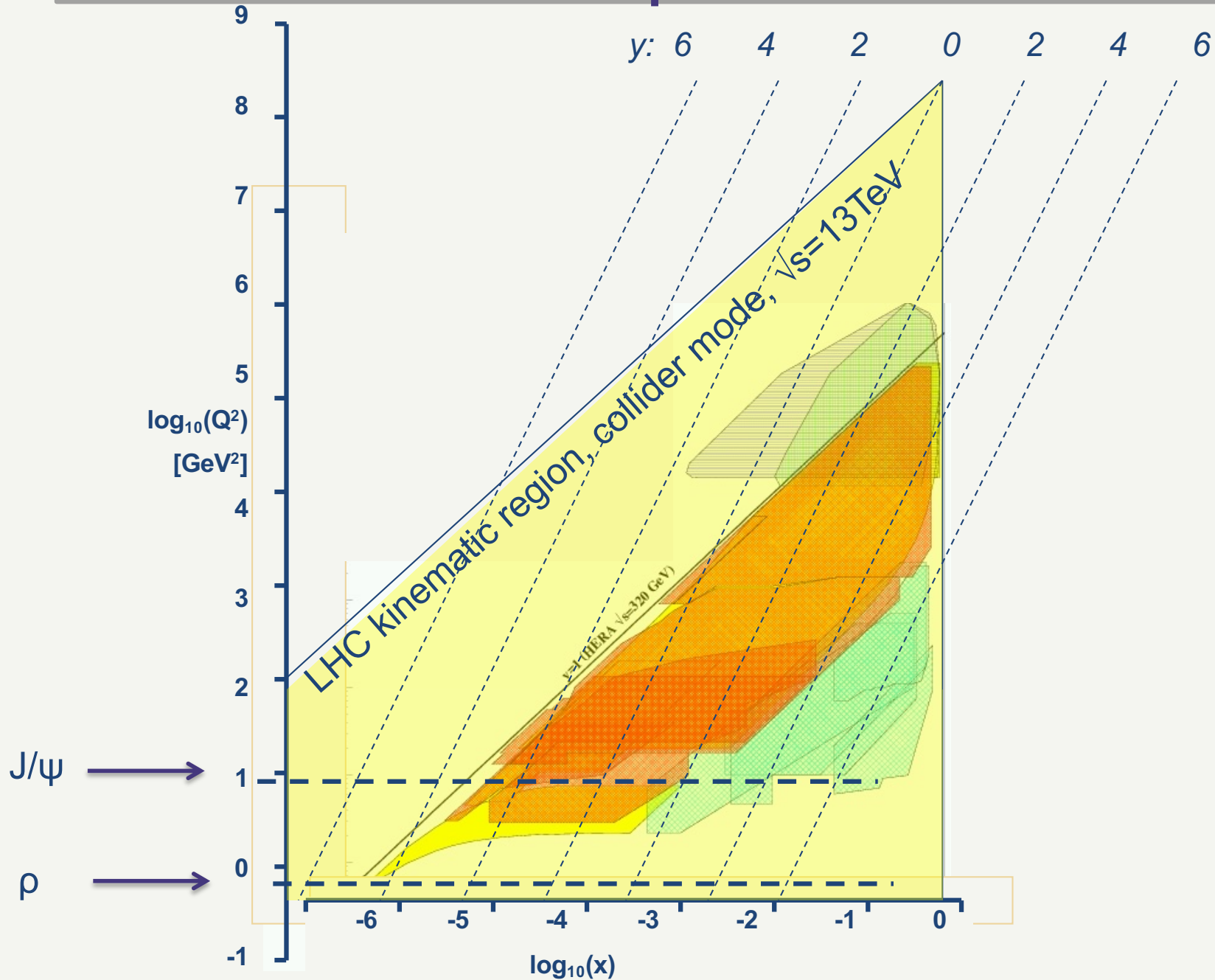


# Low-x and forward physics

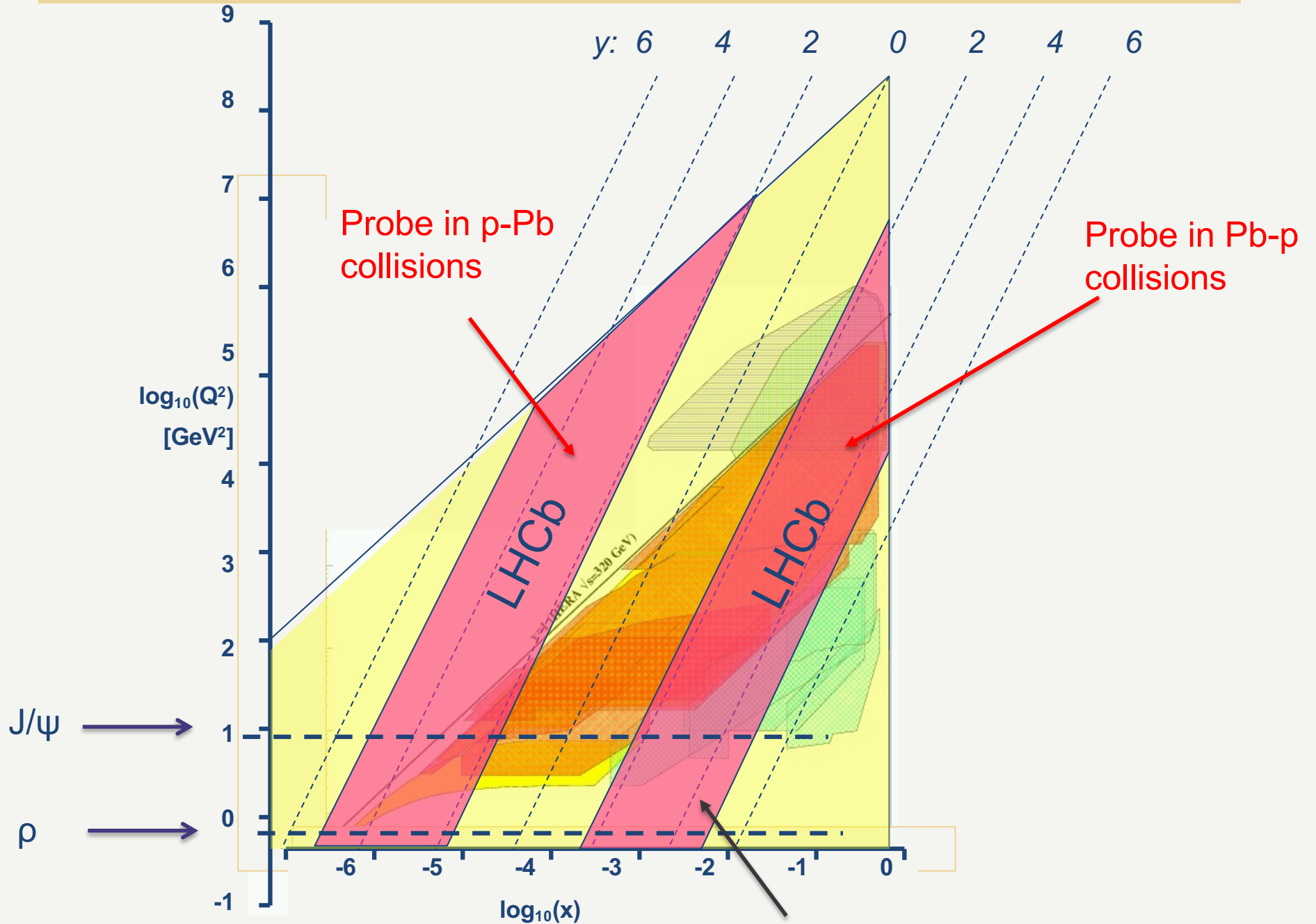
Ronan McNulty  
DIS, Grenoble, April 8-12 2024



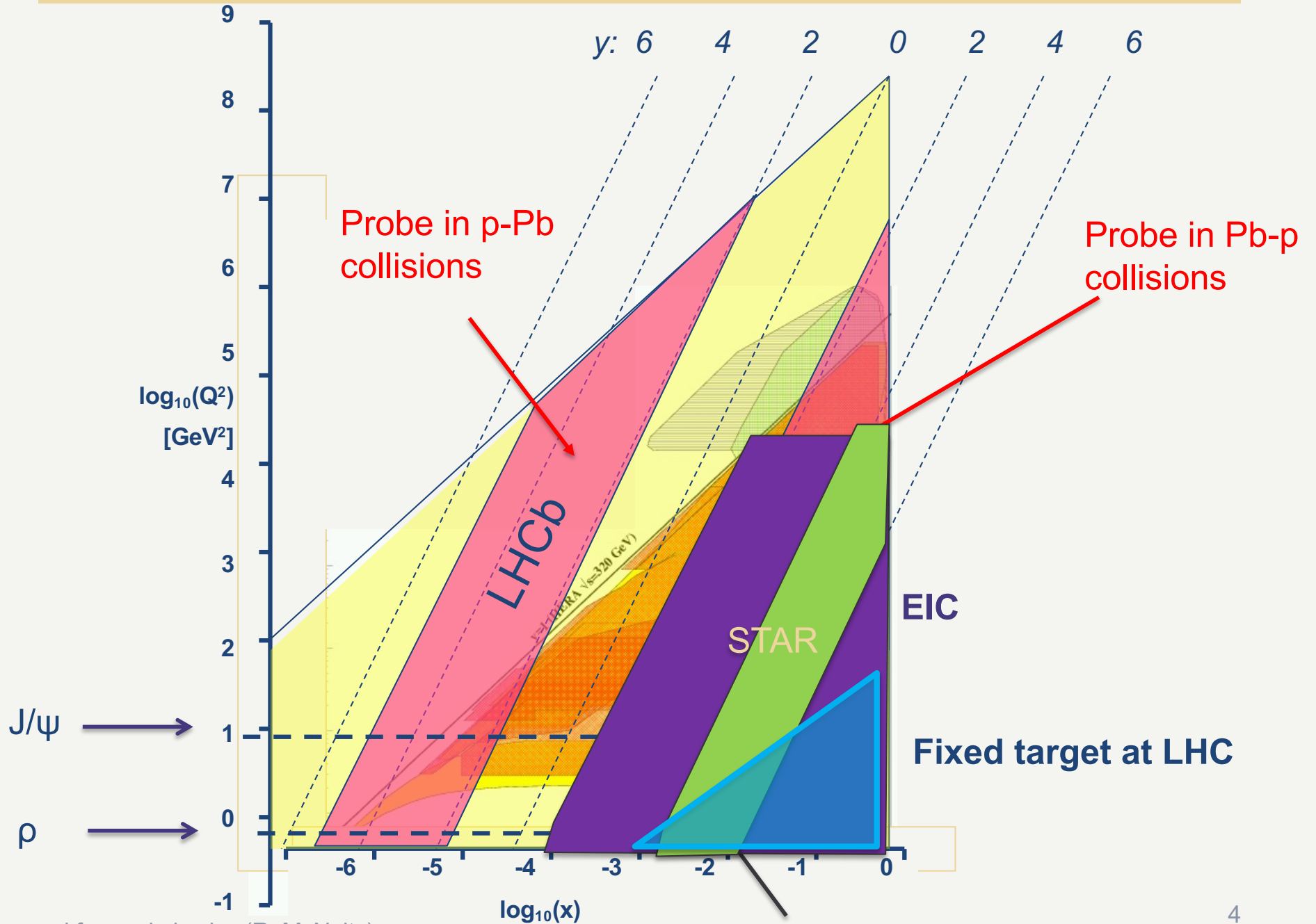
# $x$ - $Q^2$ values probed at LHC





# Low-x & Forward physics



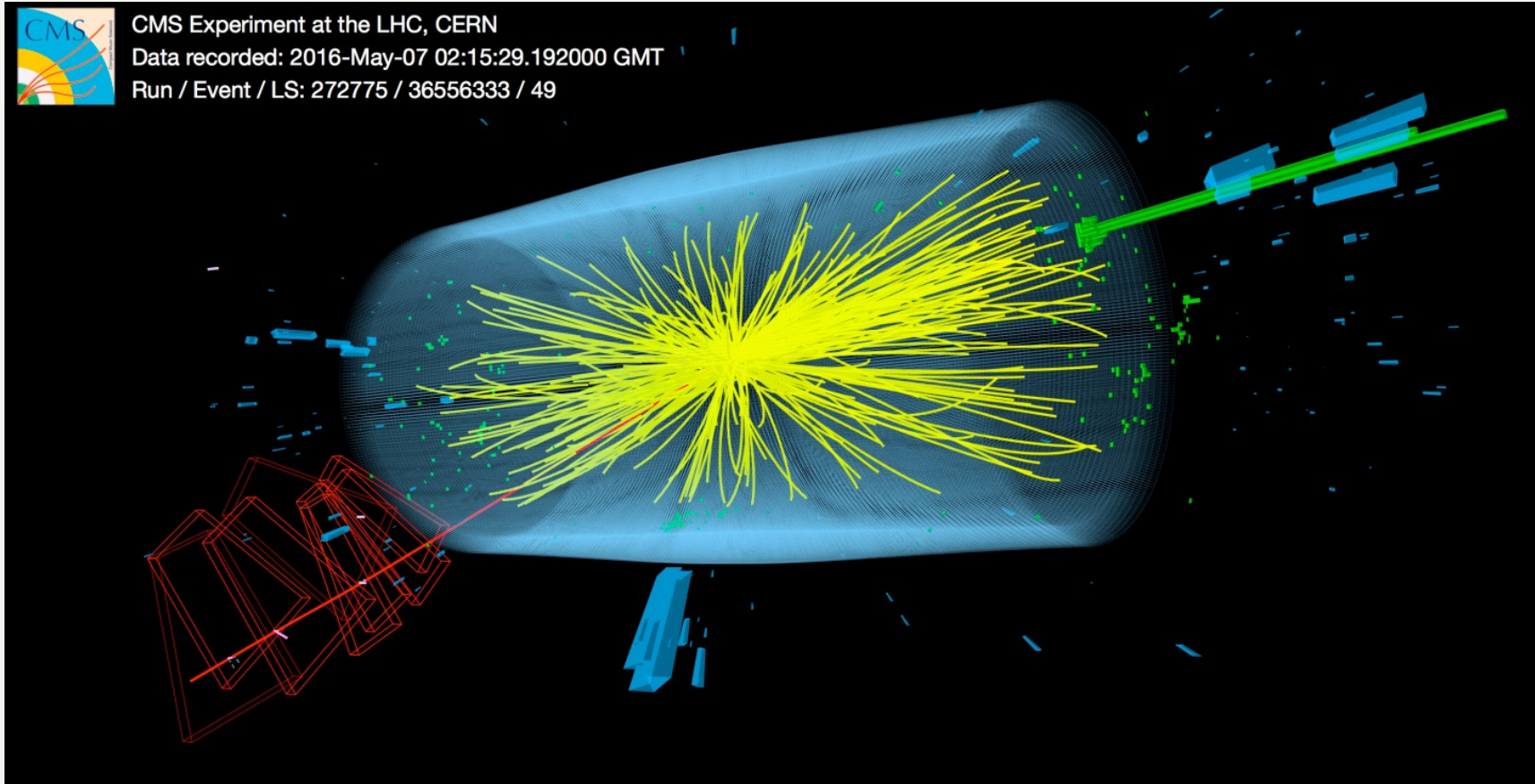
# Low-x & Forward physics



# Motivation

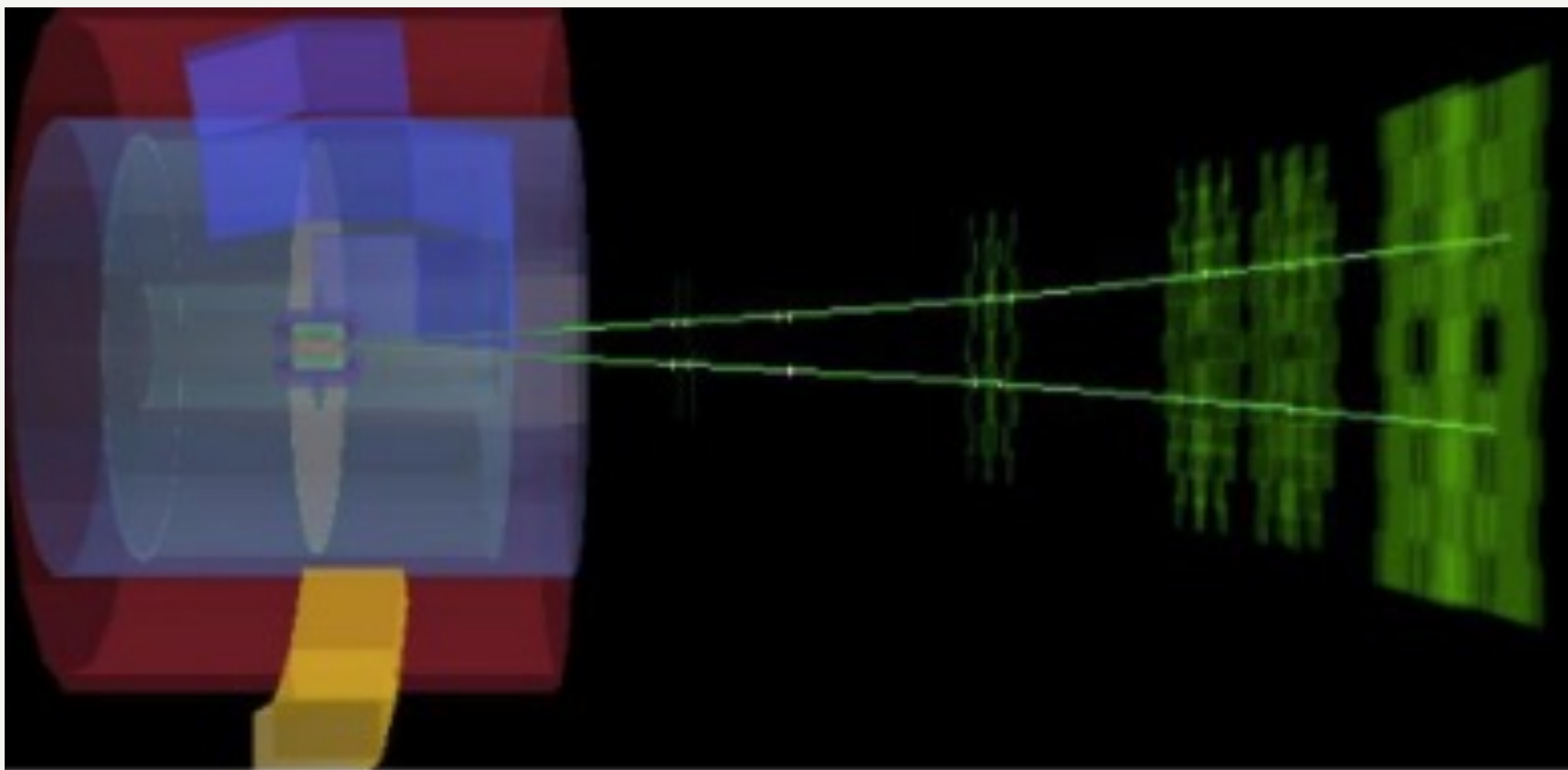
- Much to understand about QCD
  - perturbative / non-perturbative regime
  - proton and nuclear structure (PDFs GPDs)
  - saturation 
  - quark model bound states ( $\rho$ ,  $\rho'$ ,  $f_0, f_2, \dots$ )
  - beyond the naïve quark model (hybrids, tetraquarks , glueballs)
  - colourless propagators: pomerons and odderons
- Can be addressed in diffractive DIS.

# pp collision



Most collisions at the LHC, pp, pA, AA have enormous multiplicities due to colour flow. However, when colourless propagators are involved, multiplicities are low and events have large **rapidity gaps**.

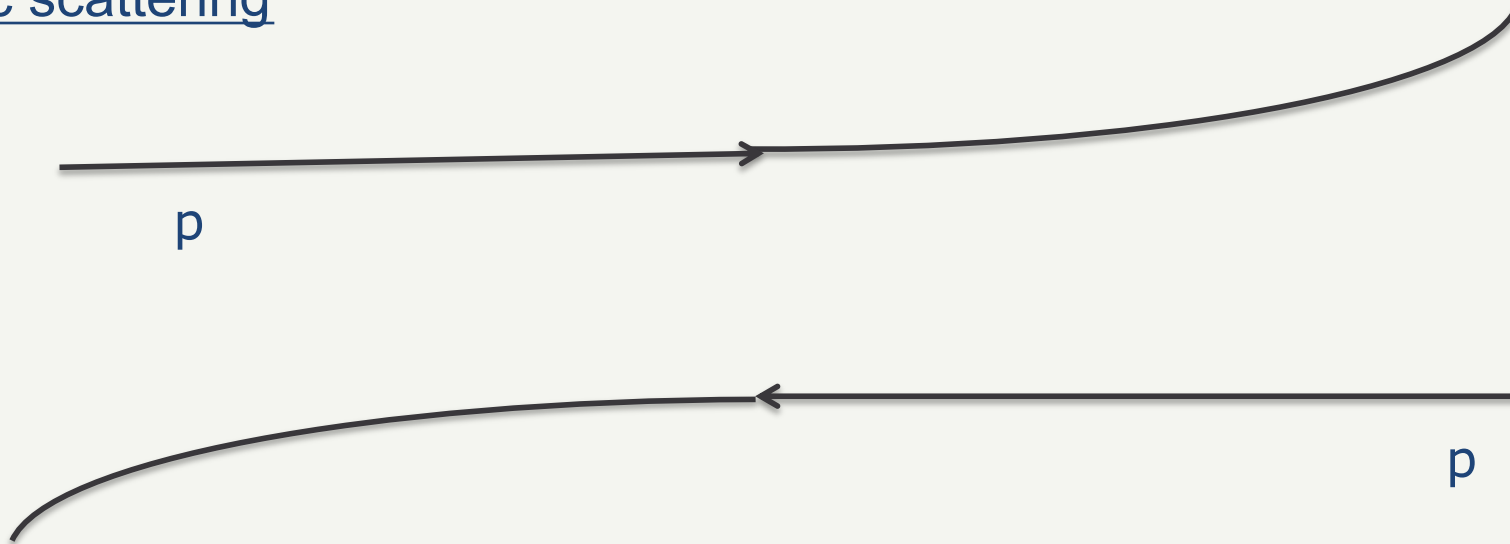
# UPC $J/\psi$ at forward rapidity in ALICE PbPb data



(from Evgeny Kryshen talk at INT workshop)

# Physics of the Vacuum

## Elastic scattering



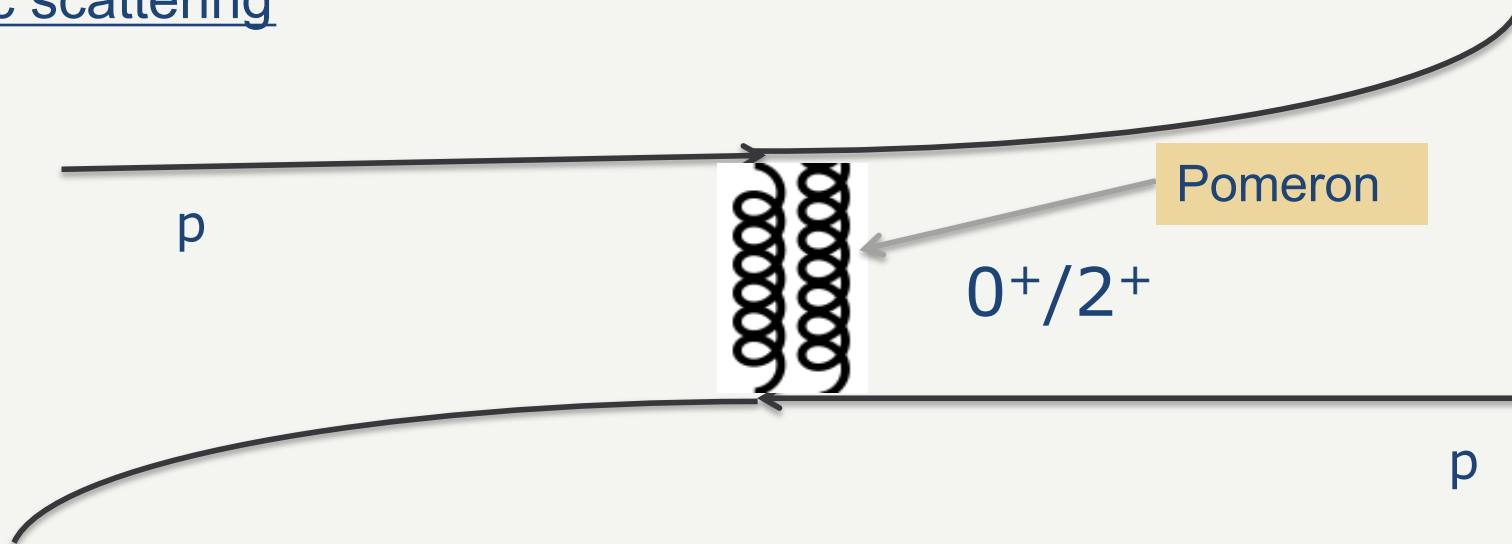
It's QCD – but not as we normally see it. It's colour-free

$\sigma_{\text{elastic}}$	$\approx 40\text{mb}$	←
$\sigma_{\text{diffractive}}$	$\approx 10\text{mb}$	
$\sigma_{\text{inelastic}}$	$\approx 60\text{mb}$	



# Physics of the Vacuum

## Elastic scattering

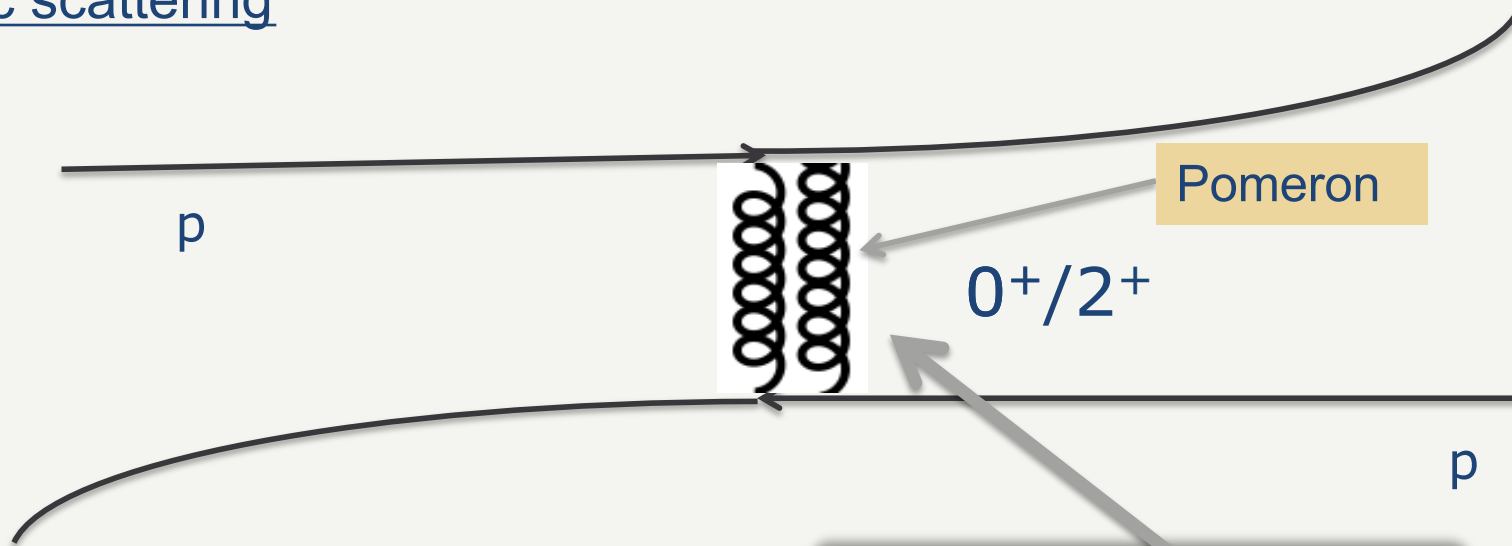


It's QCD – but not as we normally see it. It's colour-free

$\sigma_{\text{elastic}}$	$\approx 40\text{mb}$	←
$\sigma_{\text{diffractive}}$	$\approx 10\text{mb}$	
$\sigma_{\text{inelastic}}$	$\approx 60\text{mb}$	

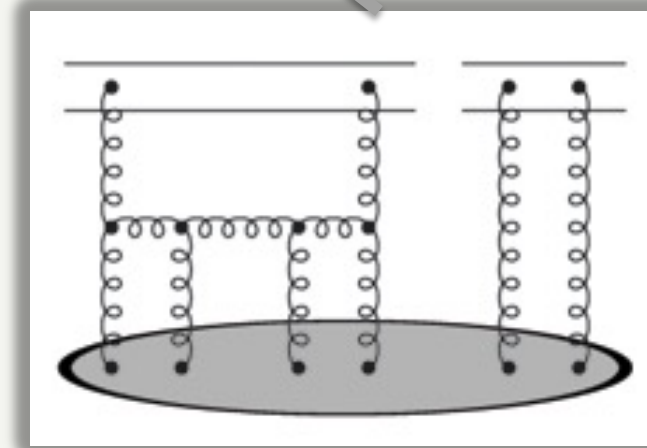
# Physics of the Vacuum

## Elastic scattering



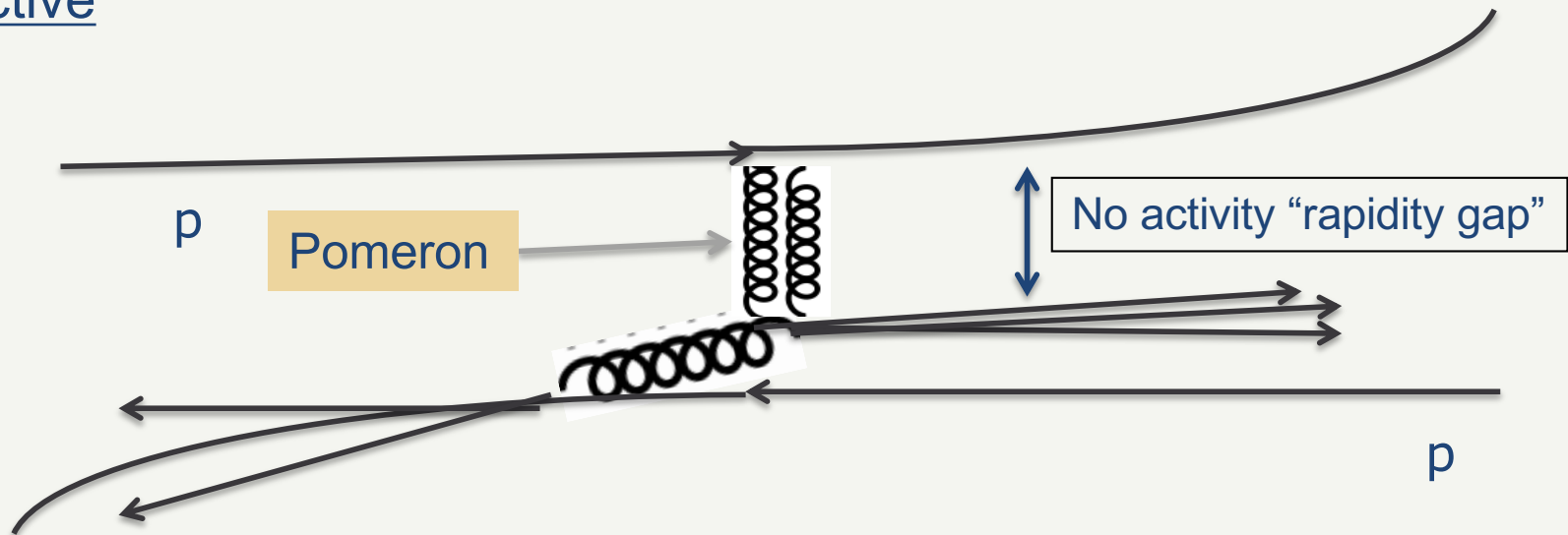
At high energy:  $A(s,t)=s^{\alpha(t)}$   
 $\alpha_P(t)=\alpha_P(0)+\alpha't$

$\sigma_{\text{elastic}} \approx 40\text{mb}$  ←  
 $\sigma_{\text{diffractive}} \approx 10\text{mb}$   
 $\sigma_{\text{inelastic}} \approx 60\text{mb}$



# Physics of the Vacuum

## Diffractive



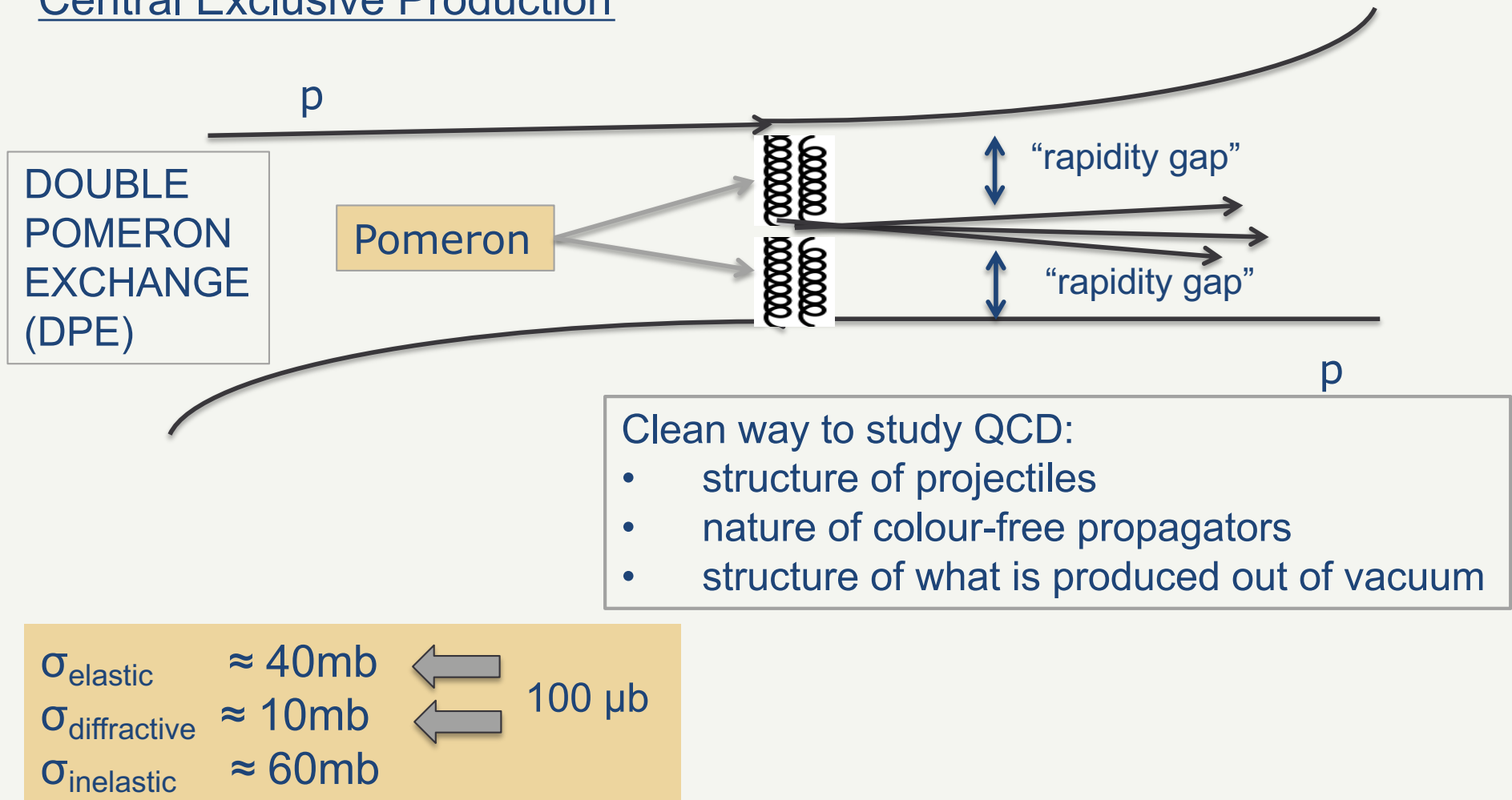
Experimental working definition of diffraction is presence of rapidity gap

$\sigma_{\text{elastic}}$	$\approx 40\text{mb}$
$\sigma_{\text{diffractive}}$	$\approx 10\text{mb}$
$\sigma_{\text{inelastic}}$	$\approx 60\text{mb}$



# Physics of the Vacuum

## Central Exclusive Production



# Physics of the Vacuum

## Central Exclusive Production

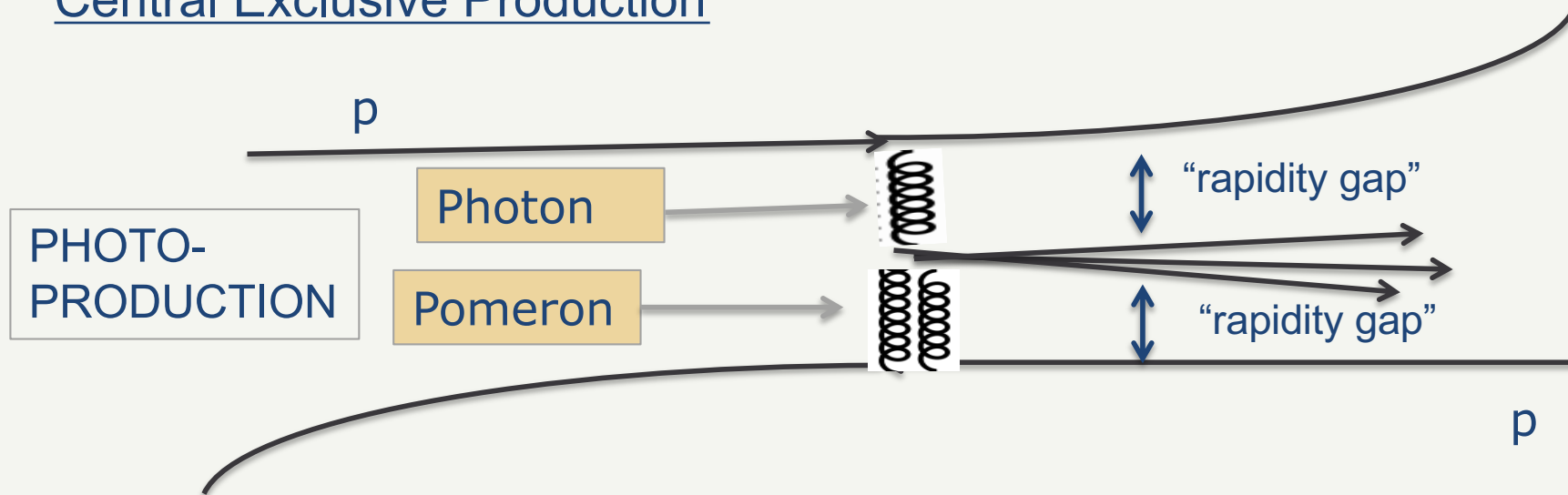


PHOTO-  
PRODUCTION

Photon

Pomeron

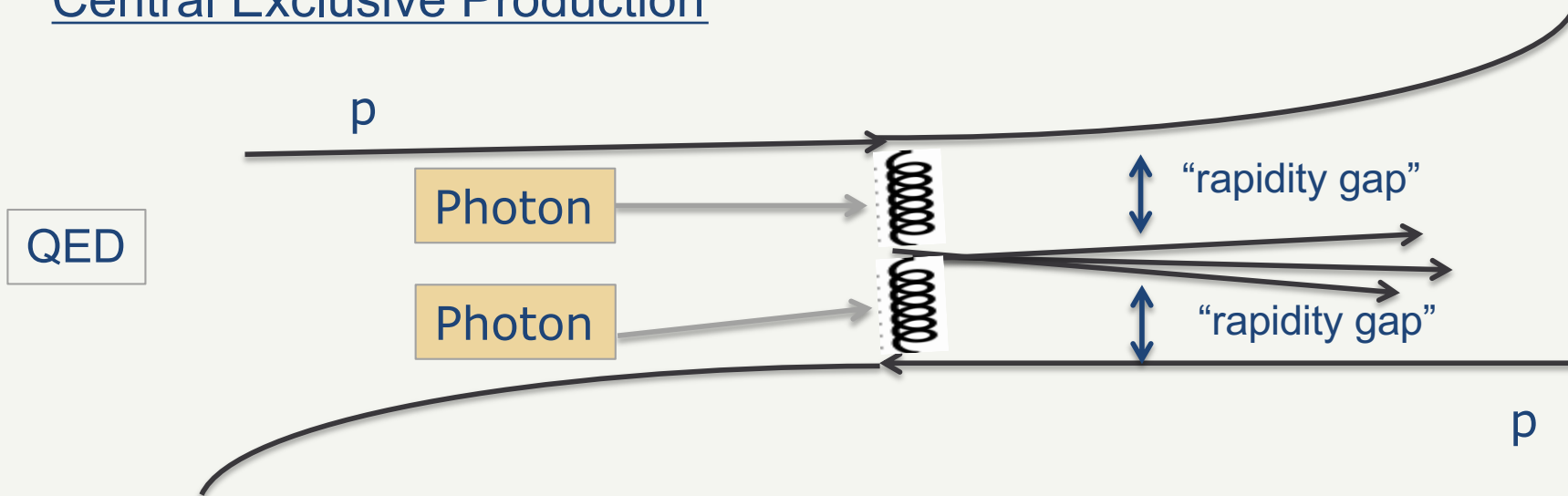
"rapidity gap"

"rapidity gap"

$\sigma_{\text{elastic}}$	$\approx 40\text{mb}$	←	100 $\mu\text{b}$
$\sigma_{\text{diffractive}}$	$\approx 10\text{mb}$	←	
$\sigma_{\text{inelastic}}$	$\approx 60\text{mb}$		

# Physics of the Vacuum

## Central Exclusive Production

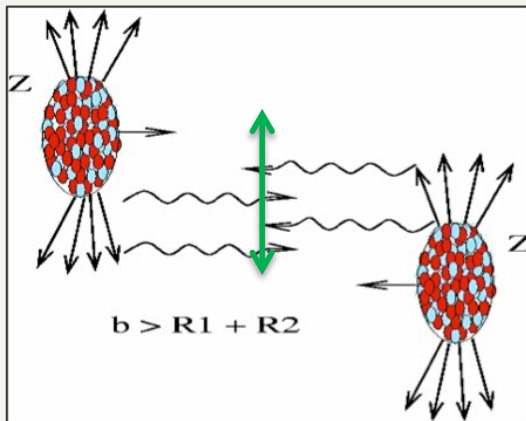
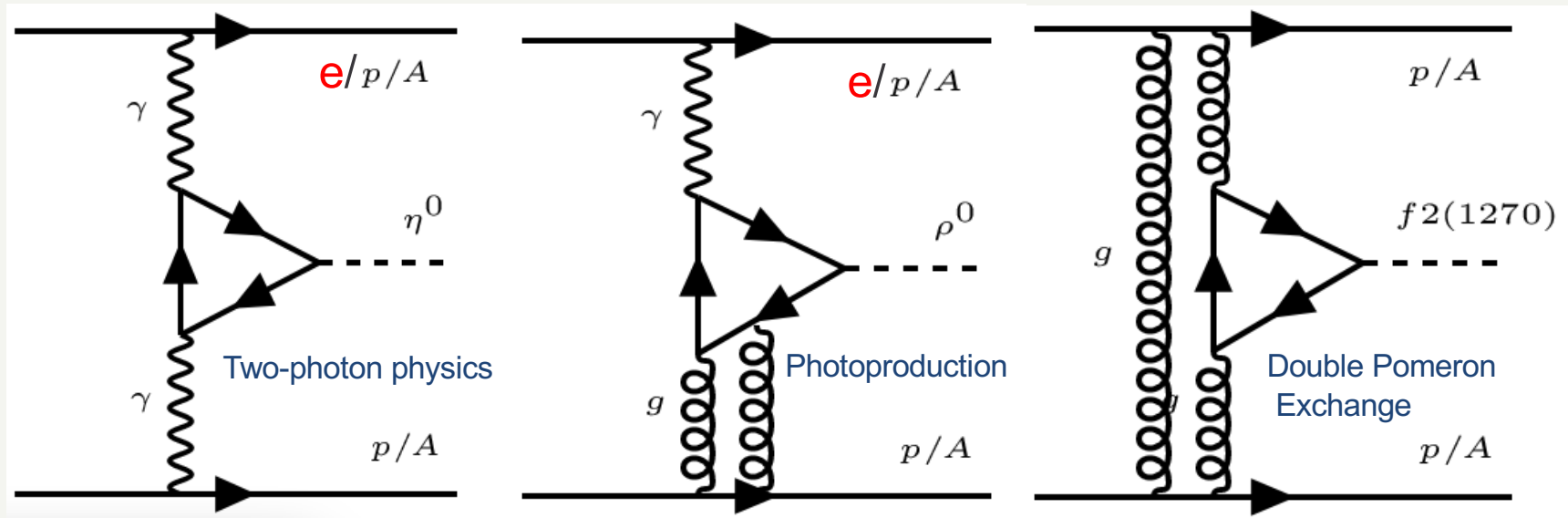


CEP is characterised by a rapidity gap all the way to the proton

Detect as large a gap as possible...

$\sigma_{\text{elastic}}$	$\approx 40\text{mb}$	←	100 pb
$\sigma_{\text{diffractive}}$	$\approx 10\text{mb}$	←	
$\sigma_{\text{inelastic}}$	$\approx 60\text{mb}$		

# Colourless propagators



## Hadron colliders:

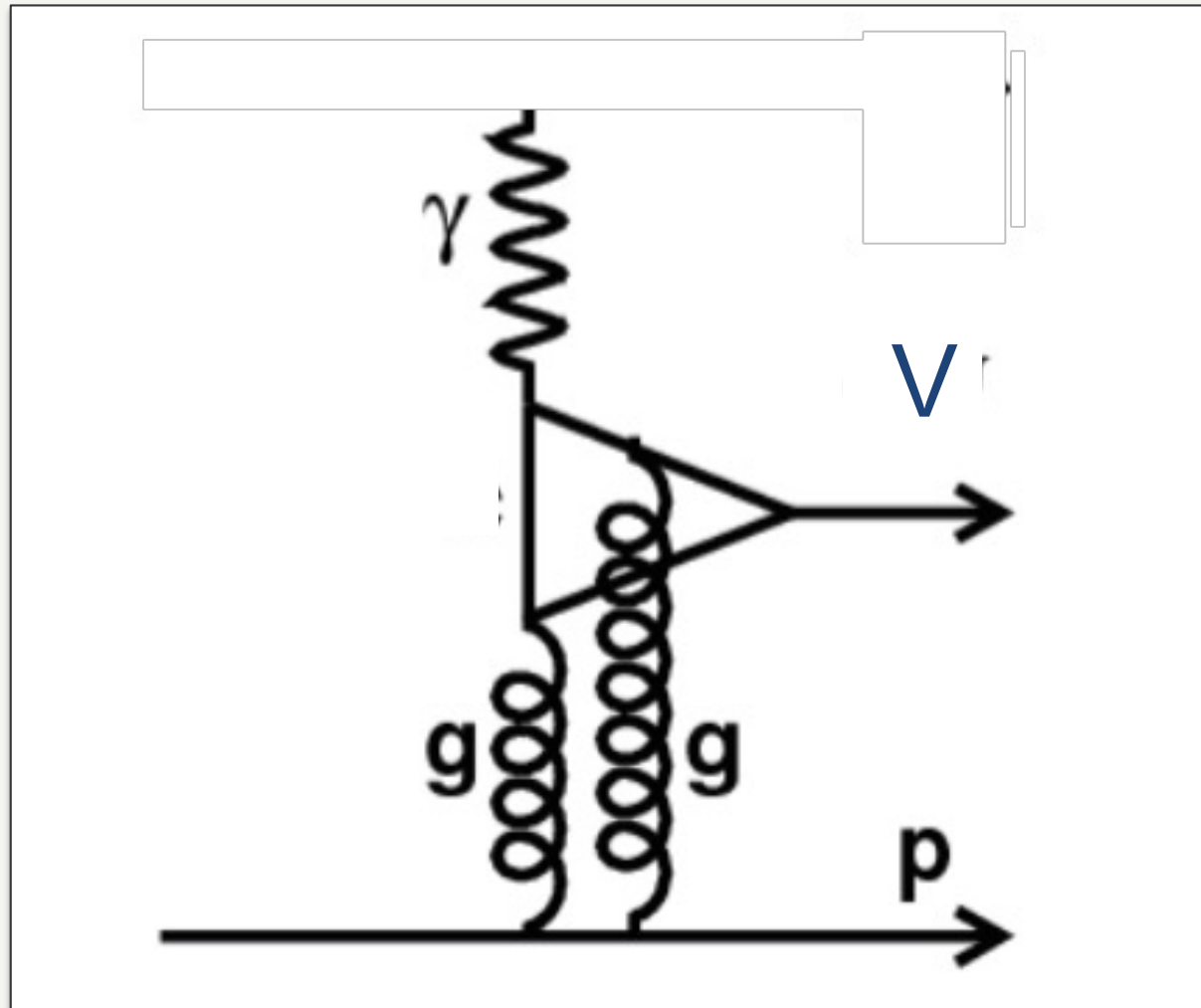
Generally, to ensure no (colourful) QCD interaction,  $d > R_1 + R_2$  (1.5 - 6 fm).

Large impact parameter  $\leftrightarrow$  Small  $p_T$

## Electron-hadron collider:

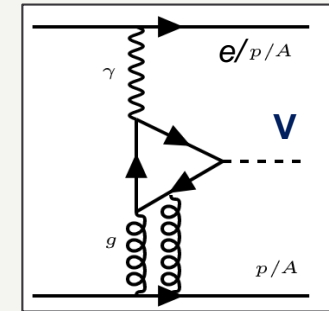
$\sim 70\%$  of total cross-section is diffractive

# Photoproduction

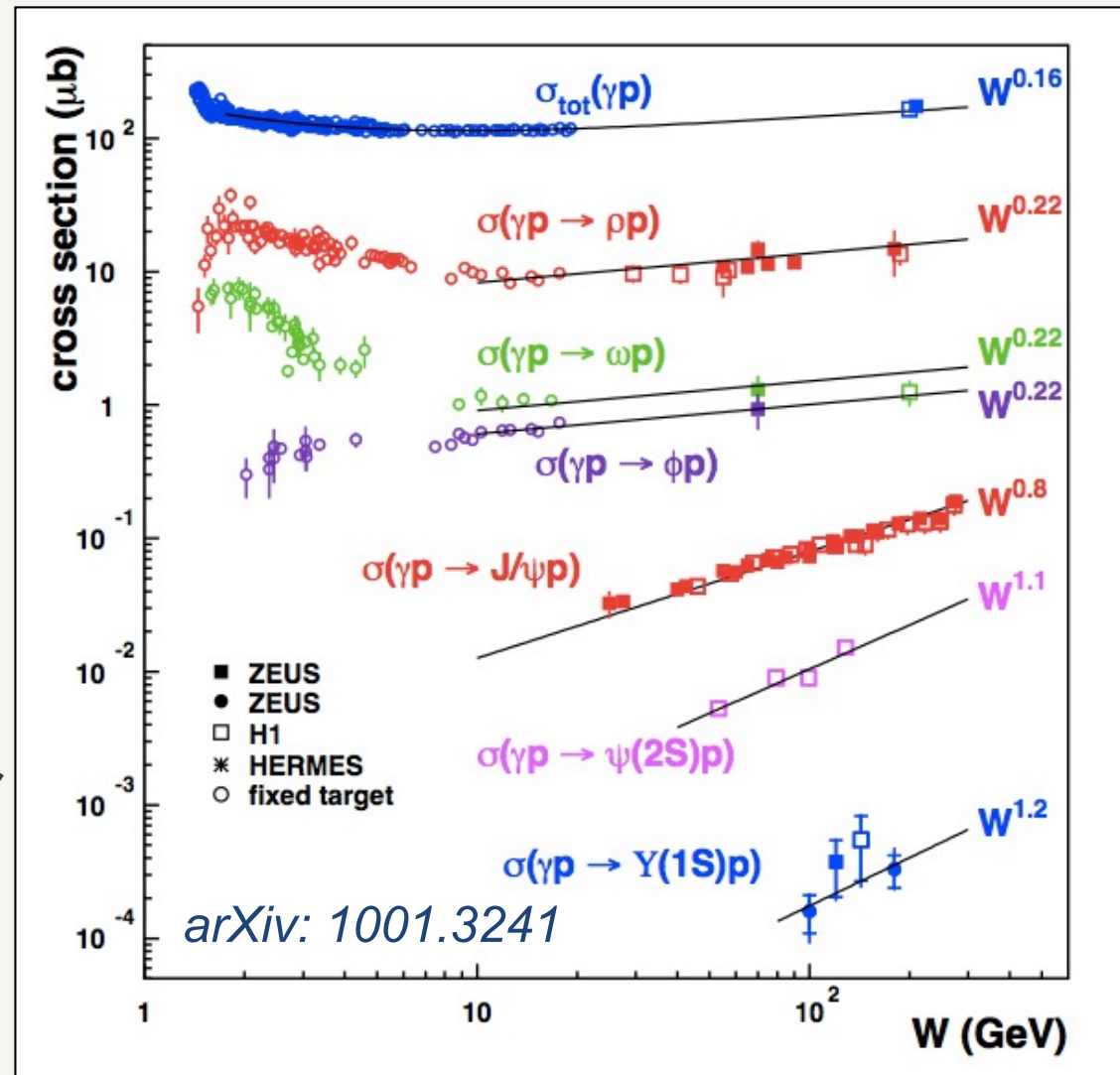


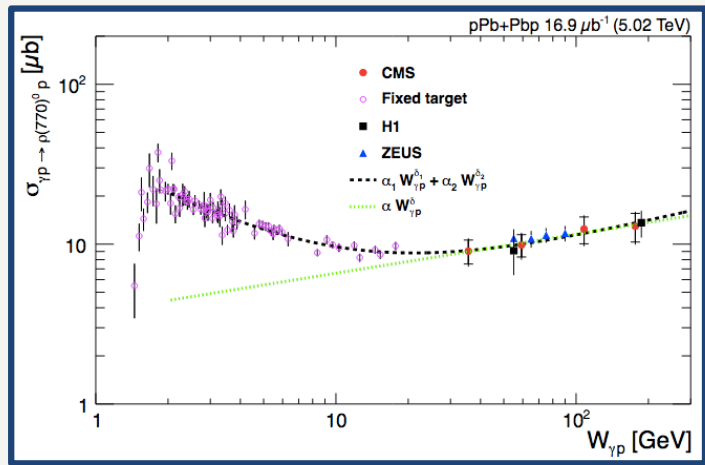


# Photoproduction



- Rise in  $\sigma$  related to Pomeron intercept
  - $\sigma \sim W^\delta$
  - $\delta = 4(\alpha_P(t) - 1)$
  - $\alpha_P(t) = \alpha_P(0) + \alpha' t$
- Compare slopes  $\rho, \omega, \phi$  to  $J/\psi, \psi', \Upsilon$
- Extract  $g(x, Q^2)$

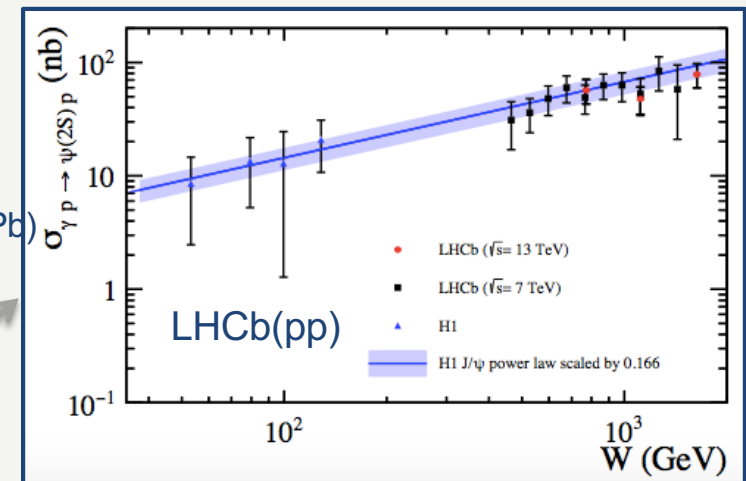
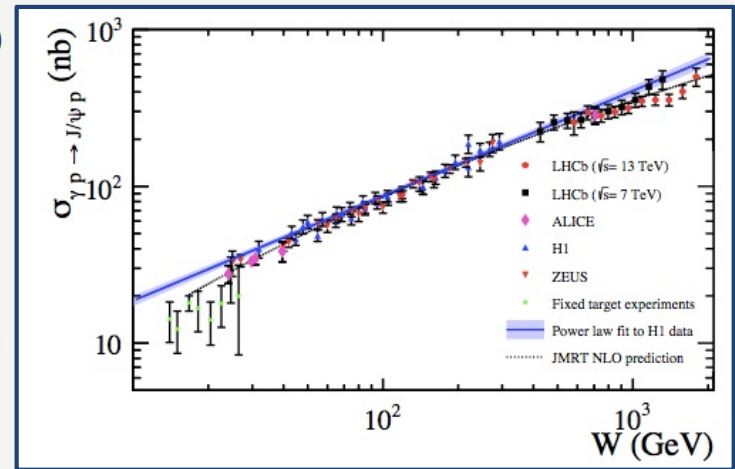




CMS (pPb) ALICE (XeXe, PbPb)

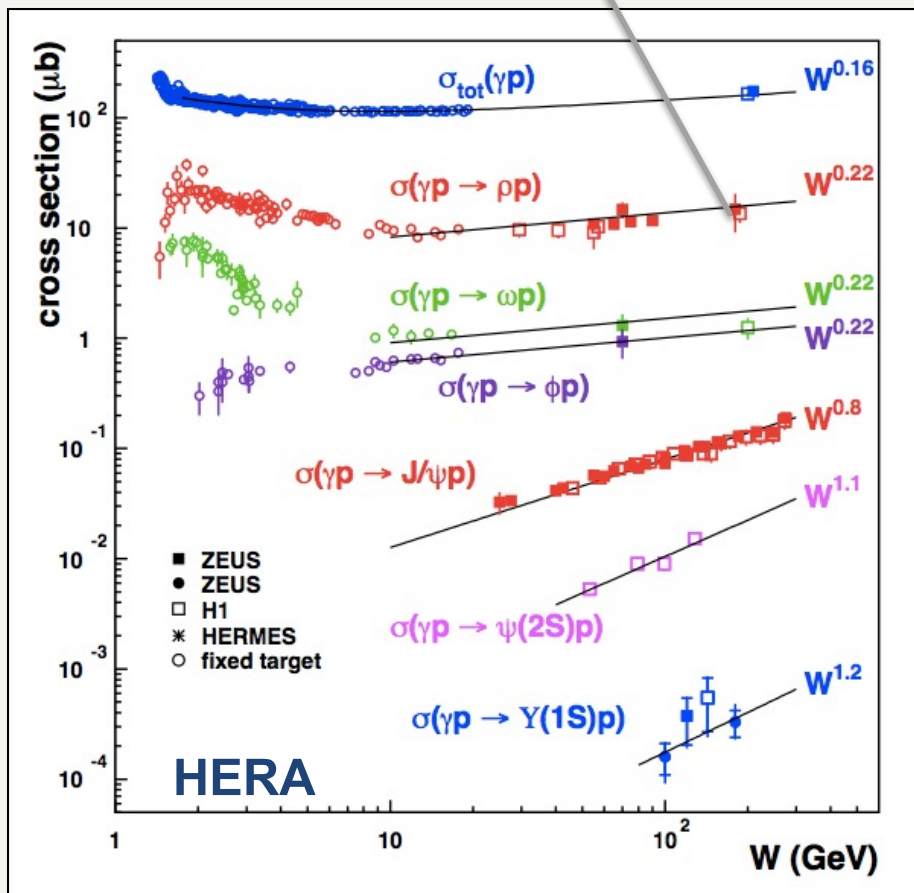
ALICE (pPb, PbPb)  
LHCb (pp, PbPb)

Central region  
 $W_{LHC} \sim W_{HERA}$   
Forward  
 $W_{LHC} \gg W_{HERA}$

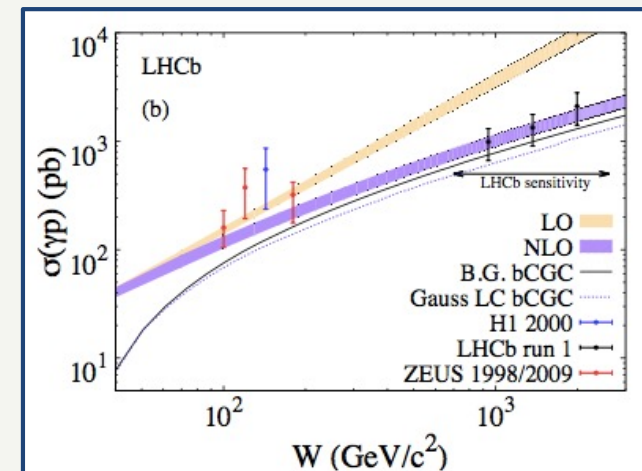


LHCb  
(pp, PbPb)

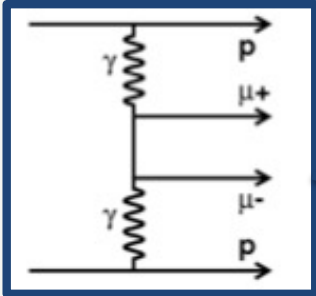
LHCb(pp)



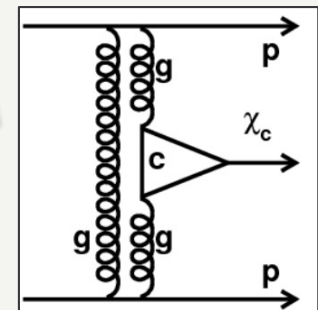
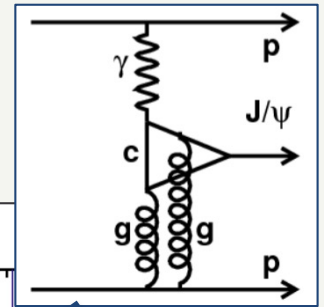
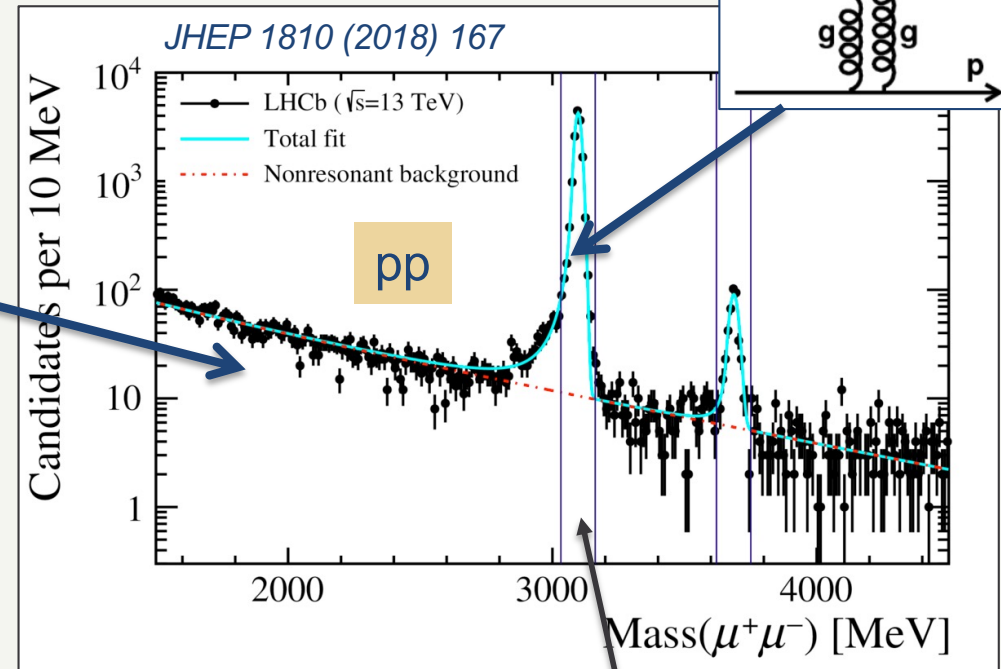
LHCb (pp)  
CMS (pPb)



# Dimuons in p(Pb)p(Pb) collisions



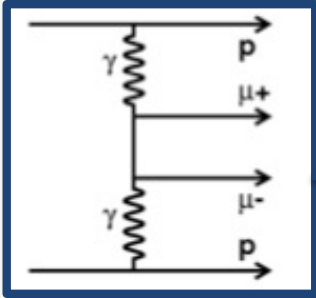
$\gamma\gamma$  events continue to detection threshold at  $\sim 600$  MeV (enhanced in PbPb)



Feed-down (not present in pA)

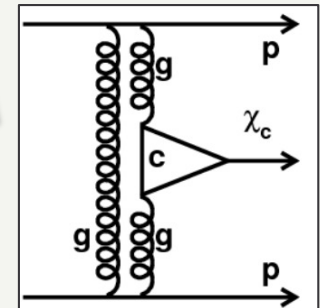
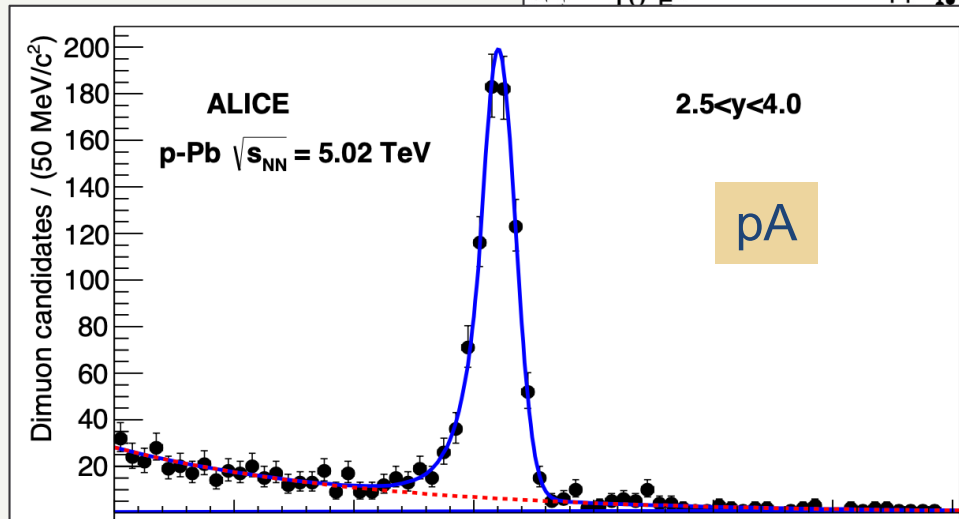
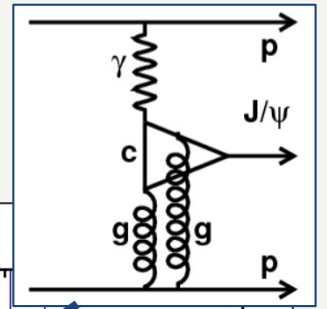
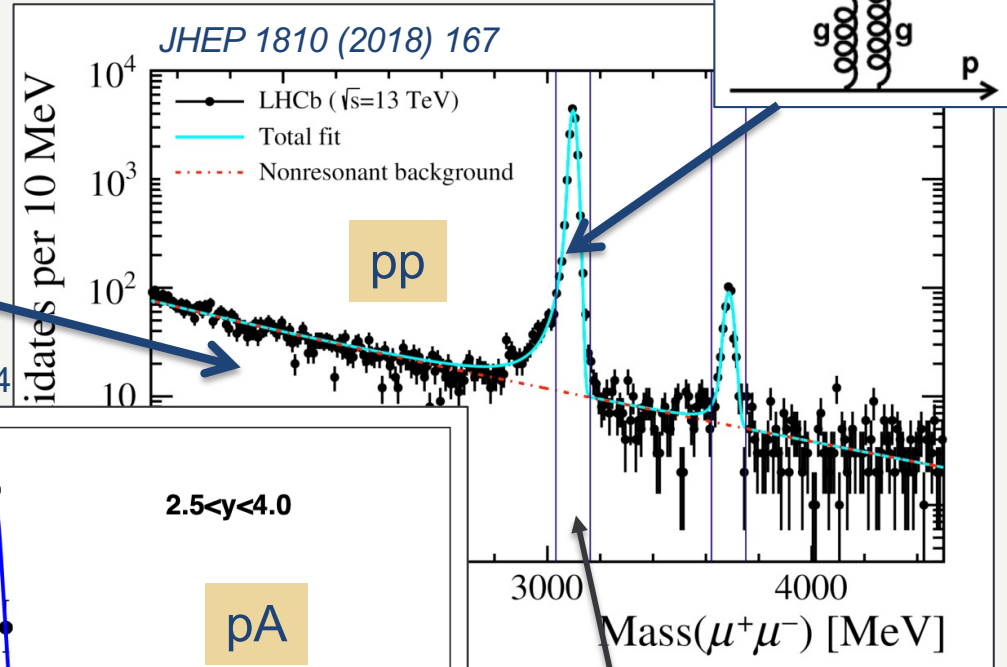
$$\chi_c \rightarrow J/\psi + \gamma$$

# Dimuons in p(Pb)p(Pb) collisions



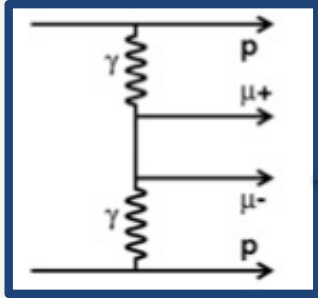
$\gamma\gamma$  events continue to detection threshold at  $\sim 600$  MeV (enhanced in PbPb)

*Phys.Rev.Lett.* 113 (2014) 23, 232504

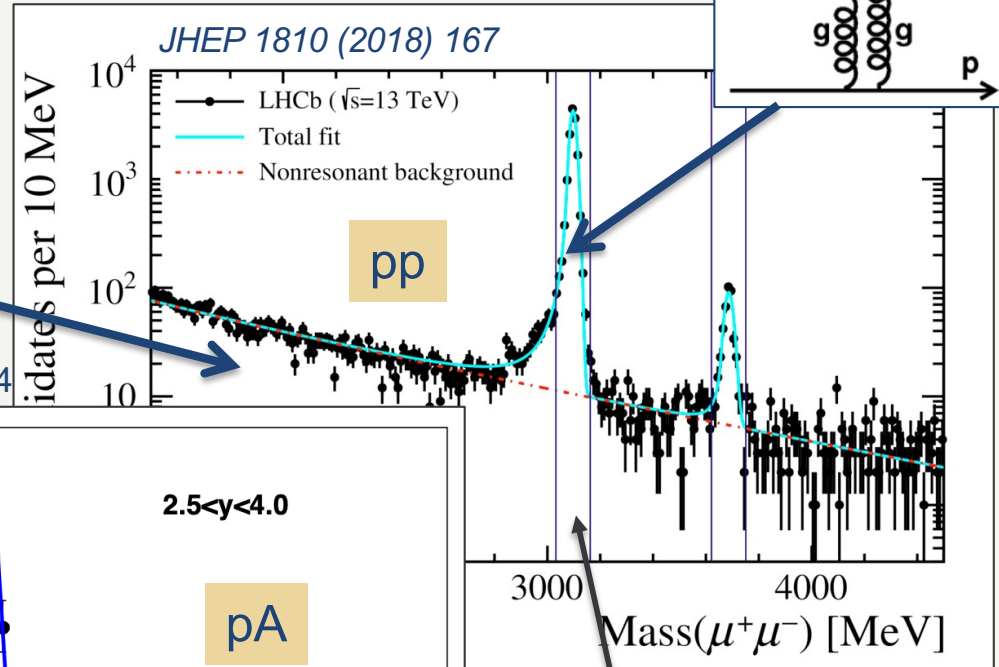
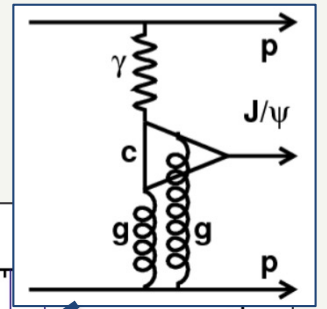


Feed-down (not present in pA)

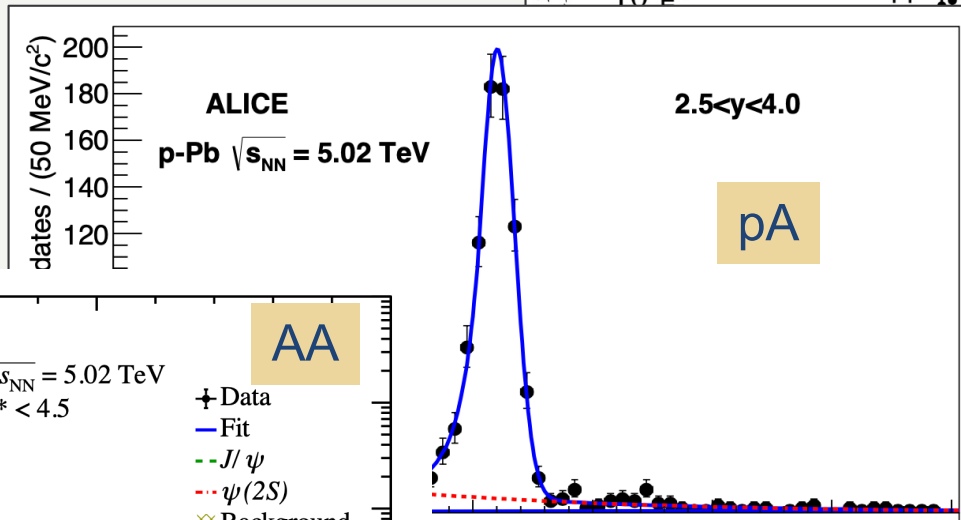
# Dimuons in p(Pb)p(Pb) collisions



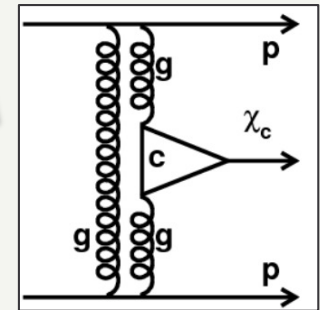
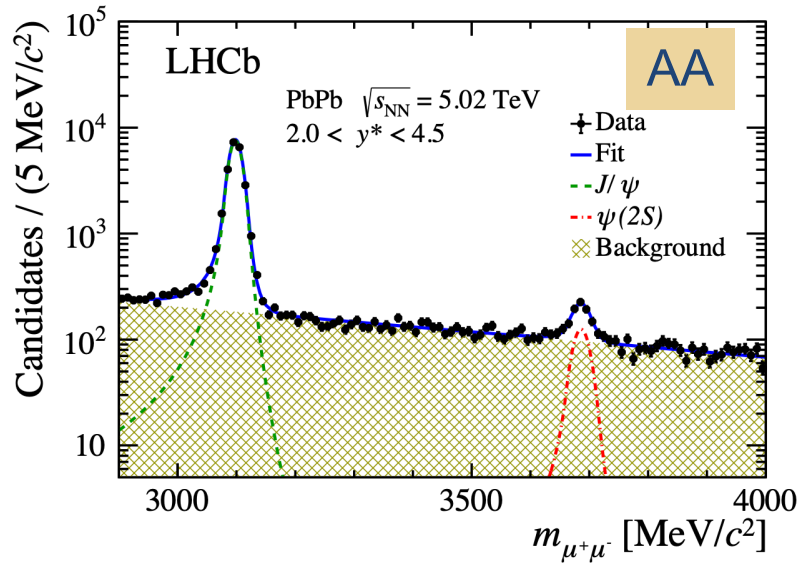
$\gamma\gamma$  events continue to detection threshold at  $\sim 600$  MeV (enhanced in PbPb)



*Phys.Rev.Lett. 113 (2014) 23, 232504*



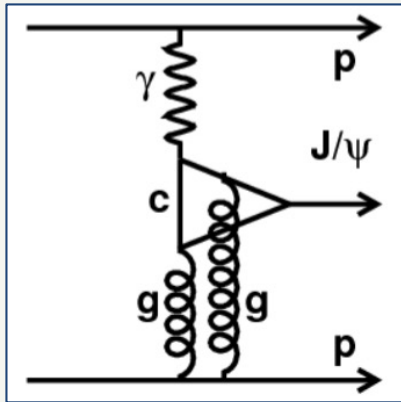
*JHEP 06 (2023) 146*



Feed-down (not present in pA)

# Implications: GPDs and PDF

Ryskin, Z. Phys. C 57 (1993) 89

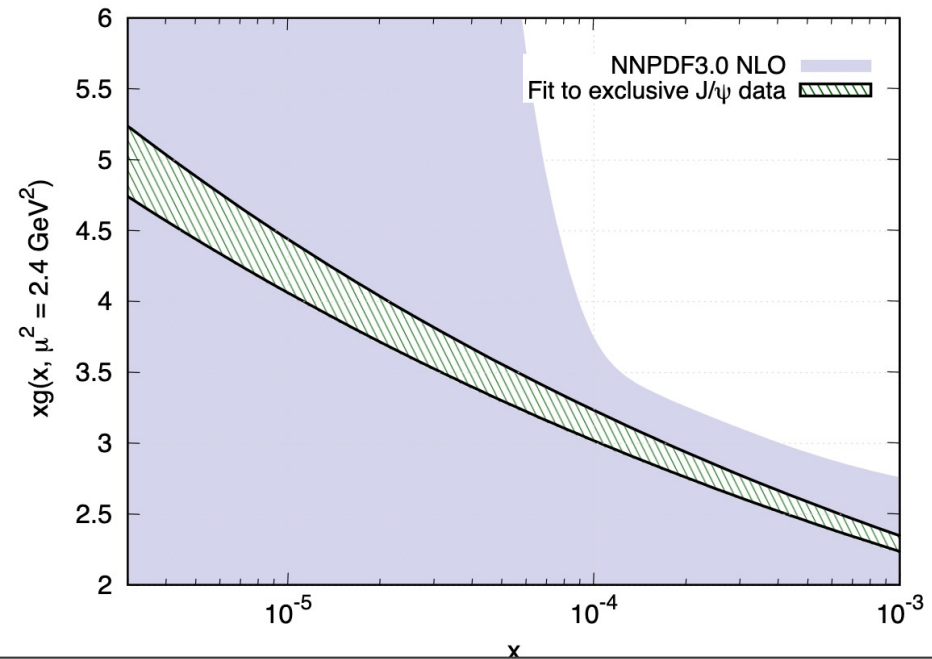
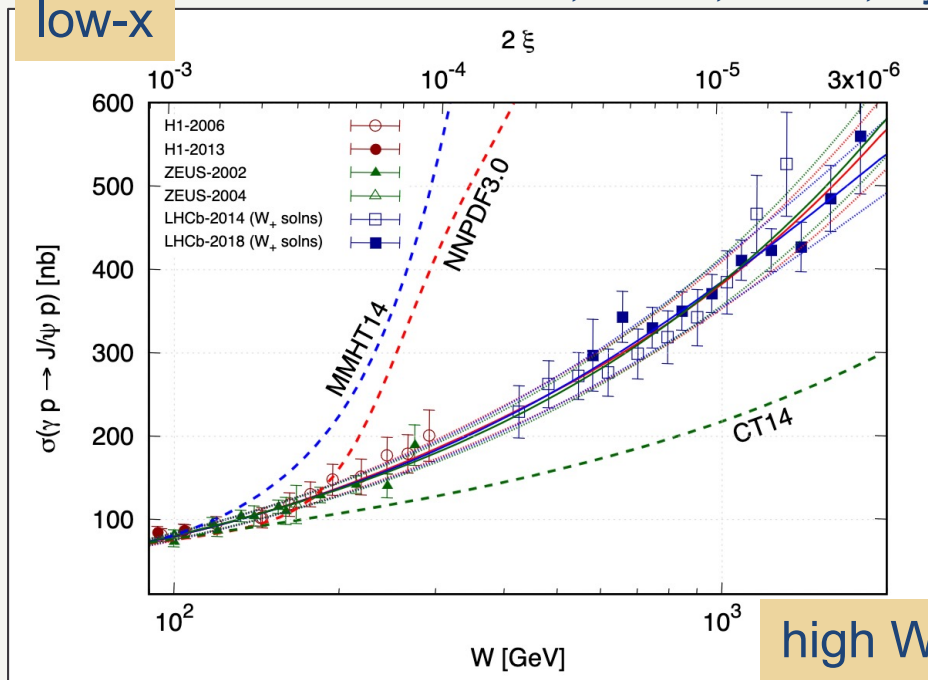


$$\frac{d\sigma}{dt} (\gamma^* p \rightarrow J/\psi p) \Big|_{t=0} = \frac{\Gamma_{ee} M_{J/\psi}^3 \pi^3}{48\alpha} \left[ \frac{\alpha_s(\bar{Q}^2)}{\bar{Q}^4} xg(x, \bar{Q}^2) \right]^2 \left( 1 + \frac{Q^2}{M_{J/\psi}^2} \right)$$

Flett, Martin, Ryskin, Teubner. Phys.Rev.D 102 (2020) 114021

Flett, Jones, Martin, Ryskin, Teubner. Phys.Rev.D 101 (2020) 9, 094011

low-x



makes use of Shuvaev transform to relate GPDs and PDFs

$$H_q(X, \xi) = \int_{-1}^1 dx' \left[ \frac{2}{\pi} \text{Im} \int_0^1 \frac{ds}{y(s)\sqrt{1-y(s)x'}} \right] \frac{d}{dx'} \left( \frac{q(x')}{|x'|} \right),$$

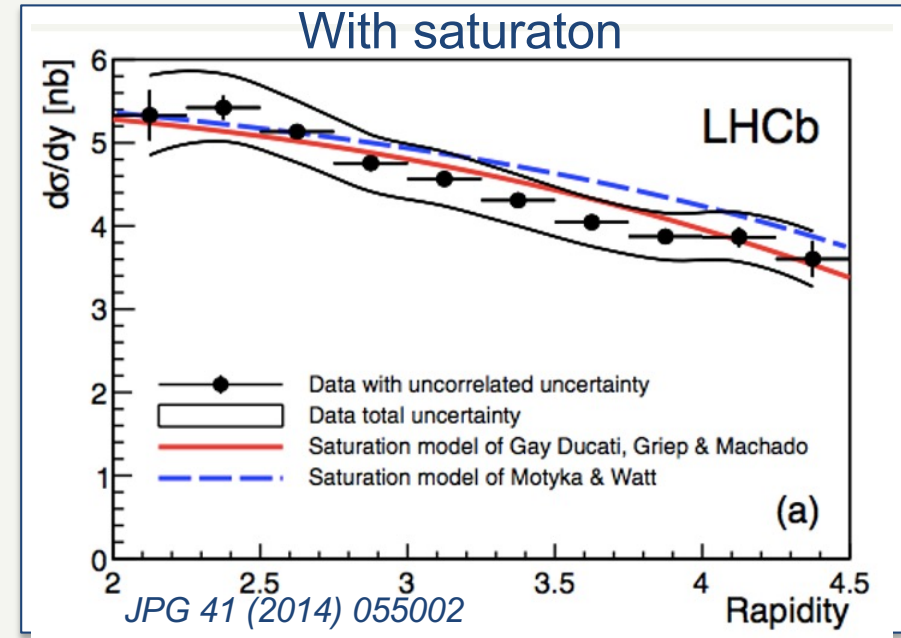
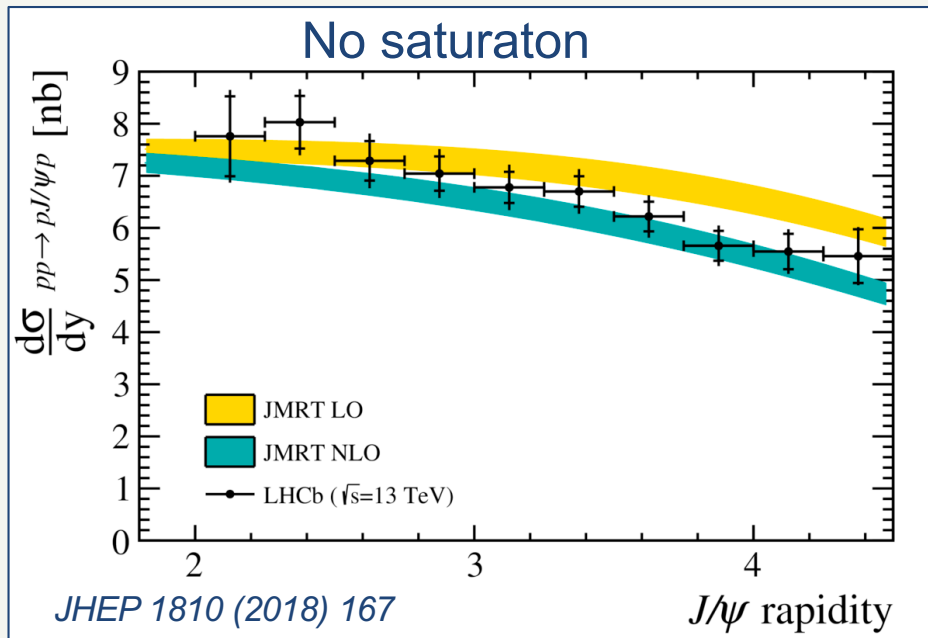
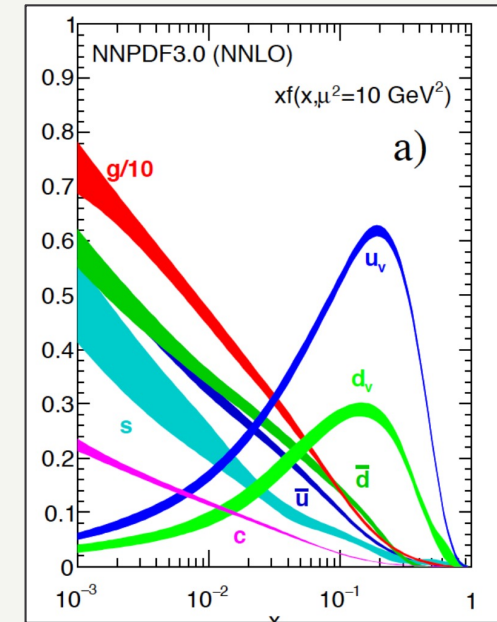
$$H_g(X, \xi) = \int_{-1}^1 dx' \left[ \frac{2}{\pi} \text{Im} \int_0^1 \frac{ds (X + \xi(1-2s))}{y(s)\sqrt{1-y(s)x'}} \right] \frac{d}{dx'} \left( \frac{g(x')}{|x'|} \right),$$

where the transform kernel,

$$y(s) = \frac{4s(1-s)}{(X + \xi(1-2s))}.$$

# Implications: Saturation

Saturation effects become visible at low-x.  
 Onset of saturation expected to scale with nucleon density  $\sim A^{1/3}$  so  
**may be easier to see in nuclear collisions**



Saturation is not inconsistent with the data, but is also not required.

# Looking for saturation in nuclear collisions

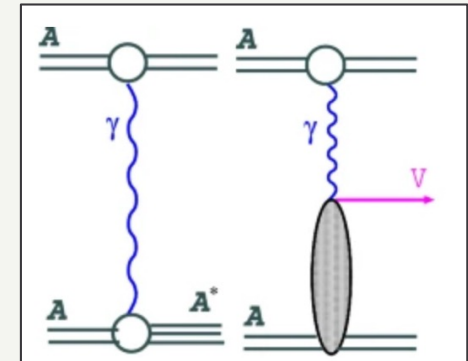
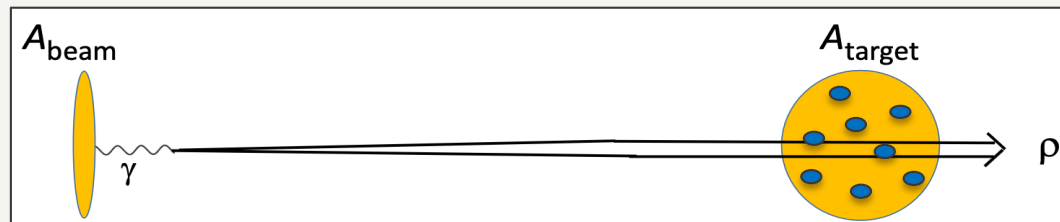
**Coherent** interaction: all nucleons behave as one.

- $b \sim 2R=13.2$  fm so  $p_T \sim 15$  MeV
- nucleus remains intact\*.

All things being equal,  $\sigma_{\gamma A \rightarrow VA} = N_A \sigma_{\gamma p \rightarrow Vp}$

Saturation would decrease cross-section at high-W (low-x)  
Nuclear suppression observed...

How much is due to saturation and how much to 'nuclear effects'?



\*additional EMD can excite or break nucleus

$$\mathcal{A}_{\rho n}(b) = i(1 - e^{-\Omega(b)/2})$$

$$\Omega(b) = T_A(b) \sigma_{\rho n} \eta$$

Glauber eikonal approx.

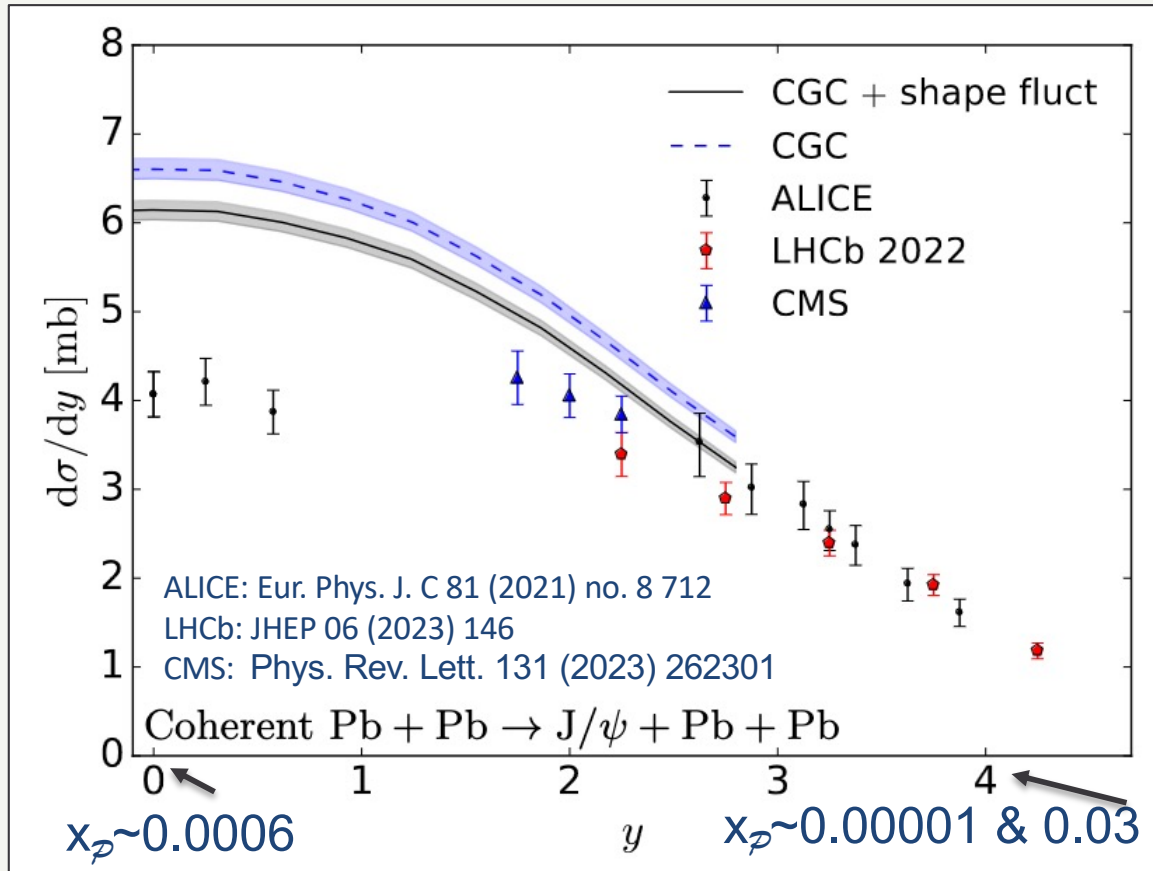
**Incoherent** interaction with nucleon or parton

- $p_T$  distribution follows  $\exp(bt)$   $b$  smaller than for coherent
- break-up is observed
- sensitive to smaller structures – saturation gives deviations from isotropy.



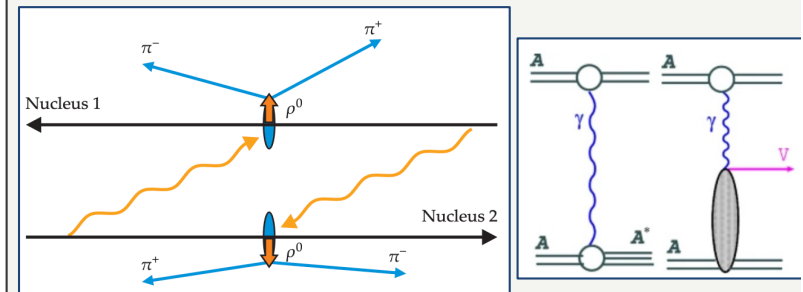
# Coherent J/ψ in PbPb

See WG2 talk by  
Heikki Mäntysaari



H. Mäntysaari, F. Salazar, B. Schenke:  
arXiv: 2312.04194

“We predict strong saturation-driven nuclear suppression at high energies, while LHC data prefers even stronger suppression.”

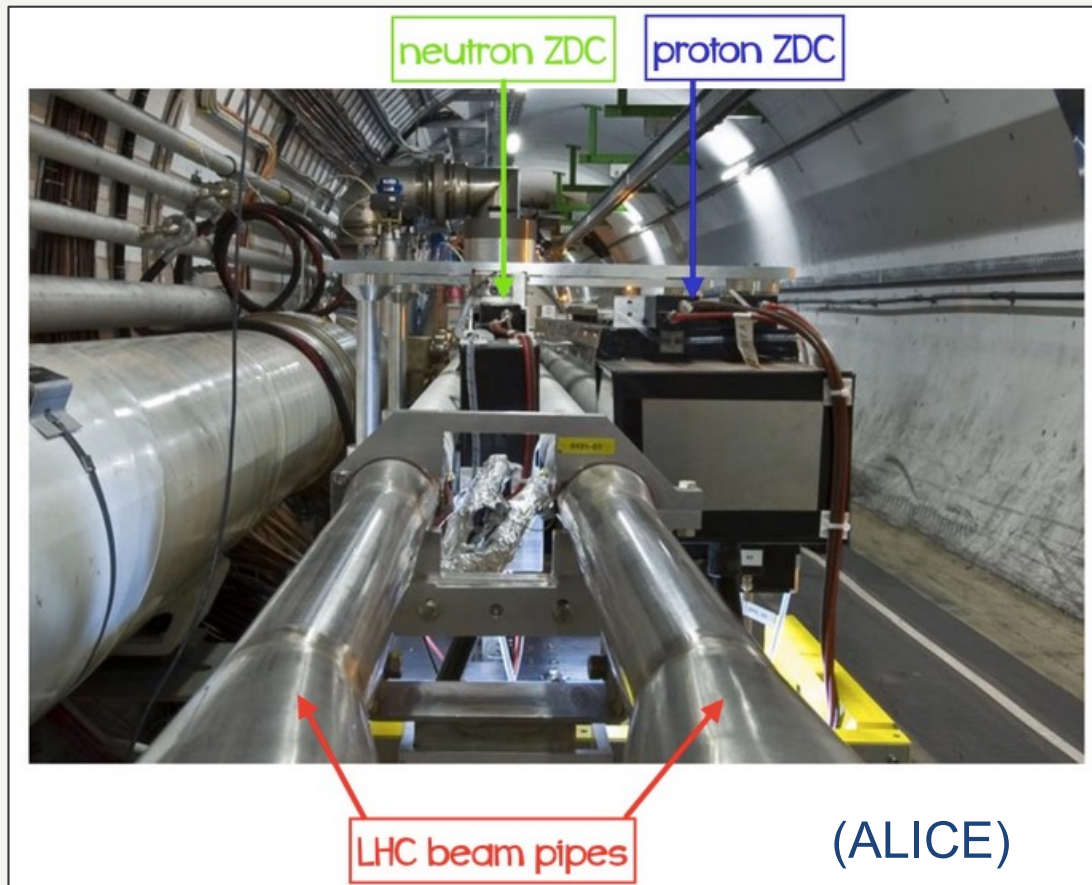


S. Klein, J. Nystrand, *Physics Today* 70, (2017) 40.

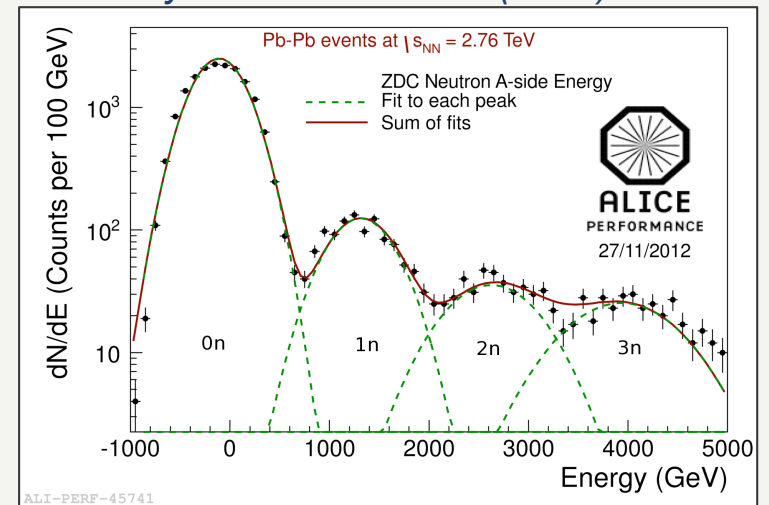
However, away from  $y=0$ , there is a two-fold ambiguity in the photon emitter and two-fold ambiguity in the value of  $W$ .

$$\frac{d\sigma_{PbPb \rightarrow PbJ/\psi Pb}}{dy} = \left( k \frac{dN_\gamma}{dk} \right)^+ \sigma_{\gamma Pb \rightarrow J/\psi Pb}(W^+) + \left( k \frac{dN_\gamma}{dk} \right)^- \sigma_{\gamma Pb \rightarrow J/\psi Pb}(W^-)$$

# ZDC calorimeters installed in CMS, ALICE, STAR



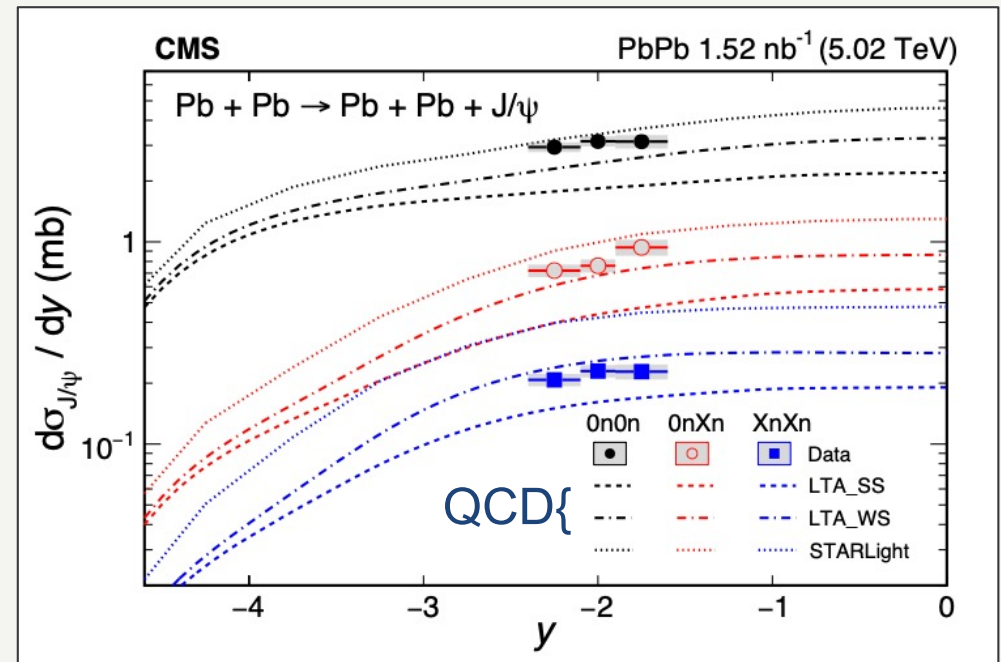
*J. Phys.: Conf. Ser. 455 (2013) 012010*



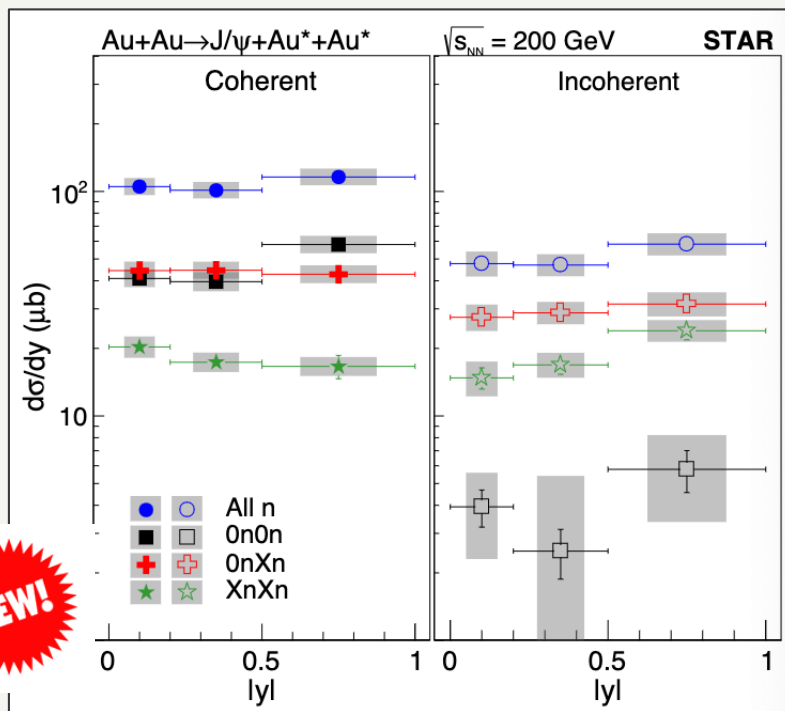
Detection of neutrons when ion breaks up allows identification of Electromagnetic Dissociation (EMD)

# Resolving the two-fold ambiguity in PbPb

EMD is more likely at small impact parameters.  
So fluxes for 0n and (X>=1)n different.

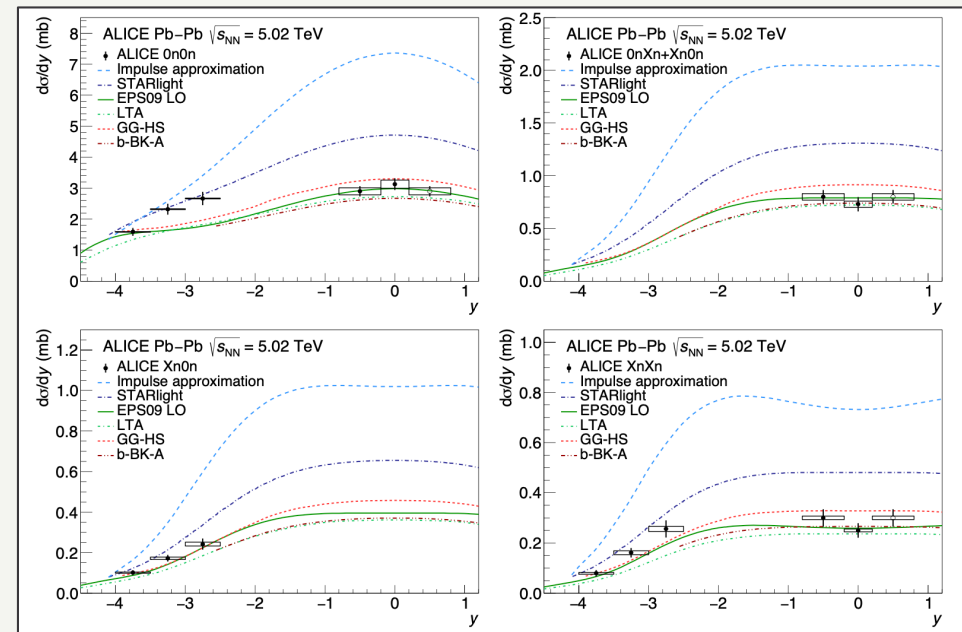


Phys. Rev. Lett. 131 (2023) 262301



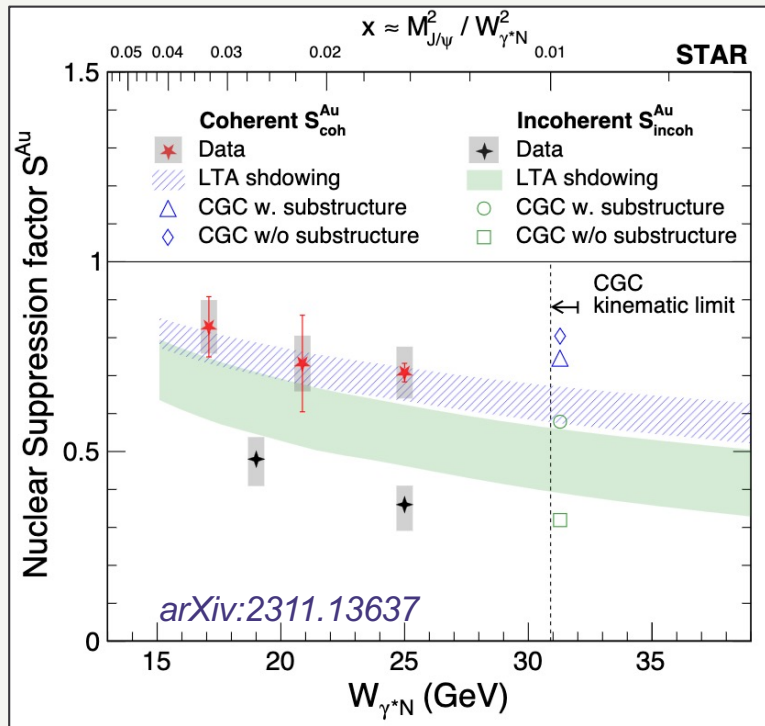
**NEW!**

arXiv: 2311.13632



IHEP 10 (2023) 119

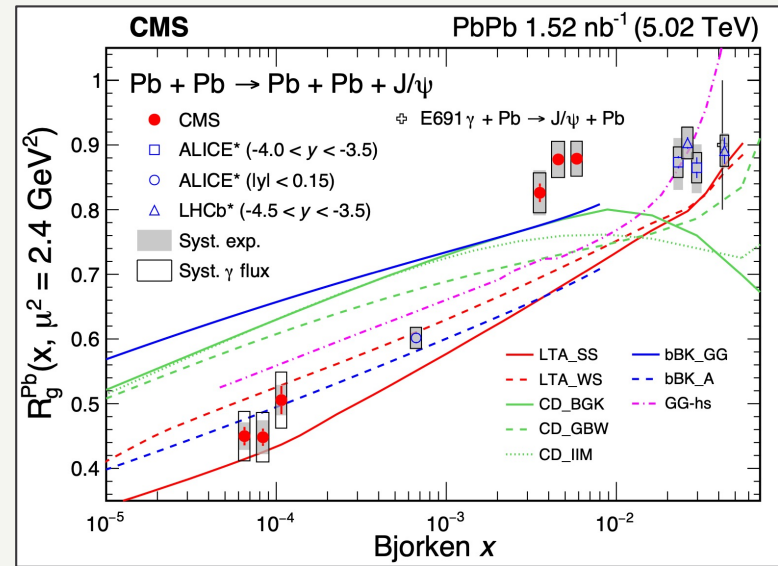
# Re-expressed in terms of nuclear suppression factors



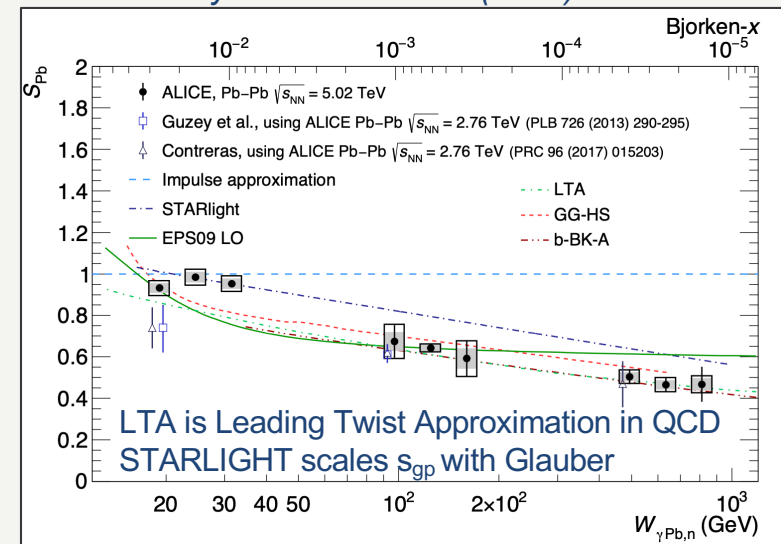
$$S_{coh} = \sqrt{\frac{\sigma_{\gamma A}}{\sigma_{IA}}}$$

IA is simple  $N_A$  scaling of  $\sigma_{\gamma p}$

None of the models does a perfect job.  
 QCD/Starlight not too bad.  
 Models with saturation also reasonable



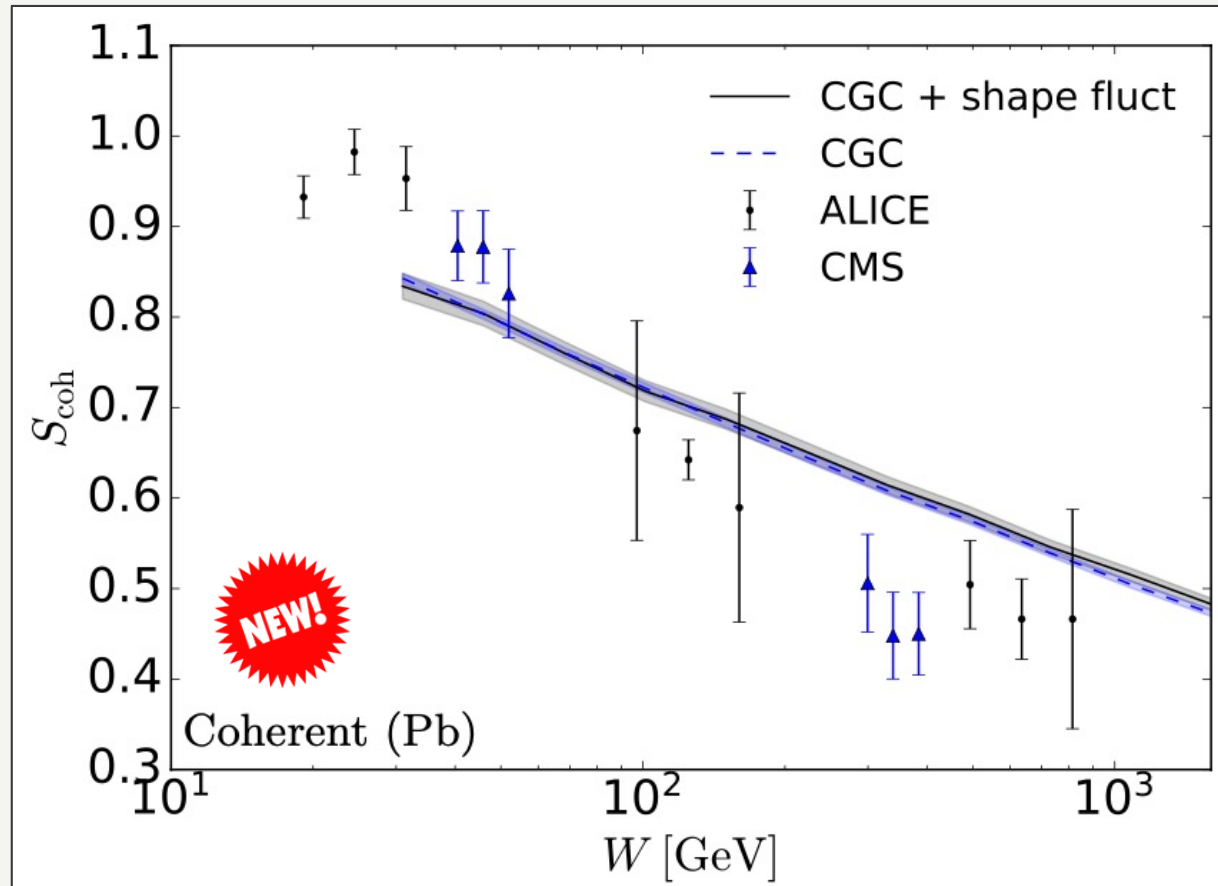
Phys. Rev. Lett. 131 (2023) 262301



JHEP 10 (2023) 119

# The case for saturation....

$$S_{\text{coh}} = \sqrt{\frac{\sigma^{\gamma A}}{\sigma^{\text{IA}}}}$$



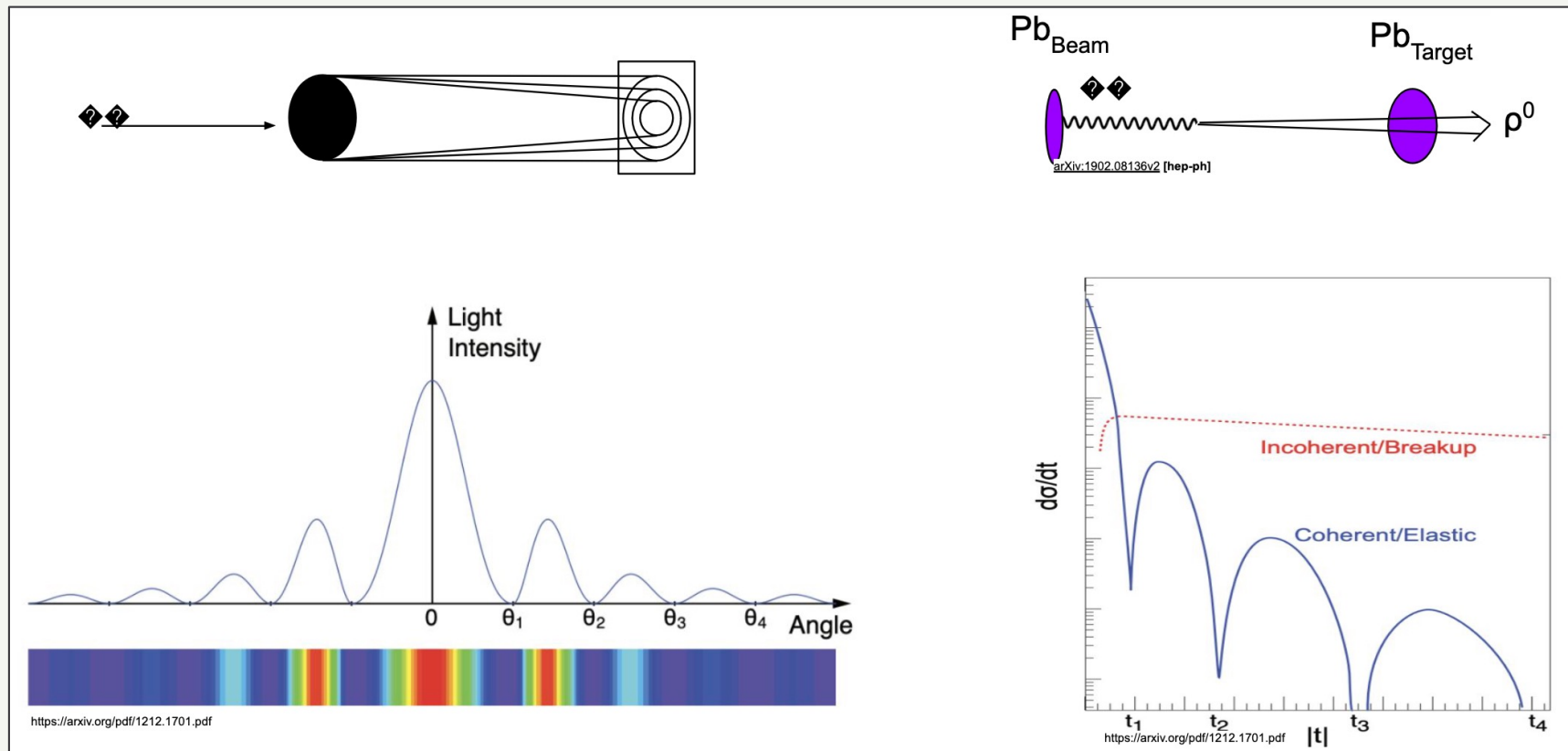
H. Mäntysaari, F. Salazar, B. Schenke: arXiv: 2312.04194

See WG2 talk by Heikki Mäntysaari

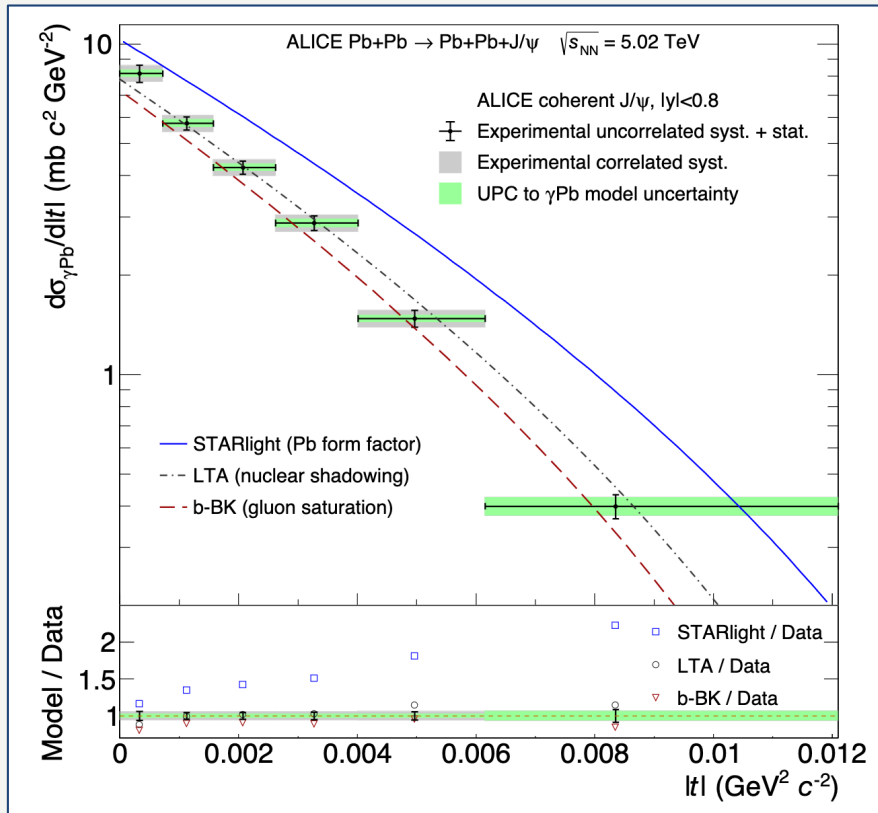
# Transverse momentum distribution

$d\sigma/dt \sim \exp(-bt)$  with  $b \sim 6 \text{ GeV}^2$  in pp collisions and  $\sim 400 \text{ GeV}^2$  in PbPb. (F.T. of ip)

$\langle p_T \rangle$  in pp and pA is  $\sim 0.5 \text{ GeV} = 1/R_p$   
 $\langle p_T \rangle$  in AA.  $\sim 0.05 \text{ GeV} \sim 1/R_A$



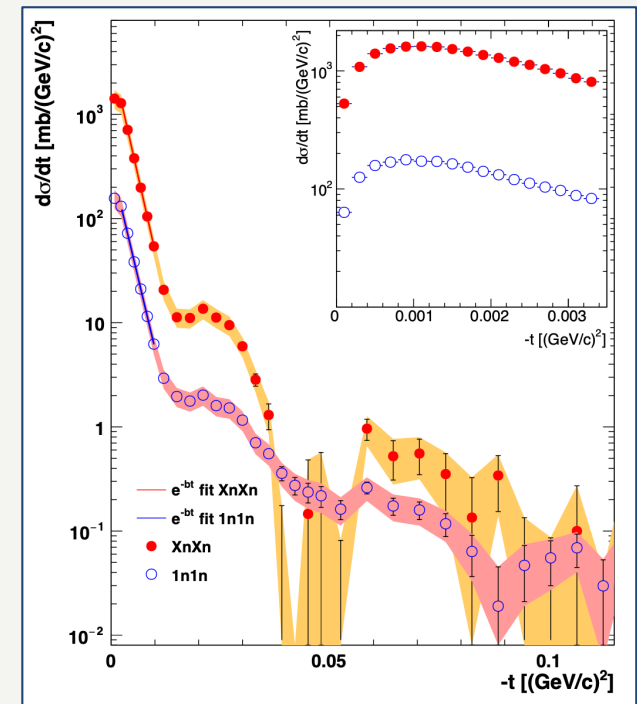
# First measurements in UPC



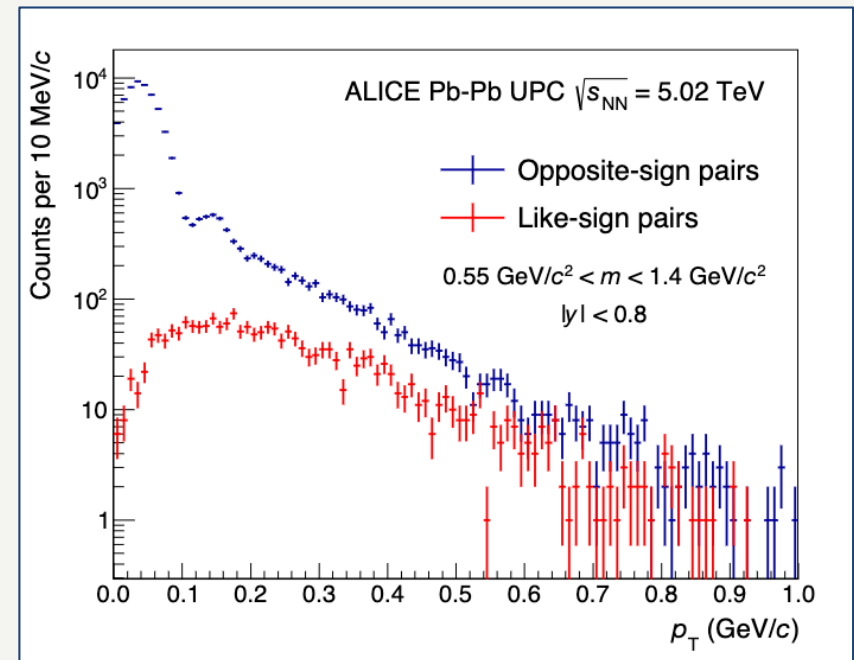
Phys.Lett.B 817 (2021) 136280

see WG2 talk by David Grund

$\rho$



Phys.Rev. C96 (2017) 054904



JHEP 06 (2020) 035

# Incoherent scatters also interesting

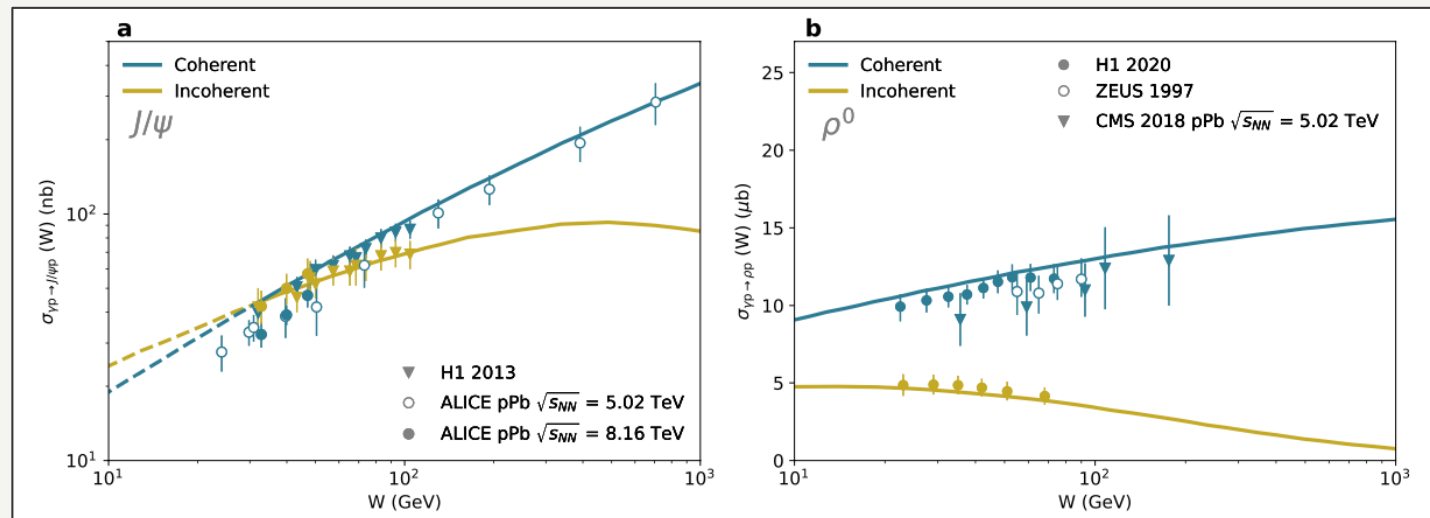
Intact target: sensitive to **average** colour  
 Breakup: sensitive to fluctuations (**rms**)

$$\frac{d\sigma^{\gamma^* p \rightarrow J/\Psi p}}{dt} = \frac{1}{16\pi} |\langle A(x_{\mathbb{P}}, Q^2, \Delta) \rangle|^2$$

$$\frac{d\sigma^{\gamma^* N \rightarrow J/\Psi N^*}}{dt} = \frac{1}{16\pi} (\langle |A(x_{\mathbb{P}}, Q^2, \Delta)|^2 \rangle - |\langle A(x_{\mathbb{P}}, Q^2, \Delta) \rangle|^2)$$

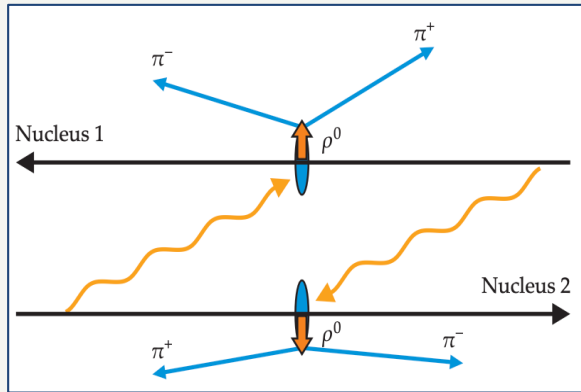
Mäntysaari, Schenke, Phys. Rev. Lett. 117, 052301 (2016)

- Original model based on gluonic fluctuations around three hot-spots (valence quarks)
- Hot-spot evolution model ( [see WG2 talk by Tobias Toll](#) )
- Energy-dependent hot-spots (J. Cepila, J. G. Contreras, J. D. Tapia Takaki Phys. Lett. B766 (2017) 186–191)
- The onset of saturation? (J. Cepilaa, J. G. Contrerasa, M. Matasa, A. Ridzikova, arXiv:2313.11320)

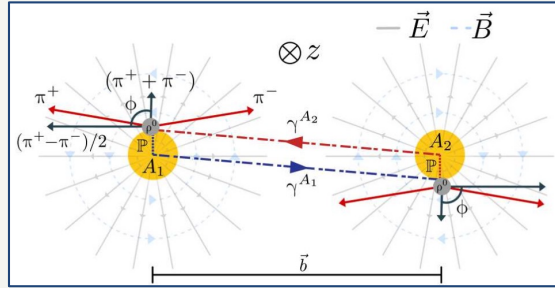




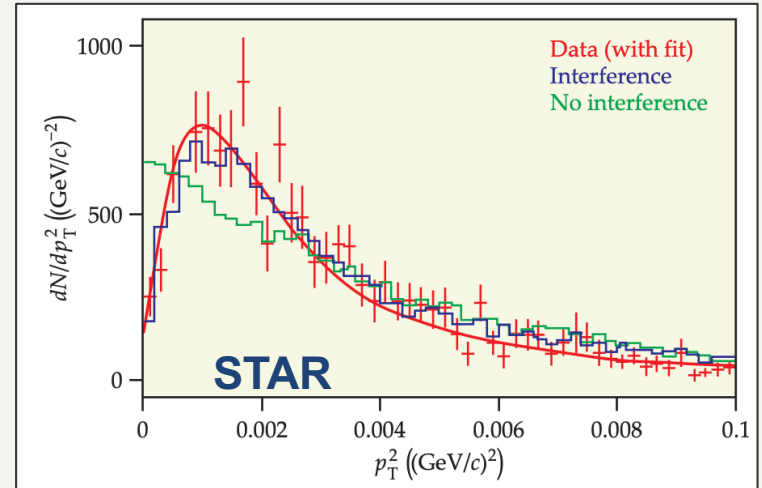
# The two-fold photon ambiguity...



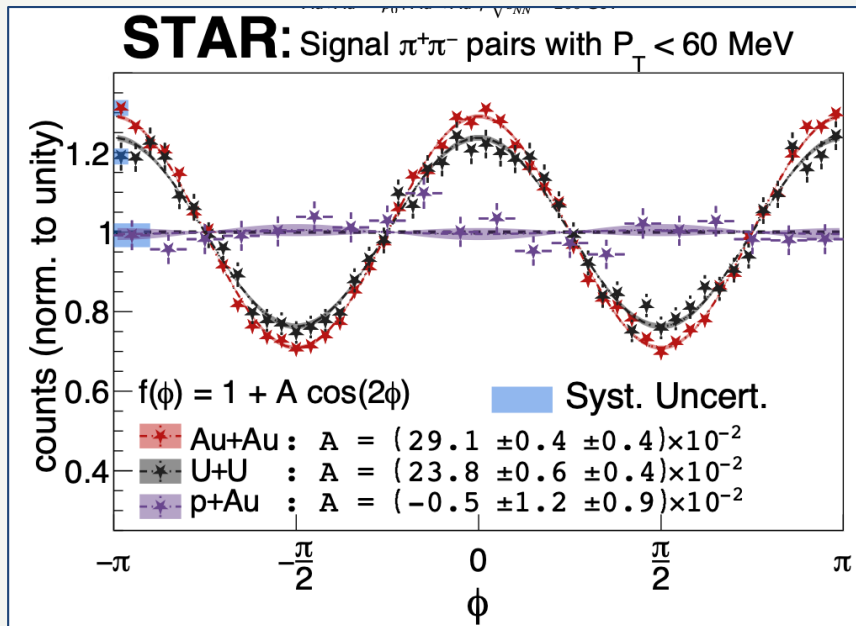
S. Klein, J. Nystrand,  
Physics Today 70, (2017) 40.



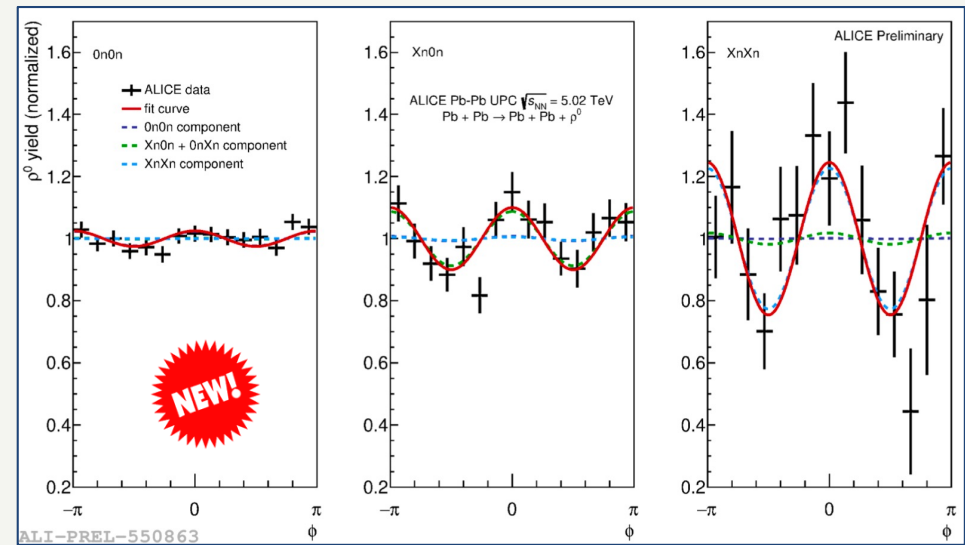
Sci.Adv. 9 (2023) eabq3903



Phys. Rev. Lett. 102, 112301 (2009)

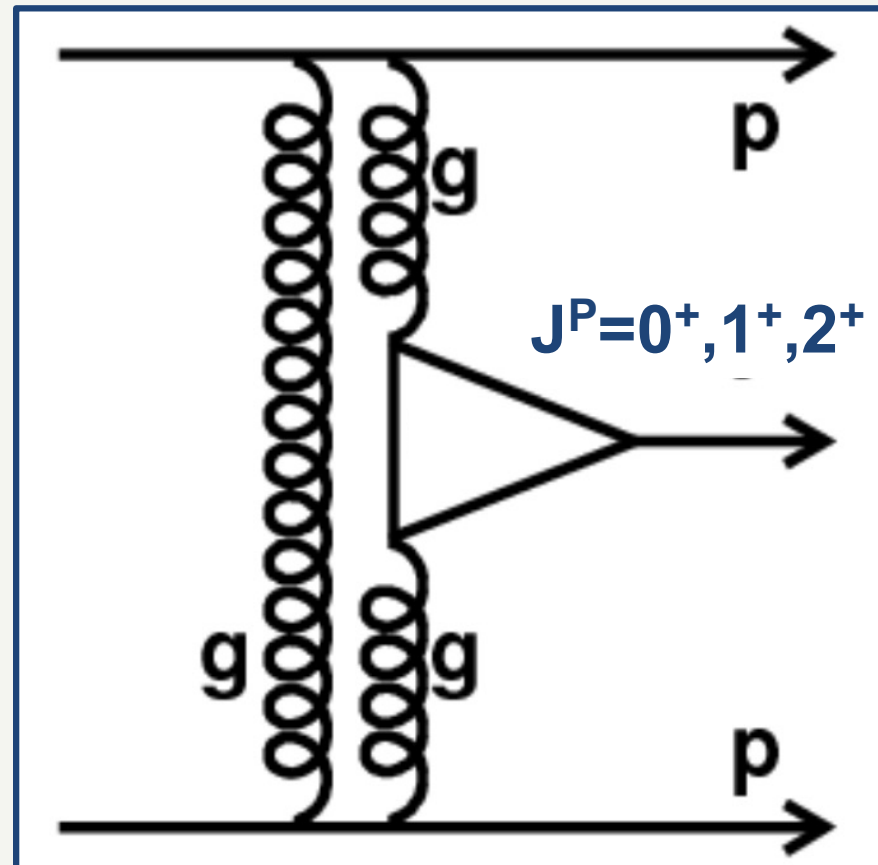


Sci.Adv. 9 (2023) eabq3903



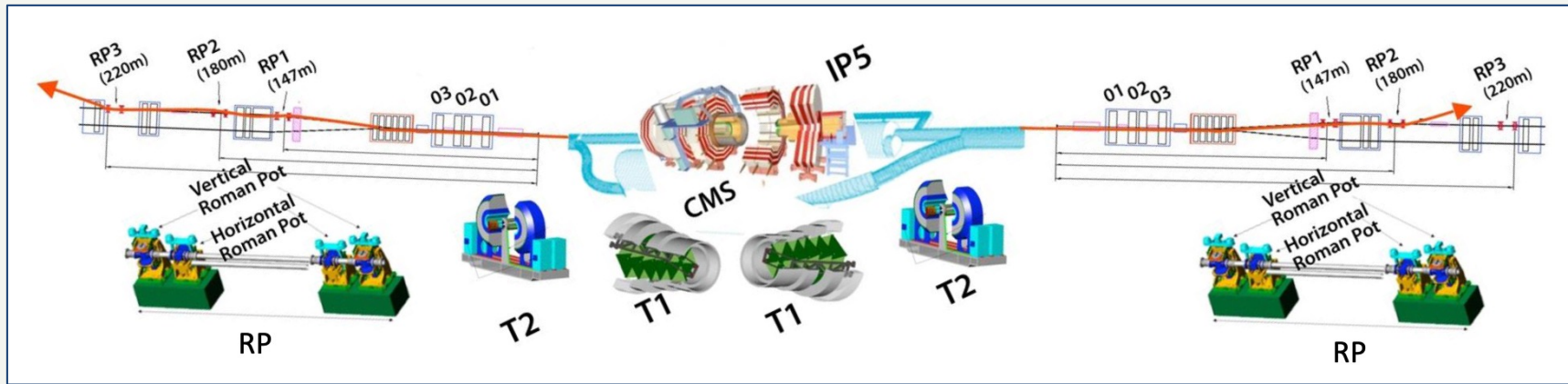
see Andrea Riffero talk WG2.

# Double Pomeron Exchange

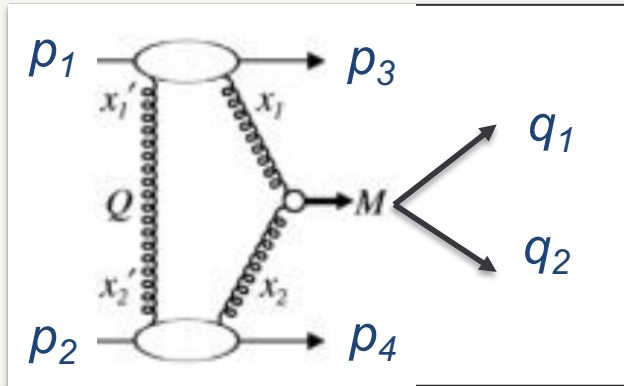


Understanding colourless strong interactions is fundamental  
Also simple environment for spectroscopy, in particular, glueballs

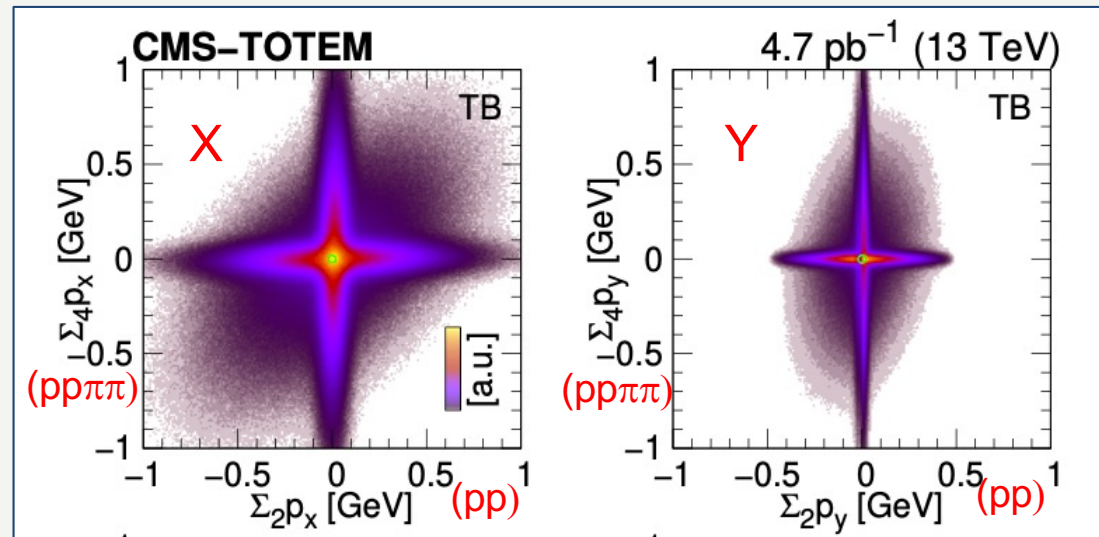
# CMS-TOTEM: Simultaneous reconstruction of central system and protons



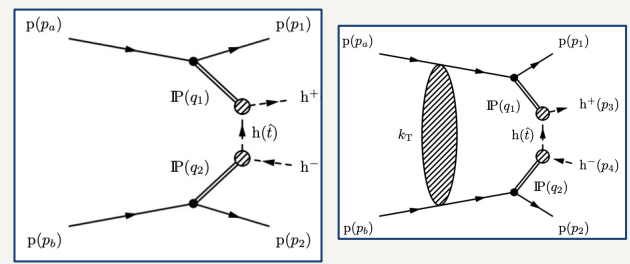
arXiv:2401.14494



$$M^2 = (q_1^\mu + q_2^\mu)^2 = (p_1^\mu + p_2^\mu - p_3^\mu - p_4^\mu)^2$$



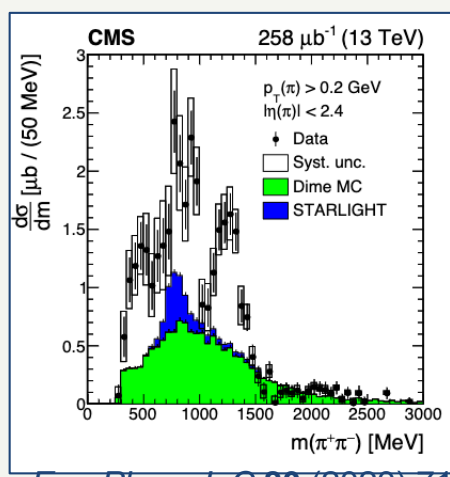
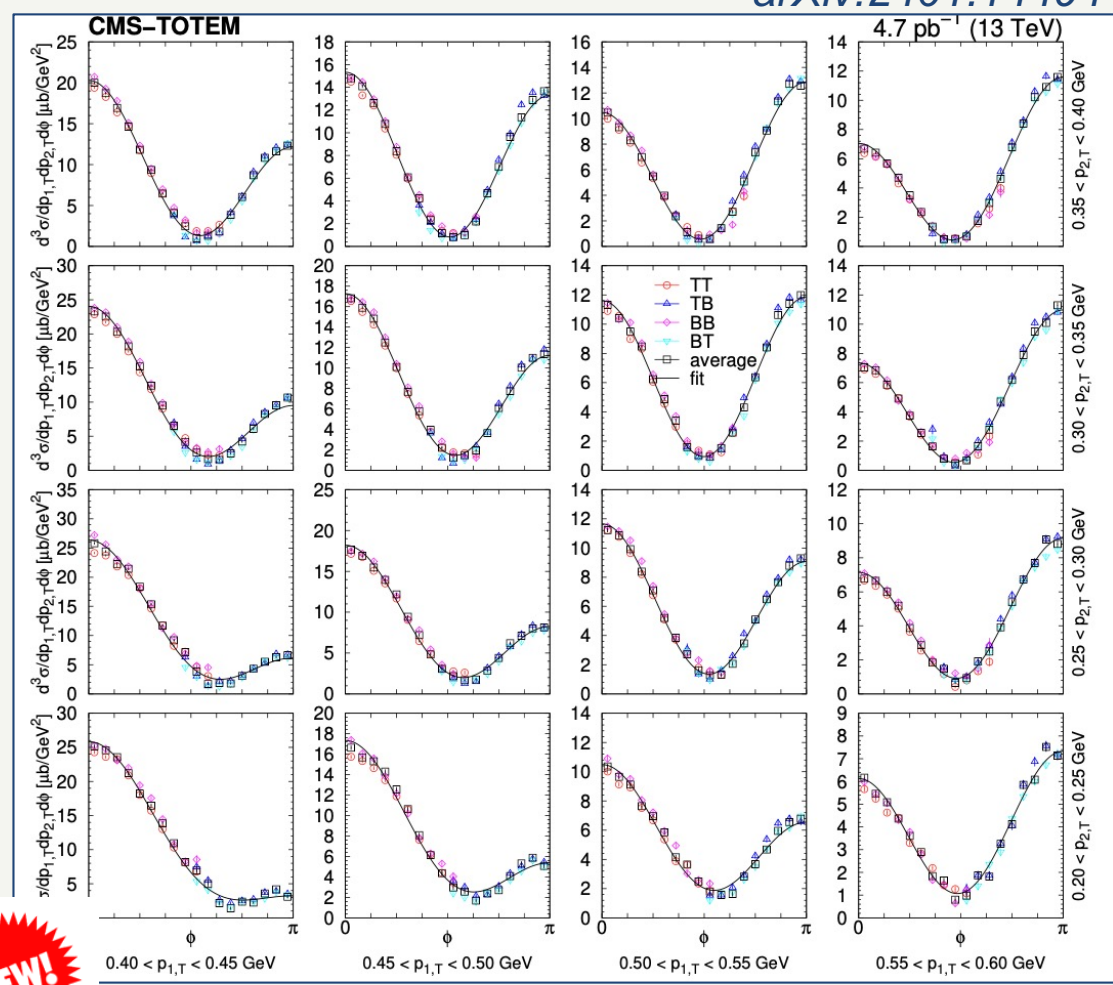
# Precision determination of pomeron interactions



arXiv:2401.14494

Shape distinctive of interference due to additional pomeron interactions

Measurements performed in non-resonant regions: (<0.7, >1.8 GeV).



Eur. Phys. J. C 80 (2020) 718

Eagerly await compelling spectroscopy of the complicated resonance region:

- photoproduced  $\rho, \omega, \rho', \rho''$
- DPE  $f_0, f_2, \dots$ , glueball candidates

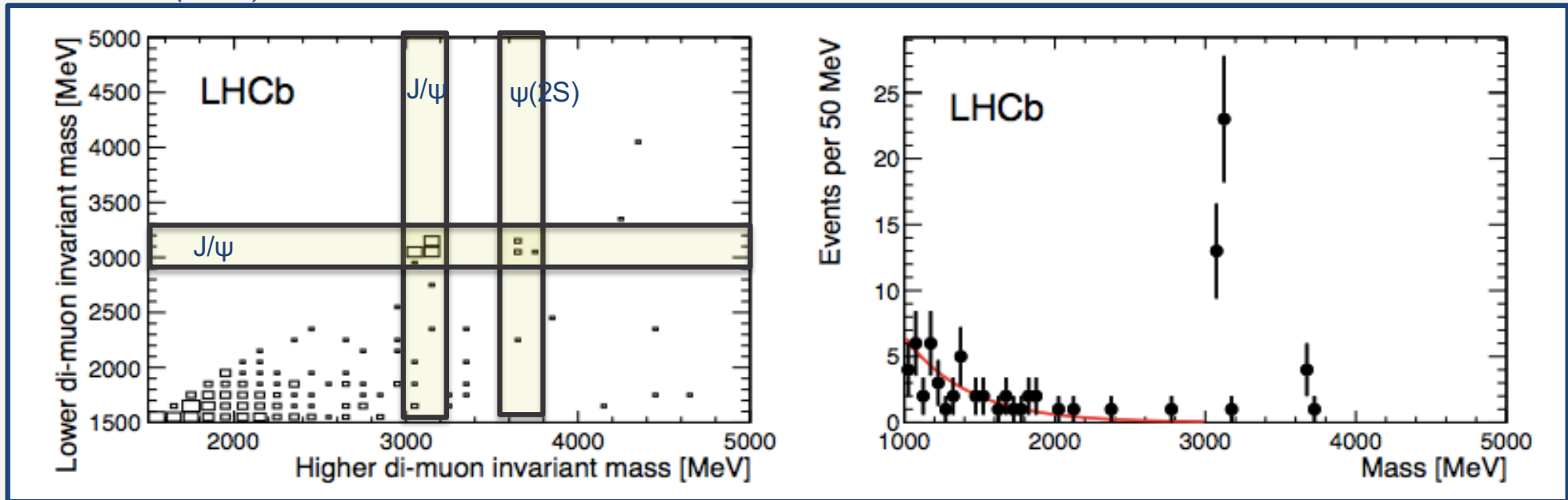


see Michael Pitt talk WG2.

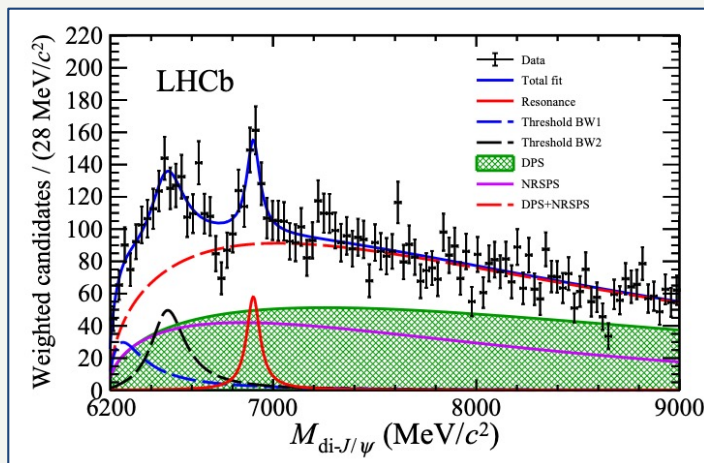
# J/ψJ/ψ: search for exotica

JPG 41 (2014) 115002

$$\begin{aligned} \sigma^{J/\psi J/\psi} &= 58 \pm 10(\text{stat}) \pm 6(\text{syst}) \text{ pb}, \\ \sigma^{J/\psi \psi(2S)} &= 63^{+27}_{-18}(\text{stat}) \pm 10(\text{syst}) \text{ pb}, \\ \sigma^{\psi(2S)\psi(2S)} &< 237 \text{ pb}, \\ \sigma^{\chi_{c0}\chi_{c0}} &< 69 \text{ nb}, \\ \sigma^{\chi_{c1}\chi_{c1}} &< 45 \text{ pb}, \\ \sigma^{\chi_{c2}\chi_{c2}} &< 141 \text{ pb}, \end{aligned}$$



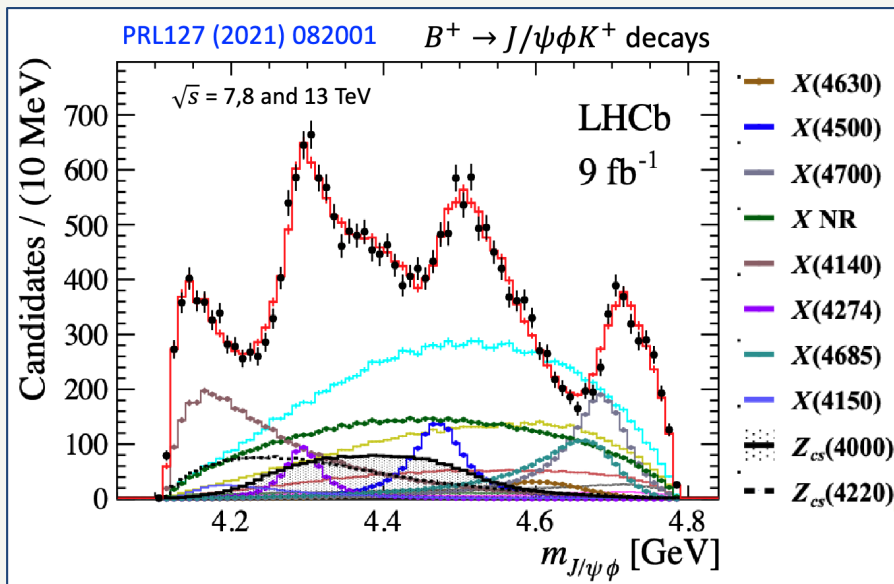
Sci.Bull. 65 (2020) 23, 1983-1993



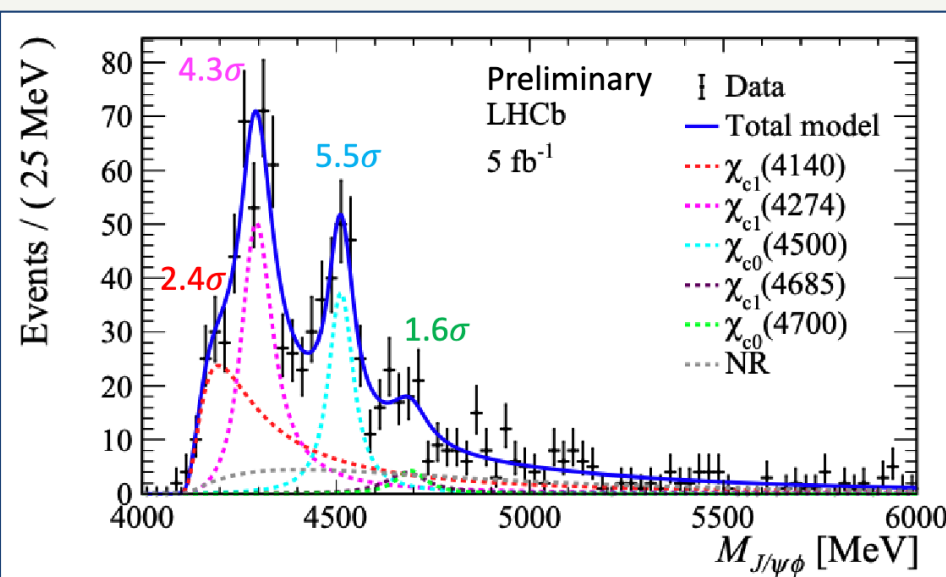
Today from inclusive measurements we know there is significant structure and tetraquark candidates

Diffraction measurements are cleaner and help identify quantum numbers

# J/ψ+φ: search for exotica



Structure seen in Inclusive production of J/ψ+φ .

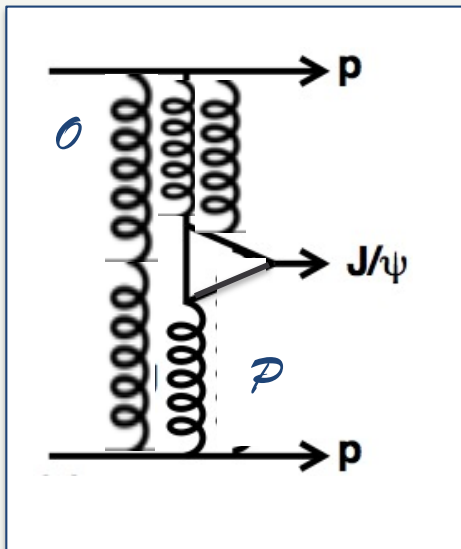


Similar and much cleaner structure now seen exclusively.

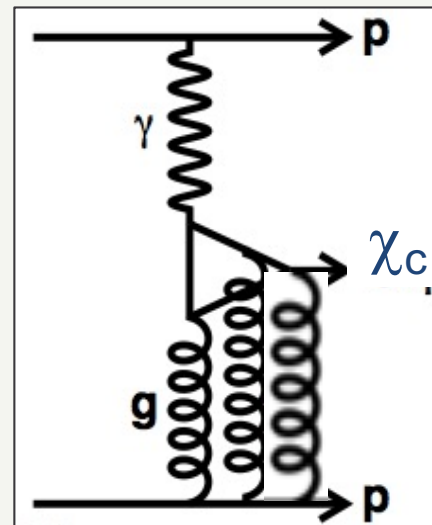
see Cesar da Silva talk WG2.



# Odderon search: partner of pomeron



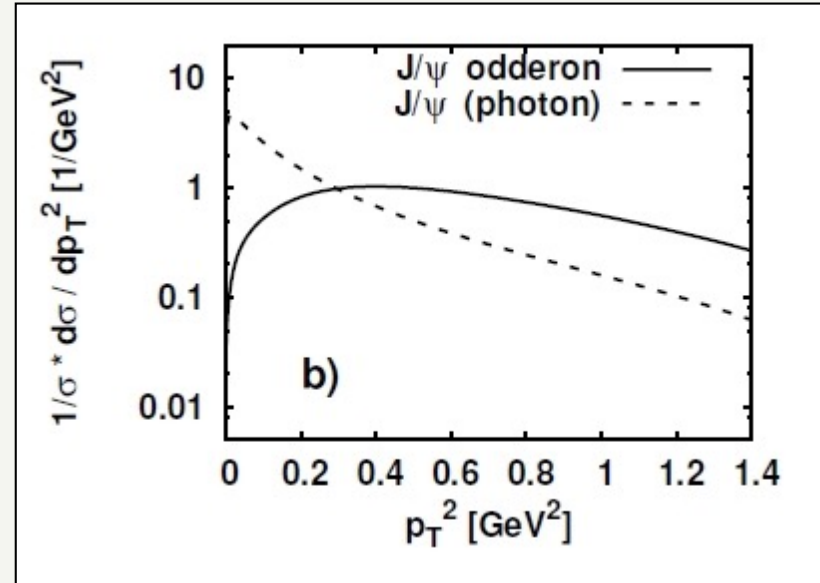
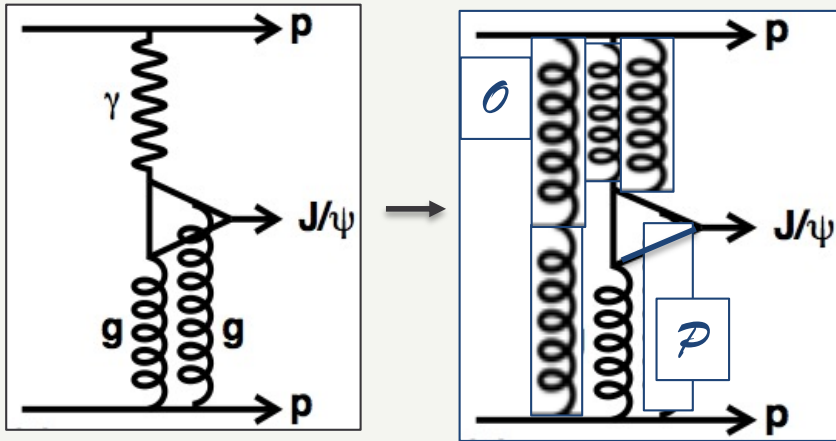
C-odd  
meson



C-even  
meson

# Method 1: High $p_T$ exclusive C- production

Replace  $\gamma$  with  $\phi$



Bzdak, Motyka, Szymanowski, Cudell PRD 75 (2007) 094023

$$\frac{d\sigma}{dt} \sim e^{bt}$$

Photoproduction:  $b \sim 6 \text{ GeV}^{-2}$   
 Proton dissociation  $b \sim 1 \text{ GeV}^{-2}$   
 Odderon  $b$  small

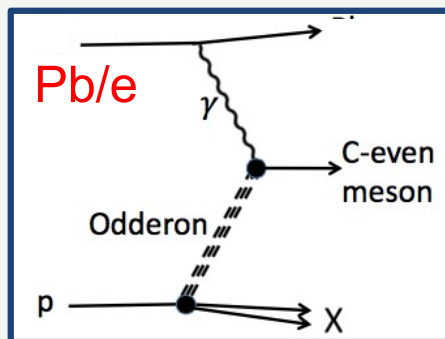
$d\sigma^{\text{corr}}/dy$	$J/\psi$	
	odderon	photon
Tevatron	0.3–1.3–5 nb	0.8–5–9 nb
LHC	0.3–0.9–4 nb	2.4–15–27 nb

Odderon contribution might be 1-10% at LHC and would dominate at high  $p_T$   
 ..... but experimentally **this is difficult to see**

Angular distribution of muons due to polarisation may also differ (R. Schnicker)

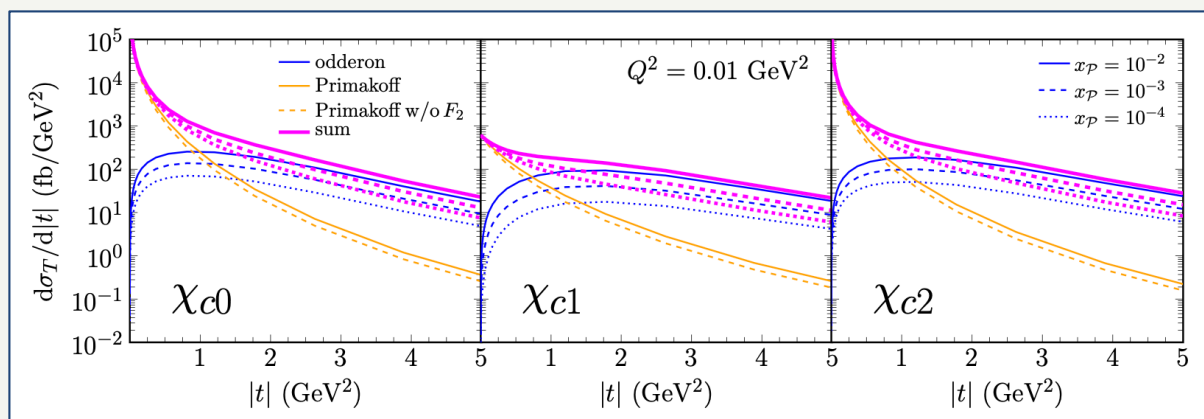
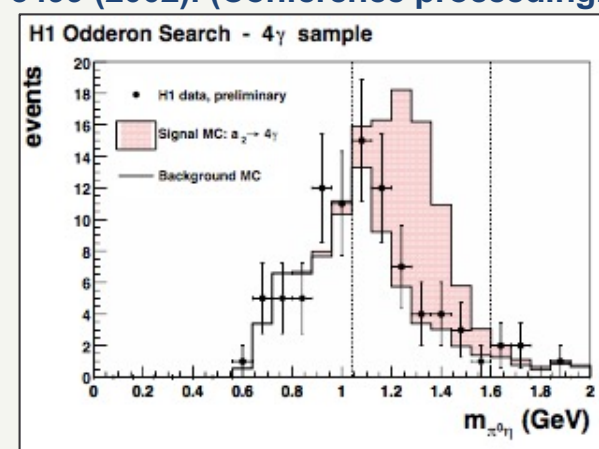
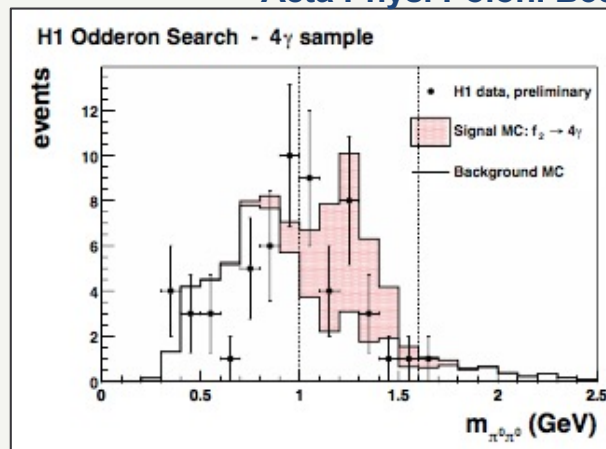
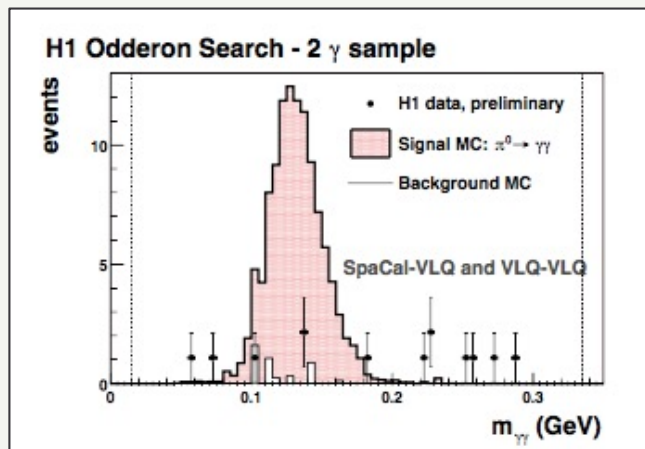


# Method 2: Photoproduction of C+



Czyzewski, Kwiecinski, Motyka, PLB398 (1997) 400.  
 Berger, Donnachie, Dosch, Kilian, Nachtmann, EPJ C9 (1999) 491.  
 Ryskin EPJ C2 (1998) 339.  
 Kilian & Nachtmann, EPJ C5 (1998) 317.  
 Harland-Lang, Khoze, Martin, Ryskin PRD 99 (2019) 3, 034011

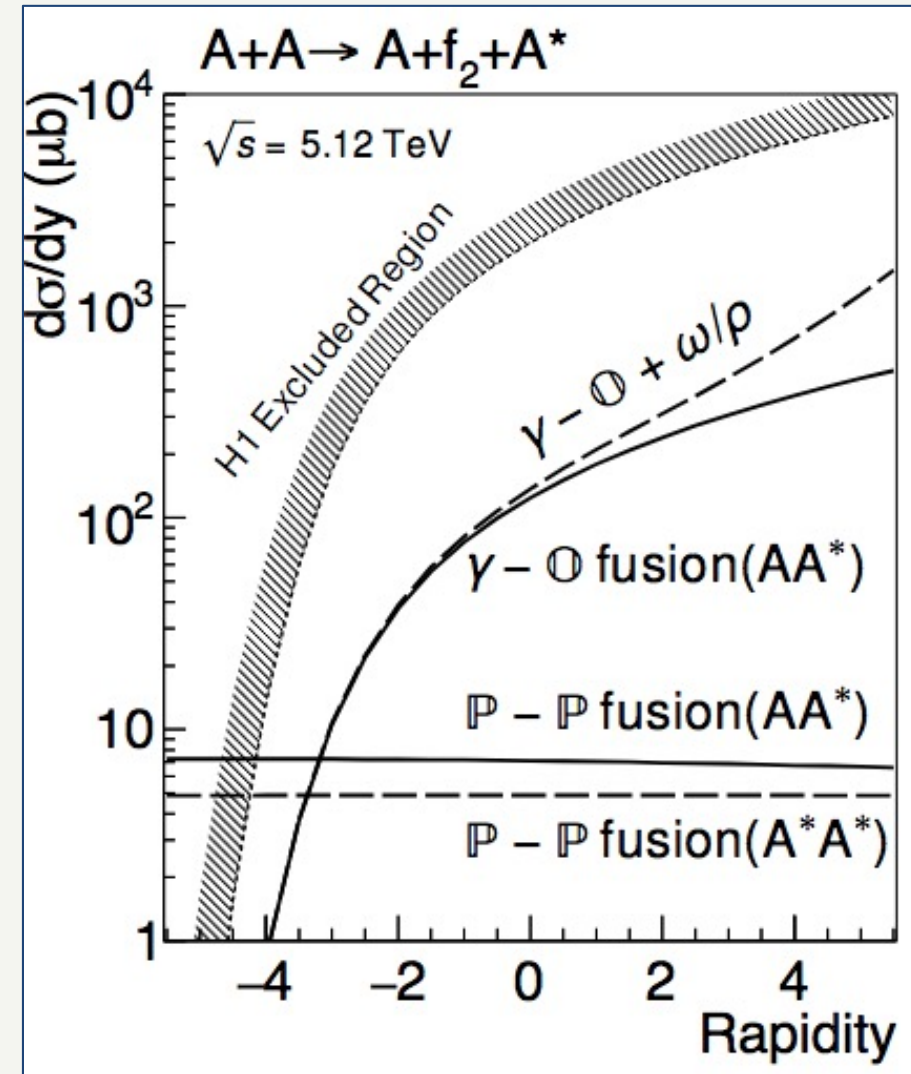
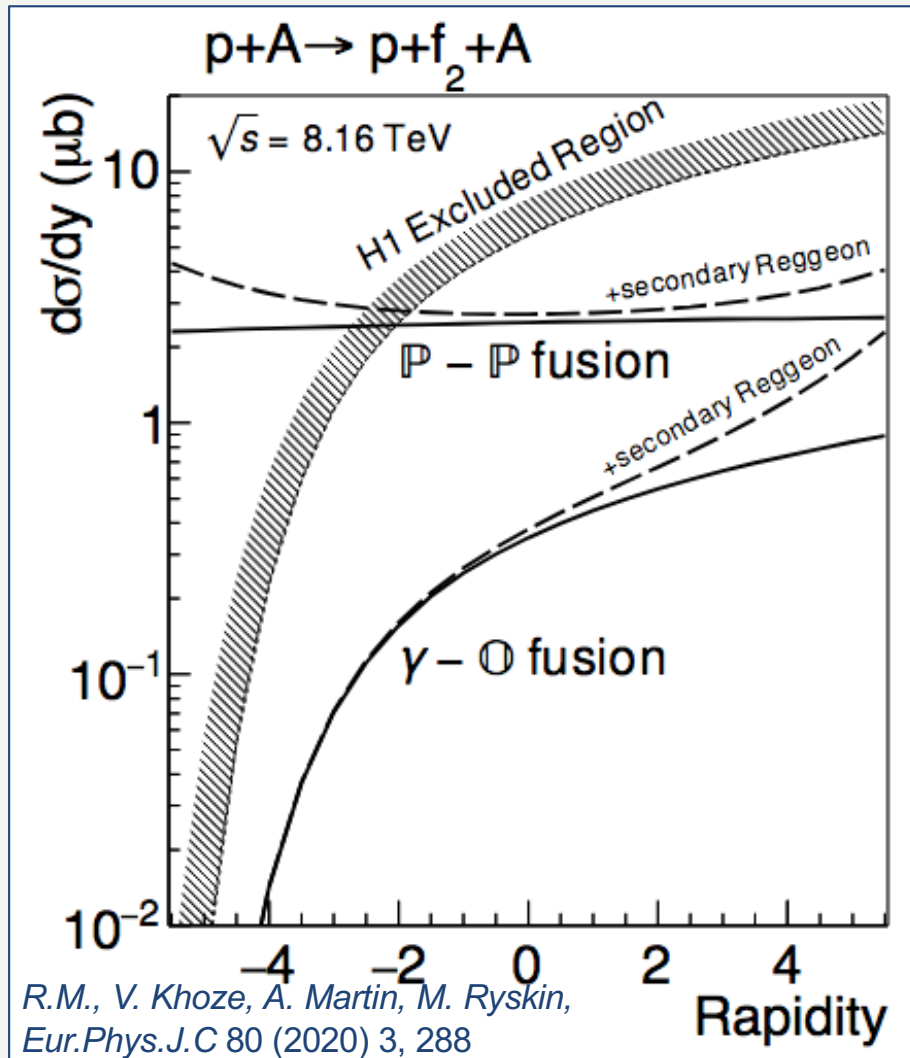
Acta Phys. Polon. B33, 3499 (2002). (Conference proceeding.)



arXiv: 2402.19134

Talk by Sanjin Benić in WG2

# Might be seen forward in p-Pb / PbPb



# Summary

- Much to understand about QCD
  - perturbative / non-perturbative regime
  - proton and nuclear structure (PDFs GPDs)
  - saturation
  - quark model bound states ( $\rho$ ,  $\rho'$ ,  $f_0, f_2, \dots$ )
  - beyond the naïve quark model (hybrids, tetraquarks, glueballs)
  - colourless propagators: pomerons and odderons
- Can be addressed in diffractive processes.