

# Measurements of $\gamma$ from tree-level decays at LHCb

A. Bertolin on behalf of the LHCb collaboration

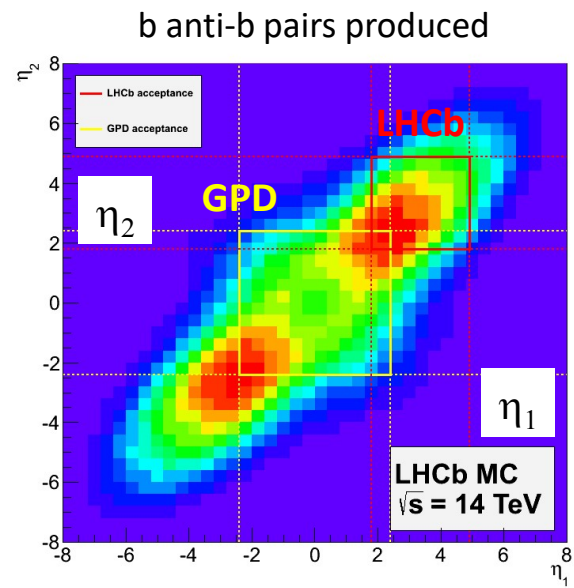
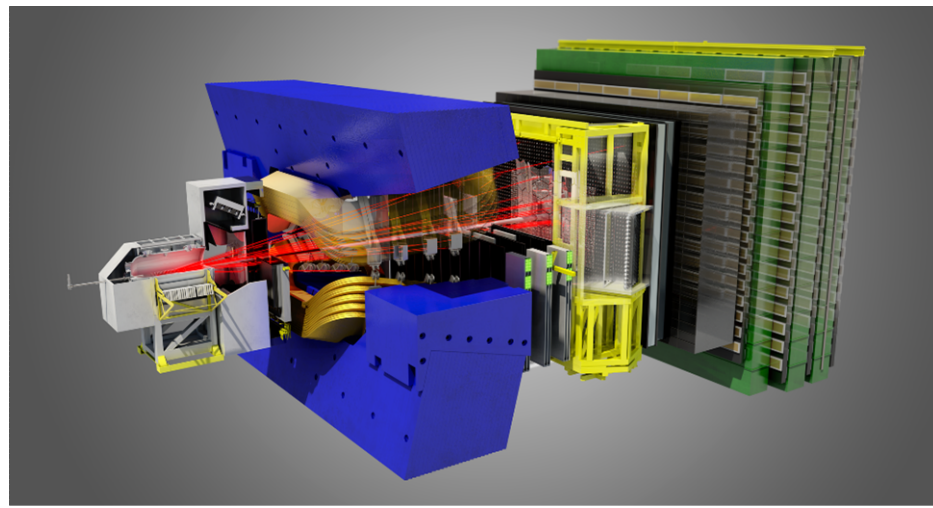


Outlook:

- ✚ short introduction to LHCb and  $\gamma$
- ✚ LHCb  $\gamma$  results:
  - the LHCb  $\gamma$  combination
  - some recent measurements not yet included in the LHCb  $\gamma$  combination
- ✚ take home message and future prospects

# LHCb: the detector and its performances

single-arm forward spectrometer at the LHC



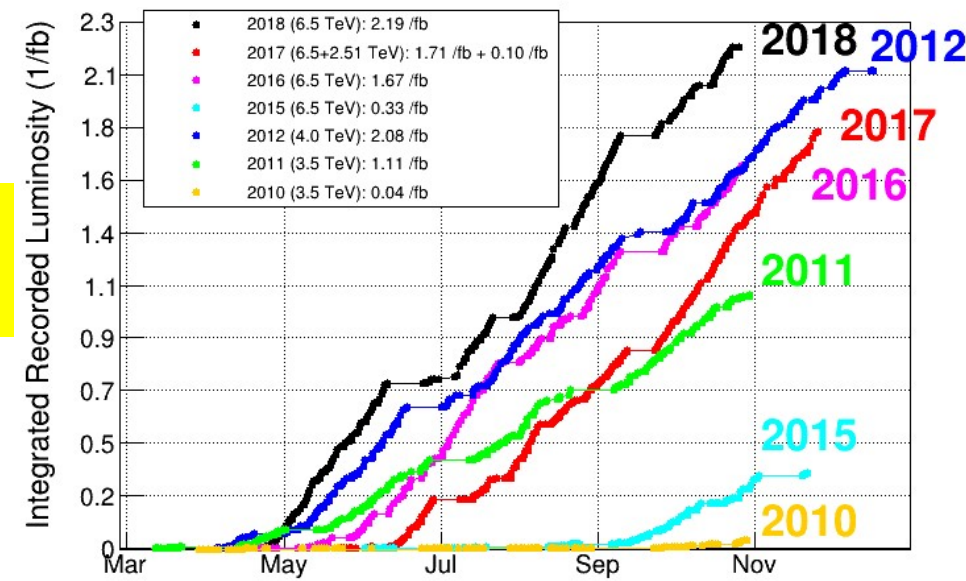
optimized for beauty and charm physics at  $2 < \eta < 5$

- detector paper: JINST 3 (2008) S08005
- Run 1 performance: Int. J. Mod. Phys. A30 (2015) 1530022
- Run 2 performance: JINST 14 (2019) P04013

## key points:

- momentum resolution ( $\sigma(p)/p \approx 0.5\%$  (low momentum) to  $1\%$  @  $200 \text{ GeV}/c$ )
- impact parameter resolution ( $\sigma(\text{IP}) \approx 15 \mu\text{m}$  at high  $p_T$ )
- primary and secondary vertices reco.
- decay time resolution ( $\sigma(t) \approx 50 \text{ fs}$ )
- 'global' PID:  $e / \mu / \pi / K$  ( $K \text{ id} \approx 95\%$   $\pi \text{ mis-id} \approx 5\%$ ,  $p < 100 \text{ GeV}/c$ )
- $\gamma$  and  $\pi^0$  reconstruction

recorded lumi.:  
 2011 → 2012 (Run 1): 3 /fb  
 2015 → 2018 (Run 2): 6 /fb



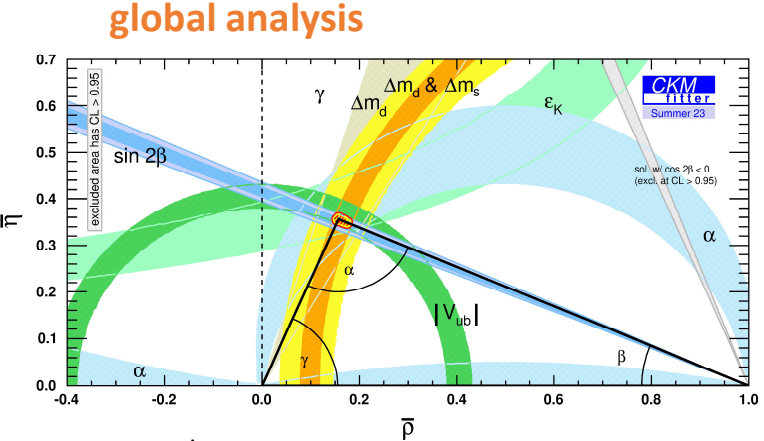
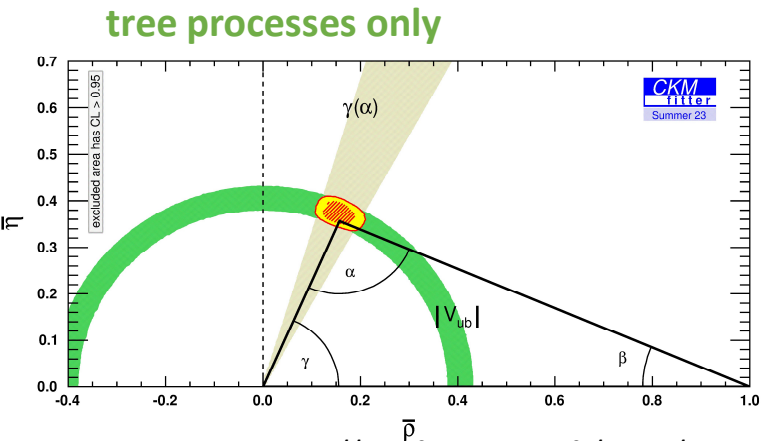
CP violation in the SM

**CP violation is one of the requirements to explain the baryon asymmetry we observe today**

a process must have been in place that took us from the equal amounts of matter - anti-matter produced in the Big Bang to the Universe dominated by matter we are living in

the SM charged current weak interactions between quarks are described by the Cabibbo-Kobayashi-Maskawa matrix  $V$ , 3 x 3, with  $V V^* = I$   
 $\Leftrightarrow$  3 angles and 1 phase or 3 reals and **1 imaginary** parameters

$V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + V_{td}V_{tb}^* = 0$  unitary condition relevant for beauty decays can be represented by a triangle in a complex plane, with angles  $\alpha$ ,  $\beta$  and  $\gamma$



[http://ckmfitter.in2p3.fr/www/results/plots\\_summer23/ckm\\_res\\_summer23.html](http://ckmfitter.in2p3.fr/www/results/plots_summer23/ckm_res_summer23.html)

$\gamma \equiv \arg \left[ -\frac{V_{ud}V_{ub}^*}{V_{cd}V_{cb}^*} \right]$   
 a.k.a.  $\phi_3$

- only CKM angle easily accessible in tree-level decays
- assuming no new physics in tree-level decays, has negligible theoretical uncertainty i.e. achievable accuracy dominated by experiments [JHEP01(2014)051]

**any disagreement between tree-level determinations and the value inferred from global CKM fits would indicate physics beyond the SM** ... due for example to new particles / mediators being exchanged in loops ...

✚ how to measure  $\gamma$

$\gamma$  can be determined by exploiting the interference between

- $b \rightarrow cW$  ( $V_{cb}$ ), favoured
- $b \rightarrow uW$  ( $V_{ub}$ ), suppressed transition amplitudes

$$A_{sup}/A_{fav} = r_B^X e^{i(\delta_B^X \pm \gamma)} \quad (- \text{ is for b-quark, + for anti-b})$$

where  $r_B^X$  and  $\delta_B^X$  are the ratio and the strong phase differences between the  $V_{cb}$  and  $V_{ub}$  transition amplitudes for the specific final state X  
these are also simultaneously determined

charm parameters can also get involved

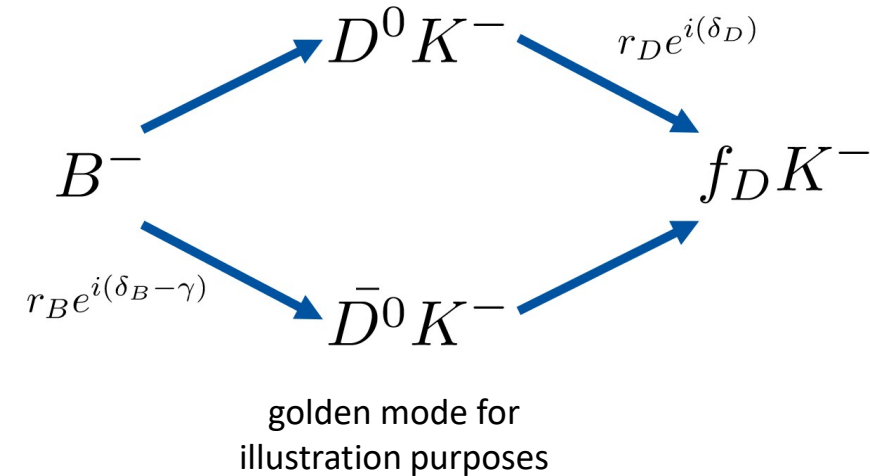
□ which/typical B meson final states ( $h=K,\pi$ ) ?

- $B^+ \rightarrow D h^+$
- $B^+ \rightarrow D h^+ \pi^- \pi^+$
- $B^0 \rightarrow D K^{*0}$
- $B^0 \rightarrow D K^+ \pi^-$

where D is a neutral charm meson mixture of the  $D^0$  anti- $D^0$  flavor eigenstates

□ which/typical D meson final states ?

- CP-eigenstates,  $D \rightarrow K^+ K^-$  and  $D \rightarrow \pi^+ \pi^-$ : GLW method [Phys. Lett. B265 (1991) 172, Phys. Lett. B253 (1991) 483]
- non CP-eigenstates,  $D^0 \rightarrow \pi^- K^+$ : ADS method [Phys. Rev. D63 (2001) 036005]
- self-conjugate multibody D meson decay, like  $K_s^0 \pi^+ \pi^-$ , with the D-Dalitz plot distributions: BPGGSZ method [Phys. Rev. D 68, 054018 (2003), Eur. Phys. J. C 47, 347 (2006)]



GLW = M. Gronau, D. London and D. Wyler

ADS = D. Atwood, I. Dunietz and A. Soni

BPGGSZ = A. Bondar, A. Poluektov, A. Giri, Y. Grossman, A. Soffer, J. Zupan

**however due to the small branching ratios the most precise determination is obtained from a combination of measurements from many decay modes**

✚ LHCb  $\gamma$  combination

measurements used in the combination, the ones denoted by (\*) include only a fraction of the Run 2 sample

$B$ decay	$D$ decay	Ref.	Dataset	Status since Ref. [14]
$B^\pm \rightarrow Dh^\pm$	$D \rightarrow h^+h^-$	[29]	Run 1&2	As before
$B^\pm \rightarrow Dh^\pm$	$D \rightarrow h^+\pi^-\pi^+\pi^-$	[30]	Run 1	As before
$B^\pm \rightarrow Dh^\pm$	$D \rightarrow K^\pm\pi^\mp\pi^+\pi^-$	[18]	Run 1&2	New
$B^\pm \rightarrow Dh^\pm$	$D \rightarrow h^+h^-\pi^0$	[19]	Run 1&2	Updated
$B^\pm \rightarrow Dh^\pm$	$D \rightarrow K_S^0h^+h^-$	[31]	Run 1&2	As before
$B^\pm \rightarrow Dh^\pm$	$D \rightarrow K_S^0K^\pm\pi^\mp$	[32]	Run 1&2	As before
$B^\pm \rightarrow D^*h^\pm$	$D \rightarrow h^+h^-$	[29]	Run 1&2	As before
$B^\pm \rightarrow DK^{*\pm}$	$D \rightarrow h^+h^-$	[33]	Run 1&2(*)	As before
$B^\pm \rightarrow DK^{*\pm}$	$D \rightarrow h^+\pi^-\pi^+\pi^-$	[33]	Run 1&2(*)	As before
$B^\pm \rightarrow Dh^\pm\pi^+\pi^-$	$D \rightarrow h^+h^-$	[34]	Run 1	As before
$B^0 \rightarrow DK^{*0}$	$D \rightarrow h^+h^-$	[35]	Run 1&2(*)	As before
$B^0 \rightarrow DK^{*0}$	$D \rightarrow h^+\pi^-\pi^+\pi^-$	[35]	Run 1&2(*)	As before
$B^0 \rightarrow DK^{*0}$	$D \rightarrow K_S^0\pi^+\pi^-$	[36]	Run 1	As before
$B^0 \rightarrow D^\mp\pi^\pm$	$D^+ \rightarrow K^-\pi^+\pi^+$	[37]	Run 1	As before
$B_s^0 \rightarrow D_s^\mp K^\pm$	$D_s^+ \rightarrow h^+h^-\pi^+$	[38]	Run 1	As before
$B_s^0 \rightarrow D_s^\mp K^\pm\pi^+\pi^-$	$D_s^+ \rightarrow h^+h^-\pi^+$	[39]	Run 1&2	As before
$D$ decay	Observable(s)	Ref.	Dataset	Status since Ref. [14]
$D^0 \rightarrow h^+h^-$	$\Delta A_{CP}$	[24] [40] [41]	Run 1&2	As before
$D^0 \rightarrow K^+K^-$	$A_{CP}(K^+K^-)$	[16] [24] [25]	Run 2	New
$D^0 \rightarrow h^+h^-$	$y_{CP} - y_{CP}^{K^-\pi^+}$	[42]	Run 1	As before
$D^0 \rightarrow h^+h^-$	$y_{CP} - y_{CP}^{K^-\pi^+}$	[15]	Run 2	New
$D^0 \rightarrow h^+h^-$	$\Delta Y$	[43] [46]	Run 1&2	As before
$D^0 \rightarrow K^+\pi^-$ (Single Tag)	$R^\pm, (x'^\pm)^2, y'^\pm$	[47]	Run 1	As before
$D^0 \rightarrow K^+\pi^-$ (Double Tag)	$R^\pm, (x'^\pm)^2, y'^\pm$	[48]	Run 1&2(*)	As before
$D^0 \rightarrow K^\pm\pi^\mp\pi^+\pi^-$	$(x^2 + y^2)/4$	[49]	Run 1	As before
$D^0 \rightarrow K_S^0\pi^+\pi^-$	$x, y$	[50]	Run 1	As before
$D^0 \rightarrow K_S^0\pi^+\pi^-$	$x_{CP}, y_{CP}, \Delta x, \Delta y$	[51]	Run 1	As before
$D^0 \rightarrow K_S^0\pi^+\pi^-$	$x_{CP}, y_{CP}, \Delta x, \Delta y$	[52]	Run 2	As before
$D^0 \rightarrow K_S^0\pi^+\pi^-$ ( $\mu^-$ tag)	$x_{CP}, y_{CP}, \Delta x, \Delta y$	[17]	Run 2	New

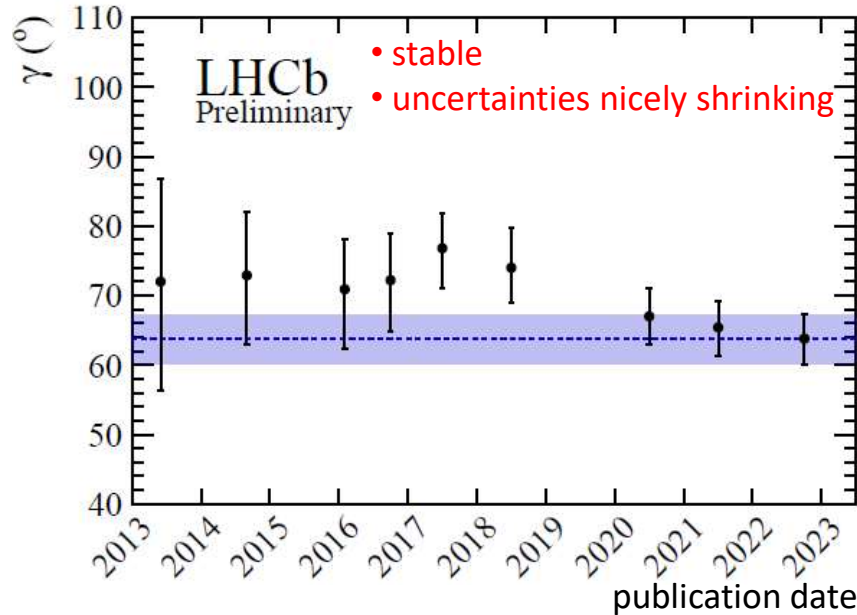
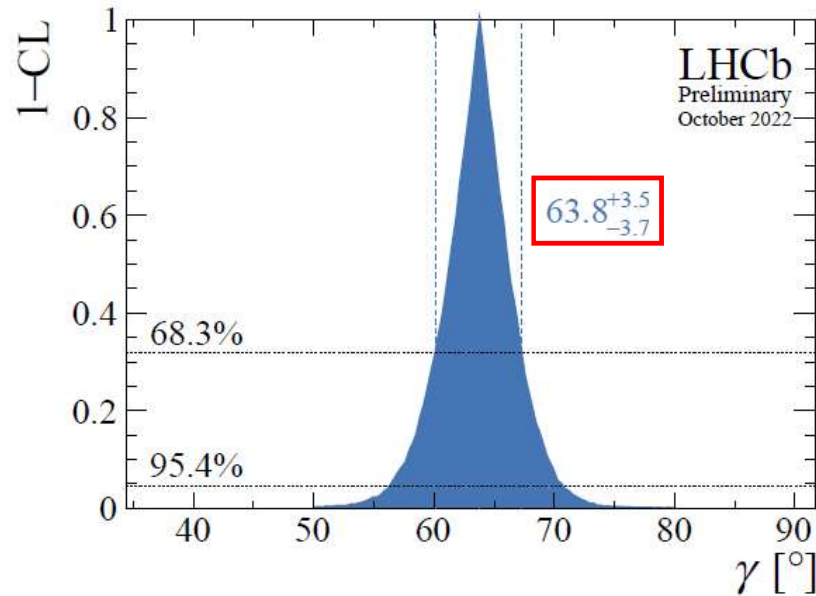
inputs from the charm system

- with so many inputs it is easy to probe the stability / strength of the final result on  $\gamma$

Simultaneous determination of the CKM angle  $\gamma$  and parameters related to mixing and  $CP$  violation in the charm sector

LHCb  $\gamma$  combination: results

- 173 input observables
- 52 free parameters
- fit quality:
  - given the  $\chi^2$  value at the best fit point and the n.d.f. the fit probability is about 80 %
  - cross checked with pseudoexperiments
- fit results:  $\gamma$ , common free parameter, ( $r_B, \delta_B$ ) for every "B", charm mixing parameters  $x$  and  $y$  and a few other floated parameters



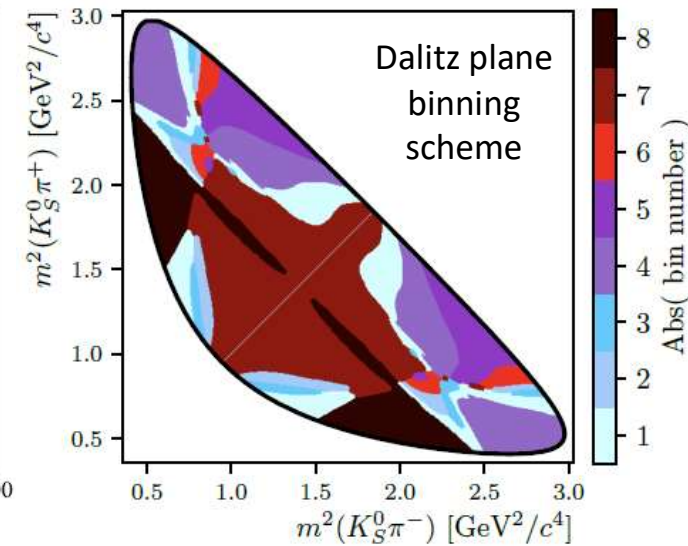
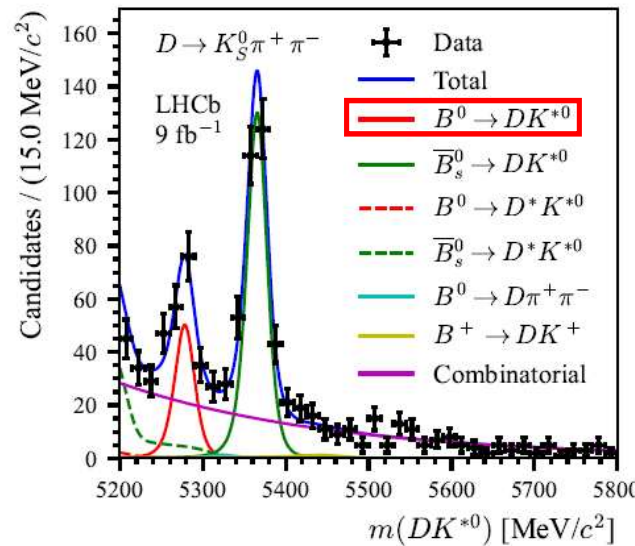
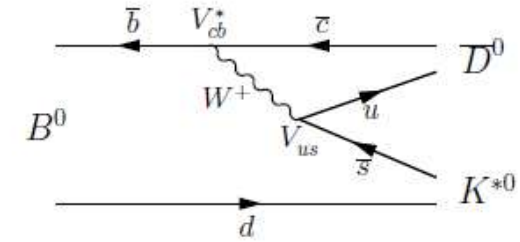
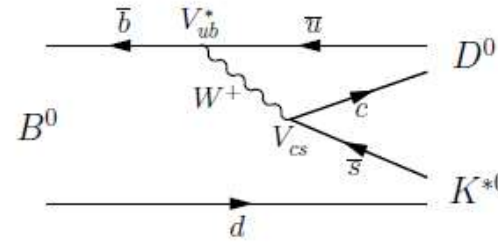
- most precise determination from a single experiment

$$x = (0.398^{+0.050}_{-0.049})\% \quad y = (0.636^{+0.020}_{-0.019})\%$$

most precise determinations to date, these were taken as auxiliary inputs from HFLAV before the 2021 measurement

✚ LHCb  $\gamma$  combination:  $B^0 \rightarrow D K^{*0}$  using self-conjugate  $D \rightarrow K_S^0 h^+ h^-$  decays

- $B^0 \rightarrow D K^*(892)^0$  has a lower BF wrt  $B^\pm \rightarrow D K^\pm$  that has the largest impact on  $\gamma$
- but the interference between the  $b \rightarrow c$  favored and  $b \rightarrow u$  suppressed amplitudes is expected to be a factor of 3 larger
- $D \rightarrow K_S^0 h^+ h^-$  with  $h = \{\pi, K\}$
- the flavor of the B meson at the point of decay is unambiguously provided by the charge of the kaon from the  $K^*(892)^0 \rightarrow K^+ \pi^-$  decay



$$x_{\pm} \equiv r_{B^0} \cos(\delta_{B^0} \pm \gamma)$$

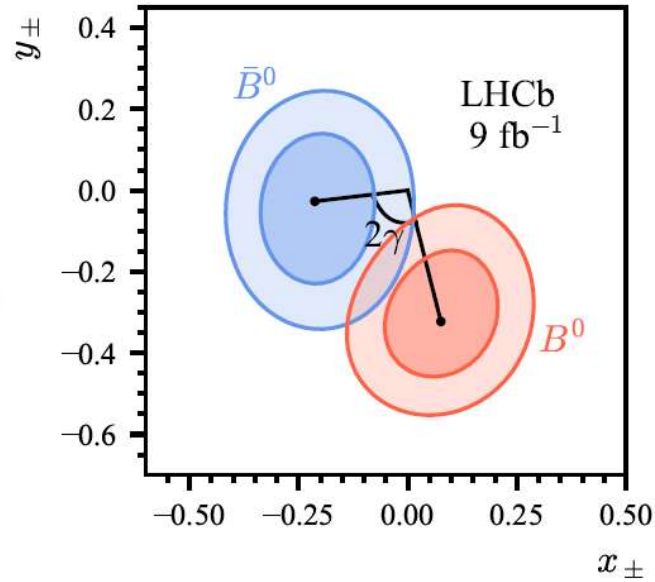
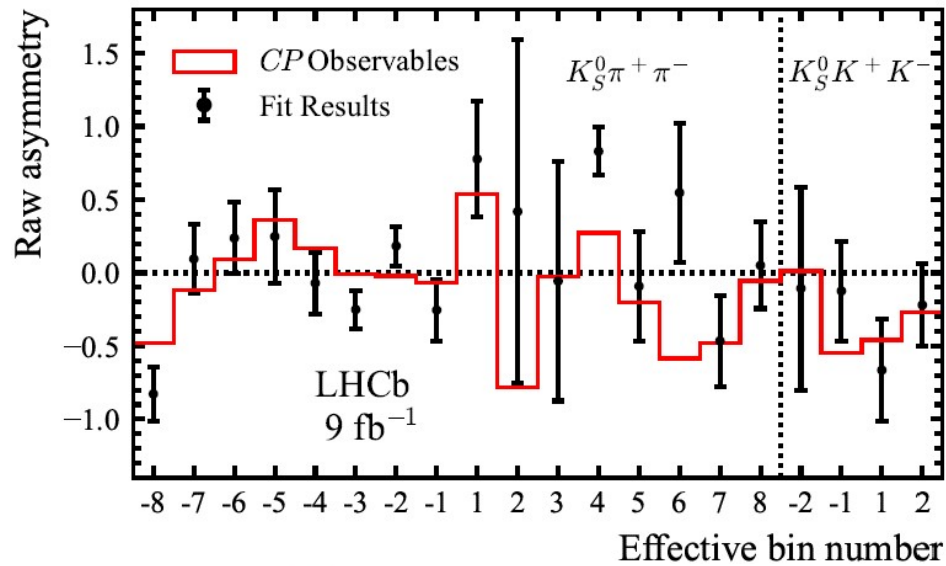
$$y_{\pm} \equiv r_{B^0} \sin(\delta_{B^0} \pm \gamma)$$

$$N_i(B^0) = h^{B^0} \left[ F_{-i} + (x_+^2 + y_+^2)F_i + 2\kappa\sqrt{F_i F_{-i}}(x_+ c_i - y_+ s_i) \right]$$

$$N_i(\bar{B}^0) = h^{\bar{B}^0} \left[ F_i + (x_-^2 + y_-^2)F_{-i} + 2\kappa\sqrt{F_i F_{-i}}(x_- c_i + y_- s_i) \right]$$

- $i$ : bin in the Dalitz plot for the D decay
- $F_i$ : fractional  $D^0$  yield in bin  $i$ , from  $B^\pm \rightarrow D \pi^\pm$ ,  $D \rightarrow K_S^0 h^+ h^-$  data
- $c_i$  and  $s_i$  are the cosine and sine of the strong-phase difference between the  $D^0$  and anti- $D^0$  decays, from BESIII and CLEO data
- $\kappa$  coherence factor diluting the interference term, fixed from data

LHCb  $\gamma$  combination:  $B^0 \rightarrow D K^{*0}$  using self-conjugate  $D \rightarrow K_S^0 h^+ h^-$  decays



$$\gamma = (49_{-19}^{+22})^\circ$$

$$r_{B^0} = 0.271_{-0.066}^{+0.065}$$



$$r_B^{DK} = 0.0964 \pm 0.0028$$

- result consistent with the LHCb  $\gamma$  combo average [LHCb-CONF-2022-003]
- confirmation of the large value of the amplitudes ratio  $r$
- external strong phase inputs from the BESIII and CLEO collaborations are not limiting the accuracy of the measurement
- the measurement is dominated by statistical uncertainties as LHCb systematics are about 1/10 of the statistical uncertainty
- expect improved accuracy from Run 3 data

JHEP04(2021)081



✚ LHCb  $\gamma$  combination:  $B^0 \rightarrow D K^{*0}$  using two- and four-body D decays

- interference in the admixture of Cabibbo-favored (D0)  $K^- \pi^+ \pi^+ \pi^-$  and doubly Cabibbo-suppressed (antiD0)  $K^- \pi^+ \pi^+ \pi^-$  decays (ADS method)
- the charge of the kaon child from the  $K^{*0}$  candidate is used to determine the flavor of the parent B meson
- $K \pi (\pi \pi) K^{*0} \equiv$  kaon child of the D candidate has the same charge as the kaon child of the  $K^{*0}$  candidate
- $\pi K (\pi \pi) K^{*0} \equiv$  the two kaons have opposite charge
- definitions of some observables:

$$\mathcal{R}_{\pi K(\pi\pi)}^+ \equiv \frac{\Gamma(B^0 \rightarrow D[\pi K(\pi\pi)]K^{*0})}{\Gamma(B^0 \rightarrow D[K\pi(\pi\pi)]K^{*0})} \quad \mathcal{R}_{\pi K(\pi\pi)}^- \equiv \frac{\Gamma(\bar{B}^0 \rightarrow D[\pi K(\pi\pi)]\bar{K}^{*0})}{\Gamma(\bar{B}^0 \rightarrow D[K\pi(\pi\pi)]\bar{K}^{*0})}$$

$$\mathcal{A}_{K\pi} \equiv \frac{\Gamma(\bar{B}^0 \rightarrow D[K\pi(\pi\pi)]\bar{K}^{*0}) - \Gamma(B^0 \rightarrow D[K\pi(\pi\pi)]K^{*0})}{\Gamma(\bar{B}^0 \rightarrow D[K\pi(\pi\pi)]\bar{K}^{*0}) + \Gamma(B^0 \rightarrow D[K\pi(\pi\pi)]K^{*0})}$$

cancellation of a large number of systematic uncertainties related to the reconstruction and selection of the signal candidates

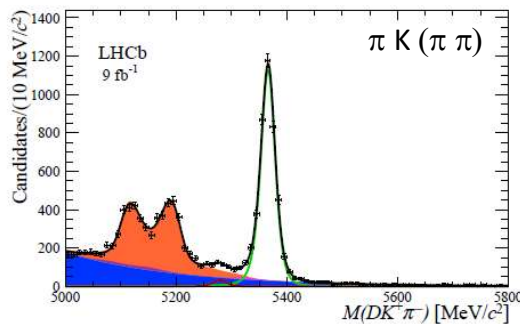
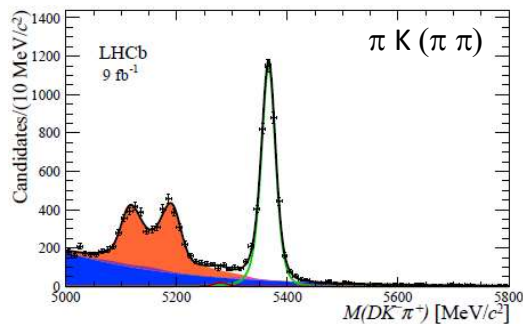
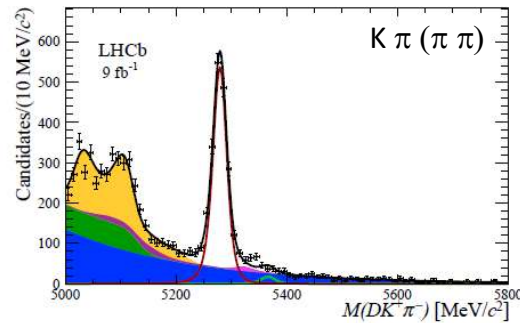
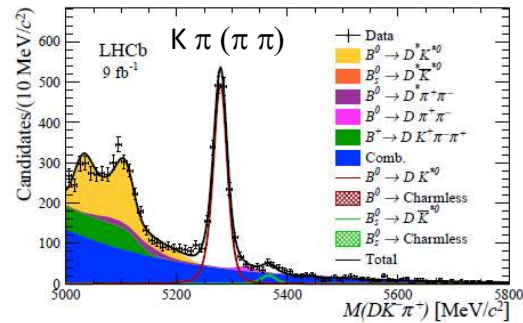
for illustration purposes, the dependence of one of the observables on  $r_\gamma$  and  $\delta$  is:

$$\mathcal{R}_{\pi K(\pi\pi)}^\pm = \frac{(r_{B^0}^{DK^*})^2 + (r_D^{K\pi(\pi\pi)})^2 + 2\kappa_{B^0} r_{B^0}^{DK^*} \kappa_D^{K\pi(\pi\pi)} \cos(\delta_{B^0}^{DK^*} + \delta_D^{K\pi(\pi\pi)} \pm \gamma)}{1 + (r_{B^0}^{DK^*})^2 + (r_D^{K\pi(\pi\pi)})^2 + 2\kappa_{B^0} r_{B^0}^{DK^*} \kappa_D^{K\pi(\pi\pi)} \cos(\delta_{B^0}^{DK^*} - \delta_D^{K\pi(\pi\pi)} \pm \gamma)}$$

fundamental parameters of interest plus some additional inputs

additional final states are included:

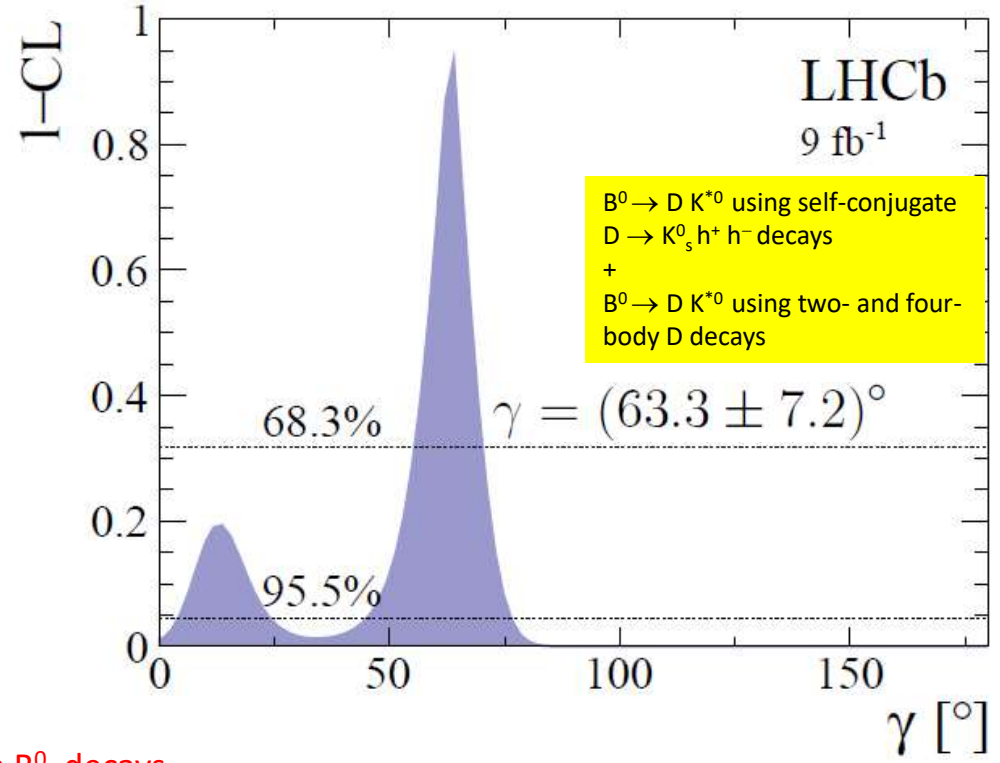
- CP eigenstates  $D \rightarrow K^+ K^-$  and  $D \rightarrow \pi^+ \pi^-$  (GLW method)
- $D \rightarrow \pi^+ \pi^- \pi^+ \pi^-$  (mostly CP even, extension of the GLW method)
- similar observables



✚ LHCb  $\gamma$  combination:  $B^0 \rightarrow D K^{*0}$  using two- and four-body D decays

LHCb-PAPER-2023-040  
<https://arxiv.org/pdf/2401.17934.pdf>

	Parameter	Value
four-body (ADS method)	$A_{K\pi}$	$0.031 \pm 0.017 \pm 0.015$
	$R_{\pi K}^+$	$0.069 \pm 0.013 \pm 0.005$
	$R_{\pi K}^-$	$0.093 \pm 0.013 \pm 0.005$
	$A_{K\pi\pi\pi}$	$-0.012 \pm 0.018 \pm 0.016$
	$R_{\pi K\pi\pi}^+$	$0.060 \pm 0.014 \pm 0.006$
	$R_{\pi K\pi\pi}^-$	$0.038 \pm 0.014 \pm 0.006$
two-body CP eigenstates (GLW method)	$R_{CP}^{KK}$	$0.811 \pm 0.057 \pm 0.017$
	$A_{CP}^{KK}$	$-0.047 \pm 0.063 \pm 0.015$
	$R_{CP}^{\pi\pi}$	$1.104 \pm 0.111 \pm 0.026$
	$A_{CP}^{\pi\pi}$	$-0.034 \pm 0.094 \pm 0.016$
four-body mostly CP even (extended GLW method)	$R_{CP}^{4\pi}$	$0.882 \pm 0.086 \pm 0.033$
	$A_{CP}^{4\pi}$	$0.021 \pm 0.087 \pm 0.016$



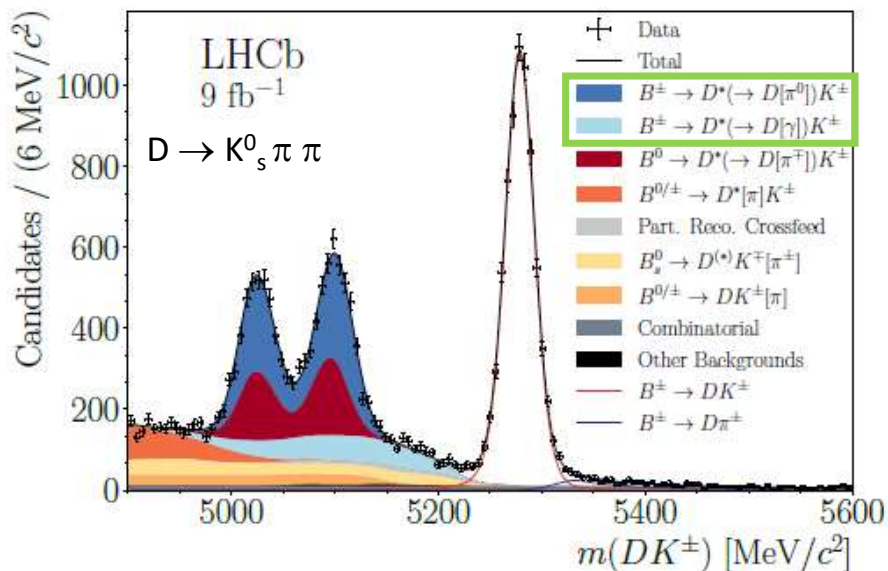
- most stringent limits to date on  $\gamma$  from  $B^0$  decays
- result is consistent with the LHCb  $\gamma$  combo average [LHCb-CONF-2022-003]
- for most of the observables the statistical uncertainties are dominant
- expect improved accuracy from Run 3 data

LHCb  $\gamma$  combination:  $B^\pm \rightarrow D^* h^\pm$  with partial  $D^*$  reconstruction

- $B^\pm \rightarrow [D \gamma / \pi^0] h^\pm$  with partial reconstruction of the  $D^*$ ,  $D \rightarrow K_s^0 hh$  and  $h = \{\pi, K\}$
- the  $B^-$  decay via  $D^0$  (anti- $D^0$ ) proceed with the favored (suppressed) amplitude, interference occur because the final state particles are identical

$$\mathcal{A}(B^- \rightarrow [D\pi^0]_{D^*} K^-) = \mathcal{A}_B(\mathcal{A}_D + r_B^{D^*K} \exp[i(\delta_B^{D^*K} - \gamma)] \mathcal{A}_{\bar{D}})$$

$$x_\pm^{D^*K} \equiv r_B^{D^*K} \cos(\delta_B^{D^*K} \pm \gamma) \quad \text{and} \quad y_\pm^{D^*K} \equiv r_B^{D^*K} \sin(\delta_B^{D^*K} \pm \gamma)$$



$D$ decay	Component	Reconstructed as:		
		$B^\pm \rightarrow DK^\pm$	$B^\pm \rightarrow D\pi^\pm$	
$D \rightarrow K_s^0 \pi^+ \pi^-$	$B^\pm \rightarrow D^*[D\pi^0]K^\pm$	$6244 \pm 12$	$2716 \pm 5$	
	$B^\pm \rightarrow D^*[D\pi^0]\pi^\pm$	$340 \pm 1$	$113170 \pm 229$	
	$B^\pm \rightarrow D^*[D\gamma]K^\pm$	$3144 \pm 6$	$1247 \pm 2$	
	$B^\pm \rightarrow D^*[D\gamma]\pi^\pm$	$166 \pm 1$	$60285 \pm 121$	
	$B^\pm \rightarrow DK^\pm$	$10398 \pm 21$	$4726 \pm 9$	
	$B^\pm \rightarrow D\pi^\pm$	$590 \pm 1$	$196804 \pm 398$	
	Other backgrounds	$10402 \pm 105$	$206664 \pm 592$	
	Combinatorial background	$1343 \pm 147$	$15177 \pm 706$	
	$D \rightarrow K_s^0 K^+ K^-$	$B^\pm \rightarrow D^*[D\pi^0]K^\pm$	$790 \pm 3$	$344 \pm 1$
		$B^\pm \rightarrow D^*[D\pi^0]\pi^\pm$	$43 \pm 1$	$14327 \pm 65$
$B^\pm \rightarrow D^*[D\gamma]K^\pm$		$397 \pm 1$	$157 \pm 1$	
$B^\pm \rightarrow D^*[D\gamma]\pi^\pm$		$21 \pm 1$	$7636 \pm 34$	
$B^\pm \rightarrow DK^\pm$		$1527 \pm 6$	$694 \pm 2$	
$B^\pm \rightarrow D\pi^\pm$		$88 \pm 1$	$29786 \pm 135$	
Other backgrounds		$1573 \pm 15$	$31278 \pm 115$	
Combinatorial background		$263 \pm 46$	$4413 \pm 261$	

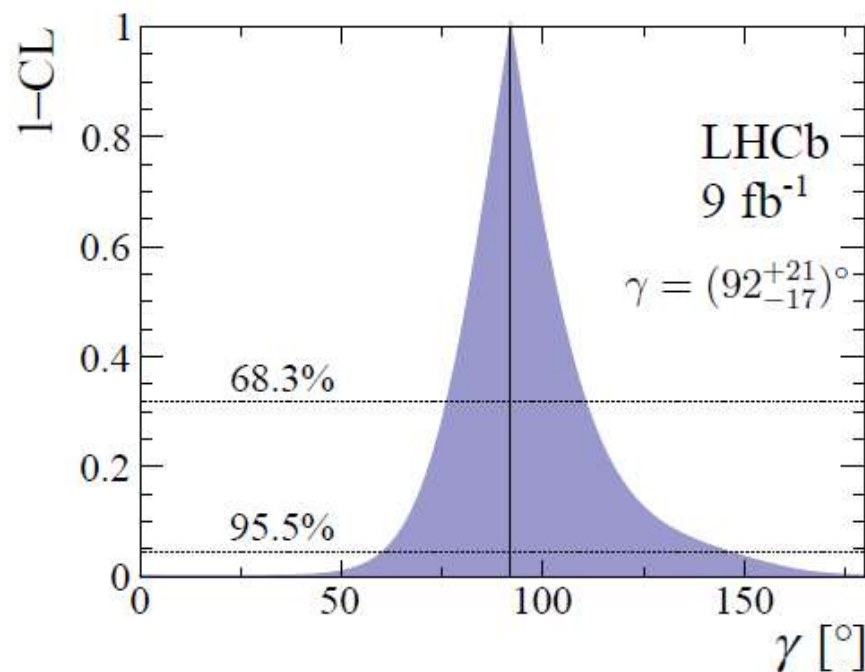
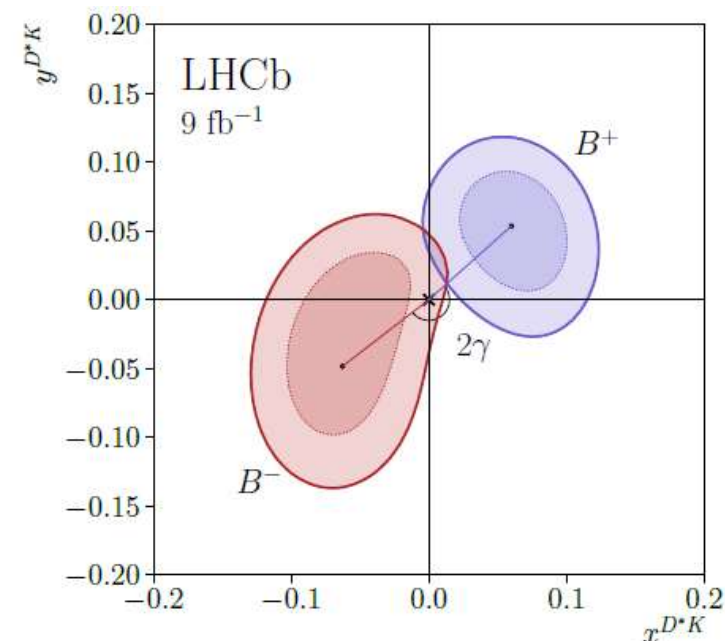
-  $i$ : bin in the Dalitz plot for the  $D$  decay

-  $F_i$   $c_i$   $s_i$  and  $\kappa$  as in pag. 7

- **unknowns:**  $F_i$ ,  $x_\pm$  and  $y_\pm$

$$N_i^+ \propto [F_{-i} + (x_+^2 + y_+^2)F_{+i} + 2\kappa\sqrt{F_{-i}F_{+i}}(c_i x_+ - s_i y_+)]$$

✚ LHCb  $\gamma$  combination:  $B^\pm \rightarrow D^* h^\pm$  with partial  $D^*$  reconstruction



- result consistent with the LHCb  $\gamma$  combo average [LHCb-CONF-2022-003]

- given the accuracy of the BESIII and CLEO data the corresponding systematic uncertainties on  $x_\pm$  and  $y_\pm$  are small ( $< \frac{1}{2}$ ) compared to the LHCb systematic uncertainties
- the total systematic uncertainty is at least a factor of 2 smaller than the statistical uncertainty
- expect improved accuracy from Run 3 data

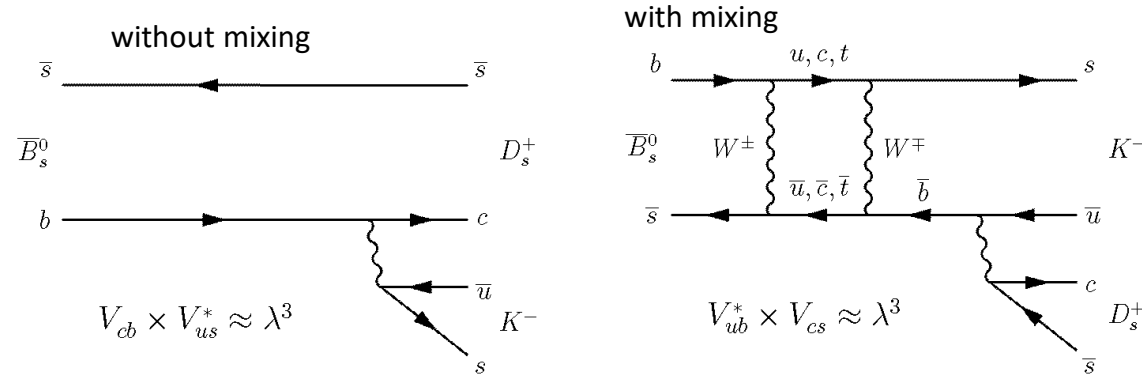
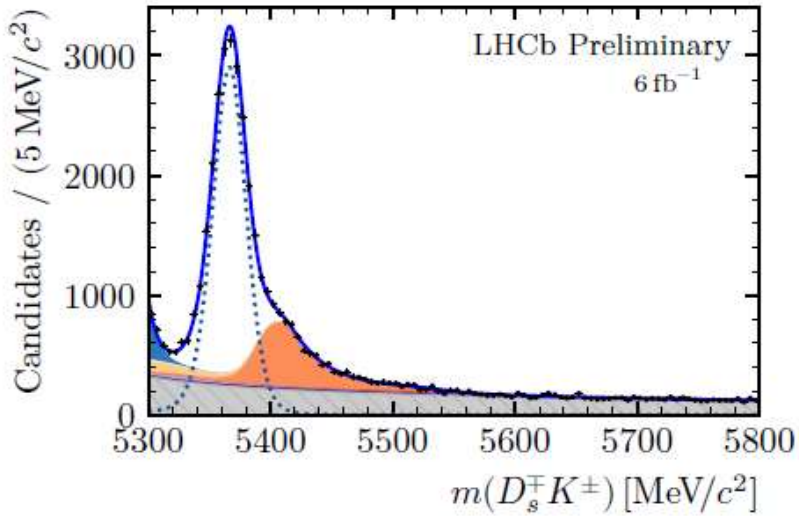
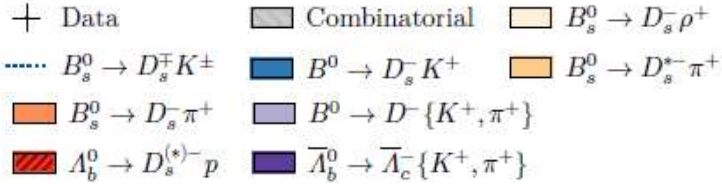
in [JHEP12(2023)013] an exclusive reconstruction of the  $D^* \rightarrow D \gamma / \pi^0$  decay is shown followed by the extraction of  $\gamma$ , the same data set is being used

the signal yield is reduced by approximately 75 %

- the extracted value of  $\gamma$  is compatible with the value obtained from the partially reconstructed approach
- the overall uncertainty on  $\gamma$  is even smaller

LHCb  $\gamma$  combination:  $B_s^0 \rightarrow D_s K$

- sensitivity to  $\gamma$  from interference of decay amplitudes with and without mixing



- time dependent analysis: decay time acceptance obtained from  $B_s \rightarrow D_s \pi$  DATA, corrected for the  $B_s \rightarrow D_s K$  to  $B_s \rightarrow D_s \pi$  MC decay time acceptance ratio (small)
- require flavor tagging

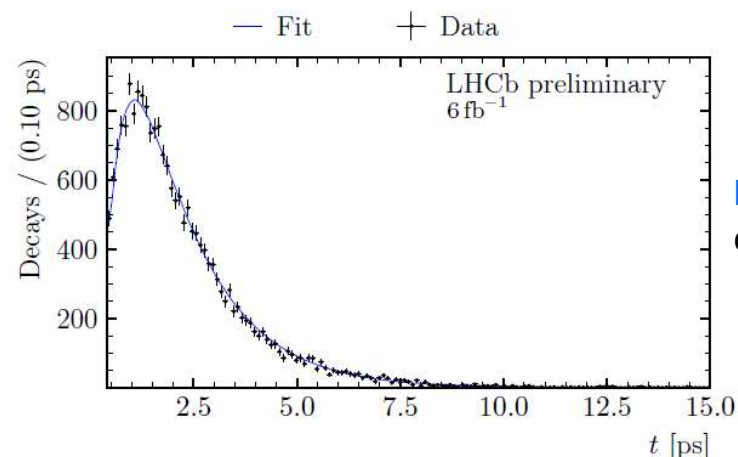
$$\frac{d\Gamma_{B_s^0 \rightarrow f}(t)}{dt} = \frac{1}{2} |A_f|^2 (1 + |\lambda_f|^2) e^{-\Gamma t} \left[ \cosh\left(\frac{\Delta\Gamma t}{2}\right) + A_f^{\Delta\Gamma} \sinh\left(\frac{\Delta\Gamma t}{2}\right) + C_f \cos(\Delta m_s t) - S_f \sin(\Delta m_s t) \right],$$

time dependent decay rates

$$\frac{d\Gamma_{\bar{B}_s^0 \rightarrow f}(t)}{dt} = \frac{1}{2} |A_f|^2 \left| \frac{p}{q} \right|^2 (1 + |\lambda_f|^2) e^{-\Gamma t} \left[ \cosh\left(\frac{\Delta\Gamma t}{2}\right) + A_f^{\Delta\Gamma} \sinh\left(\frac{\Delta\Gamma t}{2}\right) - C_f \cos(\Delta m_s t) + S_f \sin(\Delta m_s t) \right]$$

CP parameters related to  $r_B \delta_B (\gamma - 2\beta_s)$

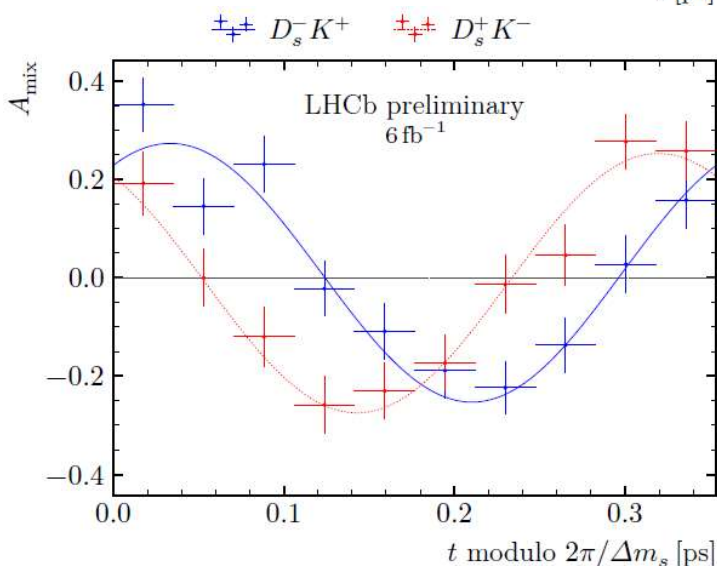
LHCb  $\gamma$  combination:  $B_s^0 \rightarrow D_s K$



blue curve: fit to the decay time distribution

$$\left\{ \begin{array}{l} C_f = 0.791 \pm 0.061 \pm 0.022 \\ A_f^{\Delta\Gamma} = -0.051 \pm 0.134 \pm 0.037 \\ A_{\bar{f}}^{\Delta\Gamma} = -0.303 \pm 0.125 \pm 0.036 \\ S_f = -0.571 \pm 0.084 \pm 0.023 \\ S_{\bar{f}} = -0.503 \pm 0.084 \pm 0.025 \end{array} \right.$$

$C_f = -C_{\bar{f}}$   
 - no CPV in decay  
 - no CPV in mixing  
 - CPV only in the interference



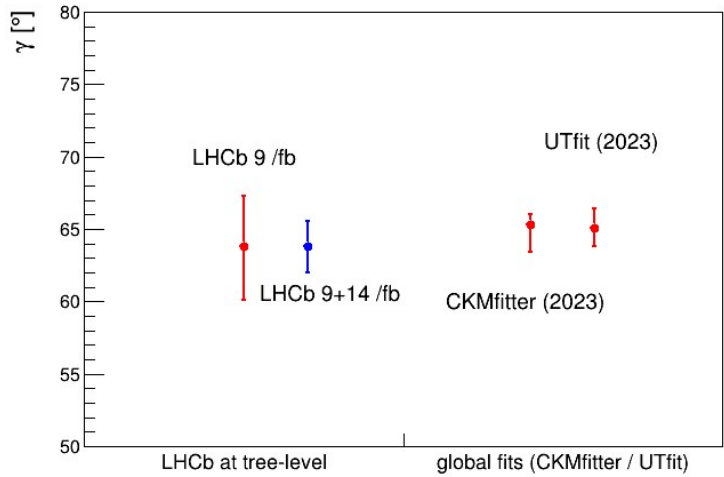
folded asymmetry plots for  $D_s^+ K^-$  and  $D_s^- K^+$   
 CP violation: non trivial phase difference for  $t = 0$  ps

• strictly speaking probing  $\gamma - 2\beta_s$ , using in addition  $\phi_s = -2\beta_s$  and  $\phi_s$  from HFLAV

- 2011  $\rightarrow$  2012 data (3 /fb), JHEP03 (2018) 059  $\gamma = (128^{+17}_{-22})^\circ$
- 2015  $\rightarrow$  2018 data (6/fb), this analysis  $\gamma = (74 \pm 11)^\circ$

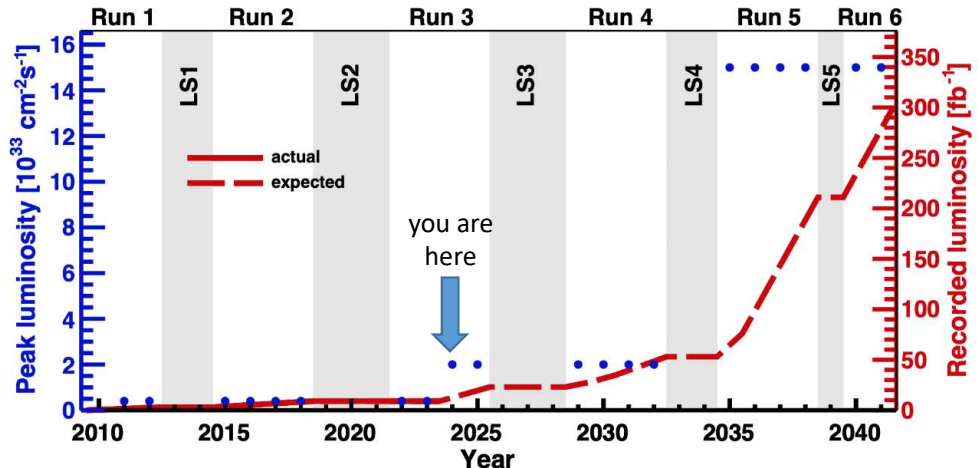
- 2011/2012 result driven by a statistical fluctuation
- 2015/2018 result closer to other channels
- will average them

✚ LHCb  $\gamma$  combination: take home message and future prospects



- 2022 LHCb  $\gamma$  combination based on 9 /fb: 63.8 [+3.5,-3.7]
- CKMfitter: 65.29 [+0.72 -1.86]

- > 10 input measurements that can cross check each other
- present uncertainty is not yet what we need
- “short term”: 9 /fb, Run 1 + Run 2, legacy combination, needed because several recent measurement are not included in the 2021 release



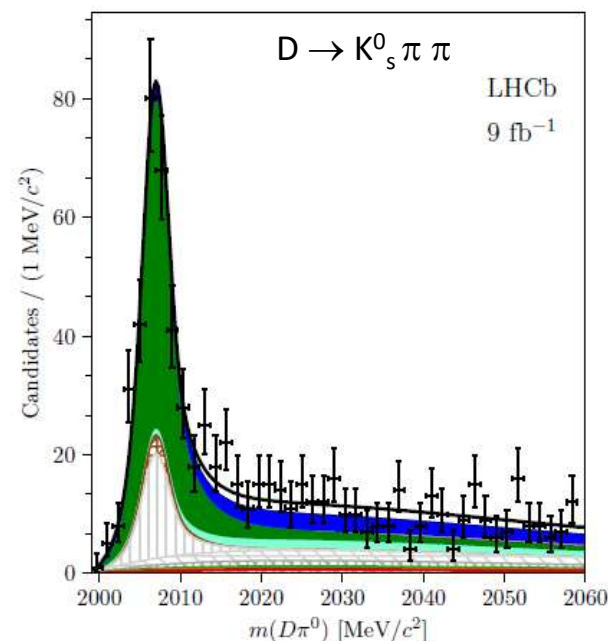
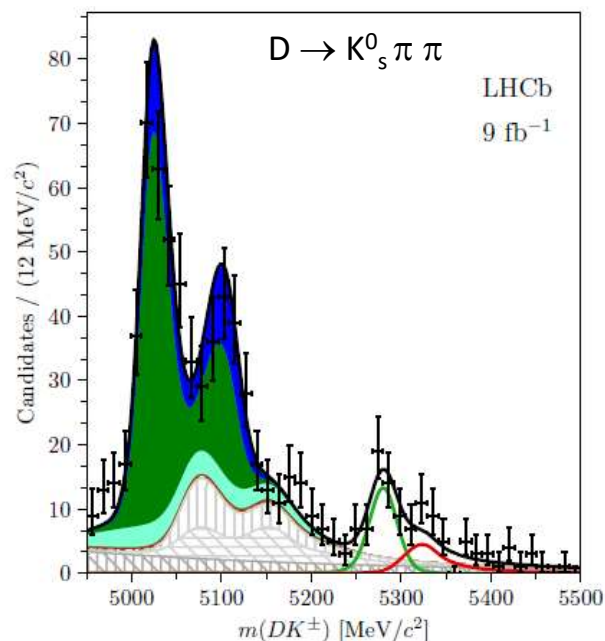
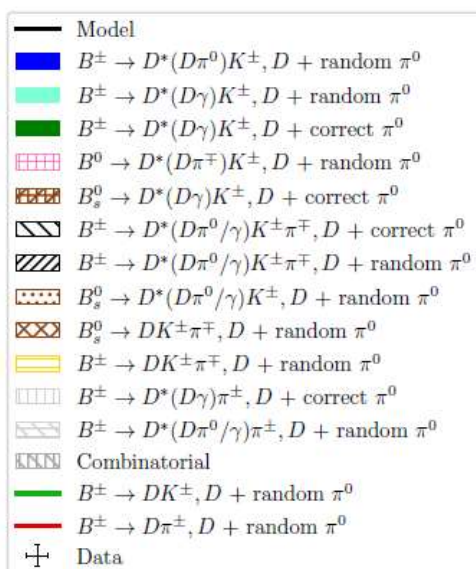
- “longer term”:
- The LHCb Upgrade II, Xuhao Yuan, WG6, Tuesday at 14:00
- fresh data from 2024 on the way, 7 /fb
- 2025 additional 7 /fb
- with 9+14 /fb we expect an accuracy in the range 1.7 - 2.3°
- if the accuracy of the external inputs will not limit the LHCb measurement could reach a 0.35° uncertainty with 300 /fb

- LHCb is doing well, with very significant improvements w.r.t. BaBar and Belle, and has excellent potentialities
- Belle II will also be able to push towards a reduction of the  $\gamma$  uncertainty, expect the same sensitivity

*Backup material*



LHCb  $\gamma$  combination:  $B^\pm \rightarrow D^* h^\pm$  with exclusive  $D^*$  reconstruction

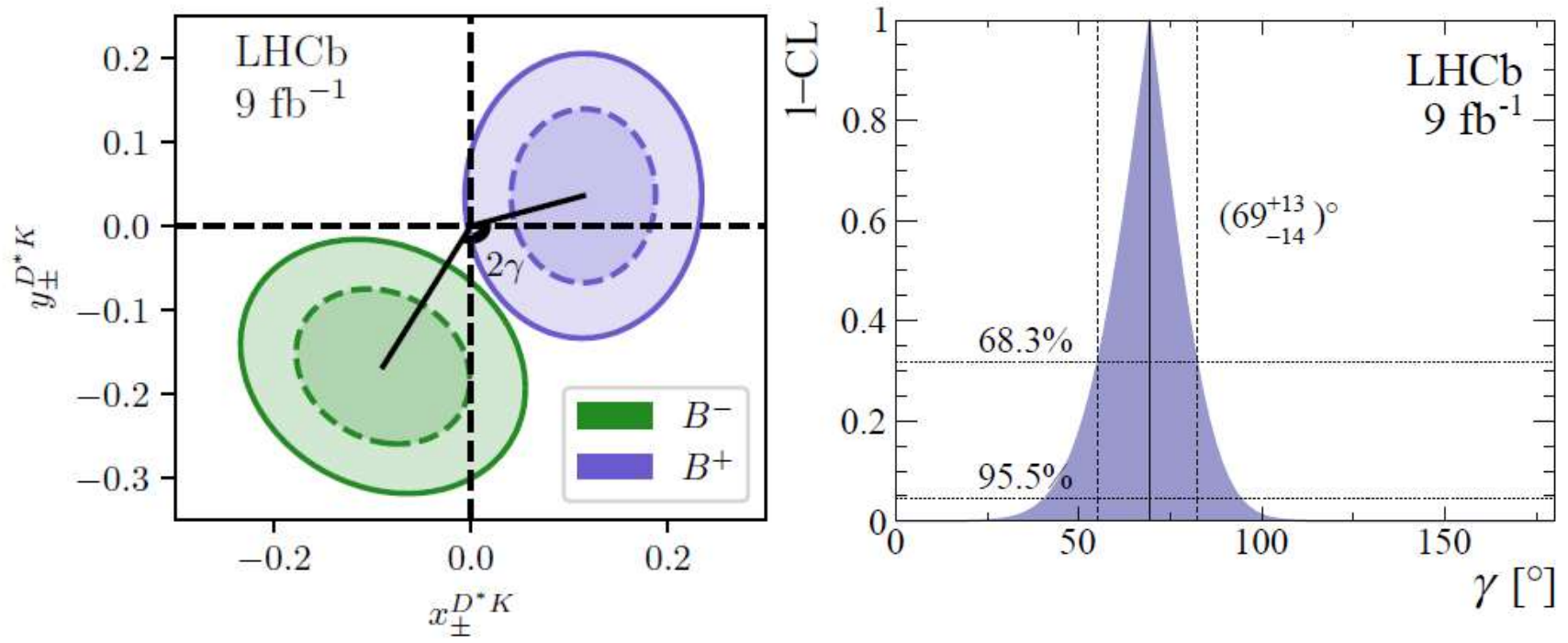


Component	Yield
$B^+ \rightarrow D^* \pi^+, D^* \rightarrow D\pi^0$	$199 \pm 13$
$B^+ \rightarrow D^* \pi^+, D^* \rightarrow D\gamma$	$782 \pm 49$
$B^- \rightarrow D^* \pi^-, D^* \rightarrow D\pi^0$	$197 \pm 13$
$B^- \rightarrow D^* \pi^-, D^* \rightarrow D\gamma$	$740 \pm 48$
$B^+ \rightarrow D^* K^+, D^* \rightarrow D\pi^0$	$13 \pm 2$
$B^+ \rightarrow D^* K^+, D^* \rightarrow D\gamma$	$69 \pm 11$
$B^- \rightarrow D^* K^-, D^* \rightarrow D\pi^0$	$13 \pm 2$
$B^- \rightarrow D^* K^-, D^* \rightarrow D\gamma$	$57 \pm 11$

invariant mass variables providing the best separation between signal and (many) backgrounds are:

- $m(Dh)$   $h=\{\pi, K\}$  not peaking at the B mass because the neutral is excluded
- $m(D\pi^0)$  or  $m(D\gamma)$  peaking at the nominal  $D^*$  mass

• LHCb can successfully use neutrals in a very busy hadron collider environments



- this result is consistent with the LHCb  $\gamma$  combo average [LHCb-CONF-2022-003]
- and dominated by statistical uncertainties

✚ LHCb  $\gamma$  combination: remark on “auxiliary” inputs

Decay	Parameters	Source	Ref.	Status since Ref. [14]
$B^\pm \rightarrow DK^{*\pm}$	$\kappa_{B^\pm}^{DK^{*\pm}}$	LHCb	[33]	As before
$B^0 \rightarrow DK^{*0}$	$\kappa_{B^0}^{DK^{*0}}$	LHCb	[53]	As before
$B^0 \rightarrow D^\mp \pi^\pm$	$\beta$	HFLAV	[13]	As before
$B_s^0 \rightarrow D_s^\mp K^\pm(\pi\pi)$	$\phi_s$	HFLAV	[13]	As before
$D \rightarrow K^+ \pi^-$	$\cos \delta_D^{K\pi}, \sin \delta_D^{K\pi}, (r_D^{K\pi})^2, x^2, y$	CLEO-c	[27]	New
$D \rightarrow K^+ \pi^-$	$A_{K\pi}, A_{K\pi}^{\pi\pi^0}, r_D^{K\pi} \cos \delta_D^{K\pi}, r_D^{K\pi} \sin \delta_D^{K\pi}$	BESIII	[28]	New
$D \rightarrow h^+ h^- \pi^0$	$F_{\pi\pi\pi^0}^+, F_{KK\pi^0}^+$	CLEO-c	[54]	As before
$D \rightarrow \pi^+ \pi^- \pi^+ \pi^-$	$F_{4\pi}^+$	CLEO-c+BESIII	[26, 54]	Updated
$D \rightarrow K^+ \pi^- \pi^0$	$r_D^{K\pi\pi^0}, \delta_D^{K\pi\pi^0}, \kappa_D^{K\pi\pi^0}$	CLEO-c+LHCb+BESIII	[55-57]	As before
$D \rightarrow K^\pm \pi^\mp \pi^+ \pi^-$	$r_D^{K3\pi}, \delta_D^{K3\pi}, \kappa_D^{K3\pi}$	CLEO-c+LHCb+BESIII	[49, 55-57]	As before
$D \rightarrow K_S^0 K^\pm \pi^\mp$	$r_D^{K_S^0 K\pi}, \delta_D^{K_S^0 K\pi}, \kappa_D^{K_S^0 K\pi}$	CLEO	[58]	As before
$D \rightarrow K_S^0 K^\pm \pi^\mp$	$r_D^{K_S^0 K\pi}$	LHCb	[59]	As before

- there are some
- whenever possible these are taken from data
  - whenever possible from LHCb data !

✚ CP violation: historical approach



2000-2008, 0.5 /ab

$$\gamma = (69^{+17}_{-16})^\circ$$

PRD 87, 052015 (2013) “legacy paper”

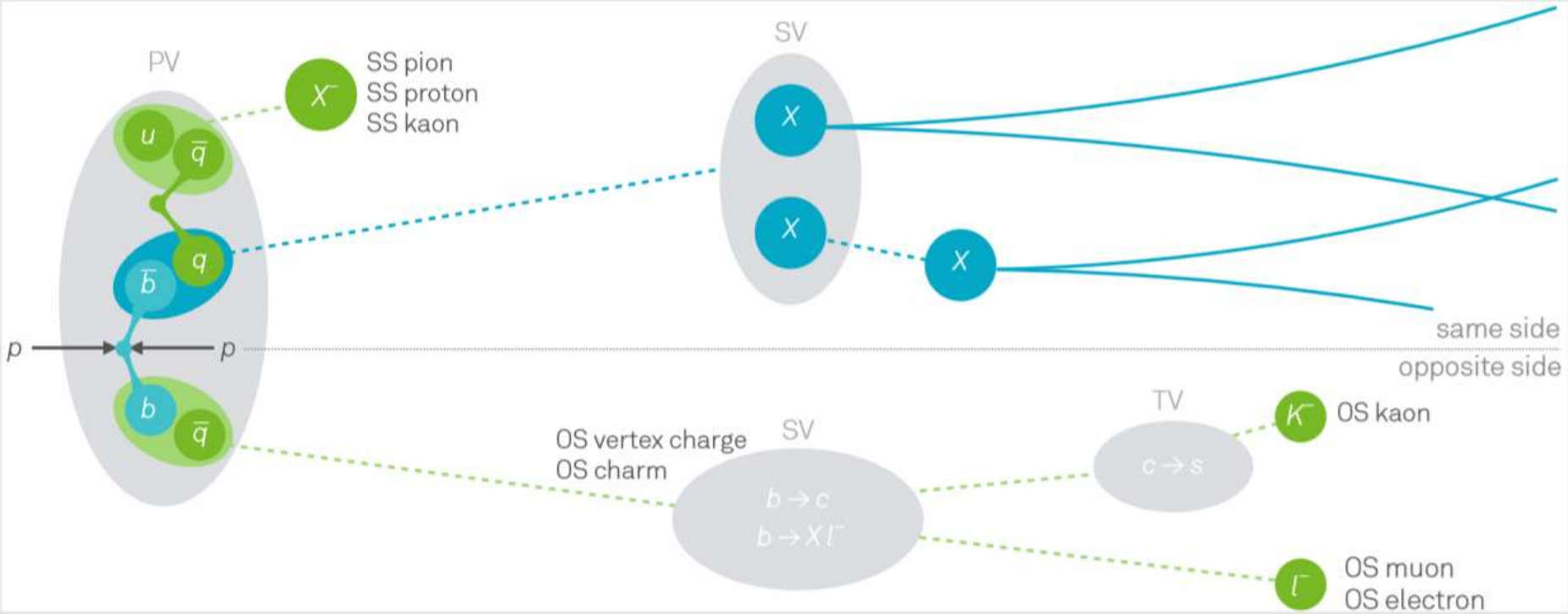


2000-2010, 0.8 /ab

so this presentation will focus on LHCb results, keeping in mind that a new player is coming into the game:



# Flavor tagging



## Tagging performances

CERN-LHCC-2018-027 or  
 LHCb-PUB-2018-009 or  
<https://arxiv.org/abs/1808.08865>

