

# Discovering (true) tauonium at colliders

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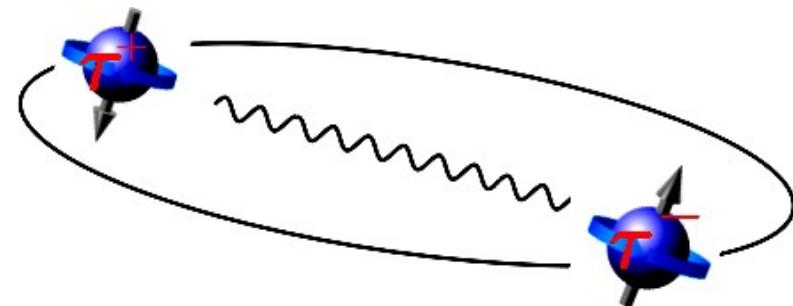
8-12 April 2024, Grenoble, France

Details: arXiv:2202.02316 [hep-ph], 2204.07269 [hep-ph], arXiv:2302.07365 [hep-ph]

# Exotic leptonium atoms

- Opposite-charge leptons ( $\ell^\pm = e^\pm, \mu^\pm, \tau^\pm$ ) can form transient “onium” bound states under their QED interaction. Out of 6 possible exotic leptonic atoms ( $e^+e^-$ ), ( $\mu^\pm e^\mp$ ), ( $\mu^+\mu^-$ ), ( $\tau^\pm e^\mp$ ), ( $\tau^\pm \mu^\mp$ ), ( $\tau^+\tau^-$ ), only the two first (positronium in 1951) and (muonium in 1960) have been observed.

→ Para- ( $J^{PC} = 0^{-+}$ ) and ortho- ( $J^{PC} = 1^{--}$ ) leptonium ground states form depending on relative spin orientation of leptons.



- Ditaunuonium  $\tau\bar{\tau}(\tau^+\tau^-)$ , barely studied, is smallest & most-bound leptonium state:

Mass:  $m_{\tau} = 2m_{\tau} + E_{\text{bind}} = 3553.6962 \pm 0.2400 \text{ MeV}$ ,  $E_{\text{bind}} = -\alpha^2 m_{\tau} / (4n^2) = -23.7 \text{ keV}$

Bohr radius:  $a_0 = 2/(\alpha m_{\tau}) = 30.4 \text{ fm}$  ( $\times 3500$  smaller than positronium)

Rydberg const ( $\gamma$  ionization):  $R_{\infty} = m_{\tau} \alpha^2 / 4\pi = 3.76 \text{ keV}$  ( $\times 3500$  larger than positronium)

- Compared to other exotic atoms, ditaunuonium can provide:

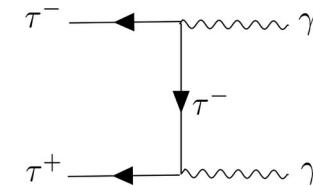
- Precision SM: Most competitive measurement of the tau mass possible.
- New tests of QED & CPT symmetries at much higher masses (smaller distances).
- Sensitivity to any BSM enhanced by  $(m_{\ell}/\Lambda_{\text{BSM}})^n$ , unaffected by hadronic uncertainties.

# Ditauonium partial widths & decays

- Para- $\tau$  decays mostly to  $\gamma\gamma$  (BR  $\approx 80\%$ ):

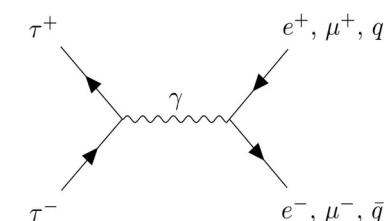
- Ortho- $\tau$  has many open channels:  $e^+e^-$ ,  $\mu^+\mu^-$ ,  $q\bar{q}$   
BR  $\approx 20\%$ , 20%, 45%

$$\Gamma^{(0)}(n^1S_0 \rightarrow \gamma\gamma) = \frac{\alpha^5 m_\tau}{2 n^3} \Big|_{n=1} = 0.018384 \text{ eV}$$



$$\Gamma^{(0)}(n^3S_1 \rightarrow e^+e^-, \mu^+\mu^-) = \frac{\alpha^5 m_\tau}{6n^3}$$

$$\Gamma^{(0)}(n^3S_1 \rightarrow q\bar{q}) = \frac{\alpha^5 m_\tau}{6n^3} R_{\text{had}}(m_{\tau_0}^2) = 2.2 \frac{\alpha^5 m_\tau}{6n^3}$$



- Weak decay of constituent  $\tau^\pm$ :  $\Gamma_{(2)\tau \rightarrow X} = 2/\tau = 0.004535 \text{ eV}$  ( $\tau \approx 290 \text{ fs}$ )  
BR<sub>eff</sub>  $\approx 19\%, 14\%$  for para-,ortho- $\tau$

- Ditauonium spectroscopy (NNLO\* non-relativ. QED):

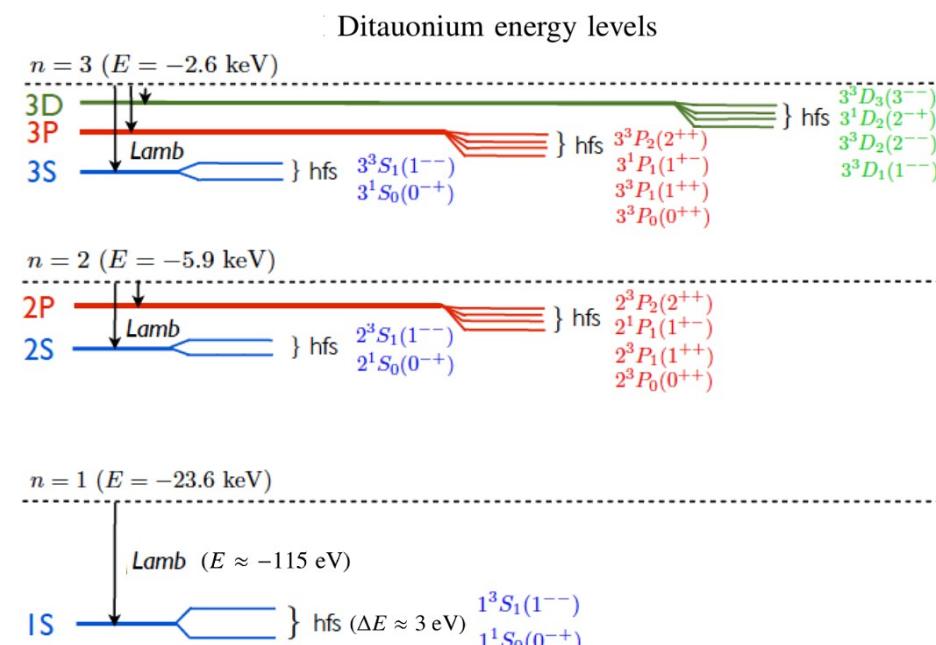
→ Lamb shifts:

$$\Delta E^{1S,2S,\dots} = -115.4, -14.4, \dots \text{ eV}$$

→ Hyperfine splittings:

$$\Delta E_{\text{hfs}}(1^1S_0, 1^3S_1, \dots) = -1.65, +1.29, \dots \text{ eV}$$

[DdE, R.Perez-Ramos, H-S. Shao:  
arXiv:2204.07269]

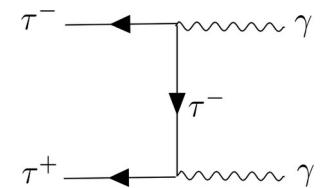


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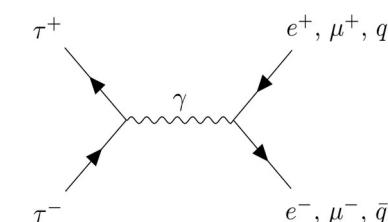
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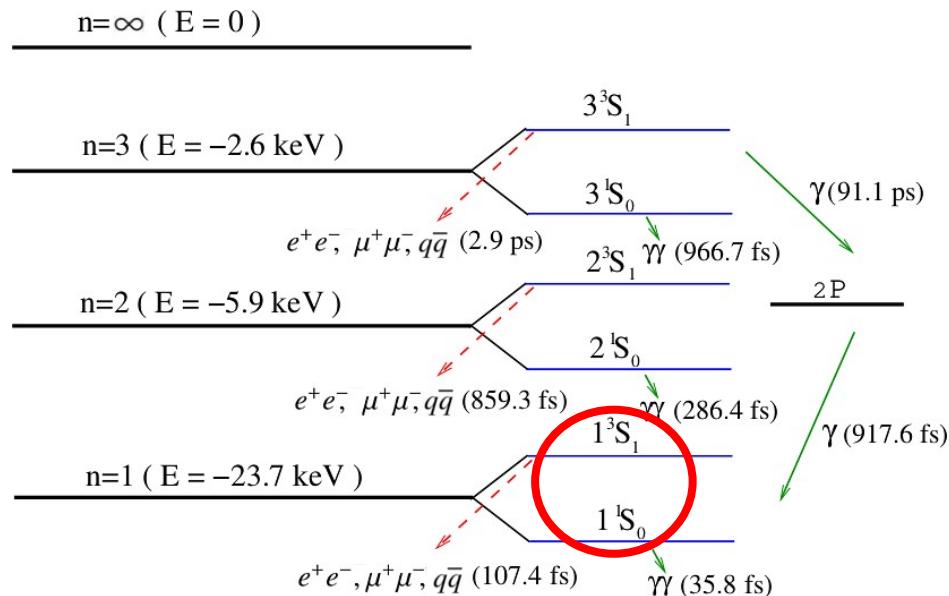
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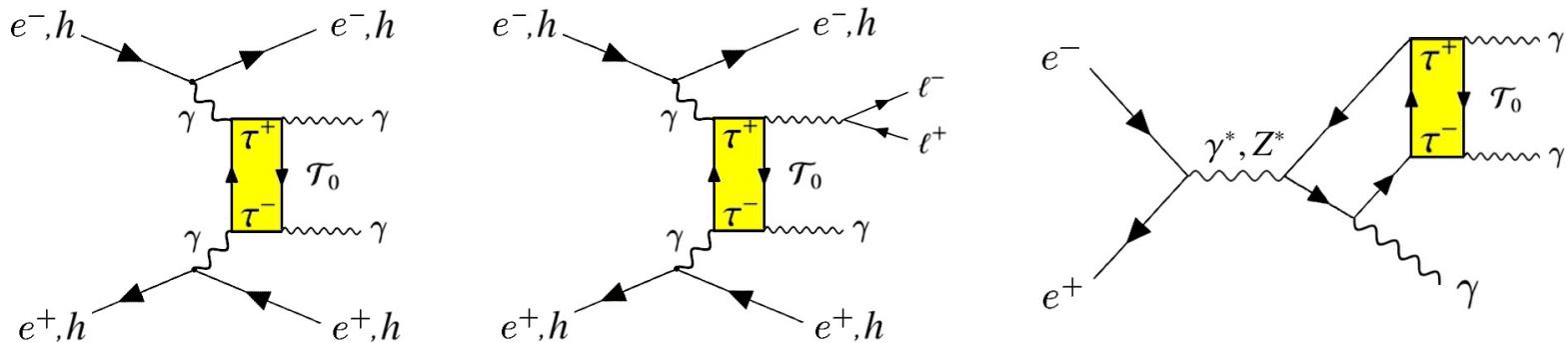
- 
- Ditauonium spectroscopy (NNLO\* non-relativ. QED):
    - Only the two lowest states ( $1^1S_0$  &  $1^3S_1$ ) have lifetimes shorter ( $\tau \approx 27.6, 20.83 \text{ fs}$ ) than the weak decay of the constituents tau's.

[DdE, R.Perez-Ramos, H-S. Shao:  
arXiv:2204.07269]

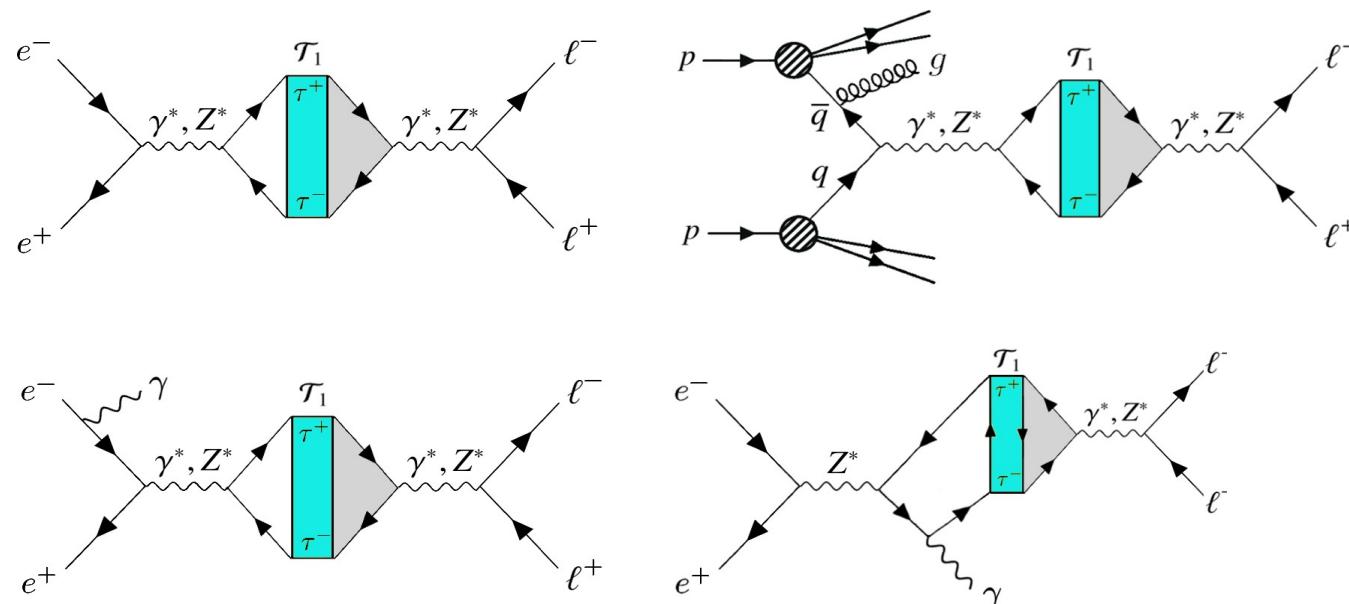


# Ditauonium production at $e^+e^-$ & hadron colliders

- 3 para-ditauonium prod./decay channels: photon-photon, s-channel+ $\gamma$

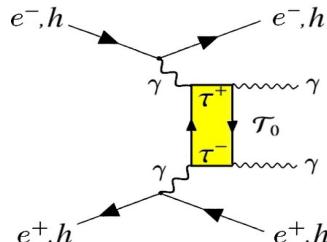


- 4 ortho-ditauonium prod./decay channels: s-channel fusion (w/ & w/o  $\gamma$ )



# Para-ditauonium via $\gamma\gamma \rightarrow \tau_0 \rightarrow \gamma\gamma$

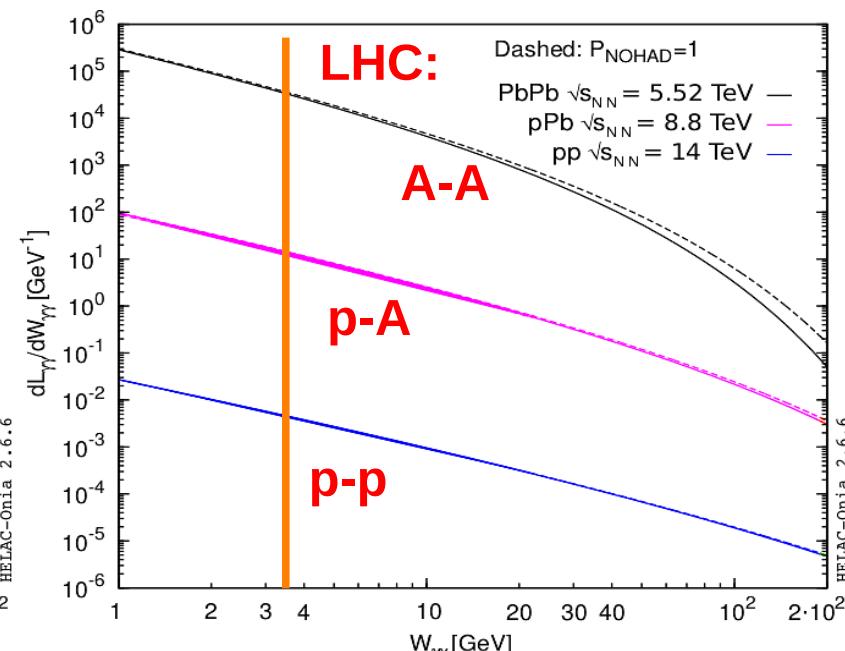
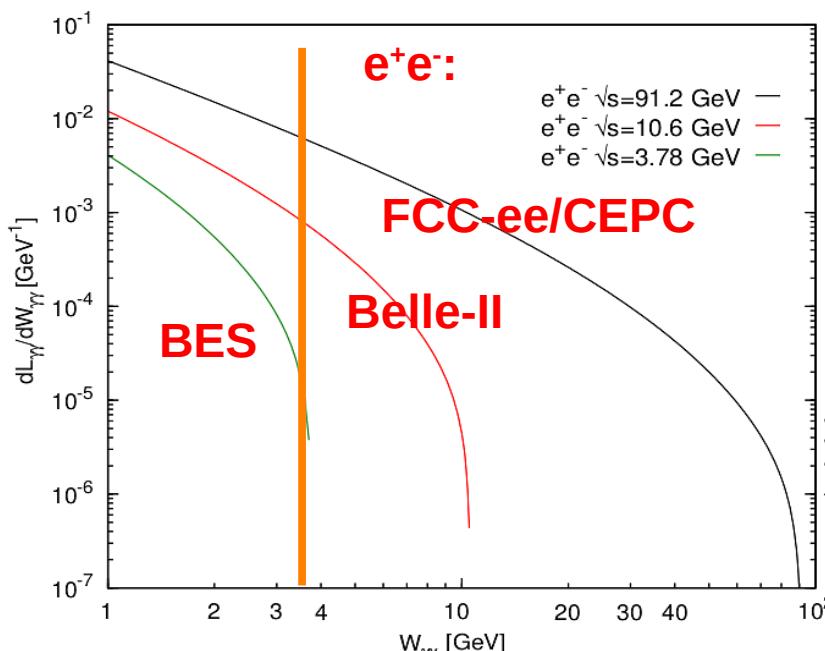
- Cross sections for signal & backgrounds computed in the Weizsäcker-Williams approximation (EPA) for  $\gamma\gamma$  collisions via gamma-UPC 2207.03012 [hep-ph]



$$\sigma(a b \rightarrow a b + X) = 4\pi^2(2J+1) \frac{\Gamma_{\gamma\gamma}(X)}{m_X^2} \left. \frac{d\mathcal{L}_{\gamma\gamma}^{(ab)}}{dW_{\gamma\gamma}} \right|_{W_{\gamma\gamma}=m_X}$$

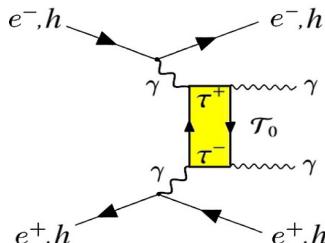
- Photon-photon luminosity for  $e^+e^-$  & ultraperipheral p-p, p-A & A-A collisions

$$\frac{d\mathcal{L}_{\gamma\gamma}^{(AB)}}{dW_{\gamma\gamma}} = \frac{2W_{\gamma\gamma}}{s_{NN}} \int \frac{dE_{\gamma_1}}{E_{\gamma_1}} \frac{dE_{\gamma_2}}{E_{\gamma_2}} \delta\left(\frac{W_{\gamma\gamma}^2}{s_{NN}} - \frac{4E_{\gamma_1}E_{\gamma_2}}{s_{NN}}\right) \frac{d^2N_{\gamma_1/\gamma_2/Z_1,Z_2}^{(AB)}}{dE_{\gamma_1} dE_{\gamma_2}}$$



# Para-ditauonium via $\gamma\gamma \rightarrow \tau_0 \rightarrow \gamma\gamma$ : Backgrounds

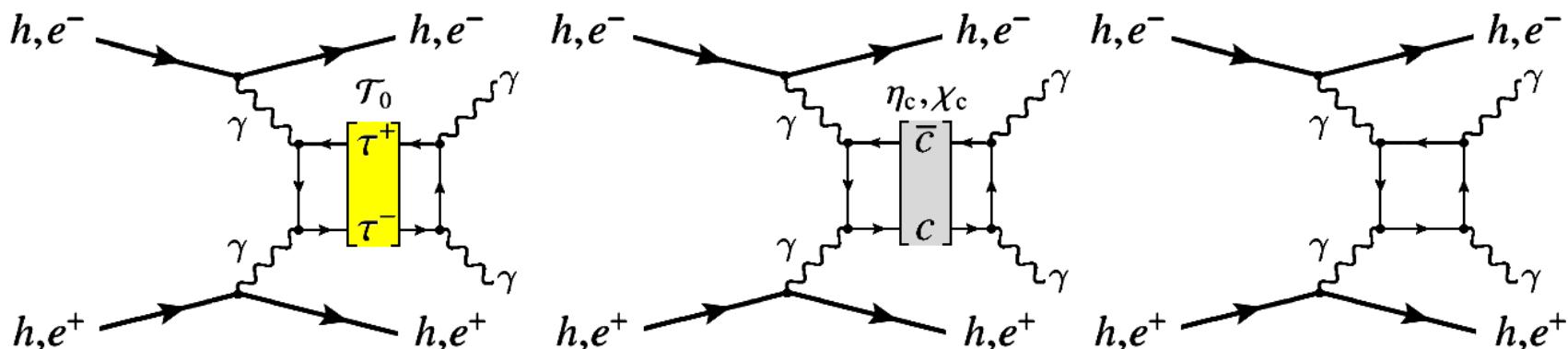
- Cross sections for signal & backgrounds computed in the Weizsäcker-Williams approximation (EPA) for  $\gamma\gamma$  collisions via gamma-UPC 2207.03012 [hep-ph]



$$\sigma(a b \rightarrow a b + X) = 4\pi^2(2J+1) \frac{\Gamma_{\gamma\gamma}(X)}{m_X^2} \left. \frac{d\mathcal{L}_{\gamma\gamma}^{(ab)}}{dW_{\gamma\gamma}} \right|_{W_{\gamma\gamma}=m_X}$$

- **Backgrounds** within  $m_{\gamma\gamma} \approx 2.9\text{--}3.7 \text{ GeV}$ :

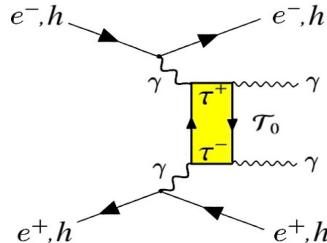
- C-even charmonium: 3  $c\bar{c}$ :  $\eta_c(2S)$ ,  $\chi_{c1,2}$  resonances within  $\sim 100 \text{ MeV}$  of  $\tau_0$
- Light-by-light scattering (LbL) continuum.



- Charmonia resonances have  $\mathcal{O}(\text{keV})$  diphoton widths:  $\mathcal{O}(10^5)$  larger than para- $\tau_0$ .  
But, diphoton BR is  $\mathcal{O}(10^4)$  larger for para- $\tau_0$  than for  $c\bar{c}$  states.

# Para-ditauonium via $\gamma\gamma \rightarrow \tau_0 \rightarrow \gamma\gamma$ : Yields

- Cross sections for signal & backgrounds computed in the Weizsäcker-Williams approximation (EPA) for  $\gamma\gamma$  collisions via gamma-UPC 2207.03012 [hep-ph]



$$\sigma(a b \rightarrow a b + X) = 4\pi^2(2J+1) \frac{\Gamma_{\gamma\gamma}(X)}{m_X^2} \left. \frac{d\mathcal{L}_{\gamma\gamma}^{(ab)}}{dW_{\gamma\gamma}} \right|_{W_{\gamma\gamma}=m_X}$$

- Results for  $e^+e^-$  and ultraperipheral p-p, p-A & A-A collisions:

Colliding system, c.m. energy, $\mathcal{L}_{\text{int}}$ , exp.	$\sigma \times \mathcal{B}_{\gamma\gamma}$					$N \times \mathcal{B}_{\gamma\gamma}$		
	$\eta_c(1S)$	$\eta_c(2S)$	$\chi_{c,0}(1P)$	$\chi_{c,2}(1P)$	LbL	$\mathcal{T}_0$	$\mathcal{T}_0$	$\chi_{c,2}(1P)$
$e^+e^-$ at 3.78 GeV, 20 $\text{fb}^{-1}$ , BES III	120 fb	3.6 ab	15 ab	13 ab	30 ab	0.25 ab	–	–
$e^+e^-$ at 10.6 GeV, 50 $\text{ab}^{-1}$ , Belle II	1.7 fb	0.35 fb	0.52 fb	0.77 fb	1.7 fb	0.015 fb	750	38 500
$e^+e^-$ at 91.2 GeV, 50 $\text{ab}^{-1}$ , FCC-ee	11 fb	2.8 fb	3.9 fb	6.0 fb	12 fb	0.11 fb	5 600	$3 \cdot 10^5$
p-p at 14 TeV, 300 $\text{fb}^{-1}$ , LHC	7.9 fb	2.0 fb	2.8 fb	4.3 fb	6.3 fb	0.08 fb	24	1290
p-Pb at 8.8 TeV, 0.6 $\text{pb}^{-1}$ , LHC	25 pb	6.3 pb	8.7 pb	13 pb	21 pb	0.25 pb	0.15	8
Pb-Pb at 5.5 TeV, 2 $\text{nb}^{-1}$ , LHC	61 nb	15 nb	21 nb	31 nb	62 nb	0.59 nb	1.2	62

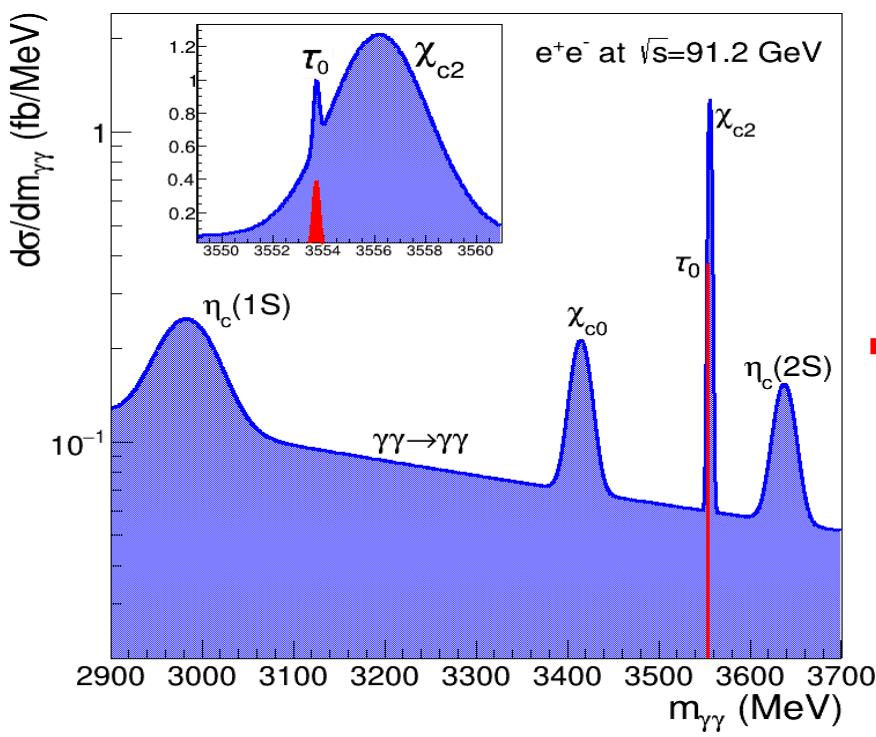
- Relative prod. x-sections:  $\eta_c(1S):\chi_{c2}(1P):\chi_{c0}(1P):\eta_c(2S):\tau_0 \approx 100:50:30:25:1$
- Para- $\tau_0$  x-sections increase with  $\sqrt{s}$  and  $Z^4$ :

Largest x-sections (0.6 nb) in PbPb UPC (but handful of evts expected at LHC)  
 Largest yields: 750, 5600 counts at Belle-II, FCC-ee thanks to  $\mathcal{L}_{\text{int}} = 50 \text{ ab}^{-1}$ .

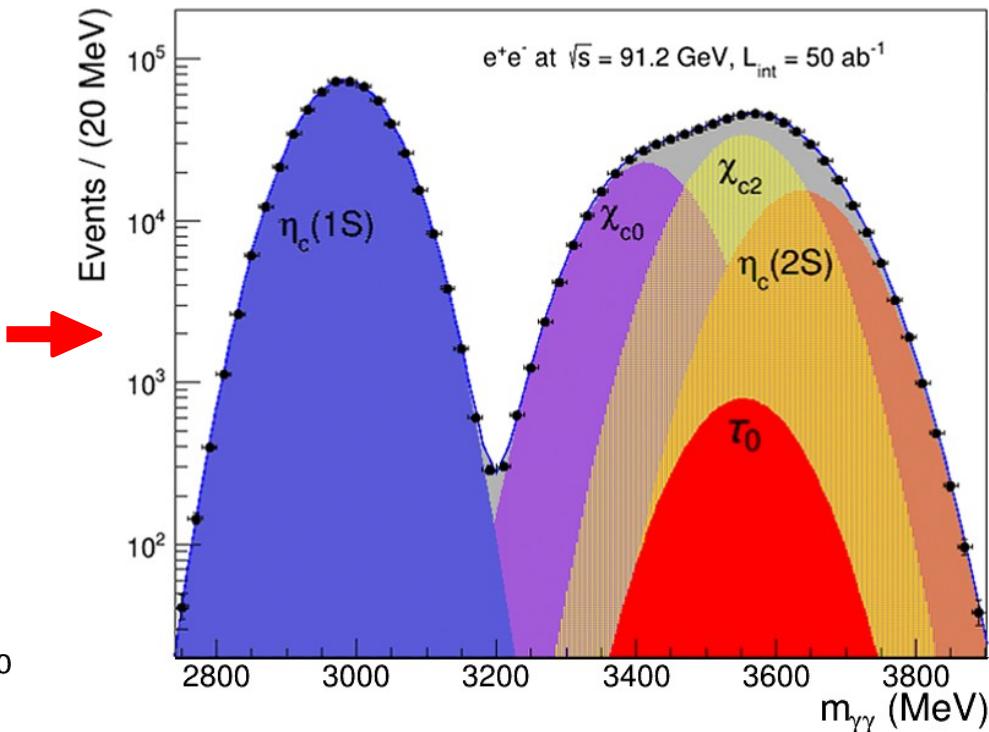
# Para-ditauonium via $\gamma\gamma \rightarrow \tau_0 \rightarrow \gamma\gamma$ (Belle II/FCC-ee)

- Trigger: Require two exclusive 1.5–2 GeV photons back-to-back with  $m_{\gamma\gamma} \approx m_{\tau_0}$
- Reco. performances (Belle-II type: high-reso low-energy crystal ECAL):  
Acceptance:  $10^\circ < \theta_\gamma < 170^\circ$ . Mass resolution: ~2%. Photon reco effic. ~100%.  
→ All diphoton resonances Gaussian-smeared with ~70 MeV widths:

Generator-level x-sections (0.1-MeV  $\tau_0$  width)



Reconstructed yields (LbL subtracted)



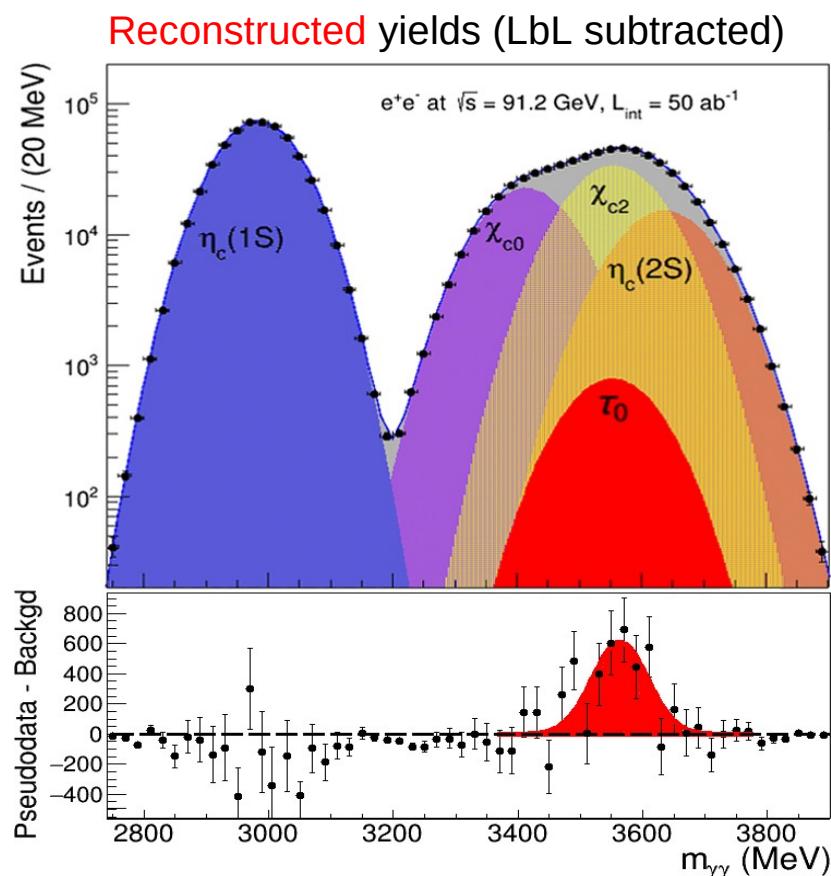
- Ditauonium signal swamped by overlapping  $\chi_{c2}(1P)$  & neighboring  $\chi_{c0}(1P)$ ,  $\eta_c(2S)$

# Para-ditauonium via $\gamma\gamma \rightarrow \tau_0 \rightarrow \gamma\gamma$ (Belle II/FCC-ee)

- 1-million events generated for signal & backgrounds. Run MVA (BDT) with 12 different single- $\gamma$  and  $\gamma$ -pair kinematic variables for signal/backgds separation:
  - (i) Strong discrimination power (factor of  $\sim 20$ ) of LbL continuum from signal.
  - (ii) No discrimination achieved for overlapping charmonia (decay- $\gamma$  angular modulation of tensor  $\chi_{c2}$  different than scalar  $\tau_0$  signal, but  $\times 50$  suppressed yields).
- Signal extracted through multi-Gaussian  $m_{\gamma\gamma}$  fit.
- Statistical significance derived from profile-likelihood of fits assuming signal presence or backgd-only, with 0.3% background syst. uncertainty:

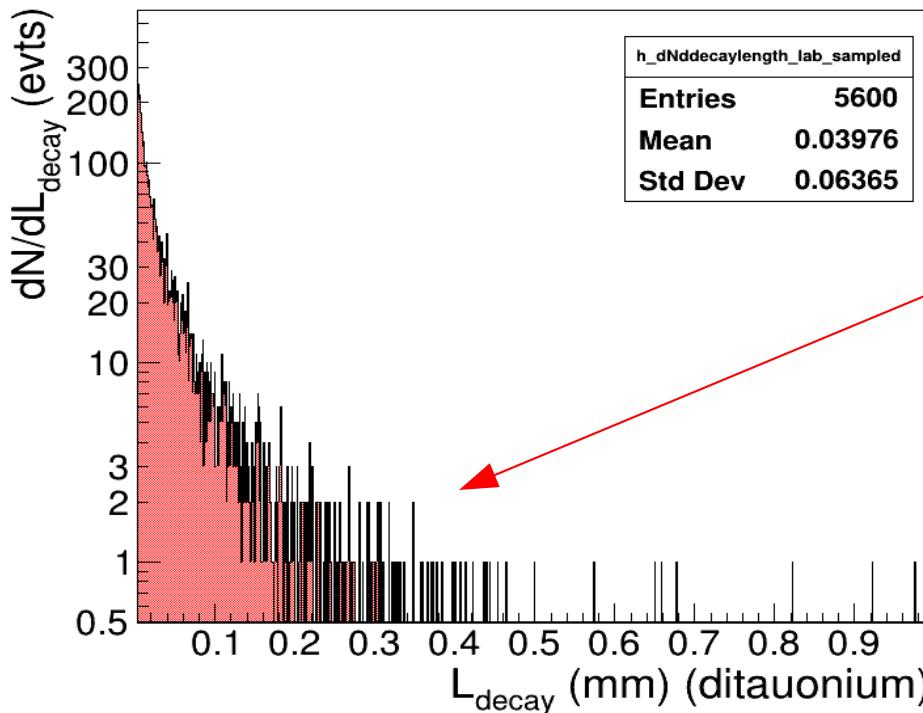
Significance (Belle-II)  $\approx 3\sigma$

Significance (FCC-ee)  $\approx 5\sigma$
- Pseudodata–null-hypothesis fit residuals:



# Para-ditauonium via $\gamma\gamma \rightarrow \tau_0$ Dalitz decays?

- Whereas background ccbar resonances decay directly from the IP, the para- $\tau_0$  has a lifetime of  $\tau \approx 28$  fs, i.e. a decay-length  $c\tau \approx 8 \mu\text{m}$ .



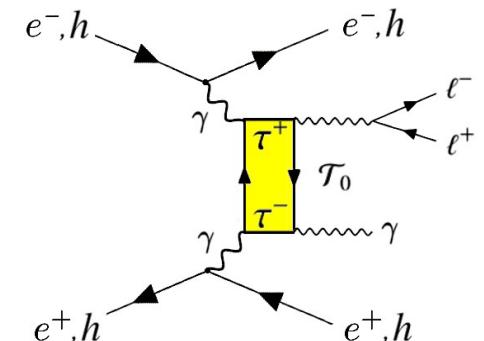
- For  $\beta\gamma \approx 3$ :  $\langle L_{\text{vtx}} \rangle \approx 25 \mu\text{m}$  tail of events up to  $\sim 1\text{-mm}$ . Any single event would be an **unambiguous  $\tau_0$  observation!**
- However, diphoton vertex pointing capabilities are much coarser: 1-cm range for LHC-type EM calos.  
Pico-second(!)  $\gamma$  ToF needed to separate  $<1\text{mm}$  distances ☹

- Displaced charged lepton vertices from Dalitz decays

$\tau_0 \rightarrow e^+e^-\gamma, \mu^+\mu^-\gamma$  with BR  $\sim 2.3\%$ ?

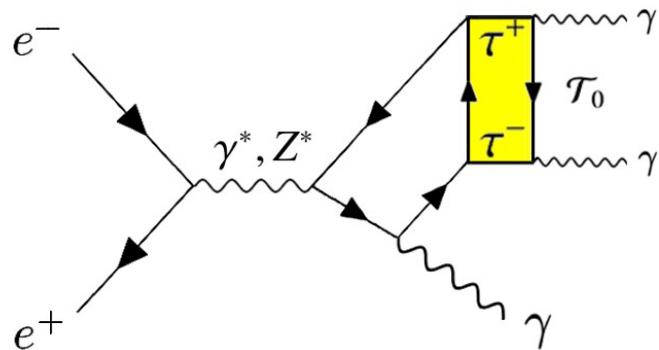
- $\mathcal{O}(150), \mathcal{O}(25)$  signal counts at FCC-ee/Belle-II...

But para- $\tau_0$  produced almost at rest ( $\beta\gamma \approx 0.06$ ) ☹



# Para-ditauonium via $e^+e^- \rightarrow \tau_0 + \gamma$ ?

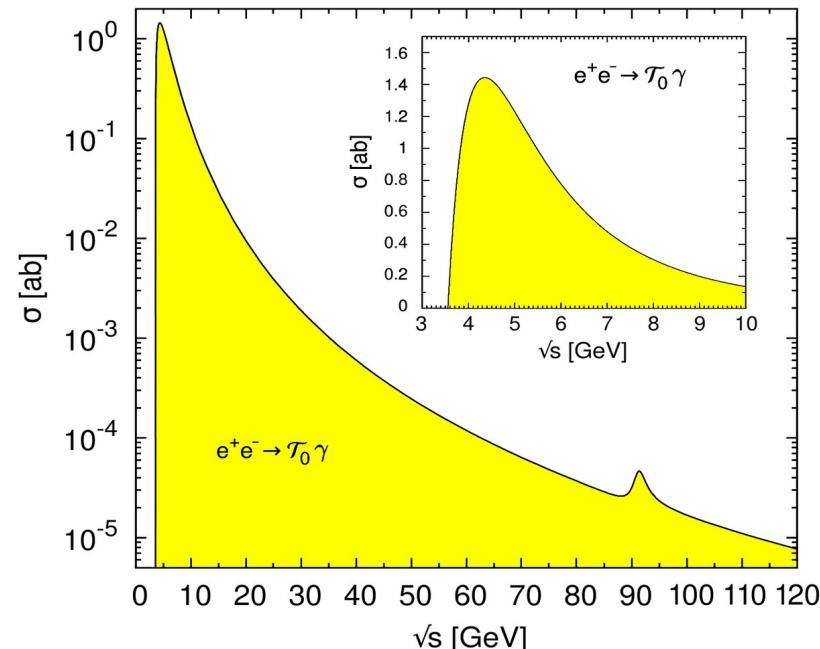
- s-channel production of para-ditauonium plus FSR in  $e^+e^-$  collisions:



$$\sigma(e^+e^- \rightarrow \tau_0 + \gamma) \approx \frac{2}{3} \frac{\pi \alpha^6}{n^3} \frac{m_\tau^2}{s^2} \left(1 - \frac{m_\tau^2}{s}\right)$$

- Tiny cross sections, in the (sub)attobarn range:

Colliding system, $\sqrt{s}$ , $\mathcal{L}_{\text{int}}$ , detector	$\sigma(\tau_0 + \gamma) \times \mathcal{B}_{\gamma\gamma}$	$N(\tau_0(\gamma\gamma) + \gamma)$
$e^+e^-$ at 4.3 GeV, 1 ab $^{-1}$ , STCF	1.1 ab	1
$e^+e^-$ at 7 GeV, 1 ab $^{-1}$ , STCF	0.37 ab	0.37
$e^+e^-$ at 3.78 GeV, 20 fb $^{-1}$ , BES III	0.69 ab	0.014
$e^+e^-$ at 10.6 GeV, 50 ab $^{-1}$ , Belle II	0.085 ab	4
$e^+e^-$ at 91.2 GeV, 50 ab $^{-1}$ , FCC-ee	$3.6 \cdot 10^{-5}$ ab	-

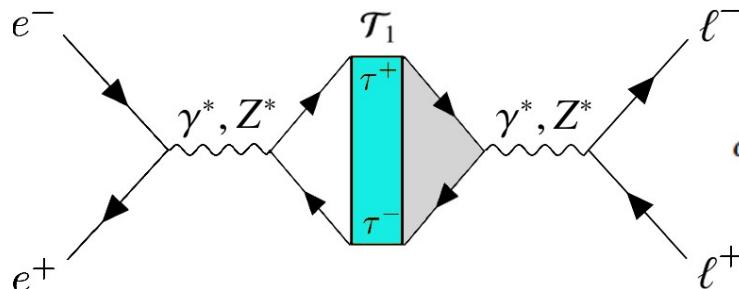


Negligible events expected, swamped by huge backgrounds:

$$\sigma(e^+e^- \rightarrow \gamma\gamma\gamma) \times \mathcal{L}_{\text{int}} = 15 \text{ pb} \times 50 \text{ ab}^{-1} = 7.5 \cdot 10^8 \text{ events at Belle II} \quad \text{:(}$$

# Ortho-ditauonium via $e^+e^- \rightarrow \tau_1$ fusion

- Resonant s-channel production of ortho-ditauonium in  $e^+e^-$  collisions:



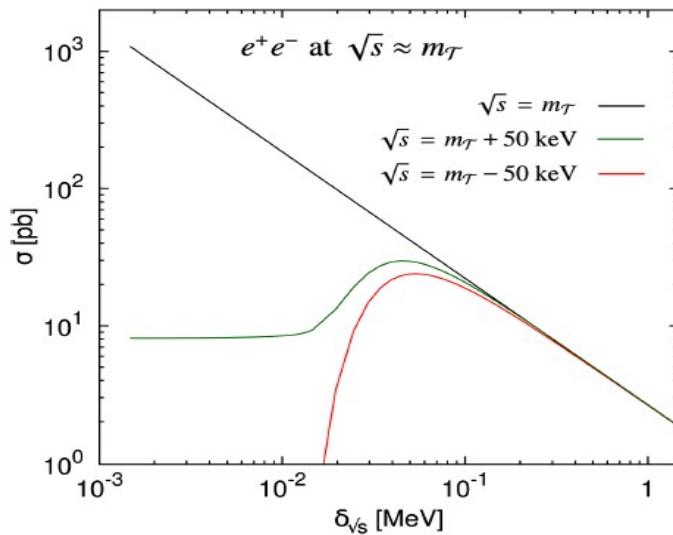
$$\sigma^{\text{ideal}}(e^+e^- \rightarrow \tau_1) = \frac{12\pi\Gamma_{\text{tot}}(\tau_1)\Gamma_{e^+e^-}(\tau_1)}{(s - m_\tau^2)^2 + \Gamma_{\text{tot}}^2(\tau_1)m_\tau^2} \stackrel{\sqrt{s}=m_\tau}{=} 236.6 \mu\text{b}$$

- Actual Breit-Wigner x-section reduced by  $>10^7$ , down to 2–20 pb, due to:
  - ISR & beam-energy spread  $\delta_{\sqrt{s}}$  (reduceable via monochromatization)
  - Accurate knowledge of  $m_\tau$  peak position required for  $\sqrt{s}$ .

$$\sigma^{\text{actual}}(e^+e^- \rightarrow \tau_1) = \frac{12\pi^2\Gamma_{e^+e^-}(\tau_1)}{m_\tau} \int_0^1 dx_1 \int_0^1 dx_2 f_{e^-/e^-}(x_1, s) f_{e^+/e^+}(x_2, s) V_2(\sqrt{x_1 x_2 s}; m_\tau, \Gamma_{\text{tot}}(\tau_1), \sqrt{x_1 x_2} \delta_{\sqrt{s}})$$

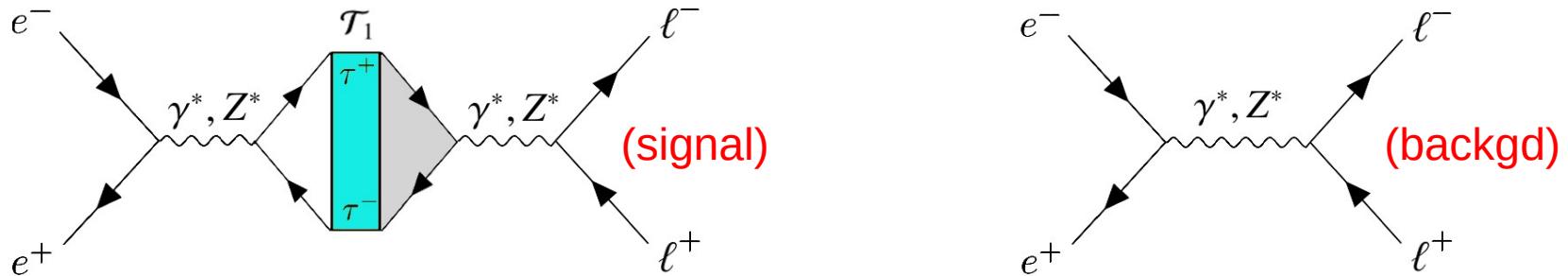
- Threshold-scan around  $\sqrt{s} = 2m_\tau$ :

Colliding system, $\sqrt{s}$ ( $\delta_{\sqrt{s}}$ spread), $\mathcal{L}_{\text{int}}$ , experiment	$\sigma$	$N$
$e^+e^-$ at 3.5538 GeV (1.47 MeV), $5.57 \text{ pb}^{-1}$ , BES III	1.9 pb	10.4
$e^+e^-$ at $\sqrt{s} \approx m_\tau$ (1.24 MeV), $140 \text{ pb}^{-1}$ , BES III	2.2 pb	310
$e^+e^-$ at $\sqrt{s} \approx m_\tau$ (1 MeV), $1 \text{ ab}^{-1}$ , STCF	2.6 pb	$2.6 \cdot 10^6$
$e^+e^-$ at $\sqrt{s} \approx m_\tau$ (100 keV), $0.1 \text{ ab}^{-1}$ , STCF	22 pb	$2.2 \cdot 10^6$



# Ortho-ditauonium observation via $e^+e^- \rightarrow \tau_1$ fusion

- Resonant s-channel production of ortho-ditauonium in  $e^+e^-$  collisions:



- Actual Breit-Wigner x-section reduced by  $>10^7$ , down to 2–20 pb, due to:

- ISR & beam-energy spread  $\delta_{\sqrt{s}}$  (reduceable via monochromatization)
- Accurate knowledge of  $m_\tau$  peak position required for  $\sqrt{s}$ .

- Threshold-scan around  $\sqrt{s} = 2m_\tau$  :  $\sigma^{\text{actual}}(e^+e^- \rightarrow \tau_1) = \frac{12\pi^2\Gamma_{e^+e^-}(\tau_1)}{m_{\tau_1}} \int_0^1 dx_1 \int_0^1 dx_2 f_{e^-/e^-}(x_1, s) f_{e^+/e^+}(x_2, s) V_2(\sqrt{x_1 x_2 s}; m_{\tau_1}, \Gamma_{\text{tot}}(\tau_1), \sqrt{x_1 x_2} \delta_{\sqrt{s}})$

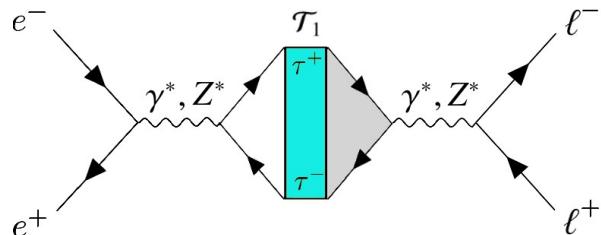
Colliding system, $\sqrt{s}$ ( $\delta_{\sqrt{s}}$ spread), $\mathcal{L}_{\text{int}}$ , experiment	$\sigma$			$N$			$S/\sqrt{B}$
	$\tau_1$	$\tau^+\tau^-$	$\mu^+\mu^-$	$\tau_1$	$\tau_1 \rightarrow \mu^+\mu^-$	$\mu^+\mu^-$	
$e^+e^-$ at 3.5538 GeV (1.47 MeV), 5.57 pb $^{-1}$ , BES III	1.9 pb	117 pb	6.88 nb	10.4	2.1	38 300	0.01 $\sigma$
$e^+e^-$ at $\sqrt{s} \approx m_{\tau_1}$ (1.24 MeV), 140 pb $^{-1}$ , BES III	2.2 pb	103 pb	6.88 nb	310	63	$9.63 \cdot 10^5$	0.06 $\sigma$
$e^+e^-$ at $\sqrt{s} \approx m_{\tau_1}$ (1 MeV), 1 ab $^{-1}$ , STCF	2.6 pb	95 pb	6.88 nb	$2.6 \cdot 10^6$	$5.3 \cdot 10^5$	$6.88 \cdot 10^9$	6.4 $\sigma$
$e^+e^-$ at $\sqrt{s} \approx m_{\tau_1}$ (100 keV), 0.1 ab $^{-1}$ , STCF	22 pb	46 pb	6.88 nb	$2.2 \cdot 10^6$	$4.5 \cdot 10^5$	$6.88 \cdot 10^8$	17 $\sigma$

- Ortho- $\tau_1$  observable at STCF (6.4 $\sigma$ ) on top of  $\mu^+\mu^-$  continuum in default run (1 ab $^{-1}$ )

- Note: Ditauonium contributes 2% of the di-tau x-section at  $\sqrt{s} = 2m_\tau$  at STCF

# Ultraprecise tau mass via $e^+e^- \rightarrow \tau_1 \rightarrow \mu^+\mu^-$

- STCF with 0.1-MeV monochromatization & 4 mass points runs (0.1  $\text{ab}^{-1}$  each) can determine very accurately peak excess of  $\mu^+\mu^-$  events corresponding to the **ortho- $\tau_1$  resonant mass point** (provided true  $m_\tau$  is known to within  $\pm 50$  keV):

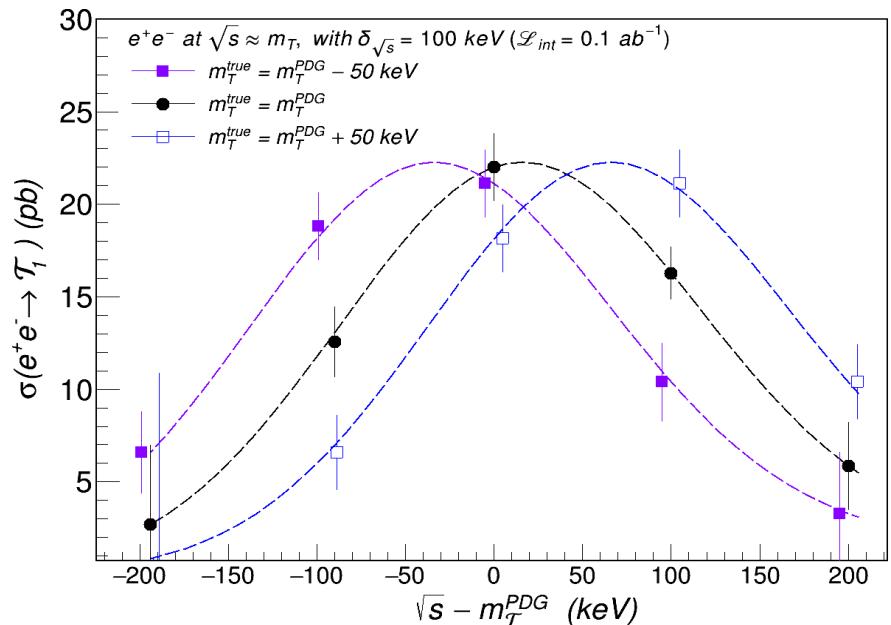


Colliding system, $\sqrt{s}$ ( $\delta_{\sqrt{s}}$ spread), $\mathcal{L}_{int}$ , experiment	$\sigma$			$N$		$S/\sqrt{B}$
	$\tau_1$	$\tau^+\tau^-$	$\mu^+\mu^-$	$\tau_1 \rightarrow \mu^+\mu^-$	$\mu^+\mu^-$	
$e^+e^-$ at $\sqrt{s} \approx m_\tau$ (1 MeV), $1 \text{ ab}^{-1}$ , STCF	2.6 pb	95 pb	6.88 nb	$5.3 \cdot 10^5$	$6.88 \cdot 10^9$	$6.4\sigma$
$e^+e^-$ at $\sqrt{s} \approx m_\tau$ (100 keV), $0.1 \text{ ab}^{-1}$ , STCF	22 pb	46 pb	6.88 nb	$4.5 \cdot 10^5$	$6.88 \cdot 10^8$	$17\sigma$

- The accuracy of the  $m(\tau_1)$  position depends only on the accuracy of the beam energy calibration:

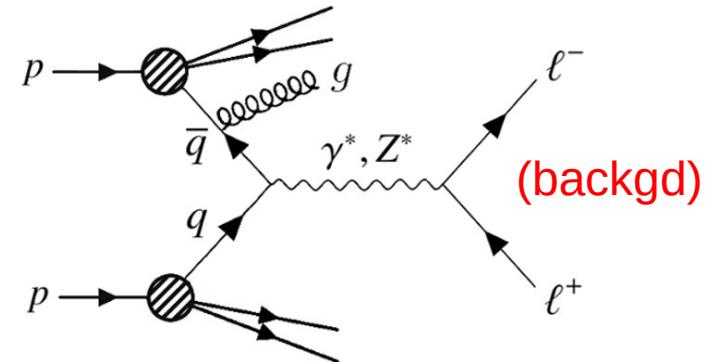
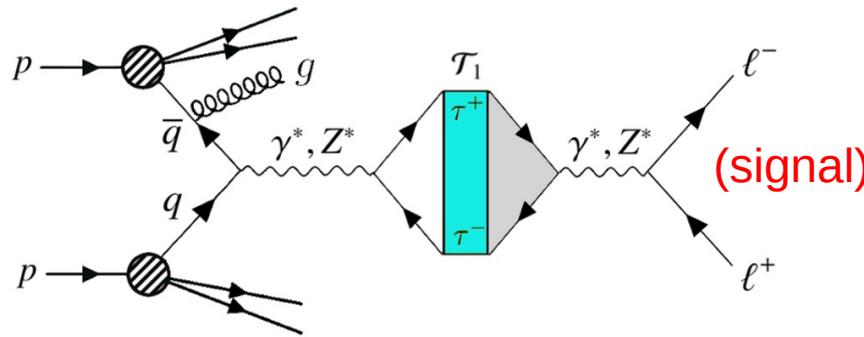
With BES-III (BEMS method):  $\Delta_{\sqrt{s}} = 10^{-5}$   
 $\Rightarrow \mathcal{O}(25 \text{ keV})$  tau mass precision

- Impact of ultraprecise  $m_\tau$ :
  - Improved LFU tests ( $\propto m_{e,\mu}^5/m_\tau^5$ )
  - CKM  $|V_{ij}|$  elements from  $\tau$  decays
  - Any other SM checks that parametrically depend on ratios of  $e, \mu, \tau$  masses



# Ortho-ditauonium via DY+j production at the LHC

- Drell-Yan production of **ortho-ditauonium + jet** in pp colls. at 14 TeV:

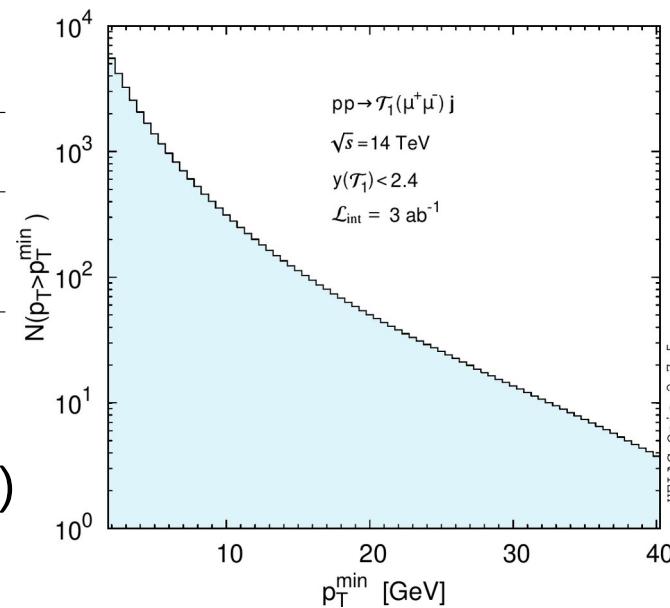


Back-to-back jet required to **boost ortho- $\tau_1$**  decay (**displaced secondary dimuon vertex**) & eliminate DY backgds. Only combinatorial heavy-Q dimuon sources left.

- Cross sections at ATLAS/CMS, ALICE/LHCb:

Colliding system, $\sqrt{s}$ , $\mathcal{L}_{\text{int}}$ , detector	$\sigma_{\text{NLO}}$		$N(\mathcal{T}_1 + j)$		with $L_{xy} > 30$ (100) $\mu\text{m}$	
	$\mathcal{T}_1 + X$	$\mathcal{T}_1 + j$	$\mathcal{T}_1 \rightarrow e^+ e^-$	$\mathcal{T}_1 \rightarrow \mu^+ \mu^-$	$\mathcal{T}_1 \rightarrow \ell^+ \ell^-$	$\mathcal{T}_1 \rightarrow \mu^+ \mu^-$
p-p at 14 TeV, $3 \text{ ab}^{-1}$ , ATLAS/CMS	$42^{+11}_{-19} \text{ fb}$	$18 \pm 9 \text{ fb}$	1100	1100	130 (10)	130 (10)
p-p at 14 TeV, $300 \text{ fb}^{-1}$ , LHCb	$42^{+11}_{-19} \text{ fb}$	$18 \pm 9 \text{ fb}$	110	110	5 (-)	5 (-)
p-p at 114.6 GeV, $10 \text{ fb}^{-1}$ , ALICE/LHCb	$2.2^{+0.3}_{-0.4} \text{ fb}$	$1 \pm 0.5 \text{ fb}$	<10	<10	-	-

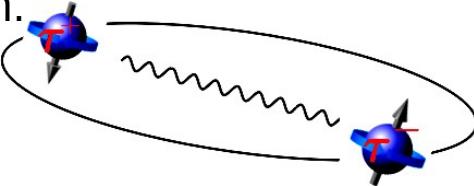
- About **130 (10)** displaced dimuon events with  $L_{xy} > 30$  (100)  $\mu\text{m}$  expected in ATLAS/CMS ( $3 \text{ ab}^{-1}$ )  
Observation feasible (even with less int. lumi)!



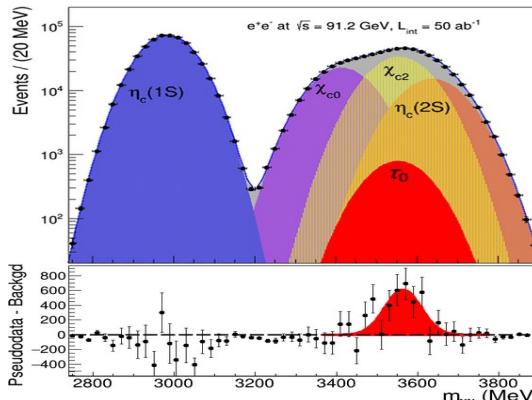
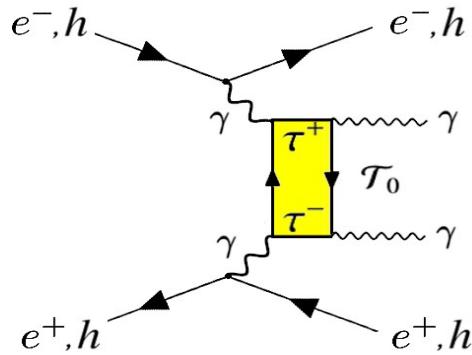
# Summary (I)

■ First-ever comprehensive study of ditaonium production/detection in the lab:

- Unobserved. Heaviest & most compact leptonic “atomic” system.
- Tests of bound QED & CPT symmetries at high-mass (BSM?).
- Ultraprecise  $\tau$  mass extraction possible via  $e^+e^- \rightarrow \tau_1 \rightarrow \mu^+\mu^-$

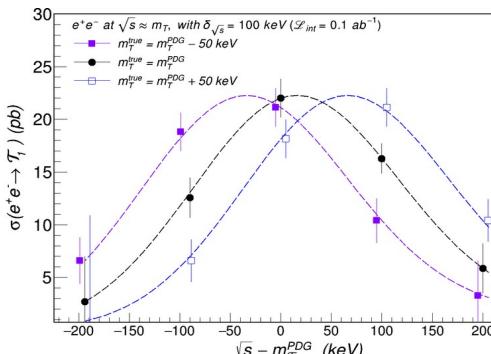
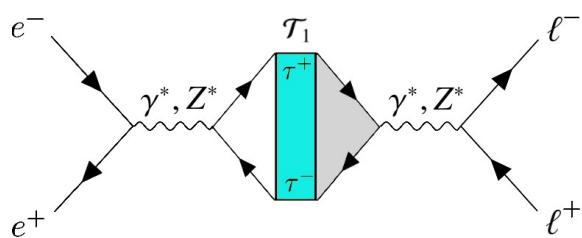


■ Para-ditaonium: Observable via  $\gamma\gamma$  fusion at high-lumi  $e^+e^-$  colliders:



- Requires accurate in-situ measure of overlapping ccbar resonances.
- Stat. significance (multi-Gaussian  $m_{\gamma\gamma}$  fit):  $S(\text{Belle-II/FCC-ee}) \approx 3\sigma, 5\sigma$

■ Ortho-ditaonium: Observable as s-channel resonance at STCF  $e^+e^-$  at  $\sqrt{s} = 2m_\tau$ :

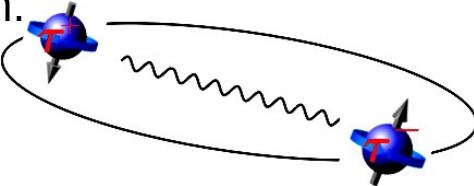


- Dimuon excess (>6 $\sigma$ ) in a nominal STFC year ( $1 \text{ ab}^{-1}$ )
- With 0.1-MeV beam monochrom. tau mass with 25-keV (or better) precision (beam calibration).

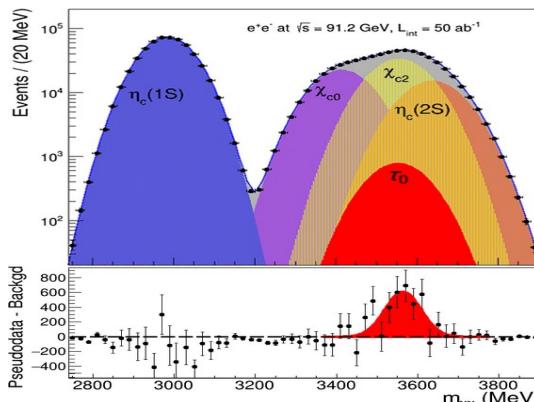
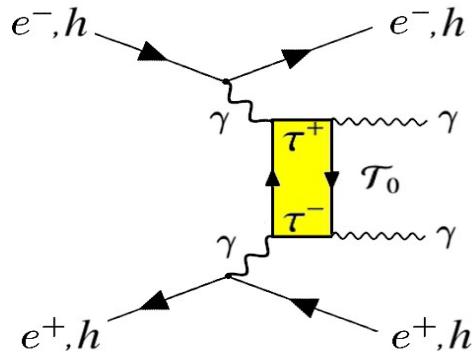
# Summary (II)

## ■ First-ever comprehensive study of ditauonium production/detection in the lab:

- Unobserved. Heaviest & most compact leptonic “atomic” system.
- Tests of bound QED & CPT symmetries at high-mass (BSM?).
- Ultraprecise  $\tau$  mass extraction possible via  $e^+e^- \rightarrow \tau_1 \rightarrow \mu^+\mu^-$

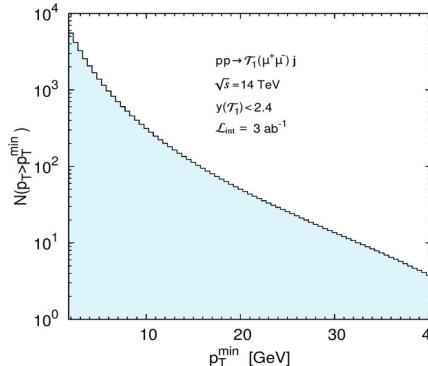
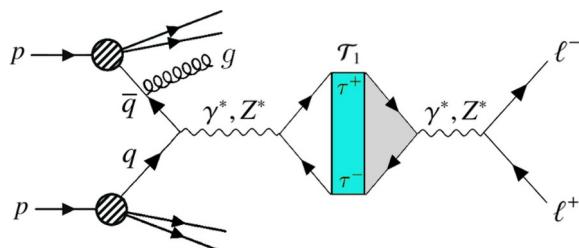


## ■ Para-ditauonium: Observable via $\gamma\gamma$ fusion at high-lumi $e^+e^-$ colliders:



- Requires accurate in-situ measure of overlapping ccbar resonances.
- Stat. significance (multi-Gaussian  $m_{\gamma\gamma}$  fit):  $S(\text{Belle-II/FCC-ee}) \approx 3\sigma, 5\sigma$

## ■ Ortho-ditauonium+jet: Observable in DY production in p-p collisions at the LHC:



- $N_{\text{evts}} = 130$  (10) displaced dimuon with  $L_{xy} > 30$  (100)  $\mu\text{m}$  at ATLAS/CMS ( $3 \text{ ab}^{-1}$ ).
- Observation feasible (even with less int. lumi)!

# Backup slides

# Ditauonium partial widths & decays

$\mathcal{T}$ state	$m_X$ (MeV)	$J^{PC}$	$\Gamma_{\text{tot}}$ (eV)	Lifetime (fs)	Decay mode	$\Gamma_X$ (eV)	$\mathcal{B}_X$
	$1^1S_0$		$3553.696 \pm 0.240$	$0^{-+}$	0.02384	27.60	
					$\gamma\gamma$	0.018533	77.72%
					$\gamma e^+ e^-$	$4.28 \cdot 10^{-4}$	1.79%
					$\gamma\mu^+\mu^-$	$1.24 \cdot 10^{-4}$	0.52%
					$\gamma q\bar{q}$	$2.20 \cdot 10^{-4}$	0.92%
					$e^+ e^- e^+ e^-$	$2.32 \cdot 10^{-6}$	0.0094%
					$e^+ e^- \mu^+ \mu^-$	$1.38 \cdot 10^{-6}$	0.0058%
					$e^+ e^- q\bar{q}$	$1.20 \cdot 10^{-6}$	0.0050%
					$\mu^+ \mu^- \mu^+ \mu^-$	$1.65 \cdot 10^{-7}$	0.00069%
					$\mu^+ \mu^- q\bar{q}$	$2.72 \cdot 10^{-7}$	0.0011%
					$q\bar{q}q'\bar{q}'$	$8.23 \cdot 10^{-8}$	0.00035%
					$(2)\tau \rightarrow X$	0.004535	19.02%
$1^3S_1$	$3553.696 \pm 0.240$	$1^{+-}$	0.03159	20.83			
					$e^+ e^-(\gamma)$	0.006436	20.37%
					o $e^+ e^-$	$2.95 \cdot 10^{-3}$	9.33%
					o $e^+ e^- \gamma$	$3.49 \cdot 10^{-3}$	11.04%
					$\mu^+ \mu^-(\gamma)$	0.006436	20.37%
					o $\mu^+ \mu^-$	$6.10 \cdot 10^{-3}$	19.30%
					o $\mu^+ \mu^- \gamma$	$3.38 \cdot 10^{-4}$	1.07%
					$q\bar{q}(\gamma)$	0.01416	44.82%
					$\gamma\gamma\gamma$	$1.62 \cdot 10^{-5}$	0.051%
					$e^+ e^- e^+ e^-$	$5.55 \cdot 10^{-6}$	0.0176%
					$e^+ e^- \mu^+ \mu^-$	$4.21 \cdot 10^{-6}$	0.0133%
					$e^+ e^- q\bar{q}$	$1.85 \cdot 10^{-6}$	0.0058%
					$\mu^+ \mu^- \mu^+ \mu^-$	$1.23 \cdot 10^{-7}$	$O(10^{-6})$
					$\mu^+ \mu^- q\bar{q}$	$7.36 \cdot 10^{-8}$	$O(10^{-6})$
					$q\bar{q}q'\bar{q}'$	$9.73 \cdot 10^{-9}$	$O(10^{-7})$
					$\nu_\tau \bar{\nu}_\tau$	$1.32 \cdot 10^{-8}$	$O(10^{-7})$
					$\nu_e \bar{\nu}_e$	$4.30 \cdot 10^{-11}$	$O(10^{-9})$
					$\nu_\mu \bar{\nu}_\mu$	$4.30 \cdot 10^{-11}$	$O(10^{-9})$
					$(2)\tau \rightarrow X$	0.004535	14.35%

# Para-ditauonium production via $\gamma\gamma$ collisions

- Cross sections for signal & backgrounds computed in the Weizsäcker-Williams approx. (EPA) for  $\gamma\gamma$  collisions (implemented in HelacOnia2.6/gamma-UPC):

$$\sigma(a b \rightarrow a b + X) = 4\pi^2(2J+1) \frac{\Gamma_{\gamma\gamma}(X)}{m_X^2} \left. \frac{d\mathcal{L}_{\gamma\gamma}^{(ab)}}{dW_{\gamma\gamma}} \right|_{W_{\gamma\gamma}=m_X}$$

- Diphoton charmonium resonances within  $m_{\gamma\gamma} \approx 2.9\text{--}3.7 \text{ GeV}$ :

Resonance	$J^{PC}$	$m_X$ (MeV)	$\Gamma_{\text{tot}}$ (MeV)	$\Gamma_{\gamma\gamma}$ (MeV)	$\mathcal{B}_{\gamma\gamma}$
$\mathcal{T}_0$	$0^{-+}$	$3553.696 \pm 0.240$	$2.28 \cdot 10^{-8}$	$1.83 \cdot 10^{-8}$	$\sim 80\%$
$\eta_c(1S)$	$0^{-+}$	$2983.9 \pm 0.5$	$32.0 \pm 0.7$	$(5.06 \pm 0.34) \cdot 10^{-3}$	$(0.0158 \pm 0.0011)\%$
$\eta_c(2S)$	$0^{-+}$	$3637.5 \pm 1.1$	$11.3 \pm 3.1$	$(2.15 \pm 1.47) \cdot 10^{-3}$	$(0.019 \pm 0.013)\%$
$\chi_{c0}$	$0^{++}$	$3414.71 \pm 0.30$	$10.8 \pm 0.6$	$(2.203 \pm 0.097) \cdot 10^{-3}$	$(0.0204 \pm 0.0009)\%$
$\chi_{c2}$	$2^{++}$	$3556.17 \pm 0.07$	$1.97 \pm 0.09$	$(5.614 \pm 0.197) \cdot 10^{-4}$	$(0.0285 \pm 0.0010)\%$

→ Charmonia resonances have  $\mathcal{O}(\text{keV})$  diphoton widths:  $\mathcal{O}(10^5)$  larger than para- $\tau_0$ .  
But, the diphoton BR is  $\mathcal{O}(10^4)$  larger for para- $\tau_0$  than for c-cbar states.

# $\gamma\gamma$ collision x-sections (signal & backgds)

- Cross sections for signal & backgrounds computed in the Weizsäcker-Williams approximation (**EPA**) for  $\gamma\gamma$  collisions via gamma-UPC: 2207.03012 [hep-ph].
- $\sigma(\text{LbL})$  computed with **MG5@NLO** (virtual box) with same photon fluxes.
- Results for  $e^+e^-$  and ultraperipheral p-p, p-A & A-A collisions:

Colliding system, c.m. energy, $\mathcal{L}_{\text{int}}$ , exp.	$\sigma \times \mathcal{B}_{\gamma\gamma}$					$N \times \mathcal{B}_{\gamma\gamma}$		
	$\eta_c(1S)$	$\eta_c(2S)$	$\chi_{c,0}(1P)$	$\chi_{c,2}(1P)$	LbL	$\tau_0$	$\tau_0$	$\chi_{c,2}(1P)$
$e^+e^-$ at 3.78 GeV, $20 \text{ fb}^{-1}$ , BES III	120 fb	3.6 ab	15 ab	13 ab	30 ab	0.25 ab	—	—
	1.7 fb	0.35 fb	0.52 fb	0.77 fb	1.7 fb	0.015 fb	750	38 500
	11 fb	2.8 fb	3.9 fb	6.0 fb	12 fb	0.11 fb	5 600	$3 \cdot 10^5$
p-p at 14 TeV, $300 \text{ fb}^{-1}$ , LHC	7.9 fb	2.0 fb	2.8 fb	4.3 fb	6.3 fb	0.08 fb	24	1290
p-Pb at 8.8 TeV, $0.6 \text{ pb}^{-1}$ , LHC	25 pb	6.3 pb	8.7 pb	13 pb	21 pb	0.25 pb	0.15	8
Pb-Pb at 5.5 TeV, $2 \text{ nb}^{-1}$ , LHC	61 nb	15 nb	21 nb	31 nb	62 nb	0.59 nb	1.2	62

(~10% uncertainties, today)

- Relative production x-sections:  $\eta_c(1S):\chi_{c2}(1P):\chi_{c0}(1P):\eta_c(2S):\tau_0 \approx 100:50:30:25:1$  driven by their different  $\Gamma^2(\gamma\gamma)/(\Gamma(\text{tot}) \cdot m^2_X)$  ratios.
- Cross sections increase with  $\sqrt{s}$  and  $Z^4$ :  
Largest x-sections (0.6 nb) in PbPb UPC (but handful of evts expected at LHCb)  
Largest yields: 750, 5600 counts at Belle-II, FCC-ee thanks to  $\mathcal{L}_{\text{int}} = 50 \text{ ab}^{-1}$ .

# Para-ditauonium signal extraction

- 1-million events generated for signal & backgrounds. Run MVA (BDT) with 12 different single- $\gamma$  and  $\gamma$ -pair kinematic variables for signal/backgds separation:
  - (i) Strong discrimination power (factor of  $\sim 20$ ) of LbL continuum from signal.
  - (ii) No discrimination achieved for overlapping charmonia (decay  $\gamma$  angular modulation of tensor  $\chi_{c2}$  different than scalar  $\tau_0$  signal, but  $\times 50$  suppressed yields)
- Signal extracted through multi-Gaussian  $m_{\gamma\gamma}$  fit, by considering:
  - $\eta_c(1S)$ : No overlap w/ signal (“std.candle”): 0.5M clean evts to fully control  $E_\gamma$  scale&res. plus exp. & theory uncertainties.
  - $\chi_{c0}$ ,  $\eta_c(2S)$ : Partial overlap with signal. Exploit  $\sim 100M \gamma\gamma \rightarrow \chi_{c0}, \eta_c \rightarrow X$  decays with  $\times 50$  larger BRs (e.g.  $X=3$ - and 4-mesons) to fully remove their contamination.
  - $\chi_{c2}$ : Full overlap with signal! Exploit alternative  $\gamma\gamma \rightarrow \chi_{c2} \rightarrow X$  decays (e.g. 11M evts. for  $X=4\pi$ ) to determine its lineshape to within  $\mathcal{O}(0.2\%)$ .

