

# Top-Bottom Interference Contribution to Fully-Inclusive Higgs Production

Based on: [Czakon et al., 2312.09896] and ongoing research

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Tom Schellenberger

DIS, April 10, 2023

In collaboration with M. Czakon, F. Eschment, M. Niggetiedt and R. Poncelet  
Institute for Theoretical Particle Physics and Cosmology

# ~~Top-Bottom Interference Contribution to Fully-Inclusive Higgs Production~~ Quark Mass Effects in Higgs Production

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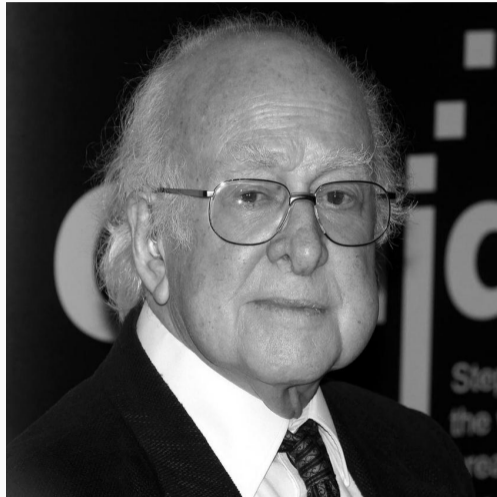
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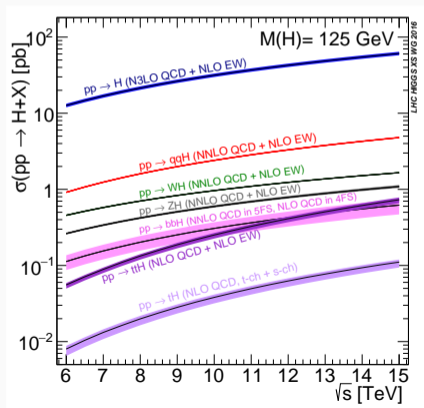
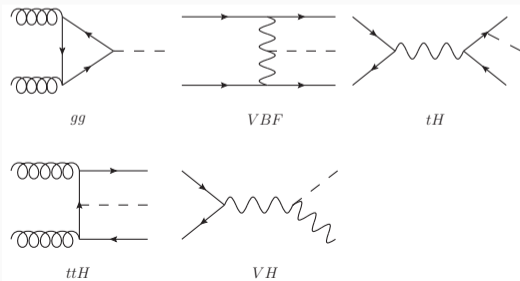
1929-2024

# Higgs Production

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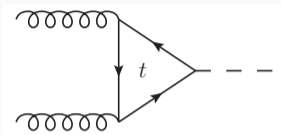
# Motivation

- With the discovery of the Higgs in 2012 we have entered a new era of precision physics.
- We need to know properties of the Higgs very accurately to be able to search for new Physics.
- An important observable here is the Higgs **production cross section**.
- The gluon fusion channel is the most dominant production channel. It is therefore the channel we must determine most precisely.

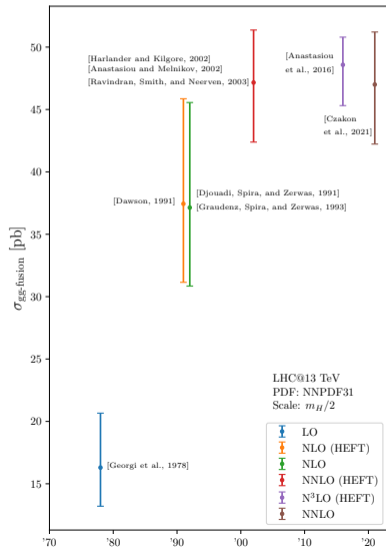


(Florian et al., 1610.07922)

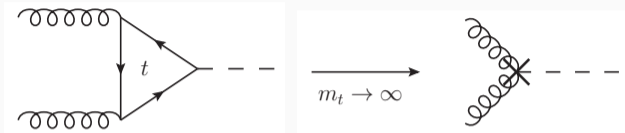
# The Gluon Fusion Channel



- gg-fusion receives large correction ( $\sim 100\%$  at NLO)



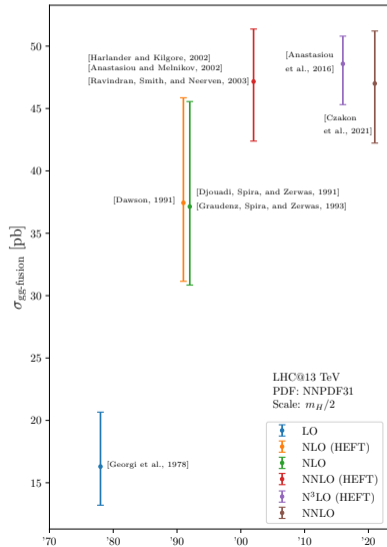
# The Gluon Fusion Channel



- gg-fusion receives large correction ( $\sim 100\%$  at NLO)
- Calculations in full QCD are hard! Instead work in the **Heavy-Top-Limit** (HTL)
- The Heavy-Top-Limit can be improved by rescaling to match the LO predictions: **Higgs-Effective-Field-Theory** (HEFT)

$$\sigma_{\text{HEFT}} = \sigma_{\text{HTL}} \frac{\sigma^{\text{LO}}}{\sigma_{\text{HTL}}^{\text{LO}}}$$

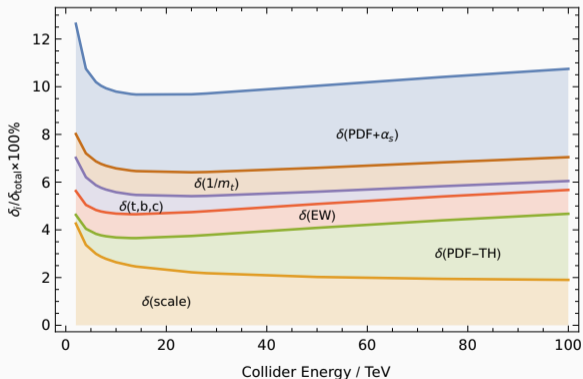
- Effective theory works surprisingly well given  $\frac{m_H^2}{4m_t^2} = 0.13 \ll 1$  is rather poor approximation



# The Gluon Fusion Channel

- Current state of the art for gluon fusion in HEFT is **N3LO** (Anastasiou et al., 1602.00695).

$$\sigma = 48.58 \text{ pb}^{+2.22 \text{ pb}(+4.56\%)}_{-3.27 \text{ pb}(-6.72\%)} (\text{theory}) \pm 1.56 \text{ pb} (3.20\%) (\text{PDF} + \alpha_s)$$



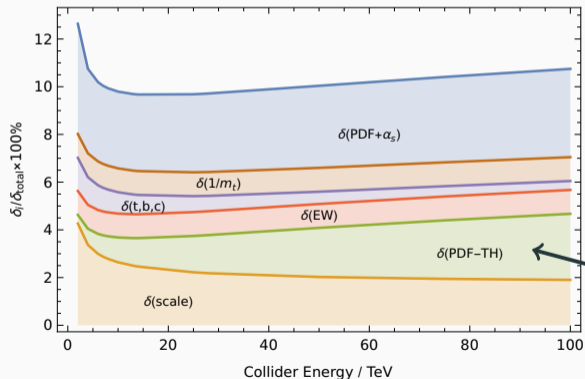
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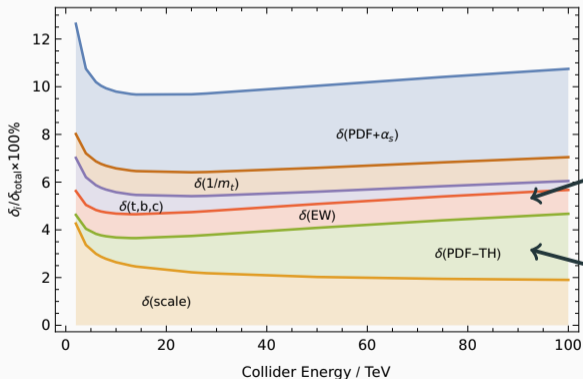
(Dulat, Lazopoulos, and Mistlberger, 1802.00827)

PDFs and their evolution only known at NNLO.  
But first moments are now available and so  
are approximate matched PDF sets  
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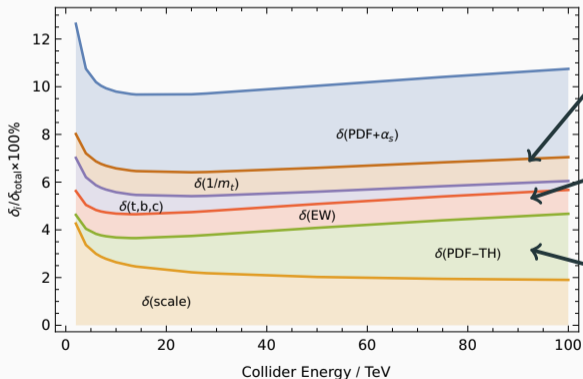
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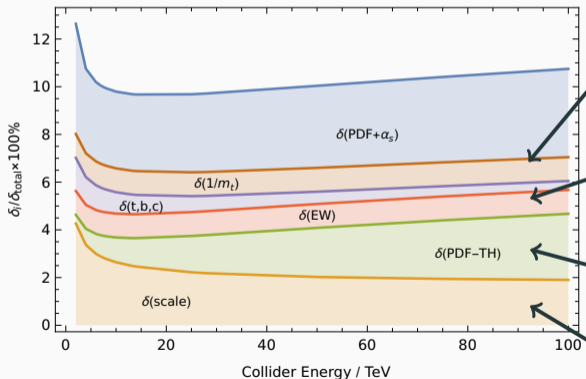
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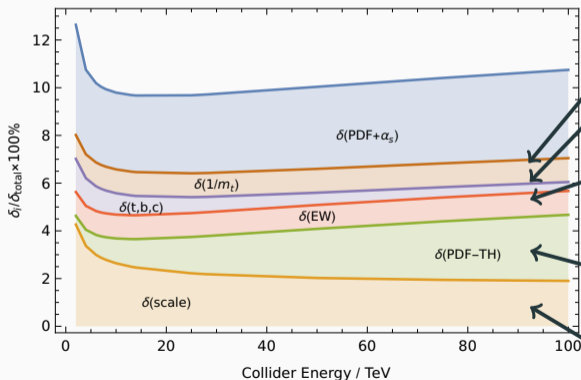
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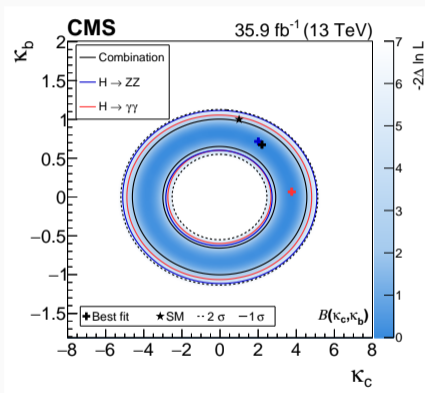
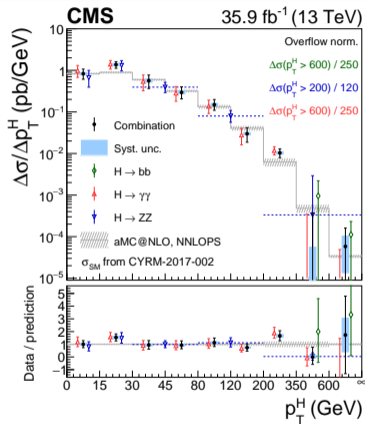
**Our Goal:** Finite  $b$ - and  $c$ -mass effects

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# Yukawa-Couplings



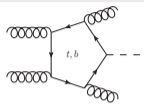
(Sirunyan et al.,  
1812.06504)

- Higgs- $p_T$  and -rapidity distributions can be used to measure the couplings to the lighter quark flavors ( $b, c$ )
- Measuring the charm-quark Yukawa coupling could become possible with HL-LHC

# Computation

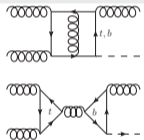
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## Real-Real



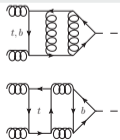
- Easy, since one-loop
- We use MCFM (Budge et al., 2002.04018)

## Real-Virtual



- Mixed contributions factorize
- Use differential equations and solve the rest numerically

## Virtual-Virtual



- Use deep expansion of the Higgs-gluon form factor (Czakon and Niggetiedt, 2001.03008) (Niggetiedt and Usovitsch, 2312.05297)

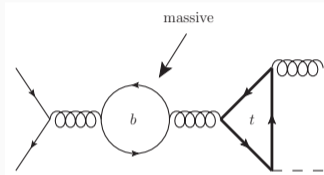
## Phase-Space Integration

- Use MC-techniques to perform phase-space integration
- Infrared divergences are dealt with using Stripper – our in-house C++ implementation of the sector improved residue subtraction scheme (Czakon, 1005.0274)



# Flavor-Schemes

- Care must be taken to have a consistent picture of heavy quarks!

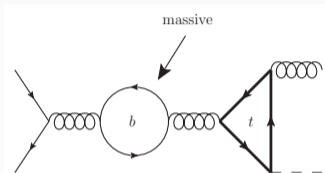


$$\xrightarrow{m_b \rightarrow 0} \ln m_b^2$$

**Mass divergence!**

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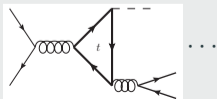


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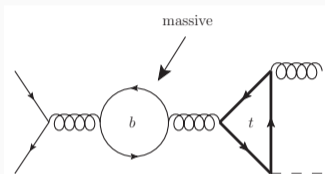
## 4 Flavor-Scheme (4FS)

- Consistently treat bottom-quark as massive
- Exclude bottom-quark from initial state
- Also consider massive  $b$ -quark radiation



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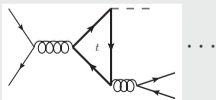


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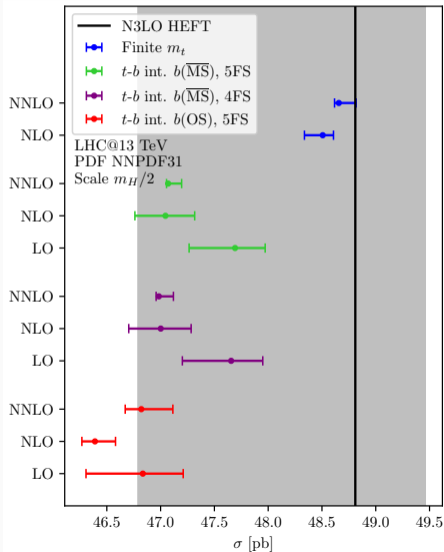
## 5 Flavor-Scheme (5FS)

- Treat bottom-quark as a massless particle
- Except for closed fermion loops that couple to the Higgs
- Equivalent to a theory with replica  $b$ -quark
  - Gauge invariant
  - Divergences cancel

## Results

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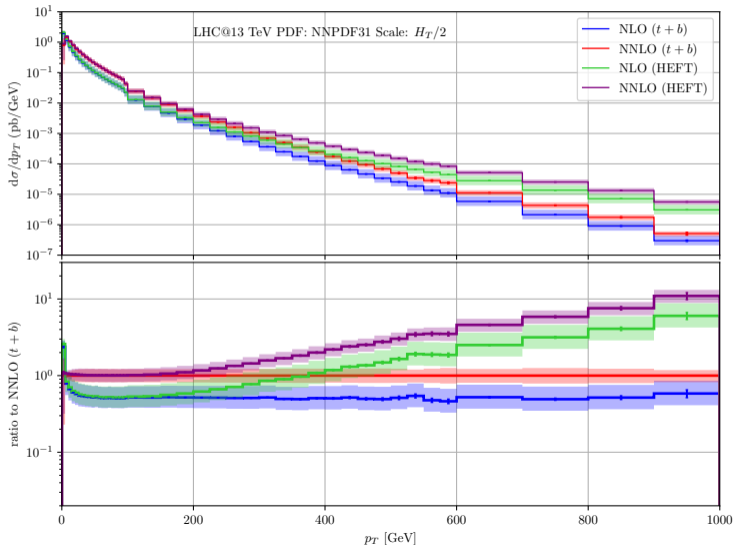
# Fully-Inclusive Cross Sections



## Preliminary

- On-Shell scheme result shows **increased** scale uncertainties
- $\overline{\text{MS}}$  scheme shows better convergence
- 4FS and 5FS show very good agreement, justifying the treatment of the bottom-quark mass in the 5FS.

# $p_T$ -distributions

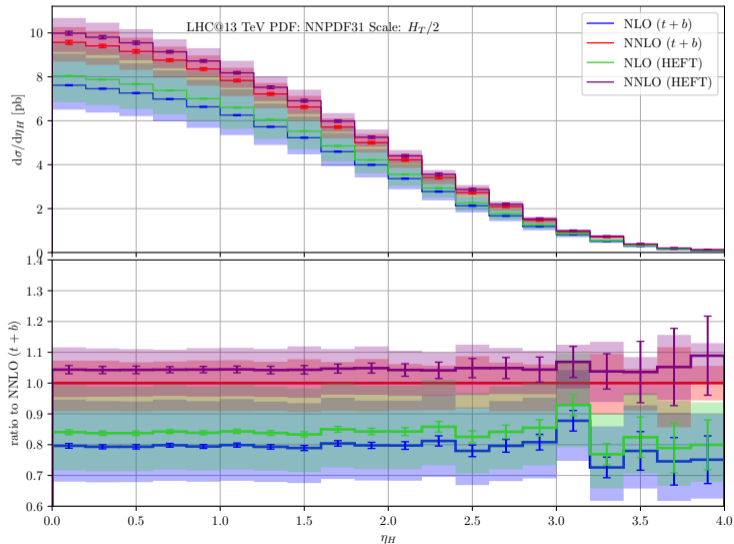


## Preliminary

$b(\overline{\text{MS}})$ , 5FS

- Notice the different scaling behavior of the HEFT ( $p_{T,H}^{-2}$ ) and full QCD ( $p_{T,H}^{-4}$ ) result (Greiner et al., 1608.01195).
- Good agreement with existing data (Lindert et al., 1703.03886), (Jones, Kerner, and Luisoni, 1802.00349)

# Rapidity-distributions



**Preliminary**

$b(\overline{\text{MS}})$ , 5FS

- Finite Quark mass effects are less pronounced for the rapidity distributions.
- Less sensitive for fitting to Yukawa couplings. This was already noticed at LO (Soreq, Zhu, and Zupan, 1606.09621)

## Conclusions

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# Conclusions

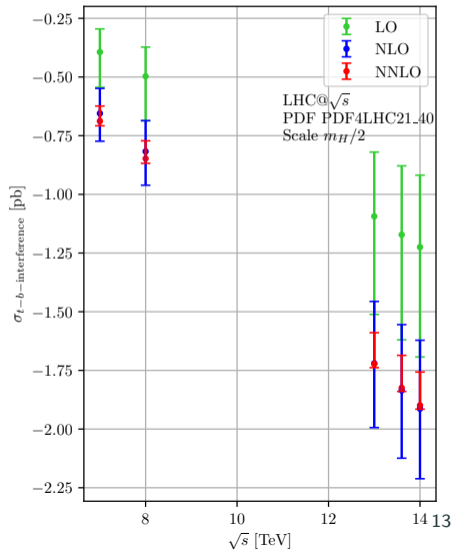
- Including finite quark mass effects for the Higgs-production cross section is crucial for precision physics at the LHC
- The  $\overline{\text{MS}}$ -scheme shows generally a better perturbative convergence
- The 4FS and 5FS showed very good agreement, i.e. the treatment of the bottom quark as a massless particle is justified

LHC@13 TeV, PDF NNPDF31, Scale  $m_H/2$

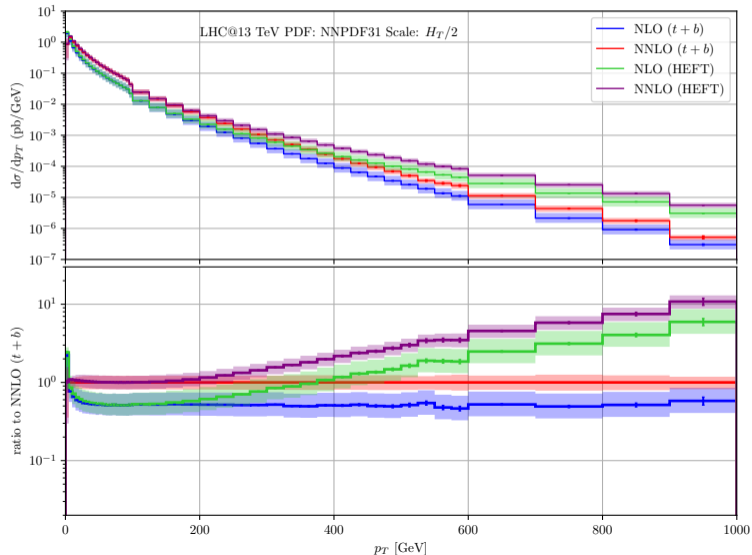
**Preliminary**

$$\sigma_{gg\text{-fusion}} = \underbrace{48.81^{+0.65}_{-2.02}}_{\text{HEFT}} \underbrace{-0.16^{+0.13}_{-0.03}}_{\text{finite } t} \underbrace{-1.74^{+0.13}_{-0.01}}_{t\text{-}b \text{ int.}} \text{ pb}$$

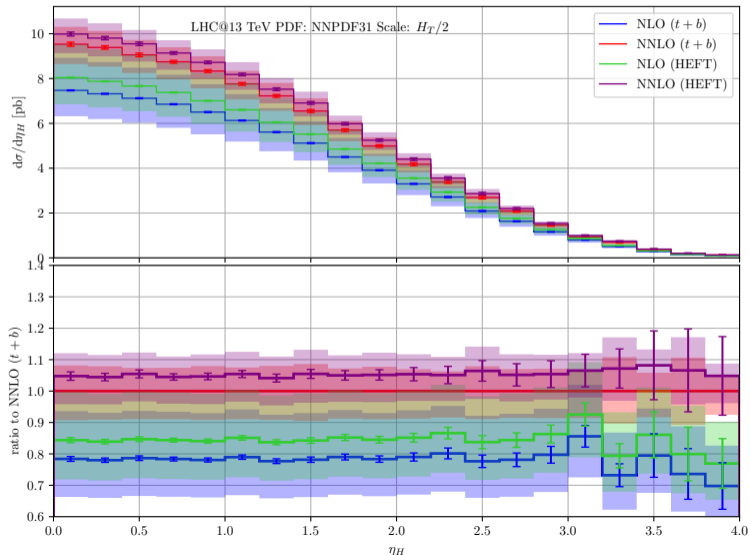
- $m_b(m_b) = 4.18 \text{ GeV}$
- $m_t = 173.05 \text{ GeV}$
- $m_H = 125 \text{ GeV}$



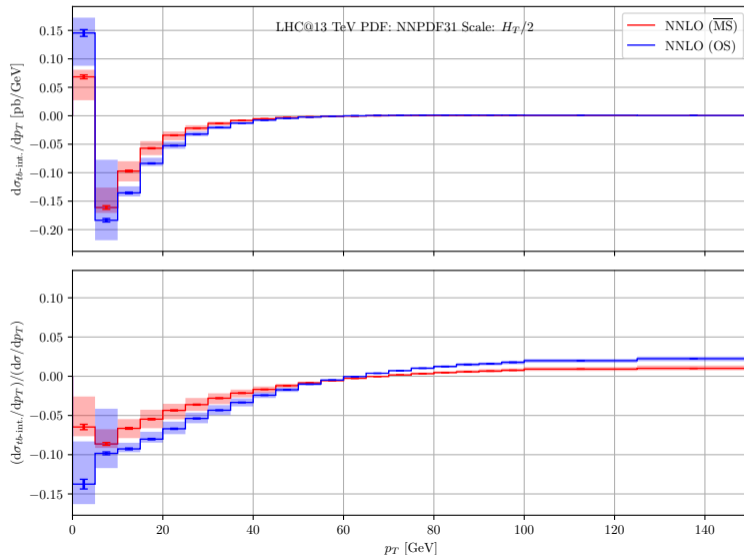
# $p_T$ -distributions (OS, 5FS)



# Rapidity-distributions (OS, 5FS)



# Top-Bottom Interference $p_T$ -Distribution



# Top-Bottom Interference Rapidity-Distribution

