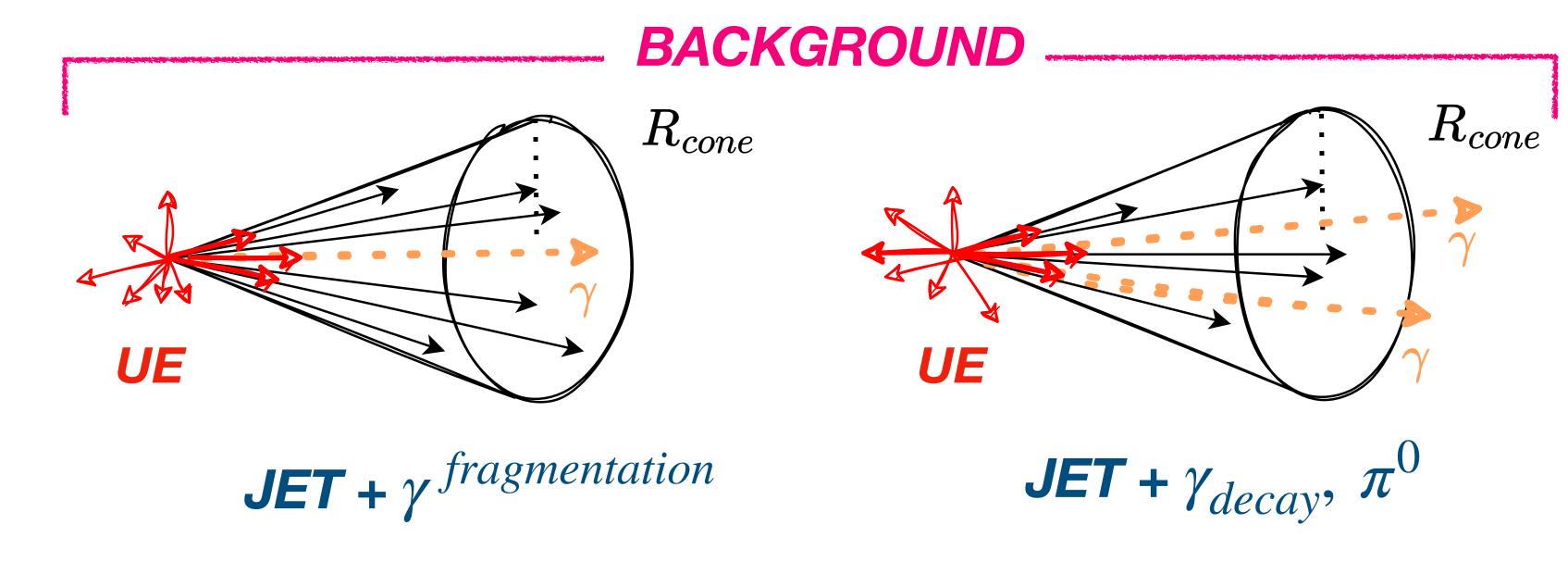
- $\gamma$  : population is mainly at  $\sigma_{long}^2 < 0.3$ ;
- $\pi^0$ : E < 30 GeV, population is mainly at  $\sigma_{long}^2 > 0.4$ ;
  - E > 30 GeV,  $\pi^0$  and  $\gamma$  bands overlap more and more

#### Photon isolation

Necessary to *reject* the largest part of *non*  $\gamma_{2\rightarrow 2}$  *photons*: *isolation criteria* 

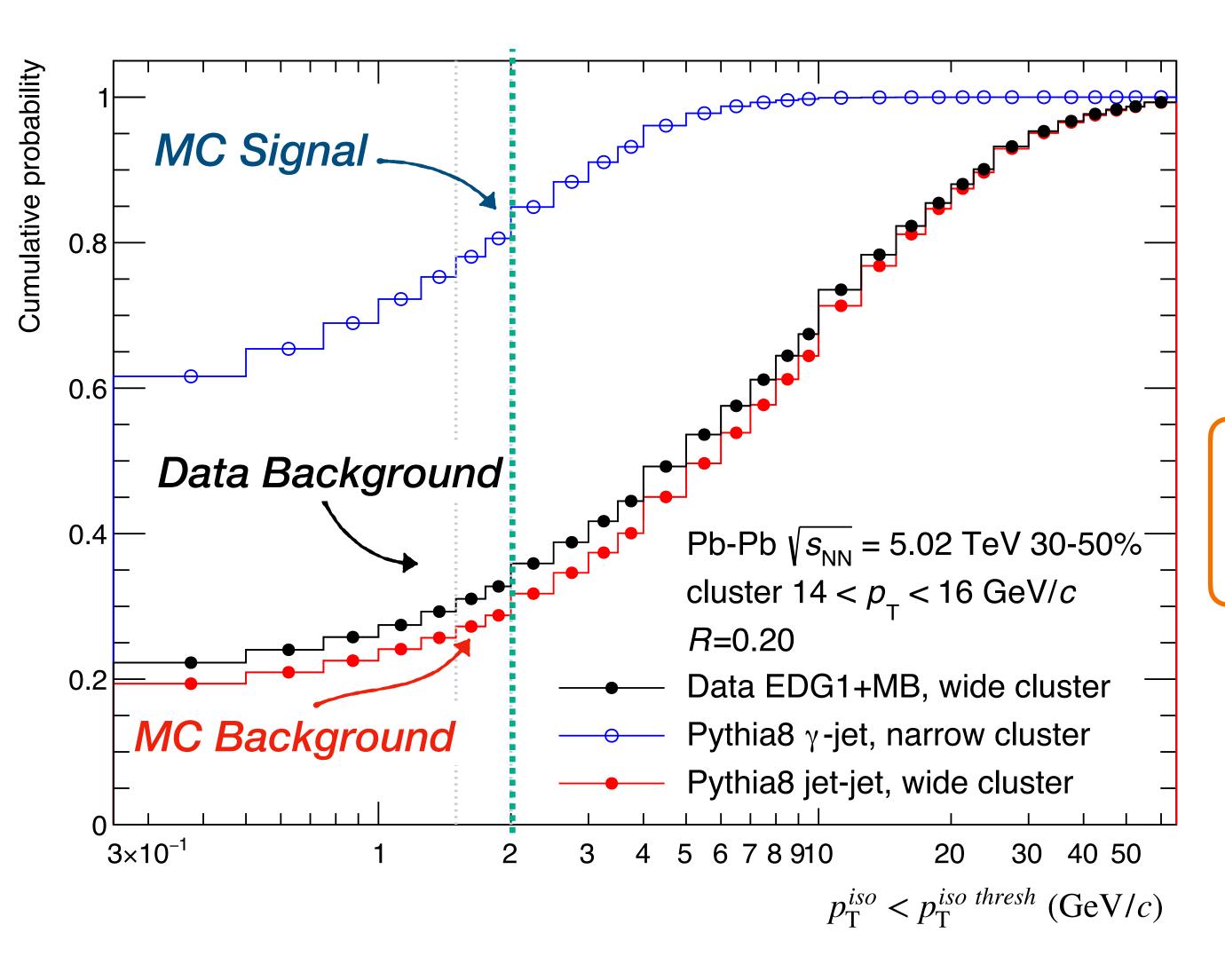
 $\gamma_{2\rightarrow 2}$  photon: Produced far from other particles (underlying event excepted)





- Define a cone radius around a candidate photon:  $R_{cone} = \sqrt{(\eta \eta_{\gamma})^2 (\varphi \varphi_{\gamma})^2} = 0.2$
- Condition on the total  $p_{\rm T}$  inside the cone:  $\sum_{cone} p_{\rm T}^i (\sum_{UE} p_{\rm T}^i)_{in\;cone} = p_{\rm T}^{iso} < 2~{\rm GeV}/c$ 
  - UE subtracted event by event, estimated in region far from our candidate gamma and jets

#### Isolation method



Check the *effectiveness* of the isolation cut, varying the threshold value

• 
$$p_{\mathrm{T}}^{iso} < p_{\mathrm{T}}^{iso \ thresh} = 2 \ \mathrm{GeV}/c$$

- ~ 85 % of the signal is accepted
- ~ 70 % of the background is rejected

Calorimeter identification



Isolation

Y<sub>2→2</sub>
trigger
selection

Photon-hadron raw azimuthal distribution



Subtraction of the UE



Purity Correction



Systematic errors

## Correlations analysis

Photon-hadron raw azimuthal distribution



Subtraction of the UE

---- Mixed Event



Purity Correction



Systematic errors

#### Mixed Event: removal of the underlying event

• Underlying event (UE): tracks uncorrelated to the hard process in a collision

**UE** produces a kind of pedestal on our azimuthal distribution

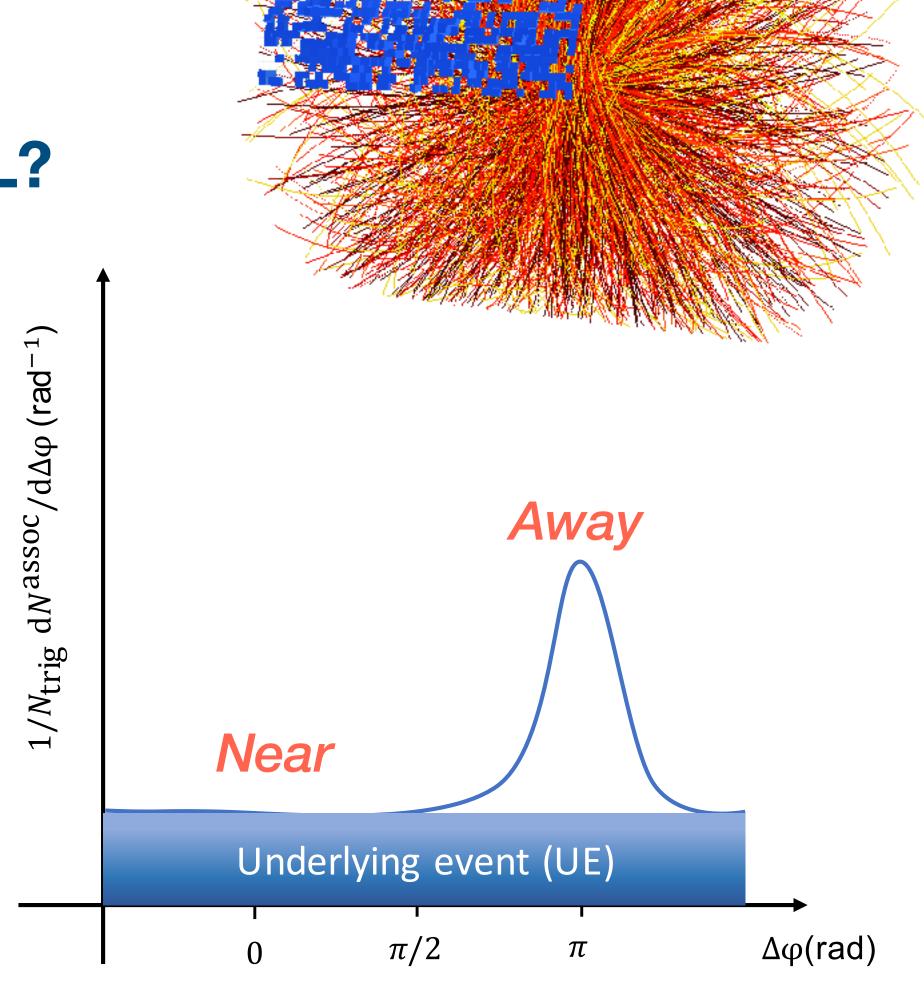
#### HOW TO REMOVE *UE* AND EXTRACT THE SIGNAL?

Mixed Event:

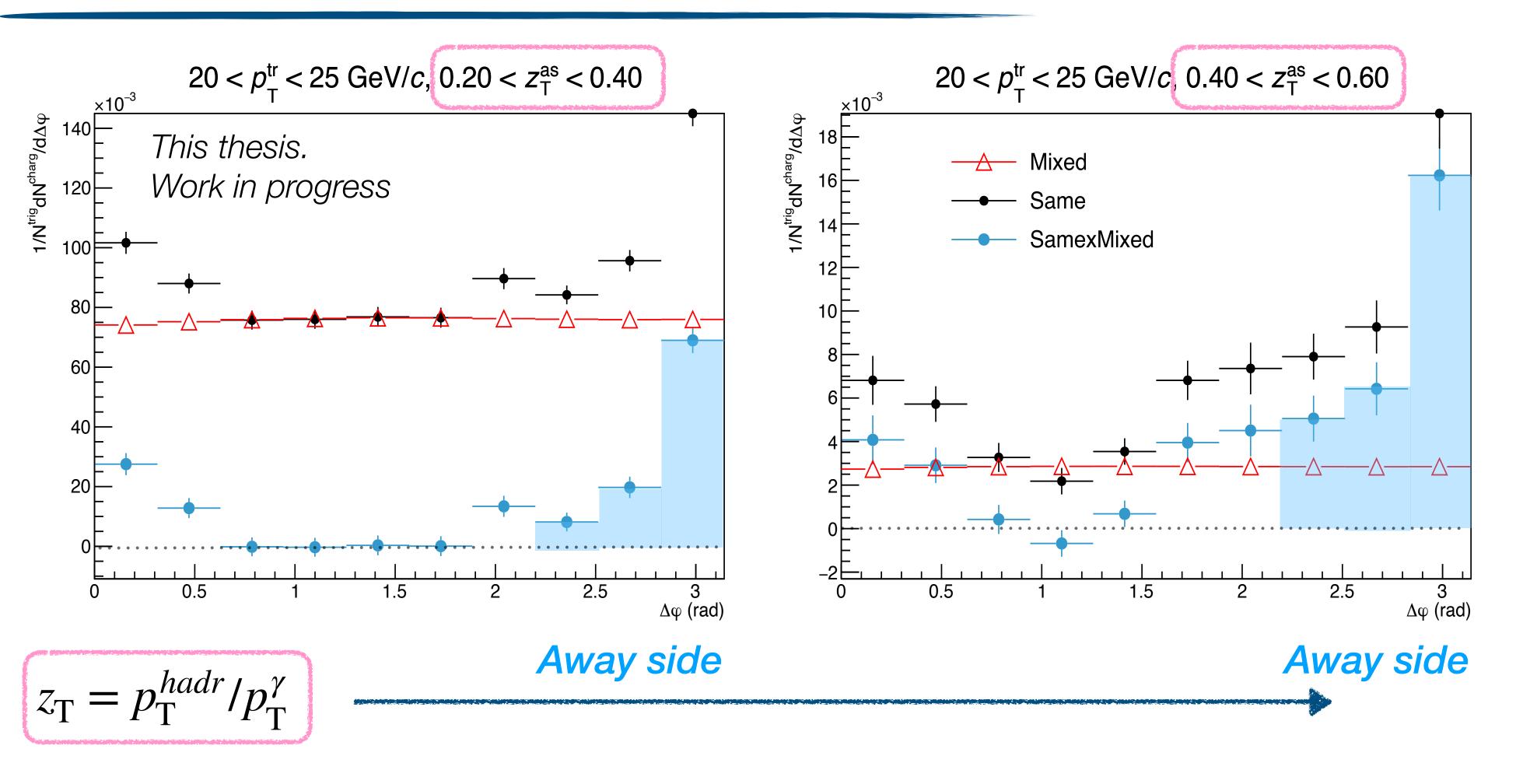
Artificial dataset created combining our trigger particle from a given collision with tracks from other collisions

Best way to subtract the UE

Based on the real data: affected by the detector status + less statistical limitation



#### UE subtraction: Mixed event method



Same - Mixed = Subtracted

- Integration of the away side peak for every  $z_{\rm T}$  interval:  $D(z_{\rm T}) = \frac{1}{N^{trig}} \frac{dN^{en}}{dz_{\rm T}}$  distribution
- But the sample is still not totally pure after the cluster selection and the removal of the UE

## Correlation analysis

Photon-hadron azimuthal distribution



Subtraction of the UE

Mixed Event



Purity Correction

The sample is still not totally pure

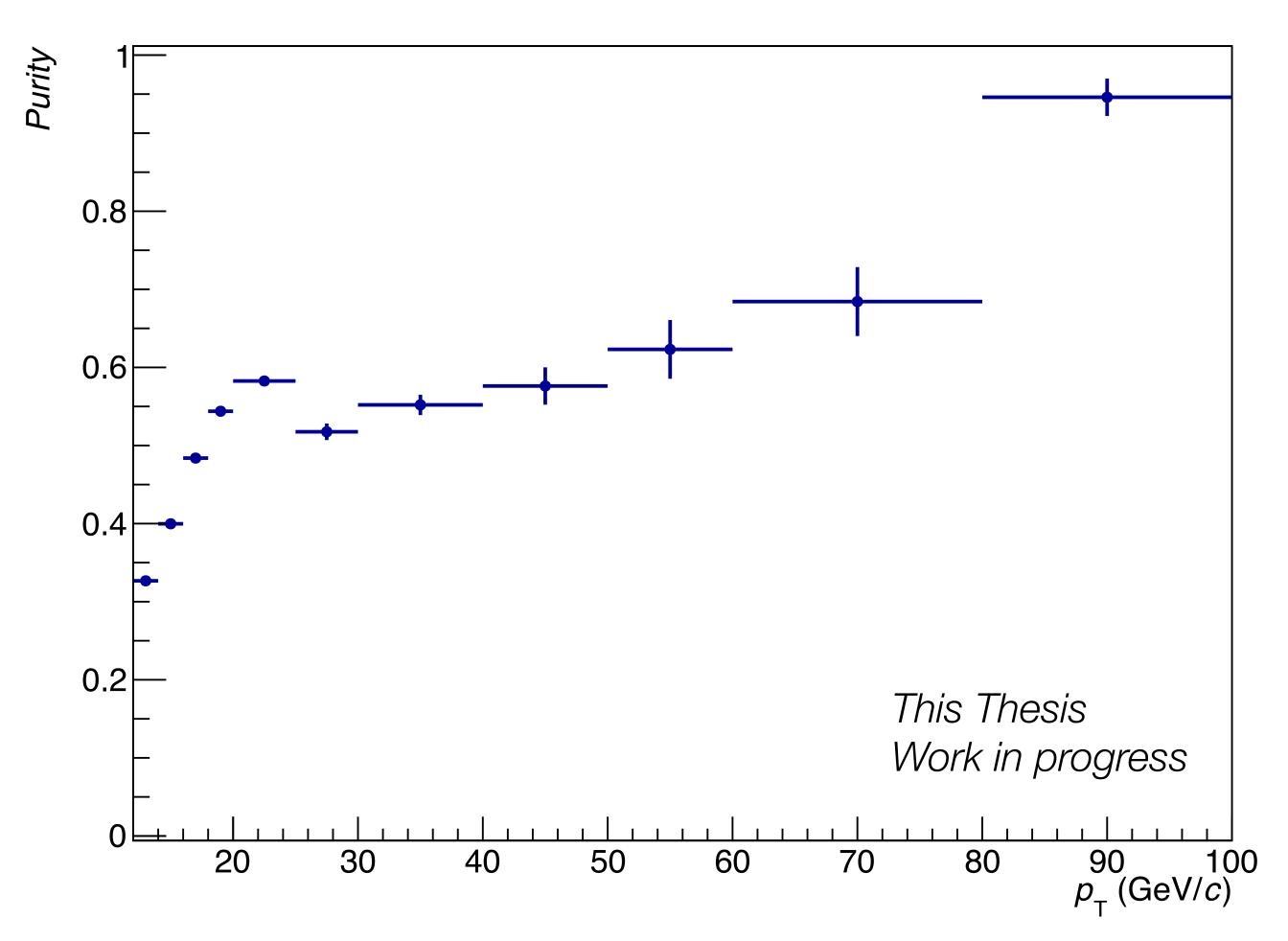


Systematic errors

#### Purity correction

The isolated  $\gamma$  candidate sample has to be corrected for  $\pi^0$  accepted

Assume correlation triggered by  $cluster_{wide}^{iso}$  equivalent to the background for  $cluster_{narrow}^{iso}$ 



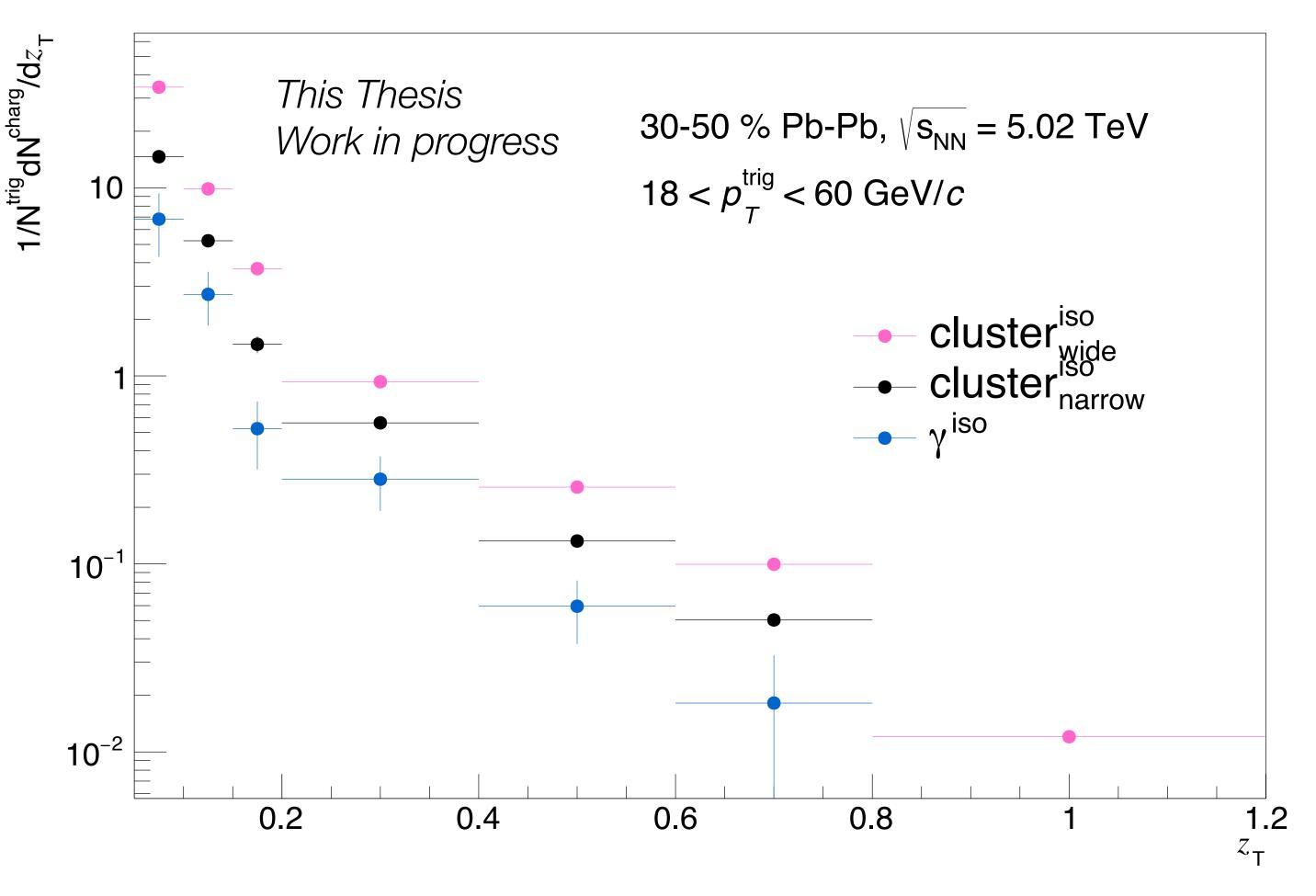
Purity 
$$P = S/N$$
,

the S signal of isolated  $\gamma$  over N = S+B

$$\gamma^{iso} = \frac{\text{cluster}_{\text{narrow}}^{iso} - (1 - P) \text{ cluster}_{\text{wide}}^{iso}}{P}$$

ALICE 2022, Internal, Gustavo Conesa's analysis

#### **Purity Correction**



$$\gamma^{iso} = \frac{\text{cluster}_{\text{narrow}}^{iso} - (1 - P) \text{ cluster}_{\text{wide}}^{iso}}{P}$$

Triggers:

For each trigger: integration of the away
 side peak for every z<sub>T</sub> interval

$$D(z_{\rm T}) = \frac{1}{N^{trig}} \frac{dN^{ch}}{dz_{\rm T}}$$
 distribution

Different 
$$p_T^{tr}$$
 bins merged:  $\frac{1}{\sum_i P_i N_i^{trig}} \sum_i P_i N_i^{trig} f(z_T)_i$  , where  $i$  is every  $p_T^{tr}$  bin

## Correlation analysis

Photon-hadron azimuthal distribution



Subtraction of the UE

----- Mixed Event



Purity Correction

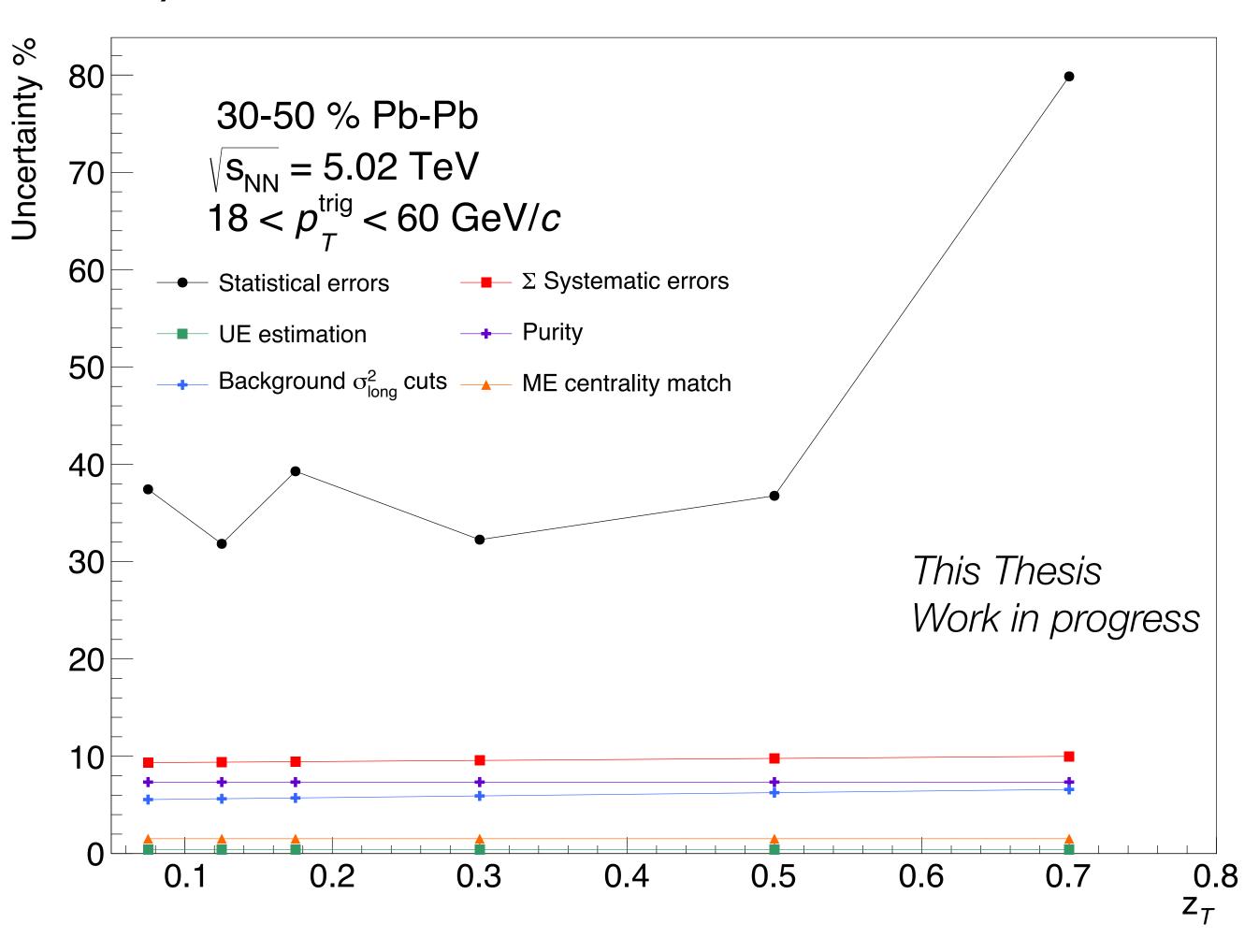
The sample is still not totally pure



Systematic errors

## Work in progress: Systematics errors

Systematic uncertainties for most sources are estimated with selection criteria variations



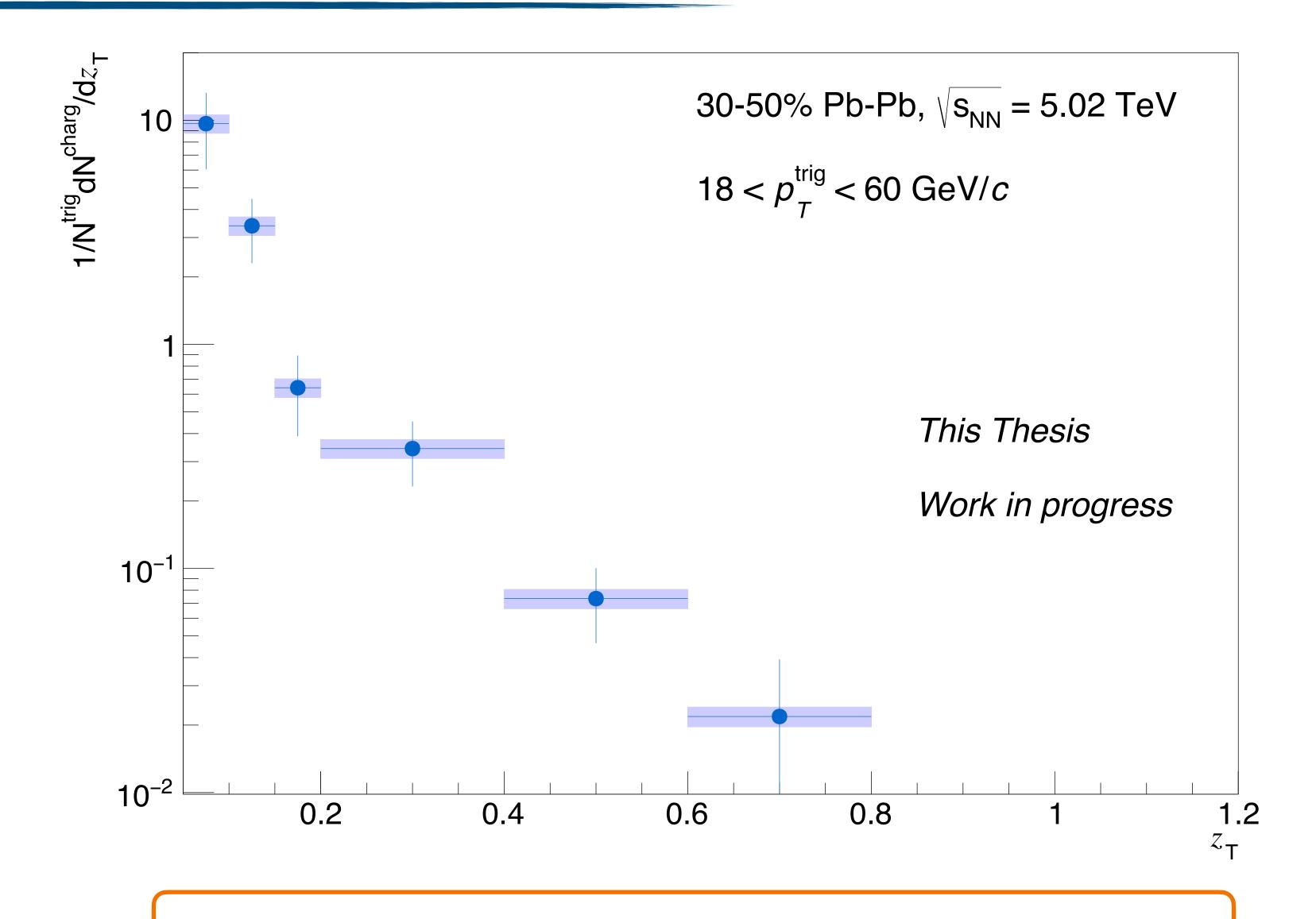
They are evaluated for the  $z_T$  function:

- UE estimation uncertainty
- Purity
- Background  $\sigma_{long}^2$  selection
- Variation of the Mixed event selection parameters
- Tracking efficiency (ONGOING)

Among all the uncertainties the statistical error is the highest one.

Statistical limitation is the major difficulty of the analysis

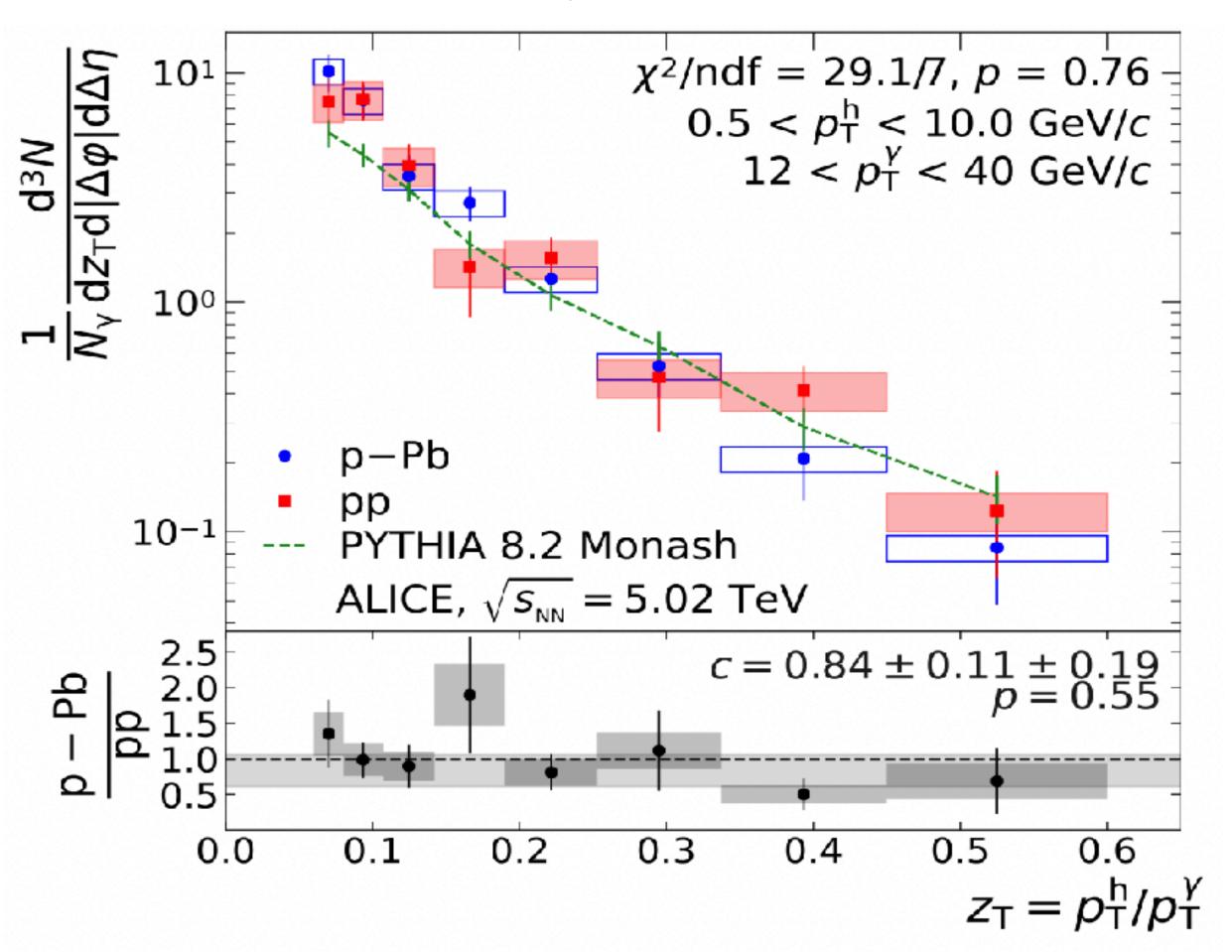
#### Z<sub>T</sub> function



The fluctuations are due to the limited statistics

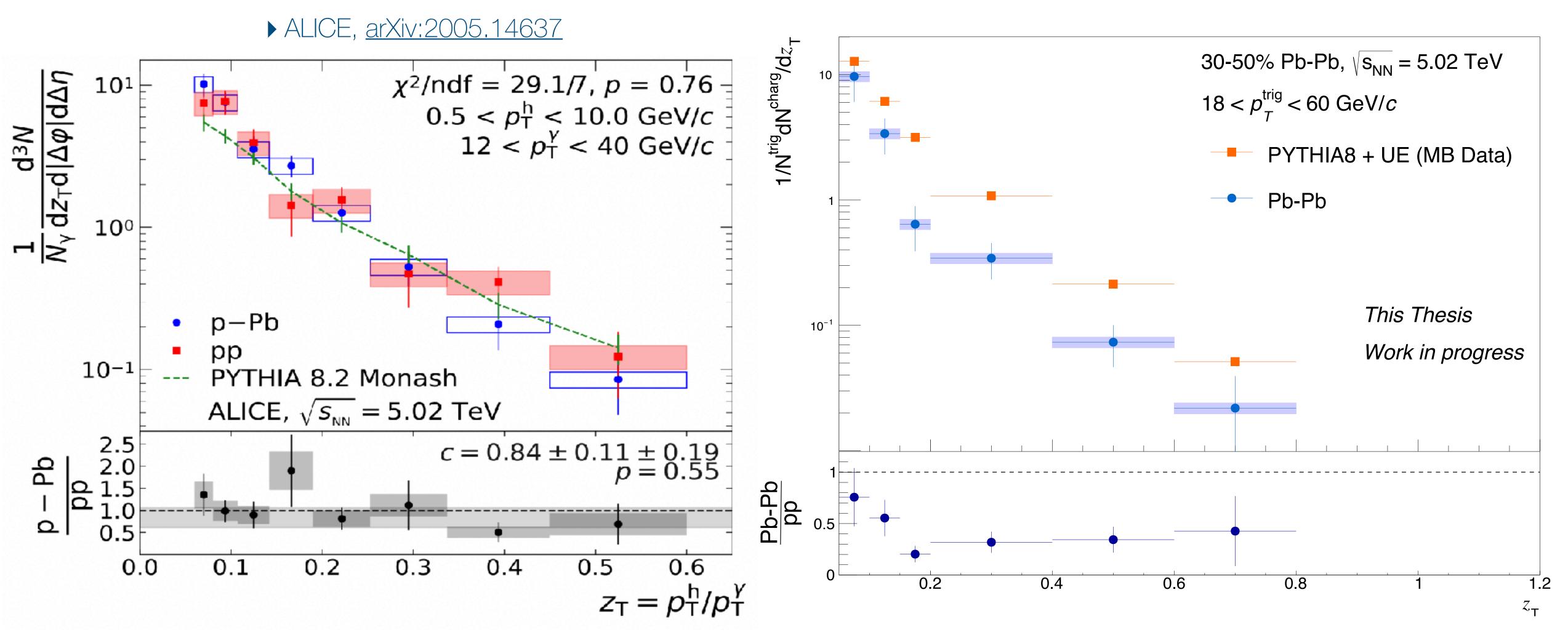
#### Z<sub>T</sub> function

#### ▶ ALICE, <u>arXiv:2005.14637</u>



In pp and p-Pb collisions → agree no QGP

#### Z<sub>T</sub> function



- In pp and p-Pb collisions → agree no QGP
- The PYTHIA8 embedded into real data Pb-Pb UE, used to represent a Pb-Pb collision wo quenching
- There seems to be a suppression in data compared to PYTHIA8, but these results are work in progress <sup>21</sup>

## Conclusions and perspectives

Challenging to make any strong conclusions due to statistical limitation and significant UE

First analysis of isolated photon-hadron correlations in Pb-Pb at  $\sqrt{s_{NN}}$  = 5.02 TeV at LHC!

#### Next Steps:

- Finish the systematics errors
- Investigate different theoretical model prediction for comparison
- Write the analysis documentation in view of the publication

Future prospect for gamma isolation and correlation analysis:

FoCal, my ALICE technical Service work

### FoCal upgrade

FoCal = Forward Calorimeter

Goals: direct photons detection in pp, p-Pb and Pb-Pb

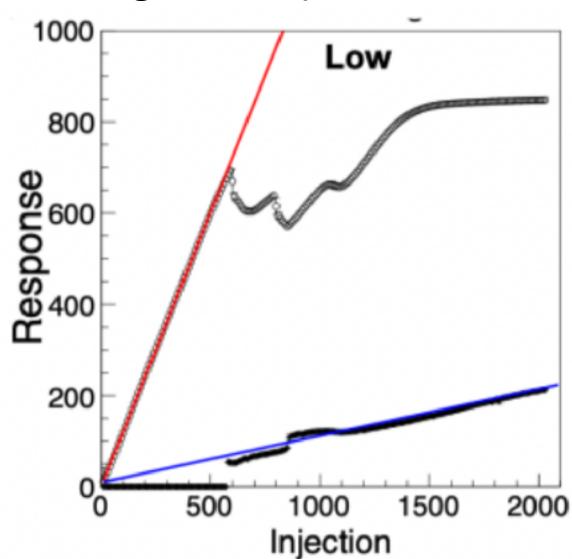
Installation in ALICE foreseen for 2029

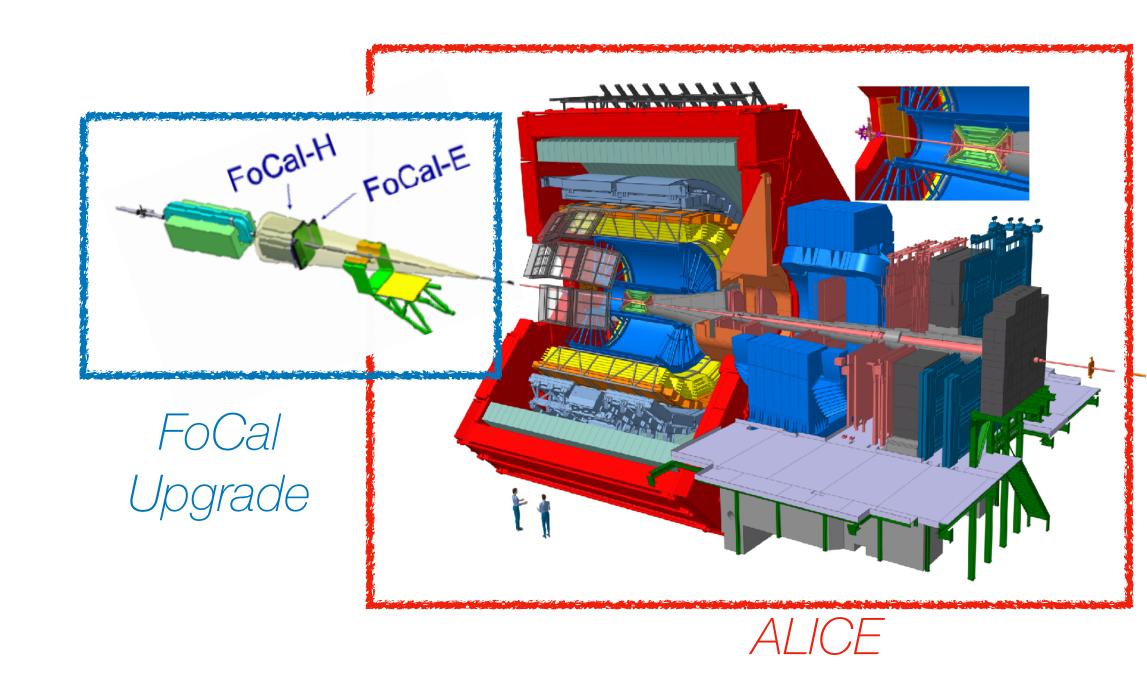
FoCal-E = Electromagnetic calorimeter

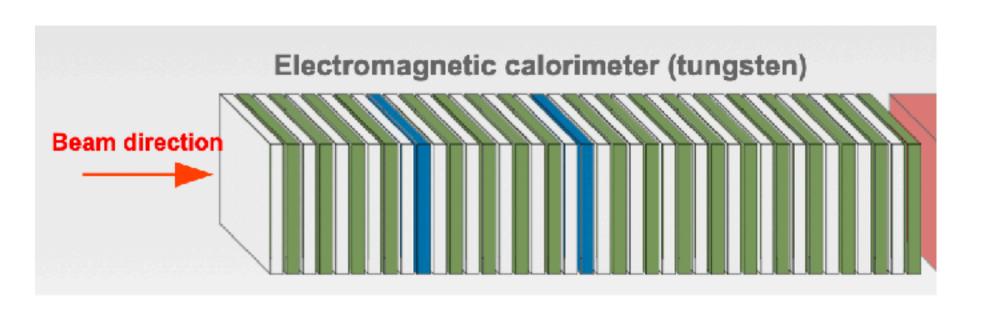
18 layers Si-W calorimeter with 72 channels

Charge deposit in pixels read by an ASIC:

- low charge ADC
- high charge ToT (Time over Threshold)







ALICE group at LPSC contribution to: readout, 18 ASICs FE, interface and DAQ

## Time Over Threshold (TOT) calibration

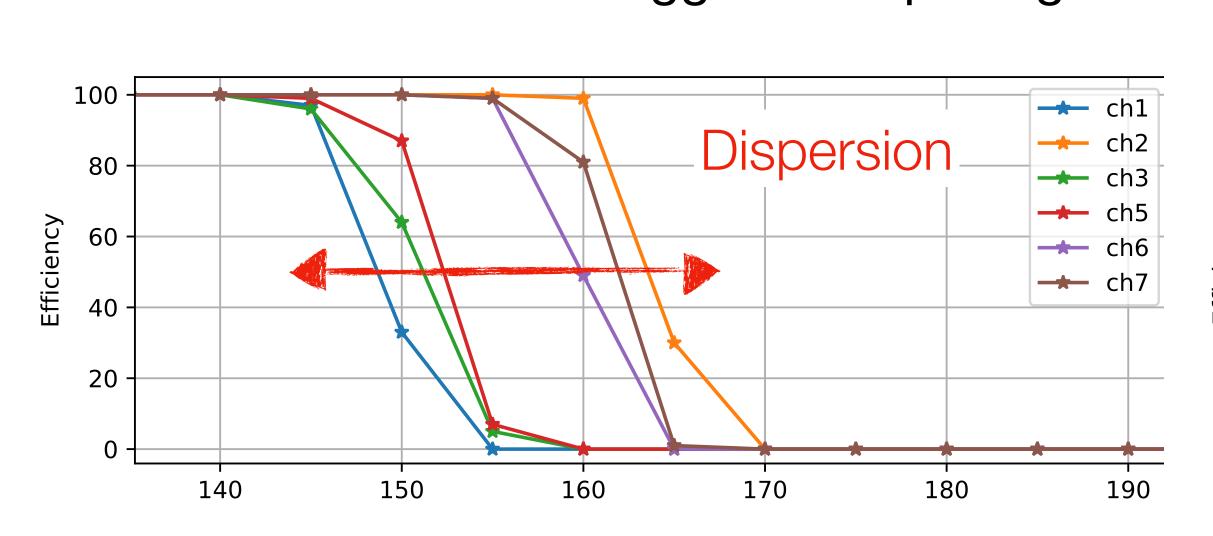
If the deposited charge is too big, the ADC saturates  $\rightarrow$  necessary to use TOT

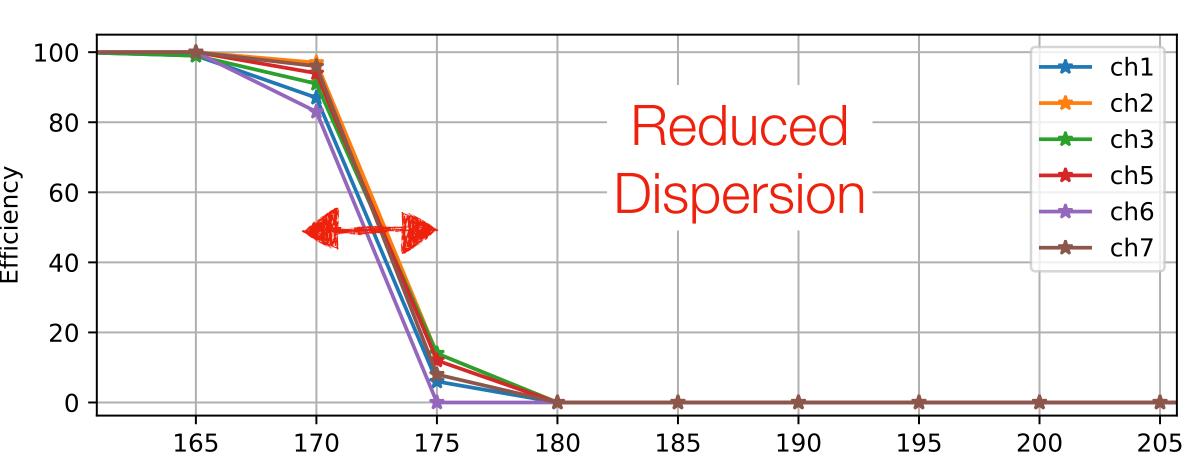
• **Discriminator:** generation of a trigger signal when the input pulse exceeds a threshold value.

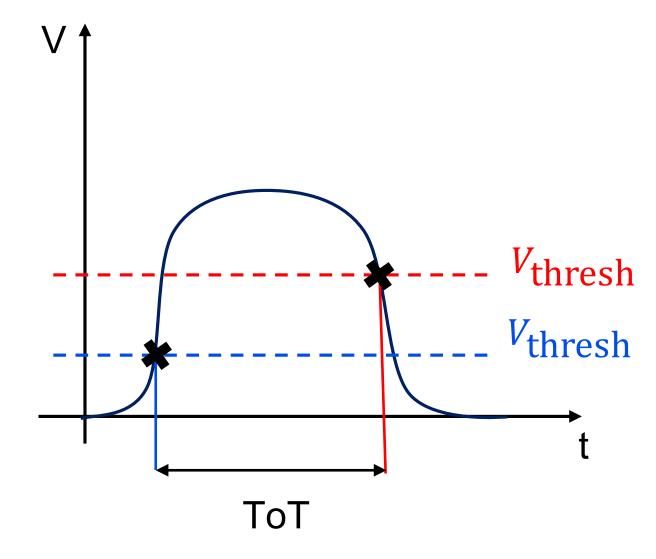
GOAL: reduce dispersion among channels

#### Calibration procedure:

- test different thresholds: check when signal is triggered
- adjust the characteristic parameters of every channel
  - \* different channels trigger the input signal similarly







#### Conclusions

#### FoCal, my ALICE technical Service work

- 6 months already done
- Calibration of the sensor prototype:
  - ADC pedestal, offset correction, ...
- Participation to 3 beam tests in CERN
- Part of these activities have been included in this publication:
  - "Prototype electronics for the silicon pad layers of the future Forward Calorimeter (FoCal) of the ALICE experiment at the LHC", arXiv:2302.13912

## Thank you for the attention!

# Backup

#### Isolation method and UE subtraction

$$R_{cone} = \sqrt{(\eta - \eta_{\gamma})^2 + (\varphi - \varphi_{\gamma})^2}$$

In Pb-Pb high multiplicity in final state →subtracting the UE

#### **How?** METHOD $\underline{\eta}$ - $\underline{\varphi}$ bands

The photon candidate is characterized by a  $\phi_{trig}$  and  $\eta_{trig}$ .

It is possible two bands  $\varphi_{band}$  and  $\eta_{band}$  with the isolation cone of radius R.

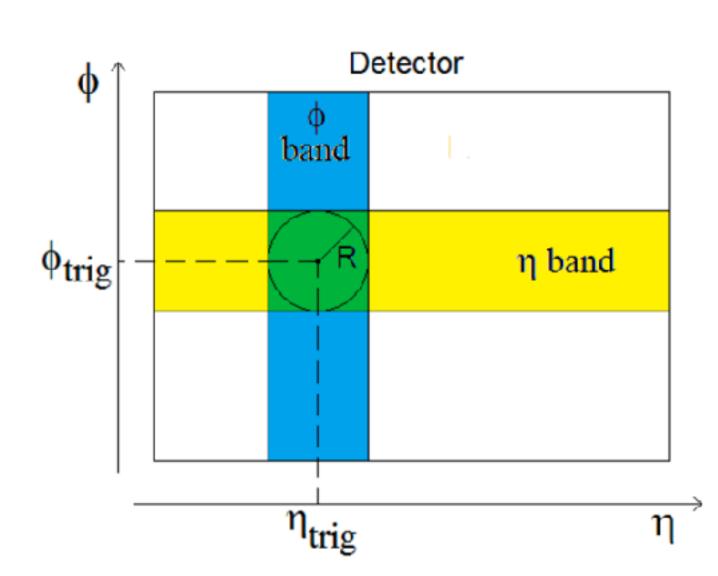
UE subtracted event by event, estimated in region far from our candidate gamma



It's removed from the total transverse momentum in the cone

The final condition on  $p_{\mathrm{T}}$  that defines a isolated  $\gamma$  is :

$$\sum_{cone} p_{\rm T}^{iso} - (\sum_{UE} p_{\rm T}^{iso})_{in\,cone} = p_{\rm T}^{iso} < 2 \text{ GeV/}c$$



#### Purity: ABCD method

Divide clusters  $\sigma_{long}^2$ - isolation energy plane into 4 regions ABCD

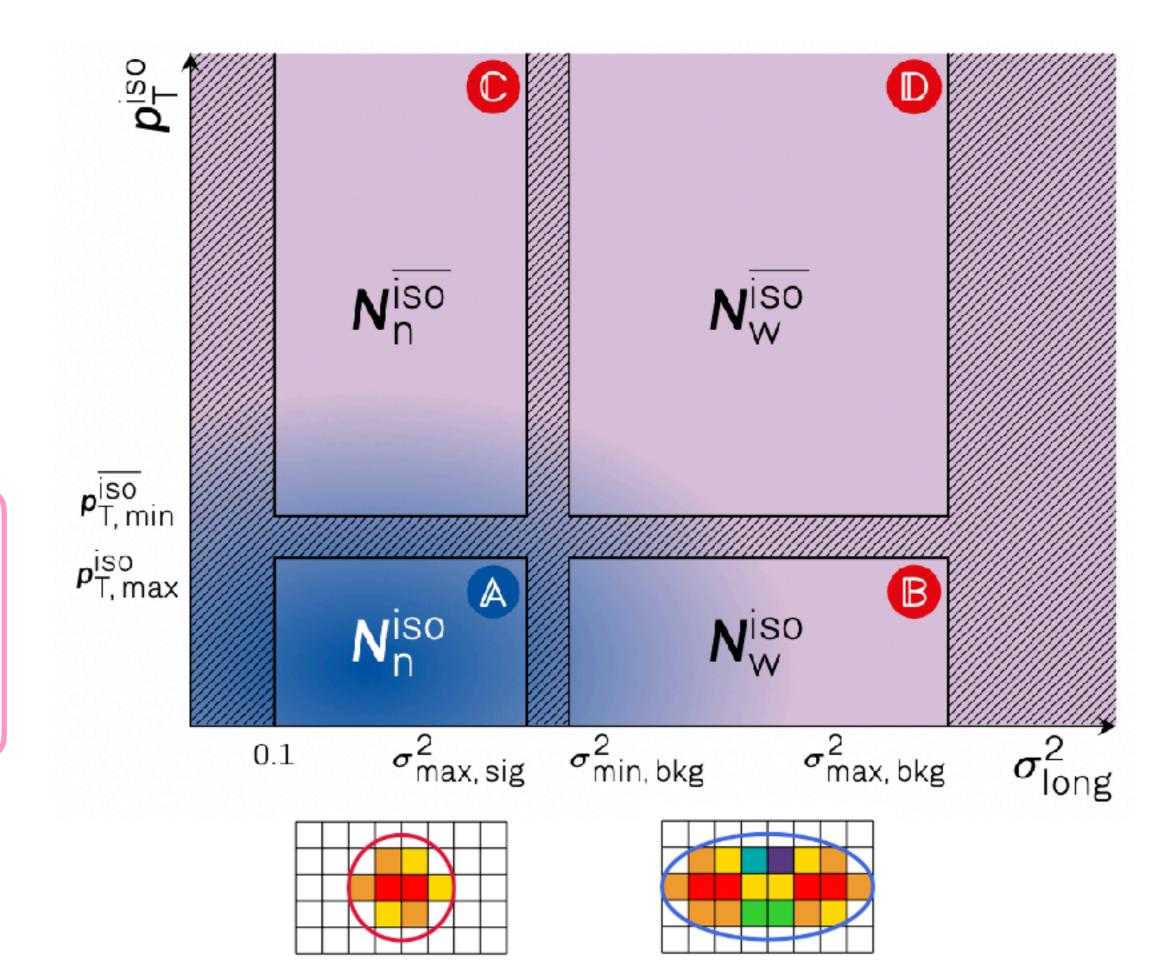
- A: signal dominated region
- B, C and D: background dominated regions

$$N = S + B = \gamma_{direct}$$
 signal + background

• *purity* = *S/N* in A region

$$purity = 1 - \left(\frac{N_n^{\overline{iso}}/N_n^{iso}}{N_w^{\overline{iso}}/N_w^{iso}}\right)_{data} \times \left(\frac{B_n^{\overline{iso}}/N_n^{iso}}{N_w^{\overline{iso}}/N_w^{iso}}\right)_{MC}$$
Data-driven purity

MC correction factor

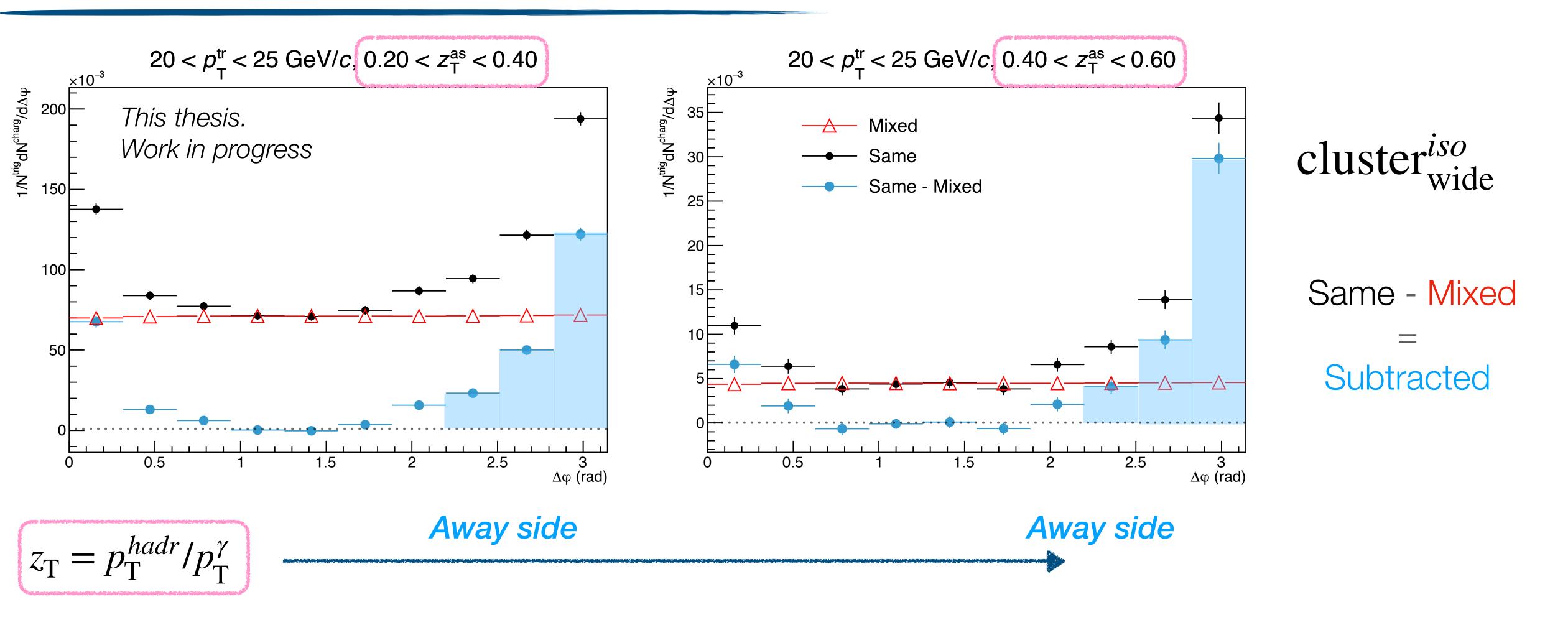


#### Corrections due to:

B isolation fraction depends on the circularity

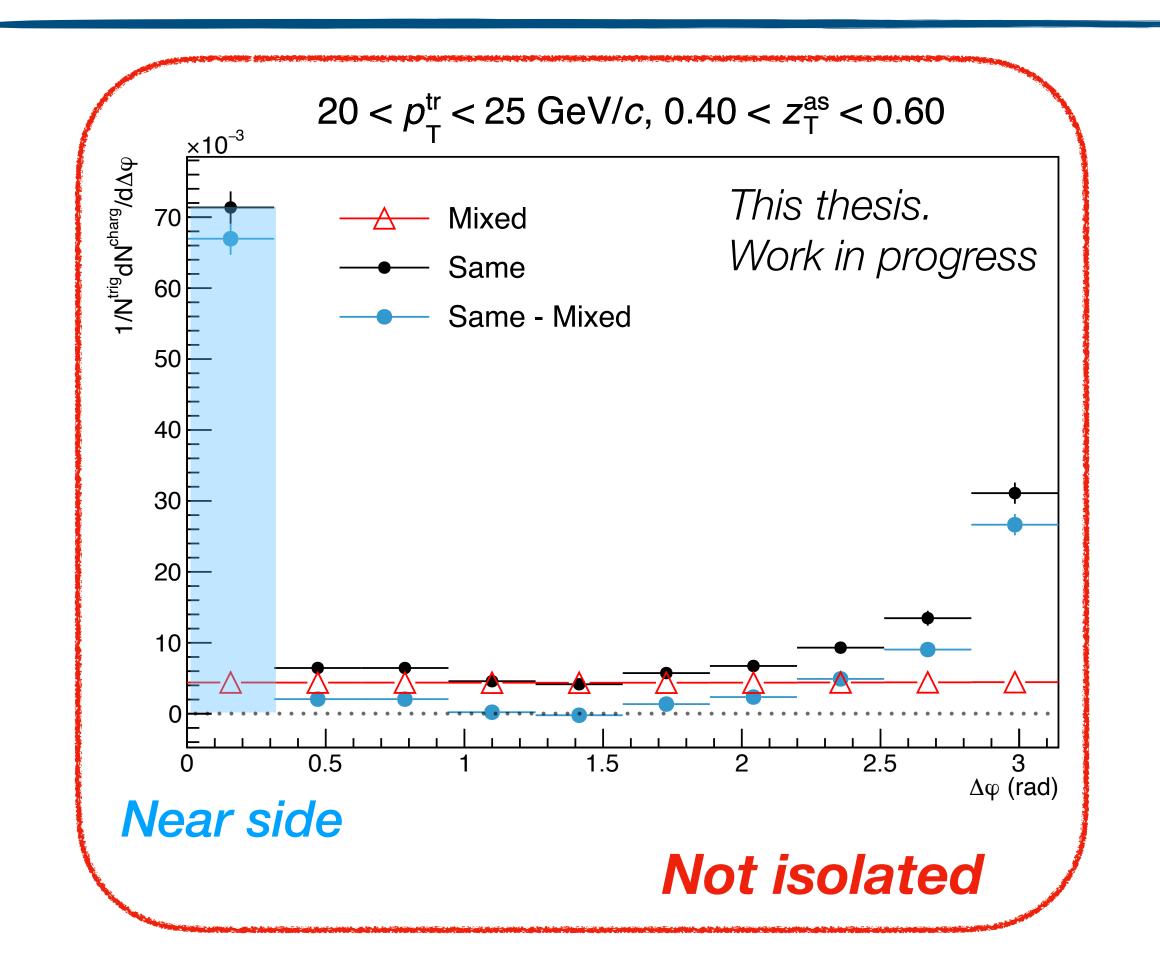
S not contained only in A. it spreads over B. C and D regions

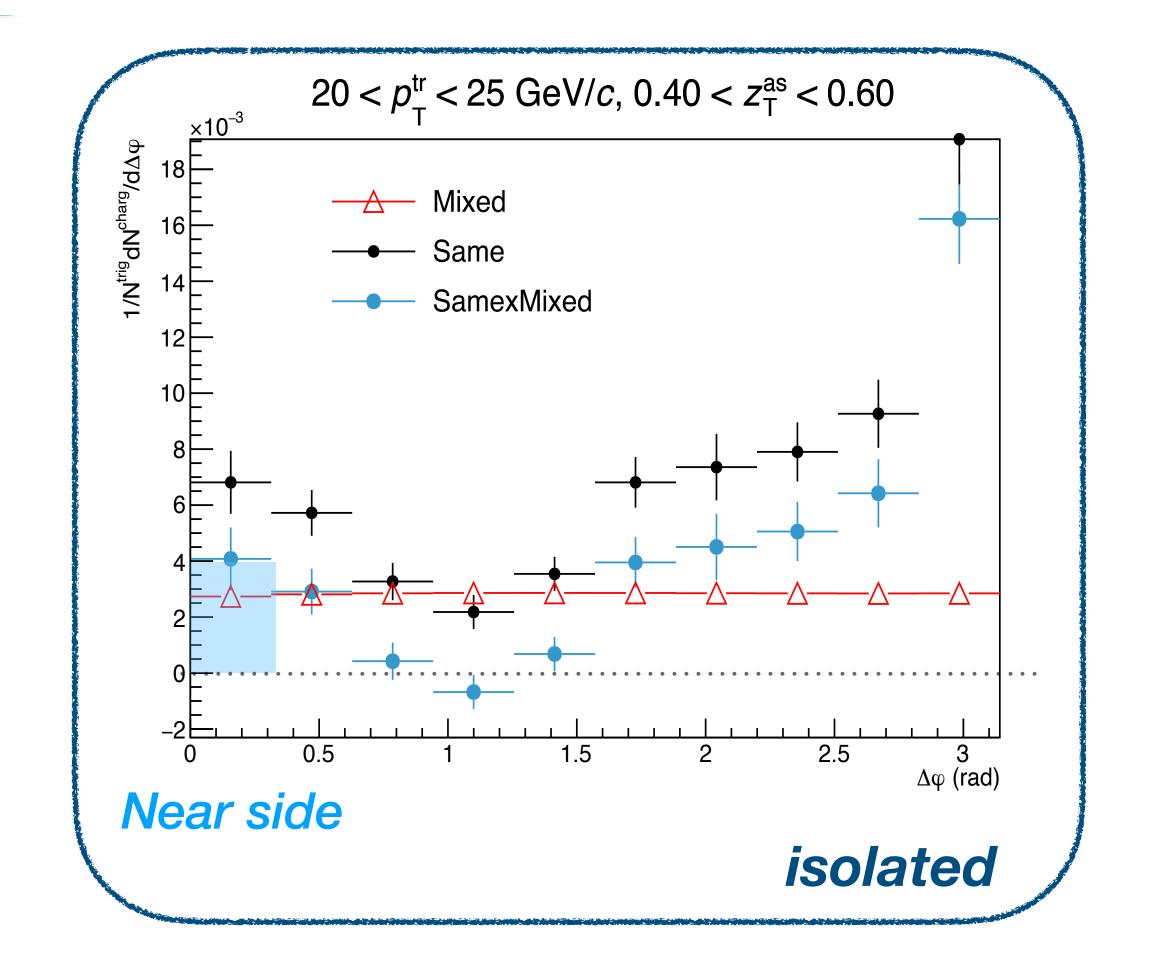
#### UE subtraction: Mixed event method



- Integration of the away side peak for every  $z_{
  m T}$  interval:  $1/N^{trig}dN^{ch}/dz_{
  m T}$  distribution
- But the sample is still not totally pure after the shower shape cut and the removal of the UE

#### UE subtraction: NOT isolated and isolated





- Not isolated: there is activity inside the cone  $\rightarrow$  2 peaks like in di-hadrons correlations
- Isolated only away side peak: it seems two peaks because the statistics fluctuates

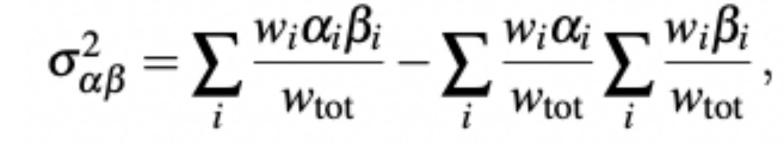
#### Photon identification with EMCal

A particle interacting with the *cell material* produces a shower spreading its energy over *neighbouring cells*.

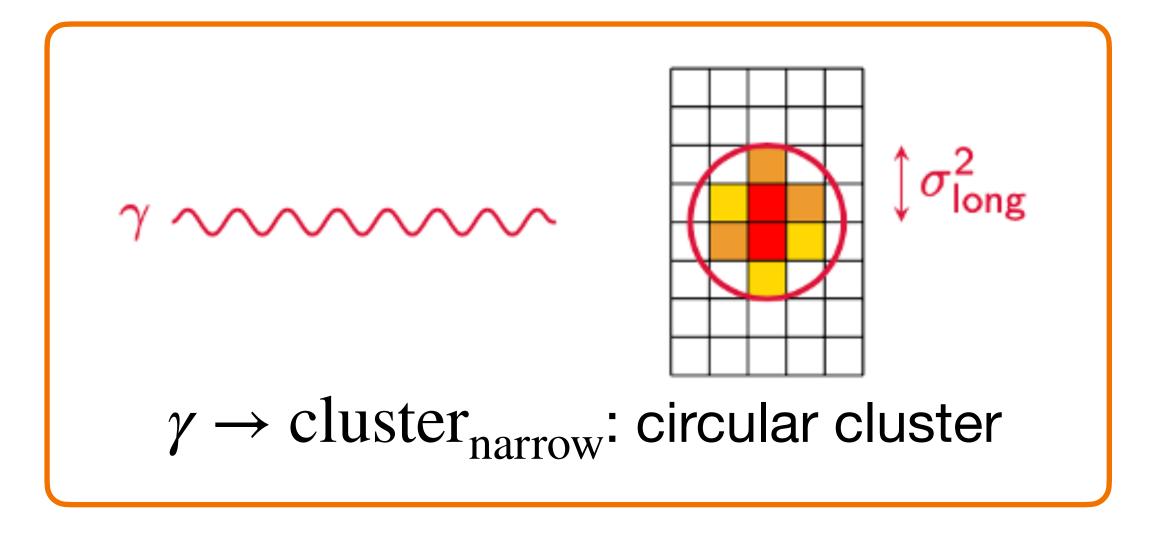
• Cluster: aggregate of cells

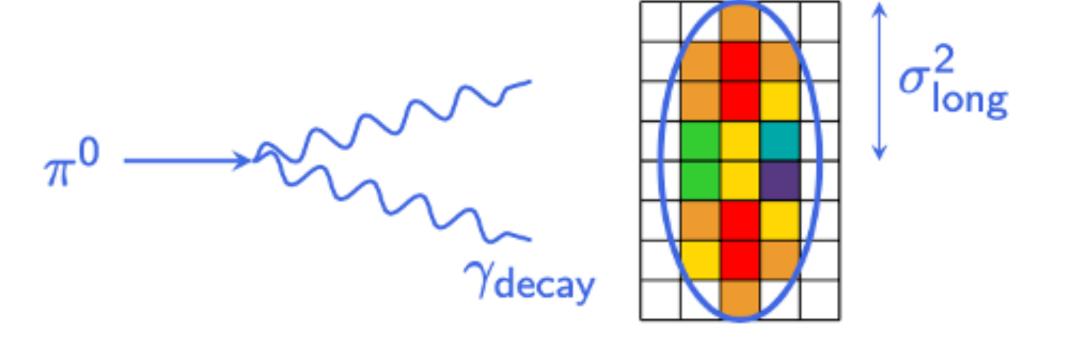
The distribution of energy within a cluster is the shower shape, describe by:

$$\sigma_{\text{long}}^2 = 0.5(\sigma_{\varphi\varphi}^2 + \sigma_{\eta\eta}^2) + \sqrt{0.25(\sigma_{\varphi\varphi}^2 - \sigma_{\eta\eta}^2)^2 + \sigma_{\eta\varphi}^2}, \qquad \sigma_{\alpha\beta}^2 = \sum_i \frac{w_i \alpha_i \beta_i}{w_{\text{tot}}} - \sum_i \frac{w_i \alpha_i}{w_{\text{tot}}} \sum_i \frac{w_i \beta_i}{w_{\text{tot}}},$$



- To discriminate between photons and overlapping decay  $\gamma^{aecay}$  from high E  $\pi^0 \rightarrow gamma$
- apply a condition on the cluster symmetry





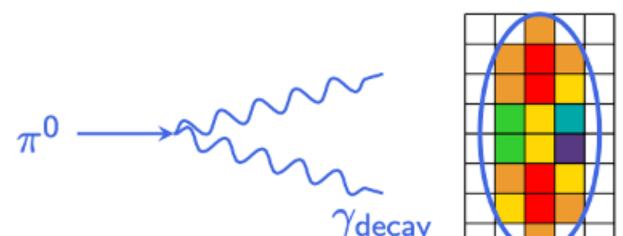
 $\pi^0 \rightarrow \text{cluster}_{\text{wide}}$ : elliptical cluster

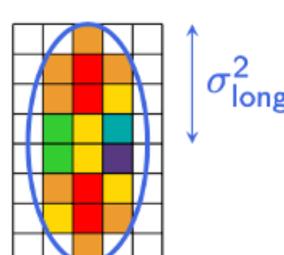
#### Photon identification with EMCal

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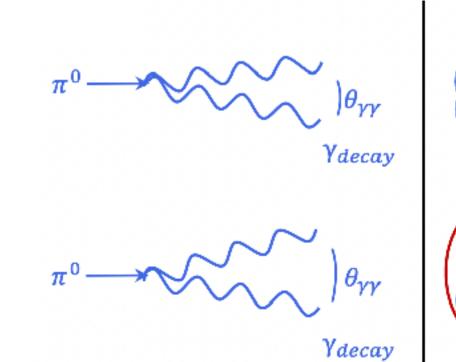
• Cluster: aggregate of cells

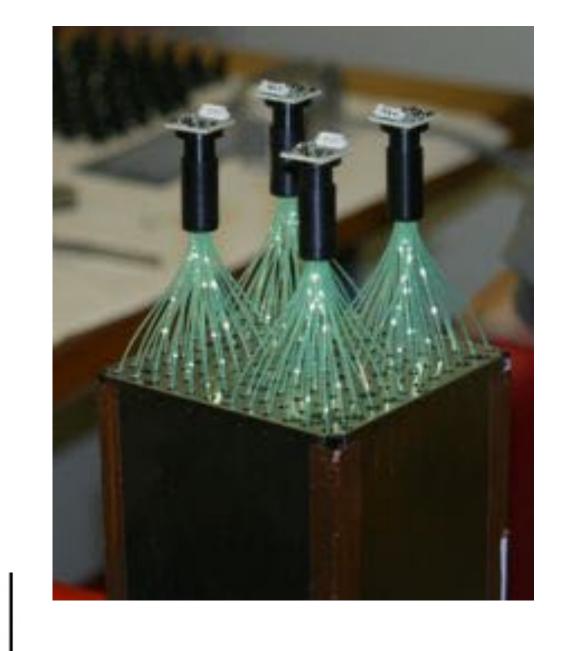
The distribution of energy within a cluster allows to discriminate between single photons  $\gamma$  shower and overlapping  $\gamma_{decay}$  from high energy  $\pi^0 \! \to \! \gamma \gamma$ 





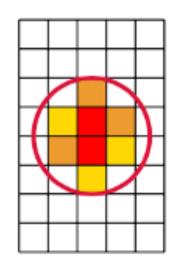
 $\begin{array}{ccc}
 & \sigma_{\text{long}}^{2} & \pi^{0} \rightarrow \text{cluster}_{\text{wide}}: \\
 & \text{elliptical cluster}
\end{array}$ 





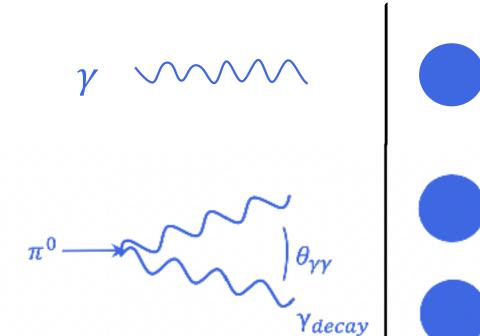


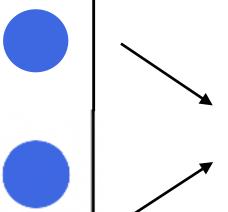




 $\sigma_{\text{long}}^2 \quad \gamma \rightarrow \text{cluster}_{\text{narrow}}$ :

circular cluster





cluster<sub>narrow</sub>

## ALICE experiment + FoCal upgrade

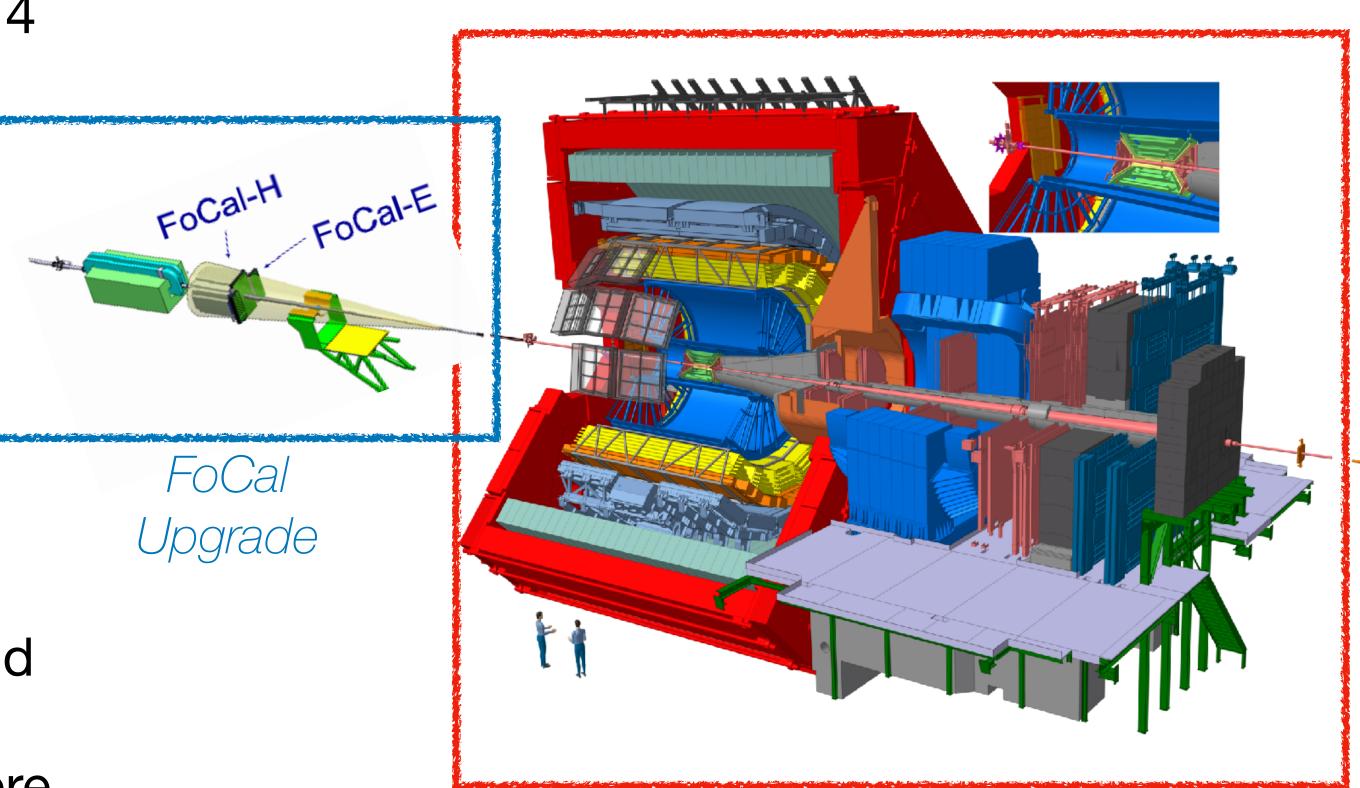
A new forward calorimeter (FoCal) developed as a detector upgrade of the ALICE experiment at LHC

 Installation in ALICE foreseen for LHC Run 4 (2029)

- FoCal unique detector composed of:
  - electromagnetic calorimeter (FoCal-E)
  - hadronic calorimeter (FoCal-H)

• Physics program:

To measure direct photon production in pp and pPb collisions, to quantify the nuclear modification of the gluon density and to explore possible nonlinear QCD evolution



ALICE

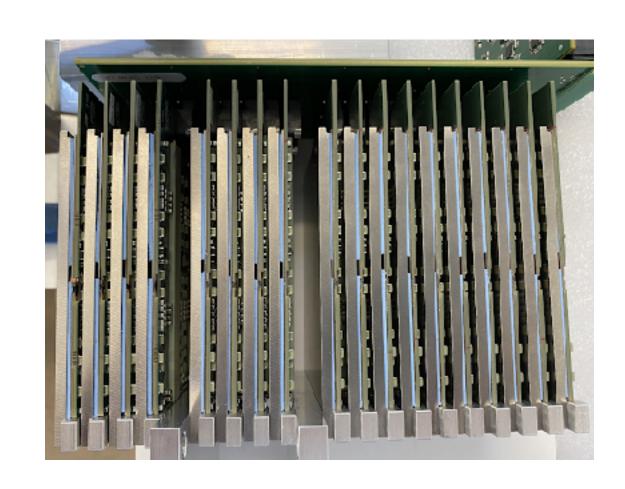
#### Electromagnetic calorimeter: Focal-E

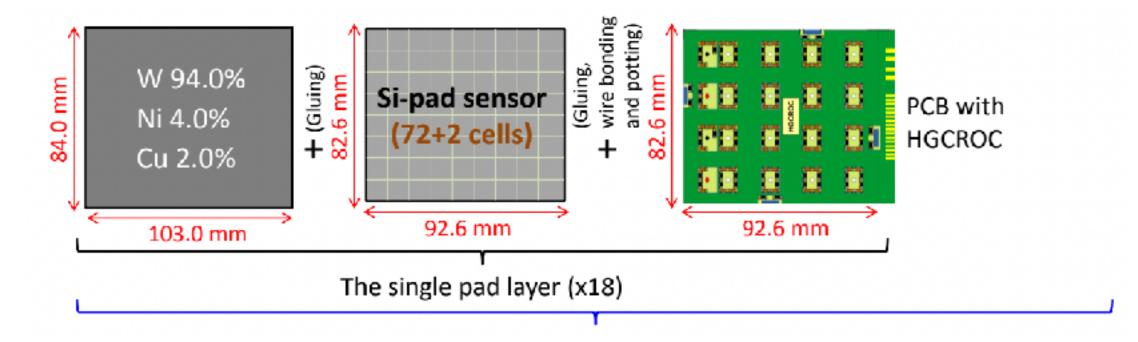
Electromagnetic calorimeter prototype: 18 Si-W layers

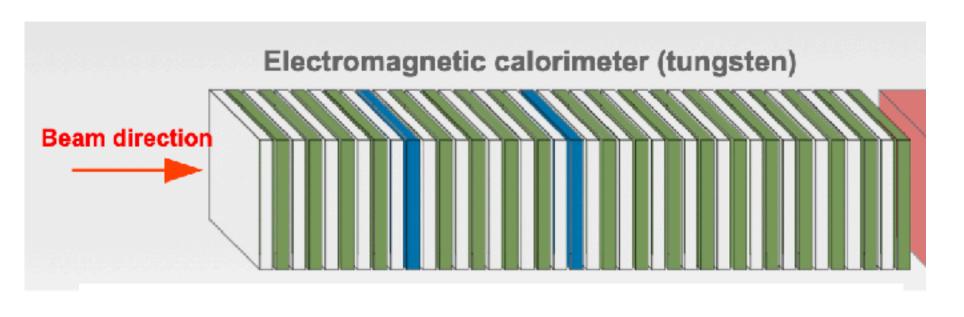
Tungsten (W): the passive absorber

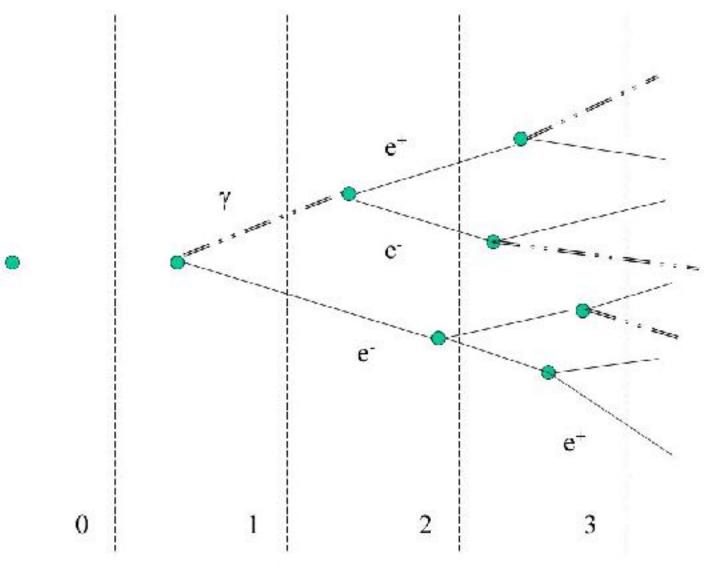
Silicon pad sensor key component of the FoCal-E:

- a good energy resolution and signal-to-noise (S/N) ratio
- a wide dynamic range for the proposed physics programs with FoCal.



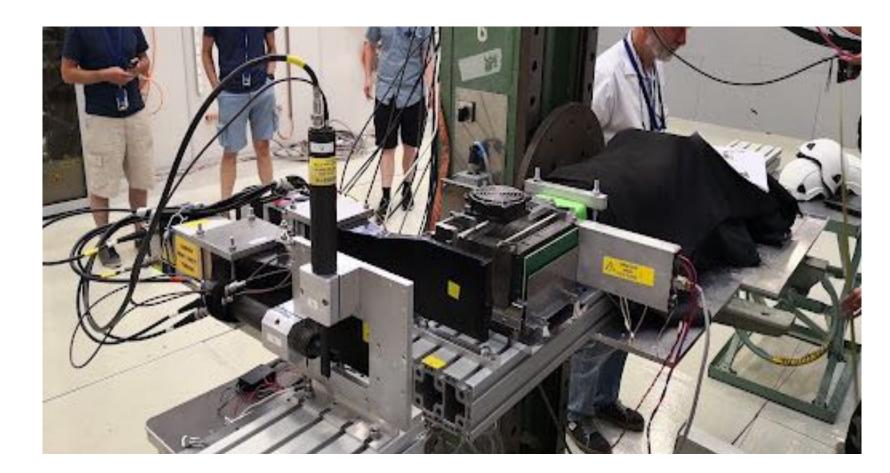


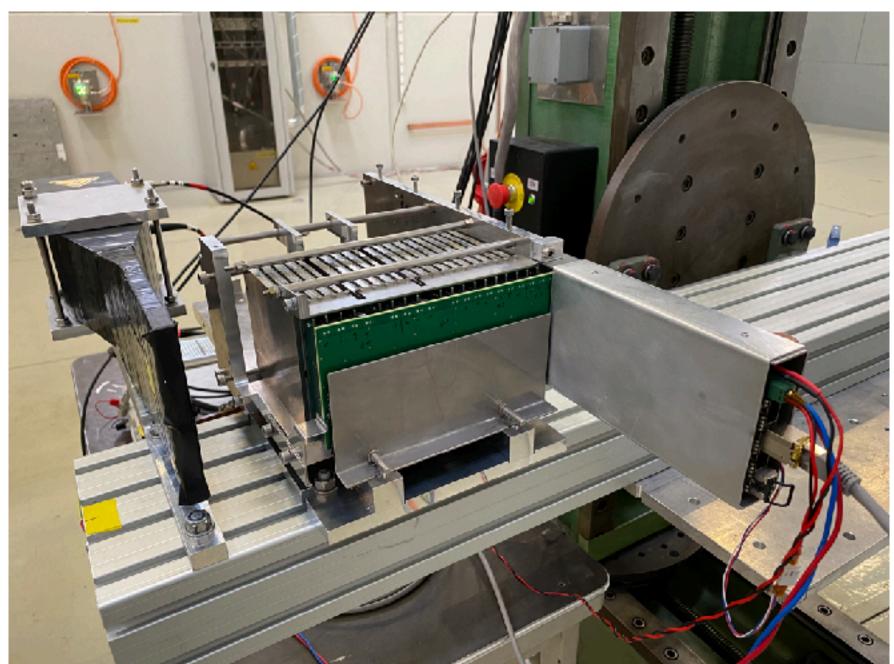




## Beam Test: configuration

June 22 @ CERN PS T9





Test of the 18 layers FoCal-E Pad prototype at CERN PS T9

Configuration:

3 PMT for the coincidence

Focal-E Pad

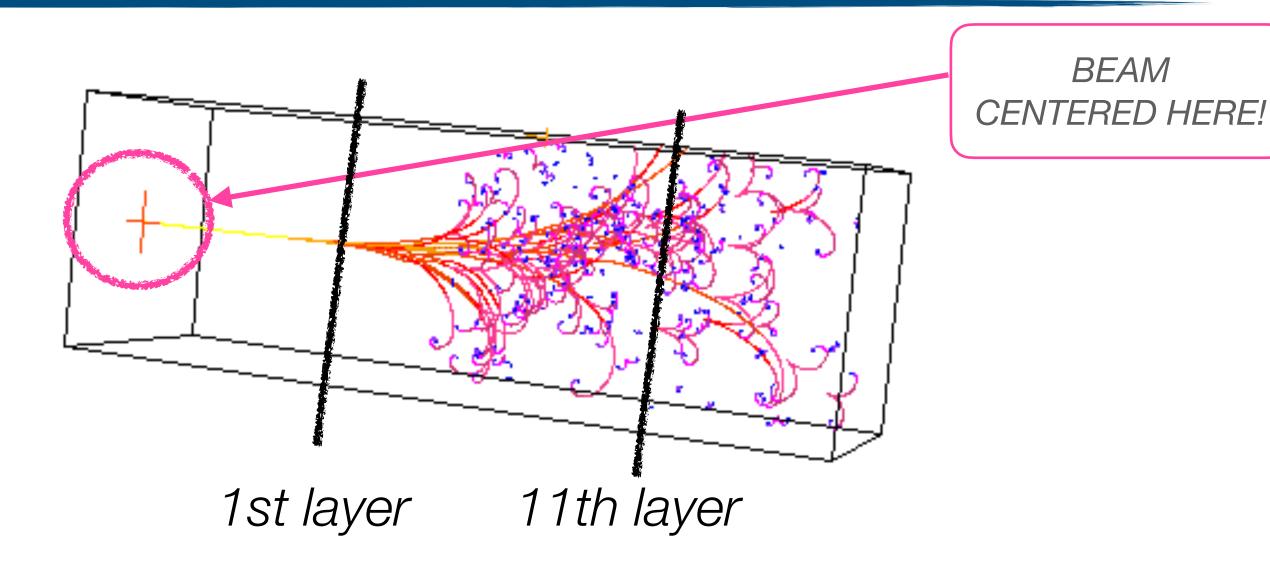
Beam characteristics:

Hadrons of 60, 50, 40, 20, 10 GeV

Objective:

Characterisation of the sensor with the measure of Minimum Ionizing Particle (MIPs)

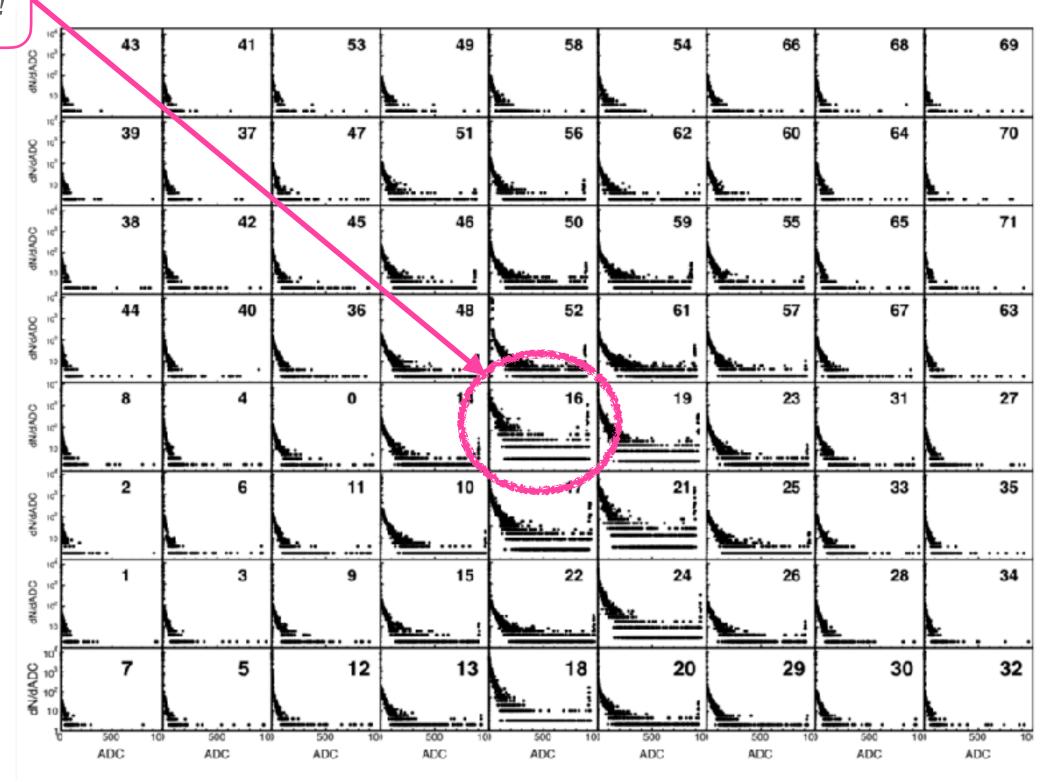
#### Beam Test measurements: Shower evolution



Check the electromagnetic shower and its evolution.

As the number of layer increases, the shower spreads more and more.

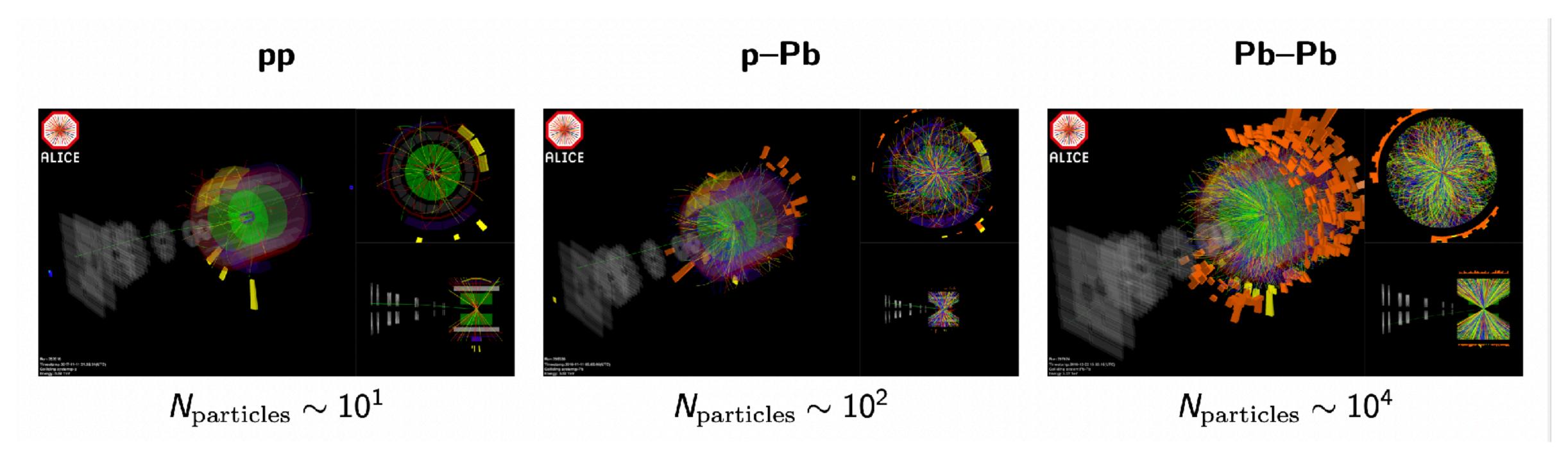
- More channels are activated
- Longitudinal evolution of the shower



11th layer, 60 GeV @ SPS beam test Sept 22

## Big questions in heavy ions physics

- How does the Quark Gluon Plasma form, evolve, and transition again into hadronic matter?
- Is there a QGP droplet formed in smaller collision systems (pp, p-Pb)?
- How do the high-energy partons originated in HIC travelling through the colored QGP medium lose energy?



#### Photon identification with EMCal

A particle interacting with the *cell material* produces a shower spreading its energy over *neighbouring cells*.

- Cluster: aggregate of cells with energy above the noise threshold
- To distinguish between neutral and charged clusters: the track information coming from ITS-TPC is used.

The distribution of energy within a cluster, referred to as "shower shape", is described using a elliptical parametrization.

$$\begin{split} &\sigma_{\text{long}}^2 = 0.5(\sigma_{\varphi\varphi}^2 + \sigma_{\eta\eta}^2) + \sqrt{0.25(\sigma_{\varphi\varphi}^2 - \sigma_{\eta\eta}^2)^2 + \sigma_{\eta\varphi}^2}, \\ &\sigma_{\text{short}}^2 = 0.5(\sigma_{\varphi\varphi}^2 + \sigma_{\eta\eta}^2) - \sqrt{0.25(\sigma_{\varphi\varphi}^2 - \sigma_{\eta\eta}^2)^2 + \sigma_{\eta\varphi}^2}, \end{split}$$

