

MSHT fit: Closure Test and Comparison of Approaches

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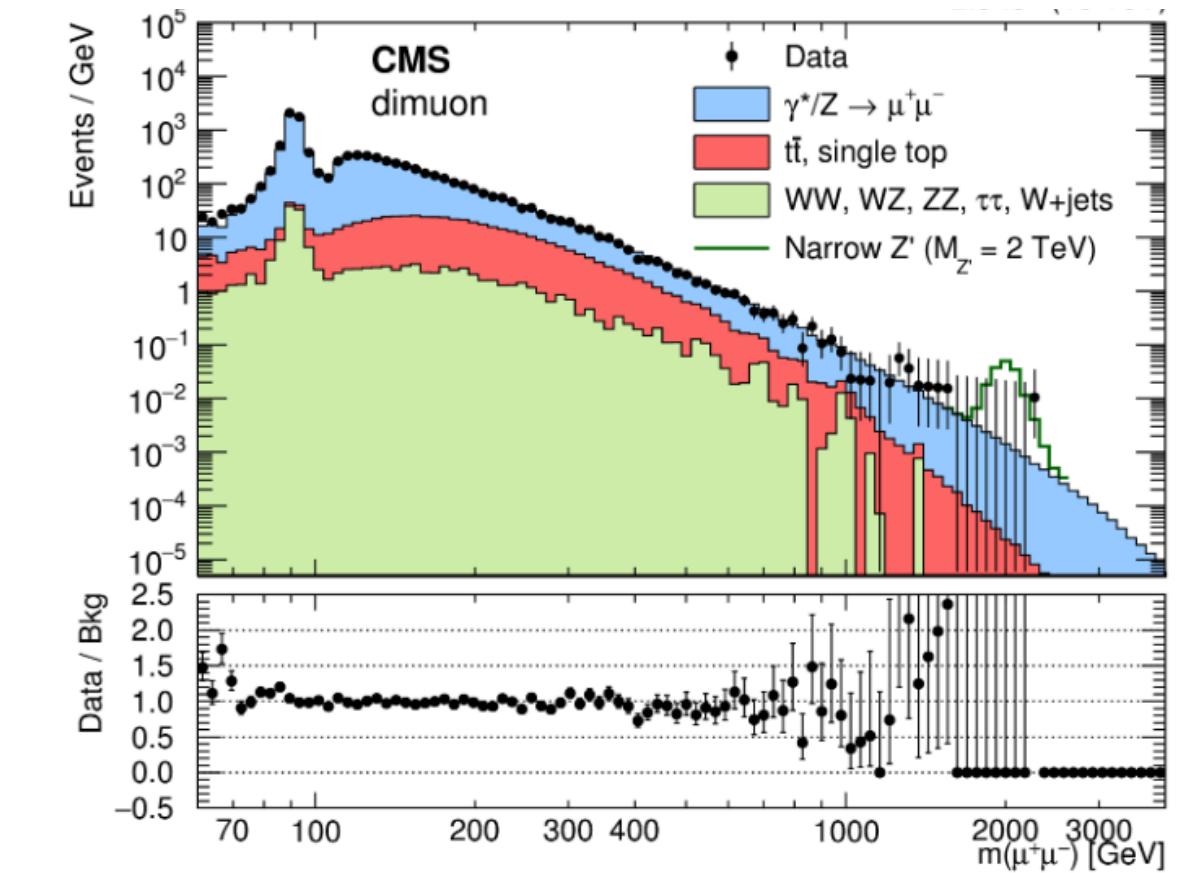
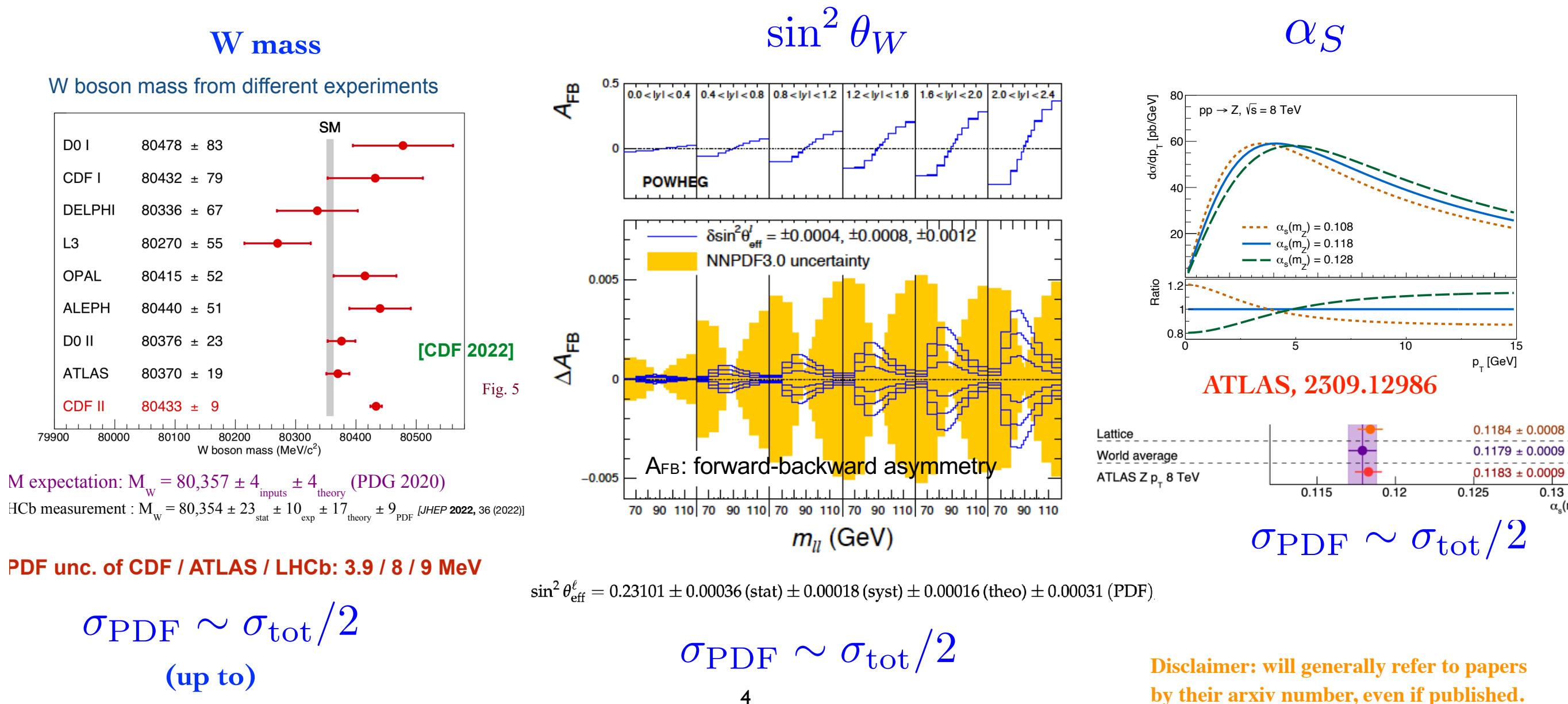
DIS 2024, Grenoble, 10 April 2024

With Robert Thorne and Tom Cridge



Introduction

- Parton distribution functions (PDFs): a key ingredient in hadron collider physics.
- Knowledge of **PDFs** and their **uncertainties** a limiting factor in LHC precision and BSM searches.



- The LHC is a BSM search machine. Often need PDFs here.
- High mass = high x , where PDFs are less well known. Key when looking for small/smooth deviations.

See my plenary talk on Monday!

- Accuracy and precision in PDF determinations essential.

Global PDF fits: parameterisation

- Two distinct methodologies on the market to parameterising PDFs: **Neural Nets** (NNPDF) or **Explicit Parameterisation** (CT, MSHT).

♦ MSHT: **52** free parameters in terms of Chebyshev polynomials.

$$u_V(x, Q_0^2) = A_u(1-x)^{\eta_u} x^{\delta_u} \left(1 + \sum_{i=1}^6 a_{u,i} T_i(y(x)) \right)$$

$$d_V(x, Q_0^2) = A_d(1-x)^{\eta_d} x^{\delta_d} \left(1 + \sum_{i=1}^6 a_{d,i} T_i(y(x)) \right)$$

$$S(x, Q_0^2) = A_S(1-x)^{\eta_s} x^{\delta_s} \left(1 + \sum_{i=1}^6 a_{S,i} T_i(y(x)) \right)$$

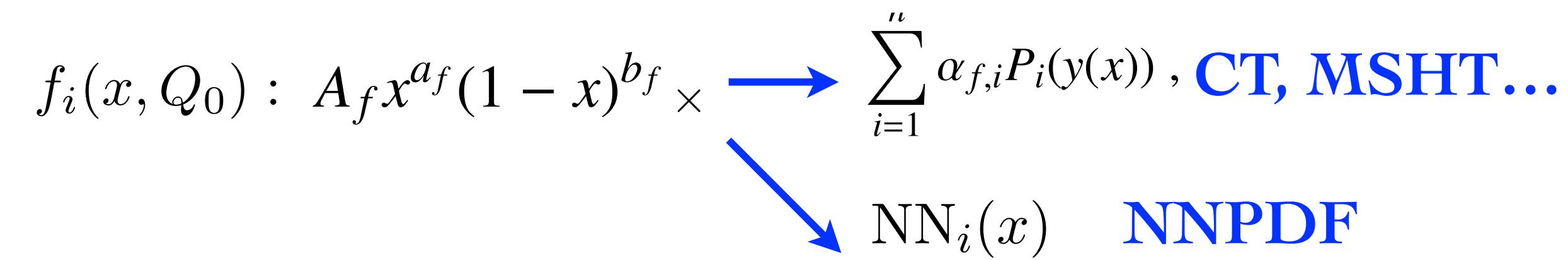
$$s_+(x, Q_0^2) = A_{s+}(1-x)^{\eta_{s+}} x^{\delta_s} \left(1 + \sum_{i=1}^6 a_{s+,i} T_i(y(x)) \right)$$

$$g(x, Q_0^2) = A_g(1-x)^{\eta_g} x^{\delta_g} \left(1 + \sum_{i=1}^4 a_{g,i} T_i(y(x)) \right) + A_{g-}(1-x)^{\eta_{g-}} x^{\delta_{g-}}$$

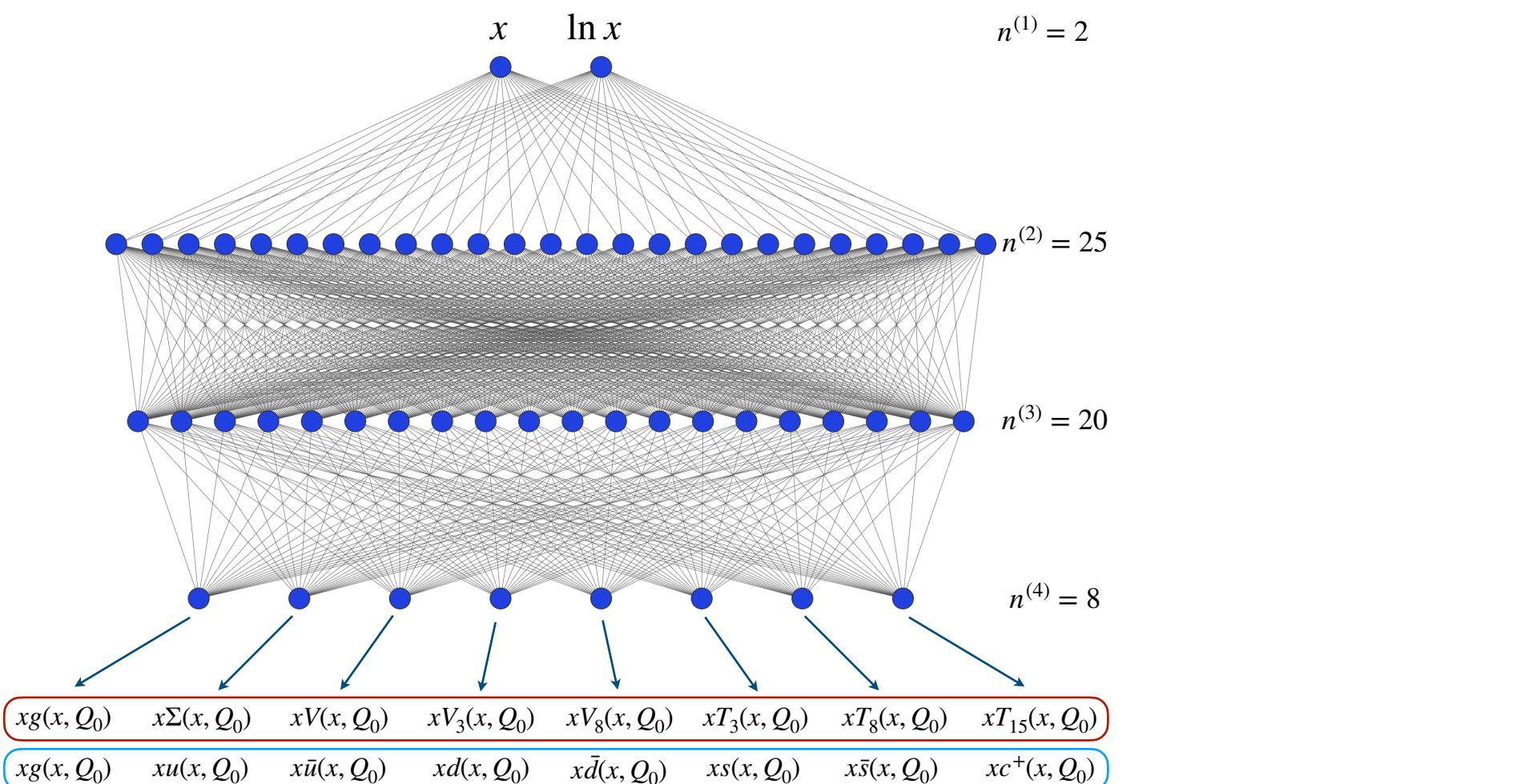
$$s_-(x, Q_0^2) = A_{s-}(1-x)^{\eta_{s-}} (1-x/x_0) x^{\delta_{s-}}$$

$$(\bar{d}/\bar{u})(x, Q_0^2) = A_\rho(1-x)^{\eta_\rho} \left(1 + \sum_{i=1}^6 a_{\rho,i} T_i(y(x)) \right)$$

♦ Less flexible in general - need to be sure flexible enough! Allows direct handle on uncertainties in Hessian framework.



★ NNPDF: **763** free parameter Neural Net.



★ Increased flexibility, but needs robust optimisation + stopping (avoid over and under fitting).

Fixed Parameterisation PDFs

- Fixed parameterisation approach:

$$\chi^2_{\text{global}} \sim \frac{(D_{\text{ata}} - T_{\text{theory}})^2}{\sigma^2}$$

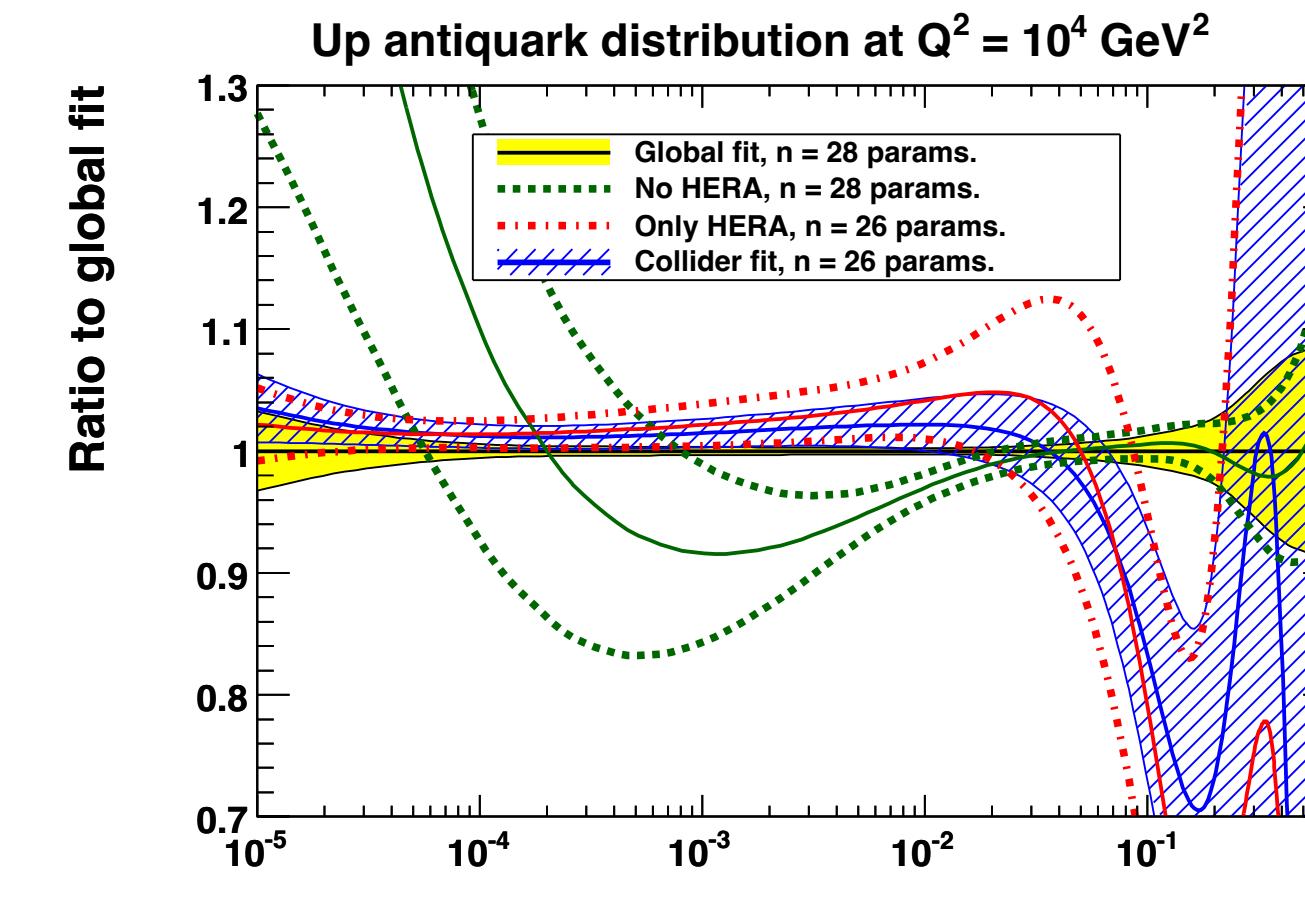
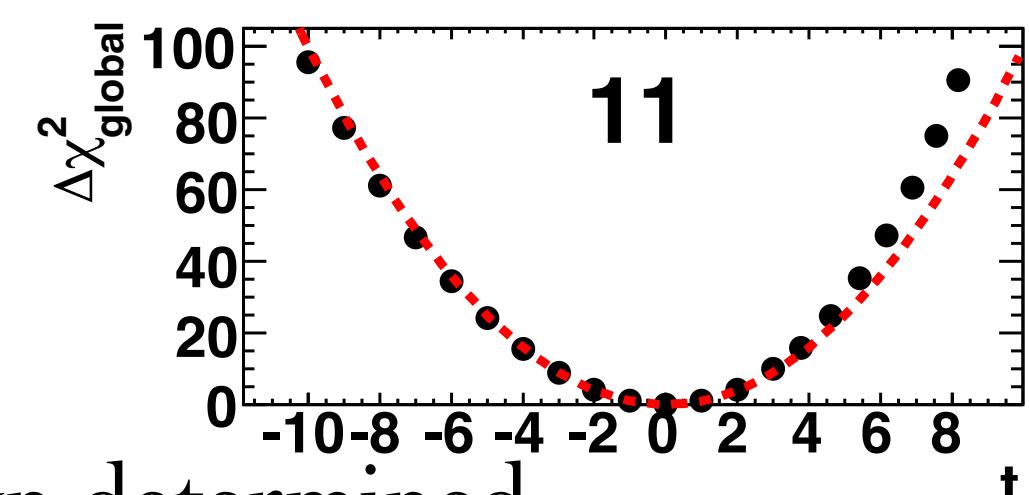
- ♦ Find global minimum of χ^2 and evaluate eigenvectors of Hessian matrix at this point.
- ♦ Parameter shifts corresponding to given $\Delta\chi^2$ criteria given in terms of these

$$a_i(S_k^\pm) = a_i^0 \pm t e_{ik}, \quad \text{with } t \text{ adjusted to give desired } T = \Delta\chi^2_{\text{global}}$$

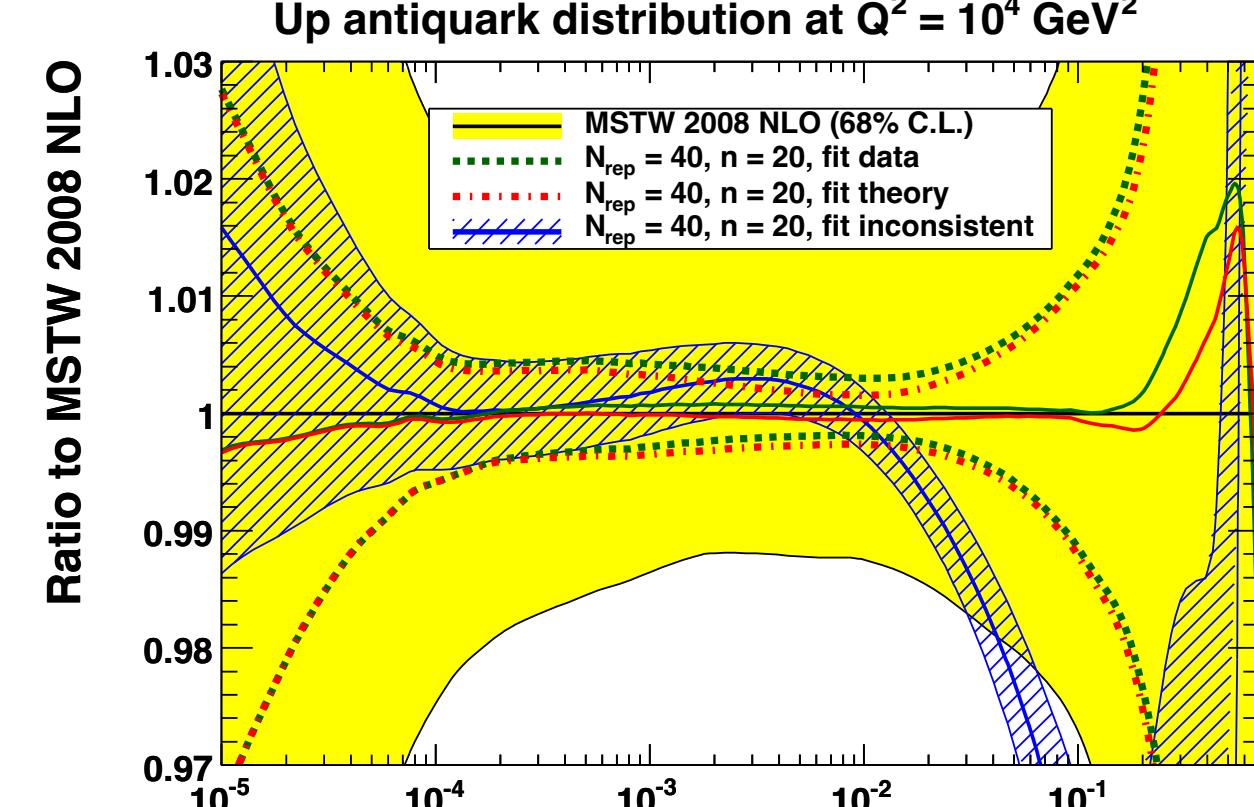
- ♦ $T = 1$: ‘textbook’ criteria for case where data matches theory perfectly up to known determined (Gaussian) errors, across entire global dataset ($N_{\text{pts}} \sim 4000 - 5000$).
- ♦ Expect to not be sufficient: fit quality poor by textbook standard, dataset tensions, theory incomplete...
- ♦ Backed up by evidence of e.g. fits to restricted datasets, or pseudodata with inconsistencies injected in.

→ Motivates an enlarged $T > 1$, either fixed or ‘dynamic’.

$$H_{ij} = \frac{1}{2} \frac{\partial^2 \chi^2_{\text{global}}}{\partial a_i \partial a_j} \Big|_{\min}$$



$$\frac{\chi^2}{N_{\text{pts}}} \gg 1 + \sigma(N_{\text{pts}}) \sim 1.02$$



G. Watt and R. Thorne, arXiv:1205.4024

See also, J. Pumplin, arXiv:0909.0268

Neural Network PDFs

- Neural network approach:
 - ♦ Generate set of MC ‘replicas’ by shifting data by errors.

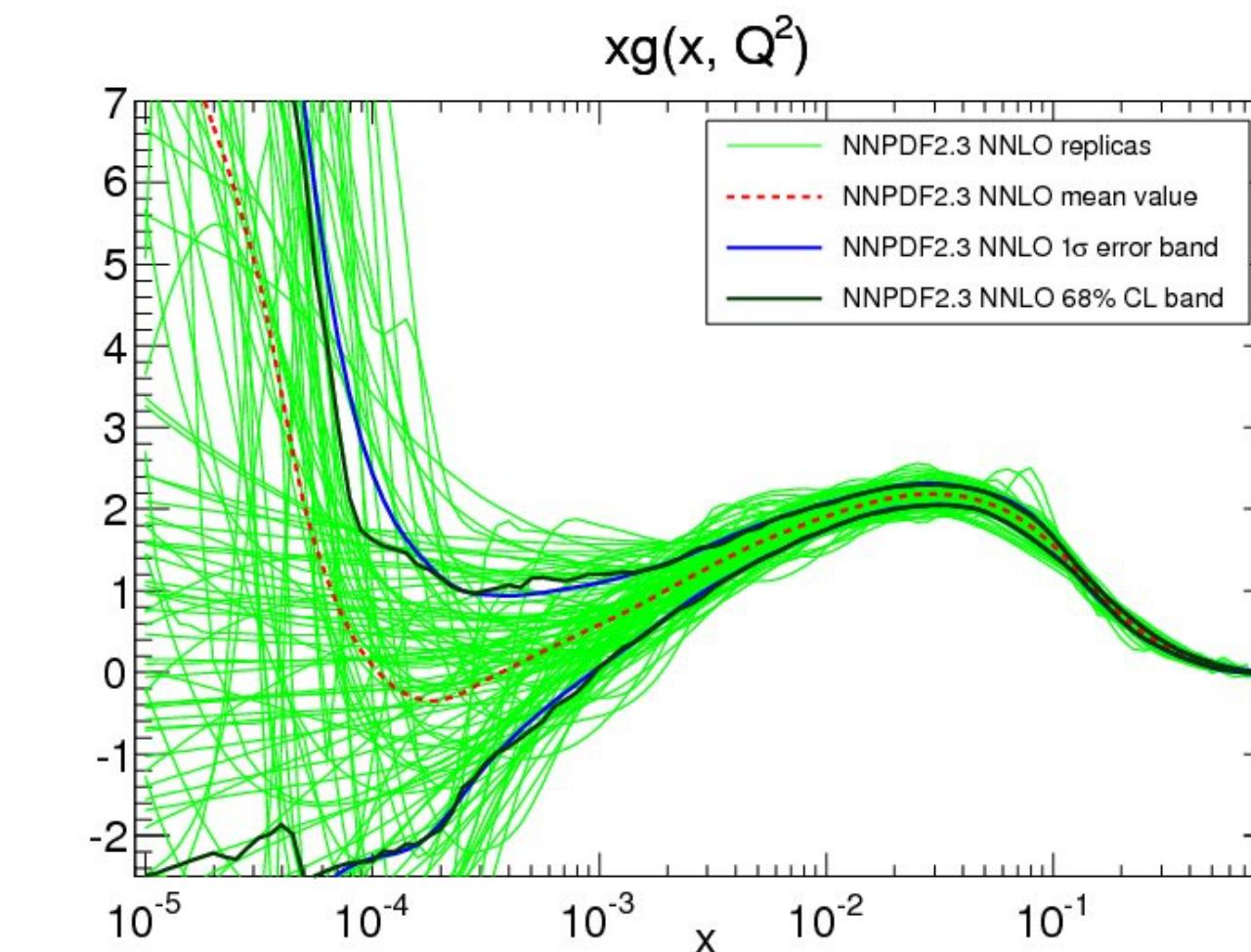
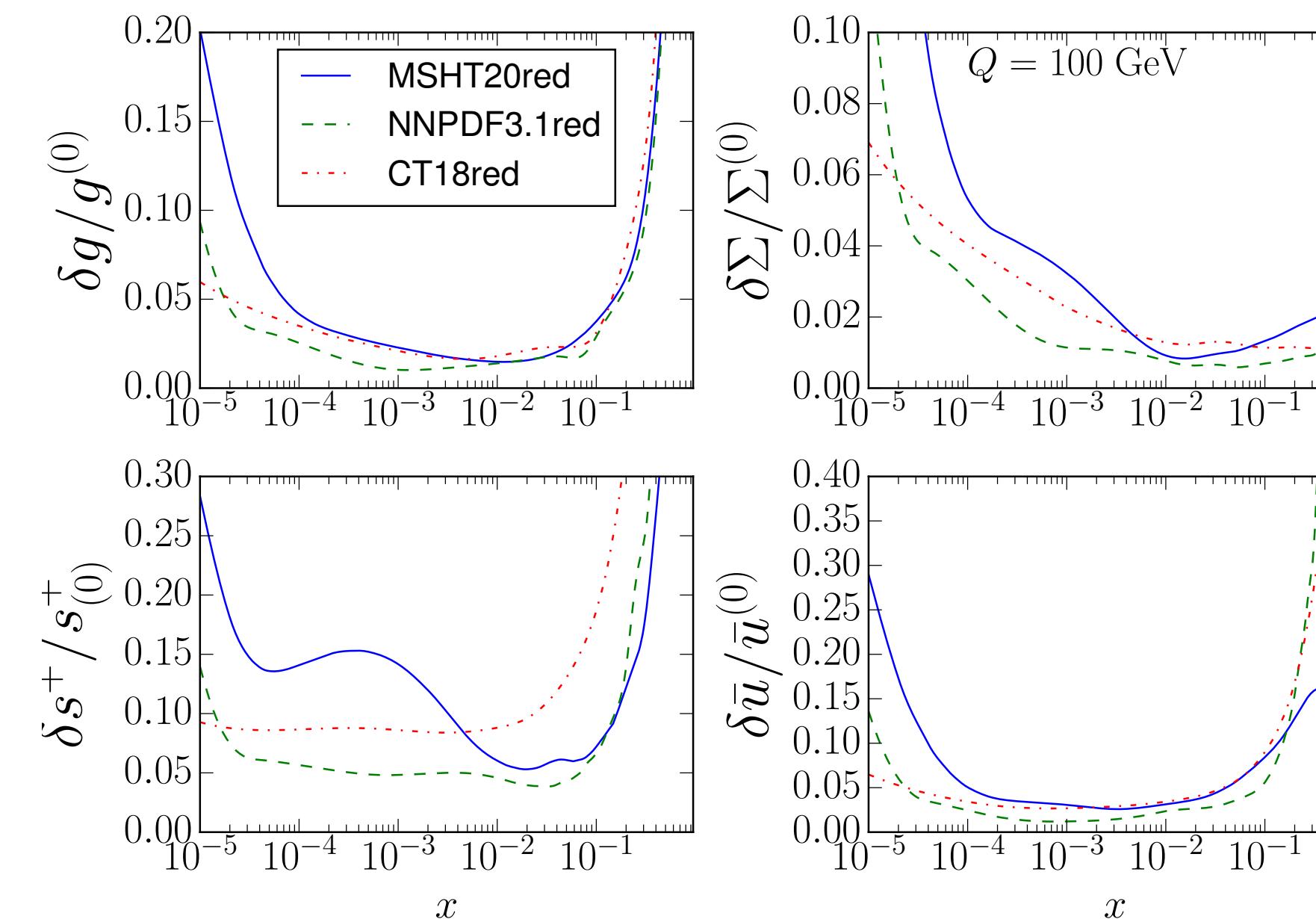
Each D_i gives f_i and from $\{f_i\} \Rightarrow$ PDF errors

G. Watt and R. Thorne, arXiv:1205.4024

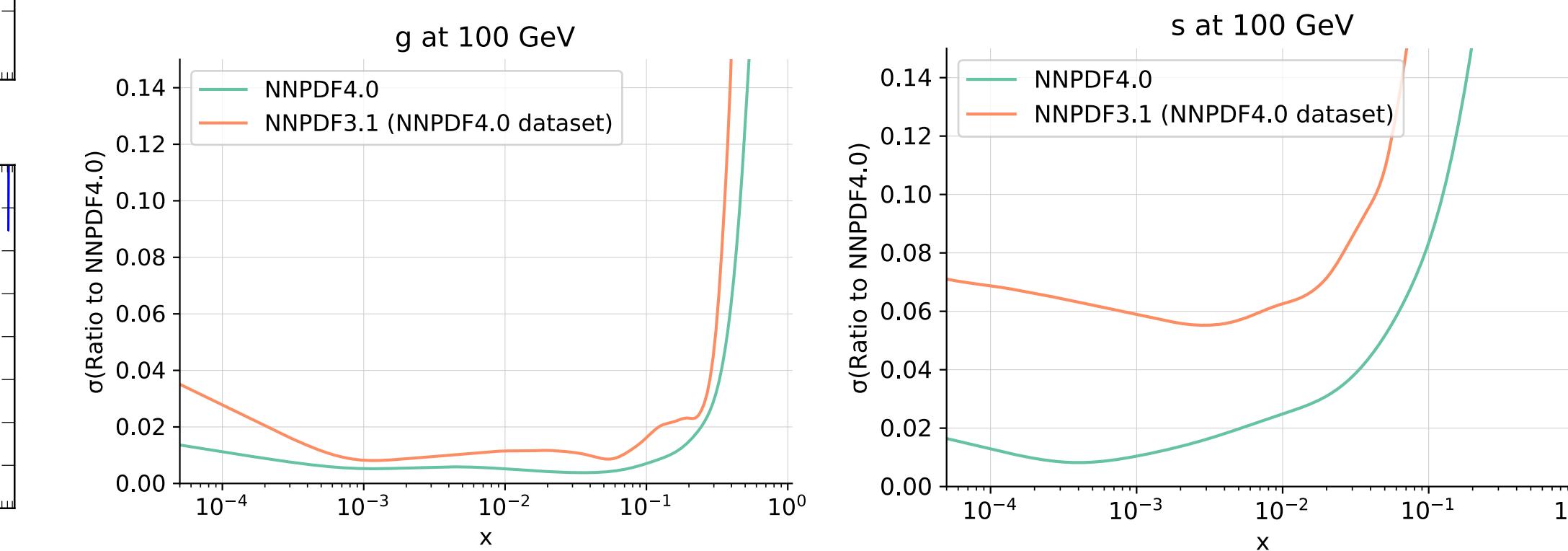
- ♦ Note not specific to NNs: can apply in fixed parameterisation as well: shown to be \sim equivalent to Hessian $\Delta\chi^2 = 1$ in that case.
- ♦ However, in NN approach direct correspondence is lost as Hessian approach does not apply.

♦ Global fits give different errors in PDF4LHC21 benchmarking. NNPDF3.1 in general **smaller errors.**

Benchmark = similar data/settings

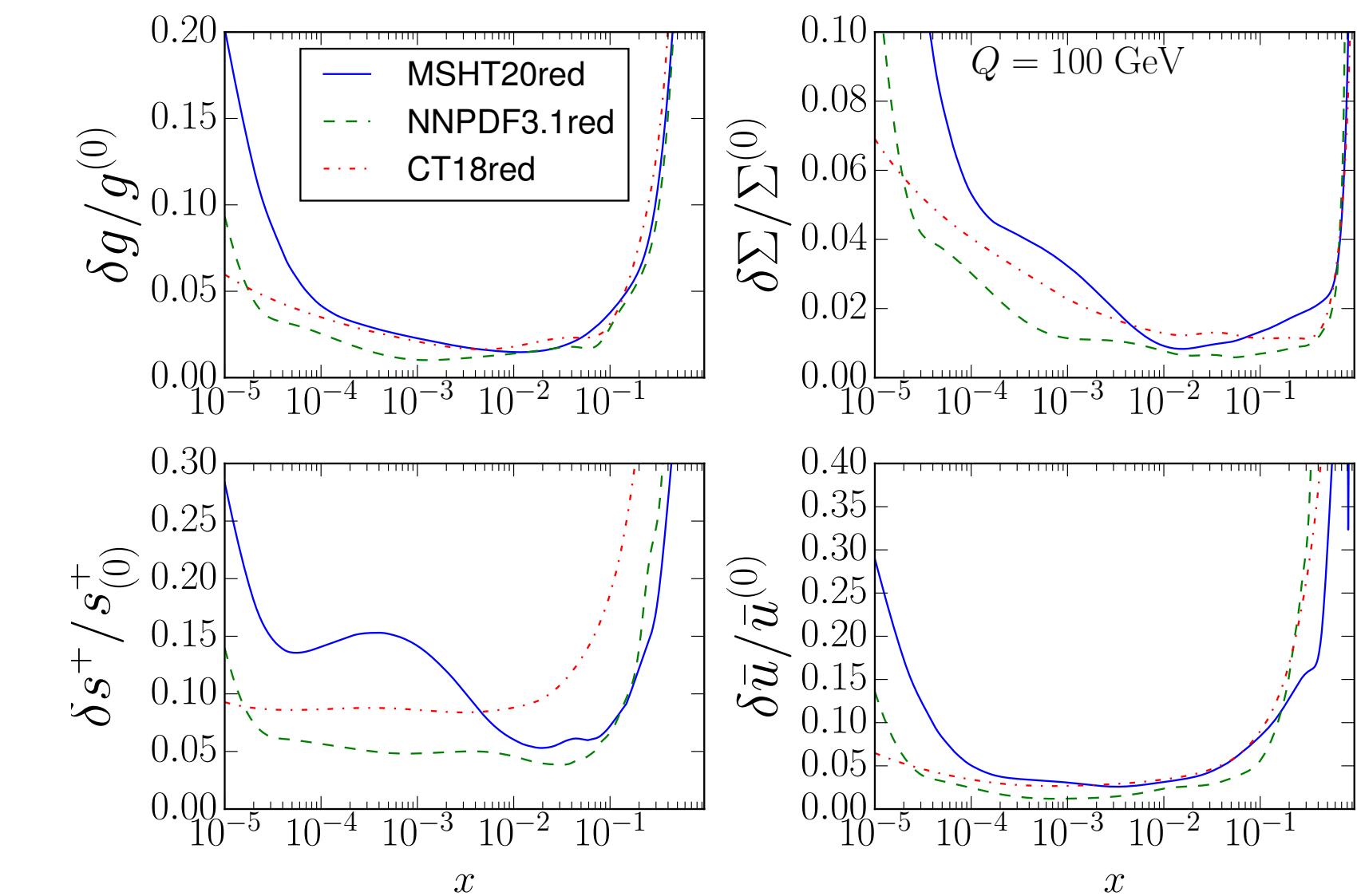


- ♦ And **4.0 methodology** gives further errors reduction.



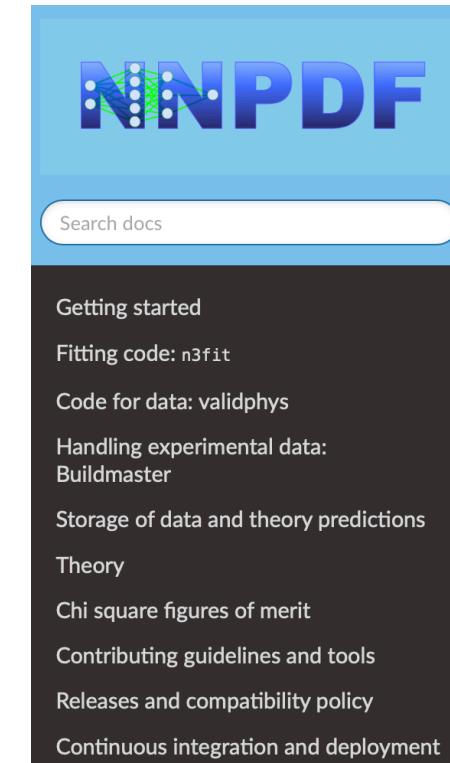
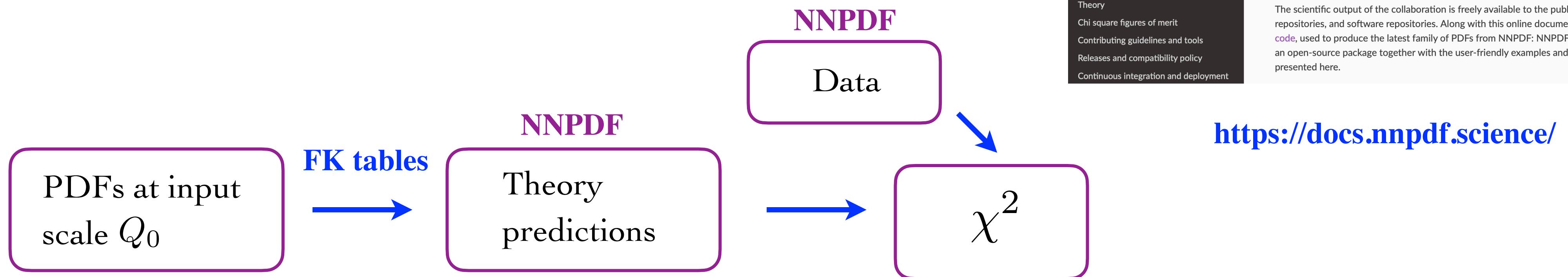
Aims of Talk

- Suggests three possibilities:
 1. **NNPDF4.0** uncertainty not conservative enough (**too small**).
 2. **MSHT (CT)** uncertainty too conservative (**too large**).
 3. **MSHT (CT)** fit less accurate, due to parameterisation inflexibility, and hence enlarged errors needed (**less precise**).
- Or some combination of the three. Finding out which clearly important for LHC precision.
- In this talk I will present results that aim to address this issue. In particular will show:
 - ★ **First** global **closure test** of fixed parameterisation (MSHT) approach: is parameterisation flexible enough to give faithful description of global pseudodata?
 - ★ **First** completely direct **comparison** between fixed parameterisation (MSHT) and NN approaches. How do these compare in full global fit?
- Study is ongoing, so all slides can be viewed as if they have a '**preliminary**' label on them!



Global Closure - set up

- How best to set up a global closure test? Will make use of publicly available NNPDF fitting code.
- Provides python libraries to load NNPDF dataset and theory predictions, given PDF set. More precisely gives:



/ The NNPDF collaboration

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The NNPDF collaboration

The [NNPDF collaboration](#) performs research in the field of high-energy physics. The NNPDF collaboration determines the structure of the proton using contemporary methods of artificial intelligence. A precise knowledge of the so-called **Parton Distribution Functions (PDFs)** of the proton, which describe their structure in terms of their quark and gluon constituents, is a crucial ingredient of the physics program of the Large Hadron Collider of CERN.

The NNPDF code

The scientific output of the collaboration is freely available to the public through the arXiv, journal repositories, and software repositories. Along with this online documentation, we release the [NNPDF code](#), used to produce the latest family of PDFs from NNPDF: NNPDF4.0. The code is made available as an open-source package together with the user-friendly examples and an extensive documentation presented here.

<https://docs.nnpdf.science/>

- Given arbitrary PDF set (grid of $\{f_i\}$ at $\{x_i\}$ and Q_0) can evaluate theory predictions + fit quality.
- This allows us to evaluate corresponding fit quality with a (MSHT) fixed parameterisation, but to NNPDF data/theory - **only difference** is **input parameterisation**. From above module can also build up optimizer in usual way to give best fit, Hessian errors etc.
- Will use for closure tests (though not essential) - but setting things up in this way will allow direct comparison at level of full fit.

Global Closure Test

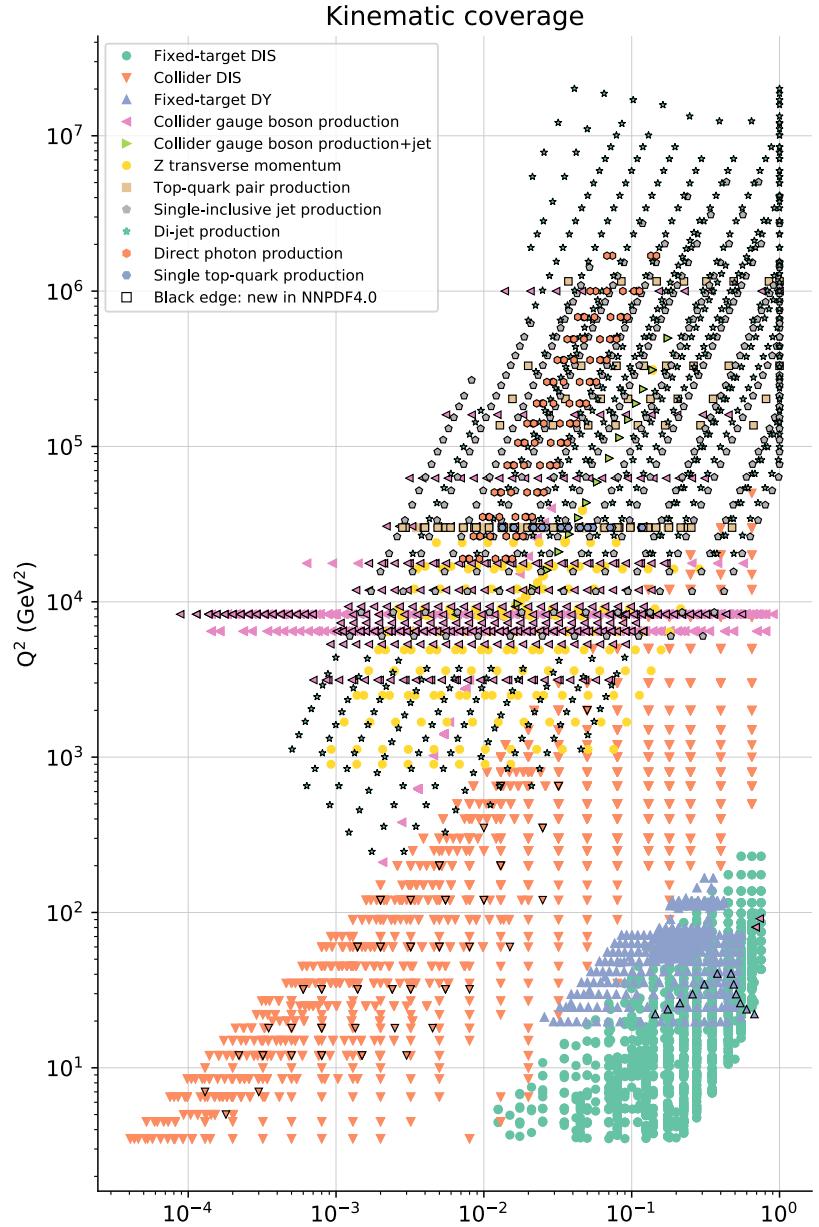
- For direct comparison will consider perturbative charm - NNPDF4.0pch set as input.
- Then generate unshifted pseudodata for 4.0 global dataset ($N_{\text{pts}} = 4627$). In principle exact agreement possible, with $\chi^2 = 0$. But will propagate data errors via Hessian approach, so $\sim \text{Level } 0 + 2$ (but not 1 yet).
- Then perform fit with default MSHT parameterisation. What do we find?

	χ^2	χ^2/N_{pts}
Fit quality:	2.4	0.0005

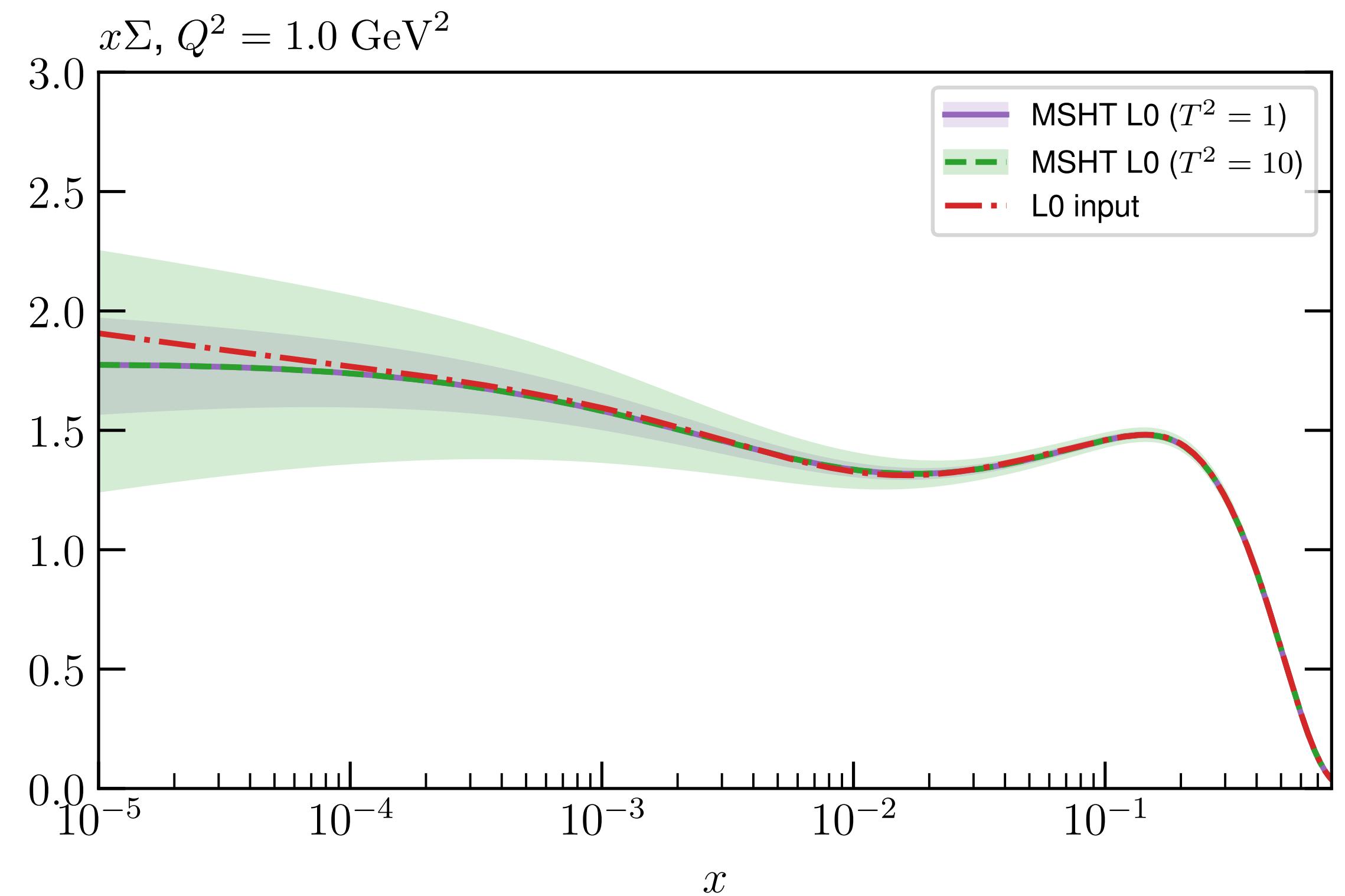
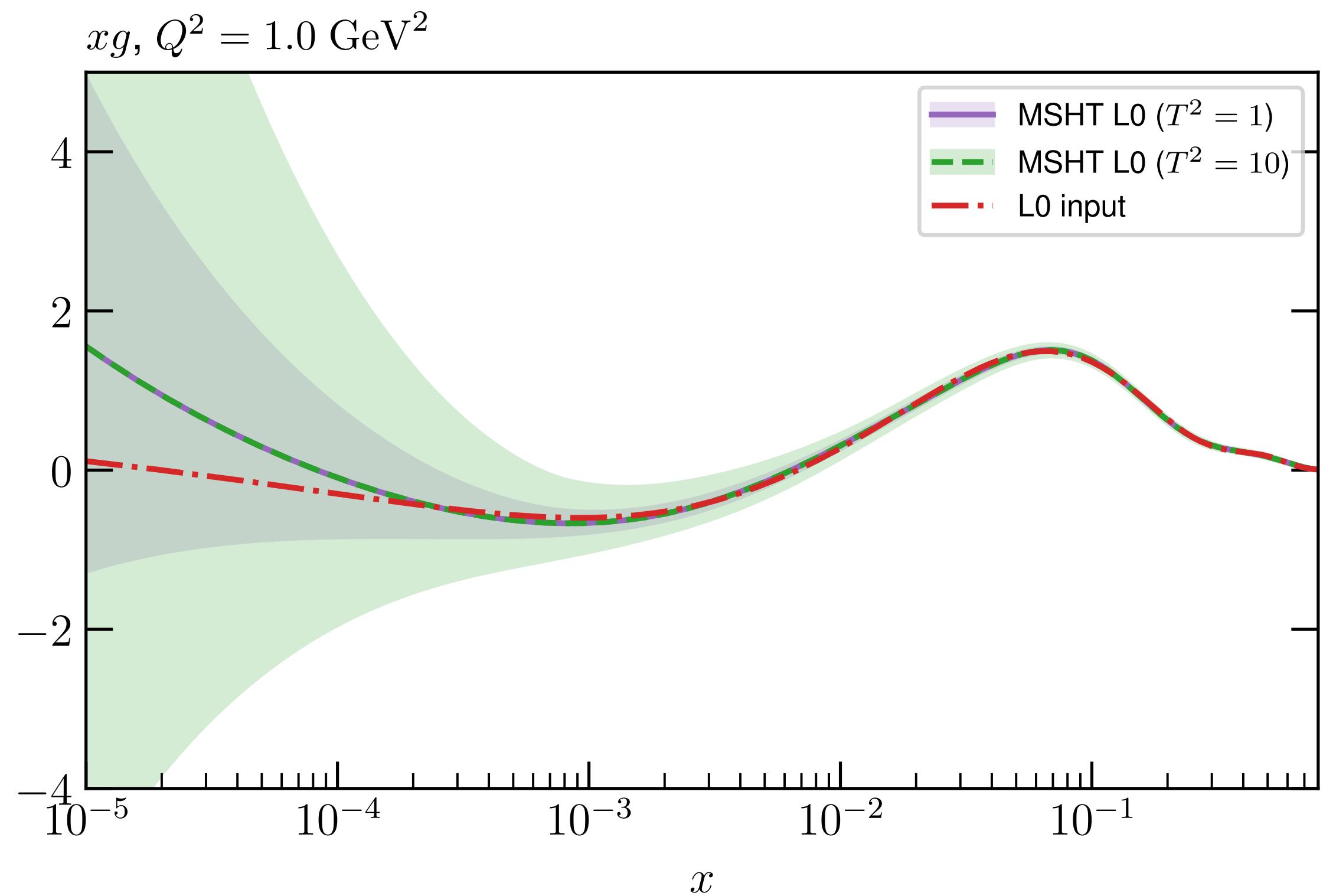
- **Remarkably good!** In fact lower than reported result of NNPDF L0 closure test.

	3.1 meth.	4.0 meth.
L. Del Debbio, T. Giani and M. Wilson, arXiv:2111.05787	χ^2/N_{pts}	0.012

- **Caveat:** only one input set, may well be different (not quite as good) for others. Trend should be similar.
- But apparently no issue with parameterisation inflexibility in this case. But what about PDFs?

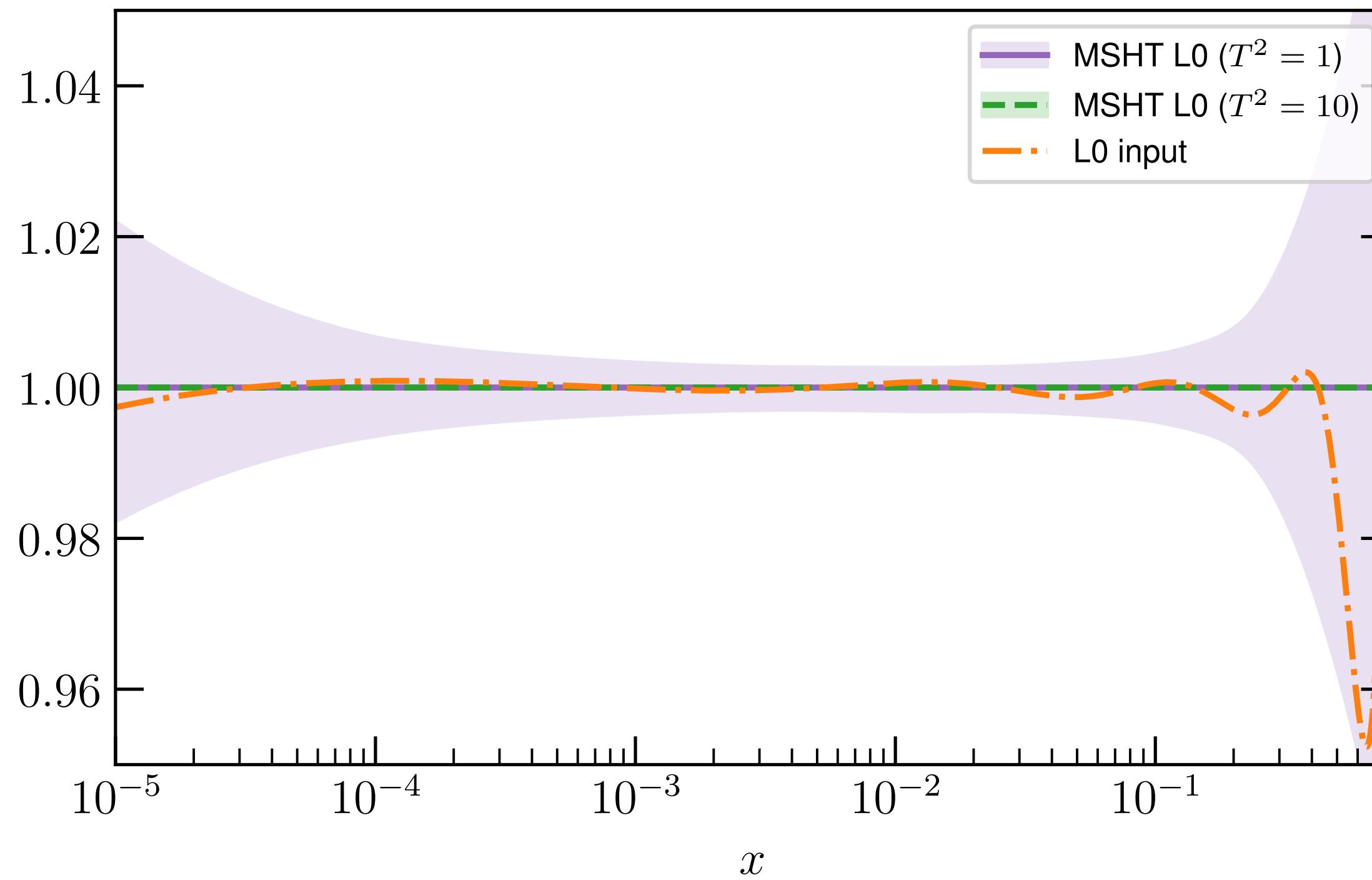


NNPDF, arXiv:2109.02653

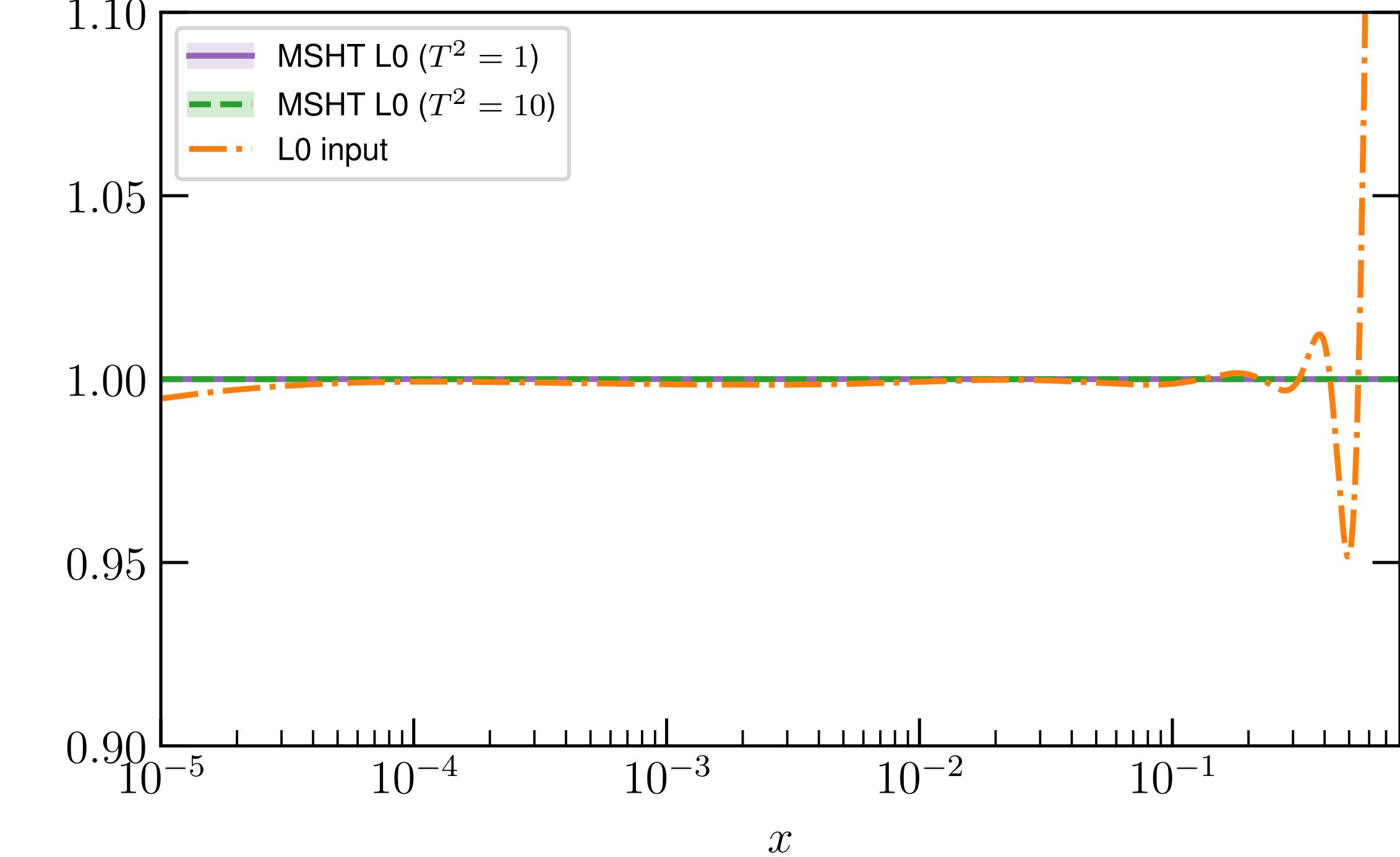


- First look: encouraging results! In more detail...

xg , PDF ratio, $Q^2 = 10^4 \text{ GeV}^2$



$x\bar{u}$, PDF ratio, $Q^2 = 10^4 \text{ GeV}^2$



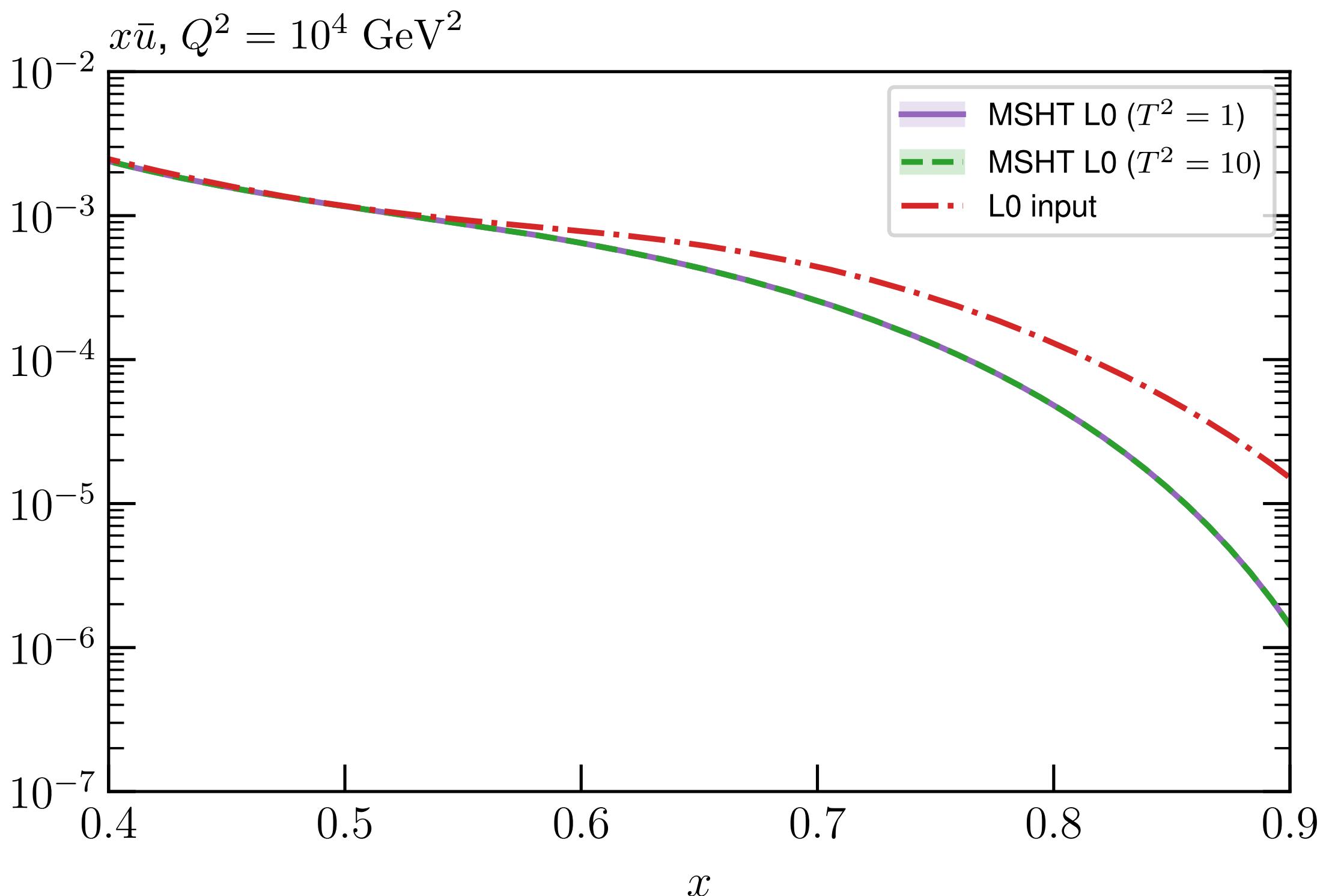
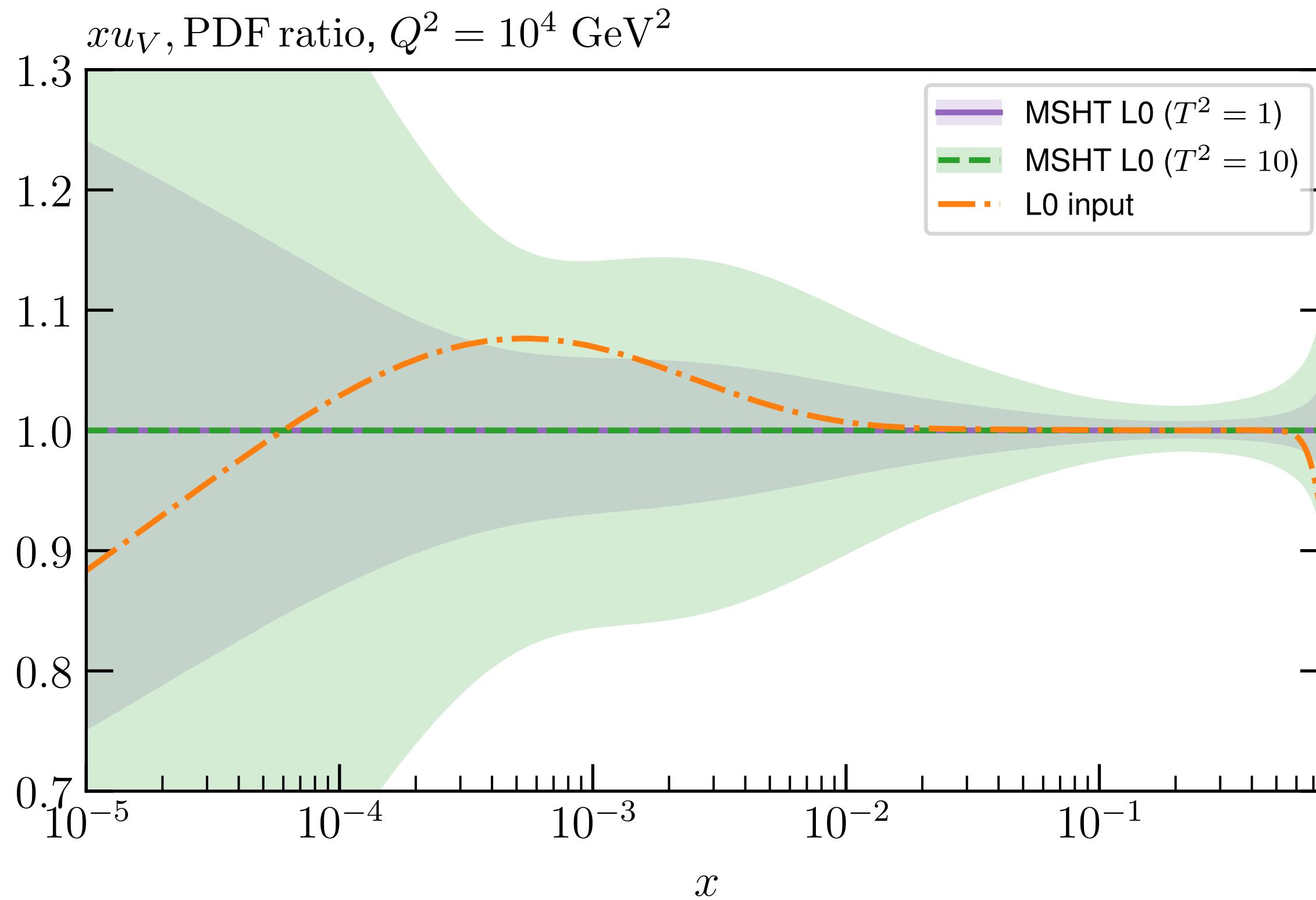
- Ratio of (NNPDF4.0pch) L0 input to fit result, including PDF uncertainties with $T^2 = 1$ and 10 that come from the closure test fit. Latter is \sim result of dynamic tolerance used in MSHT20 (checked here).

★ **Deviation** in general (in data region) **per mille** level and well within the $T^2 = 1$ uncertainties.

★ More precisely, deviation is $\sim 10\%$ or less of $T^2 = 1$ uncertainty, and a factor of $\sim 2 - 5$ lower a fraction of the $T^2 = 10$ uncertainty.

→ In **data region** L0 input PDF matched very well, and much better than $T^2 = 1$ uncertainties. **No evidence** that the increased tolerance is driven by **parameterisation inflexibility** for MSHT.

Similar results for
other quarks - see
backup



- In less well constrained regions deviation larger, e.g for u_V, d_V at low and high x and the \bar{u}, \bar{d} at high x .
- Hence in extrapolation region L0 input not always consistent within uncertainties
- As \sim outside data region not inconsistent (errors driven by data), but indicates more conservative error definition in these regions may be desirable (as tends to happen in NN approach).
- Though arguably no ‘right’ answer in true extrapolation region (too conservative vs. over-conservative).
- So far only considered L0 test, though L1 underway. Would not expect to change picture dramatically.

Comparison to NNPDF uncertainties

- Can do first comparison of **MSHT** vs. **NNPDF** PDF **uncertainties**. Not completely direct as MSHT is in closure test, and NNPDF result of full fit. But theory and underlying datasets same. Find:

★ Quark flavour decomposition: $\sigma(\text{NNPDF}) \sim \sigma(\text{MSHT}, T^2 = 1)$

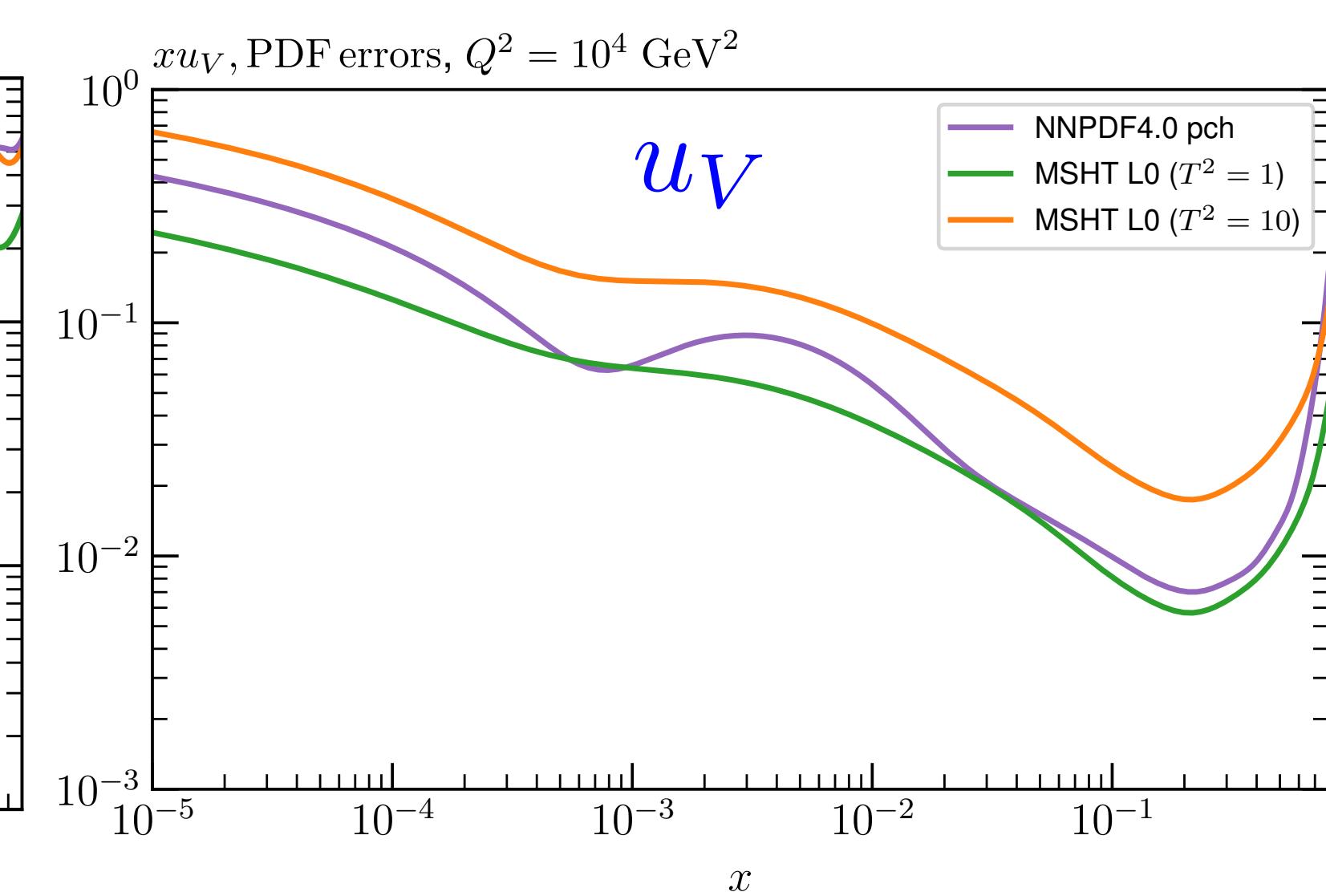
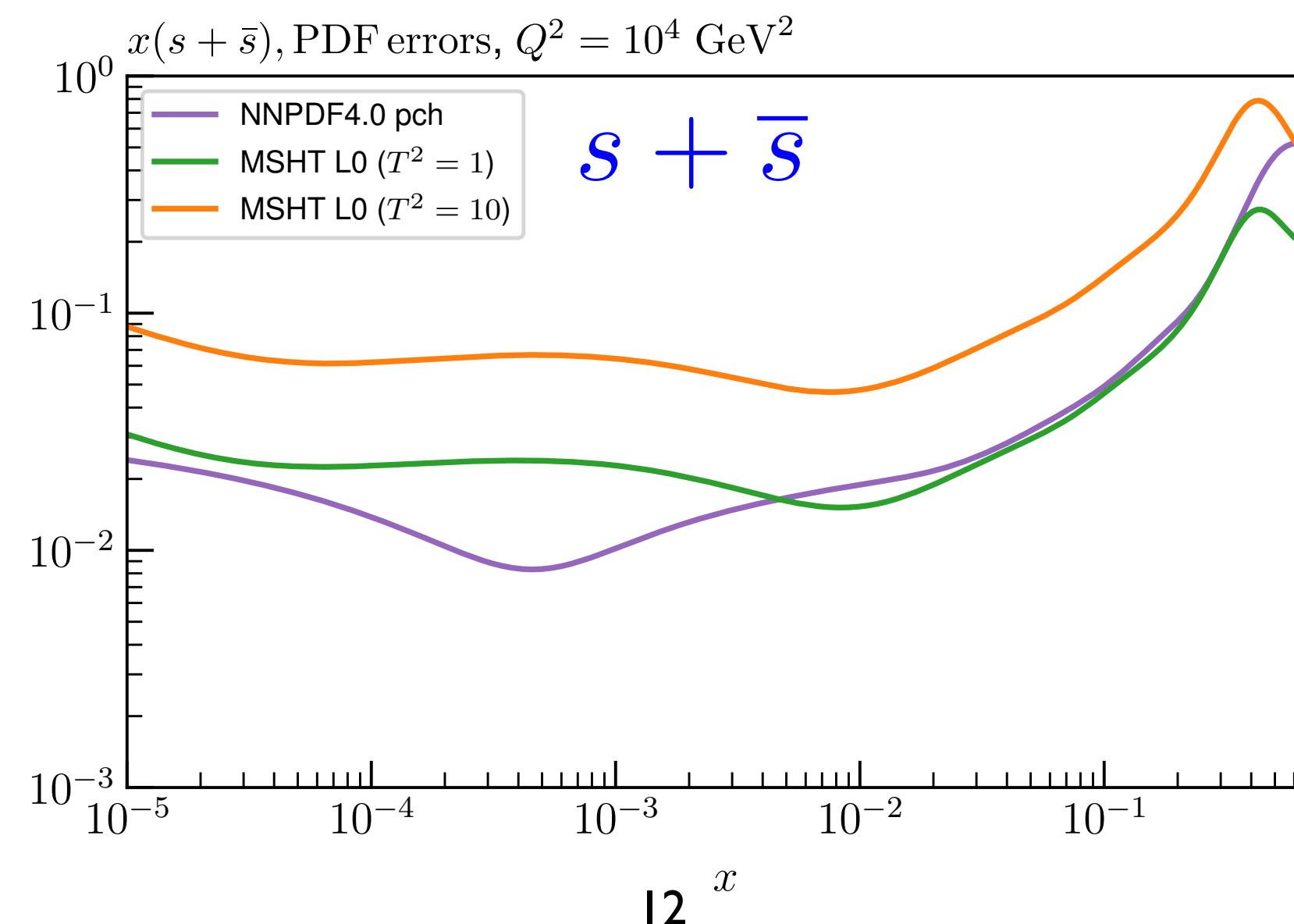
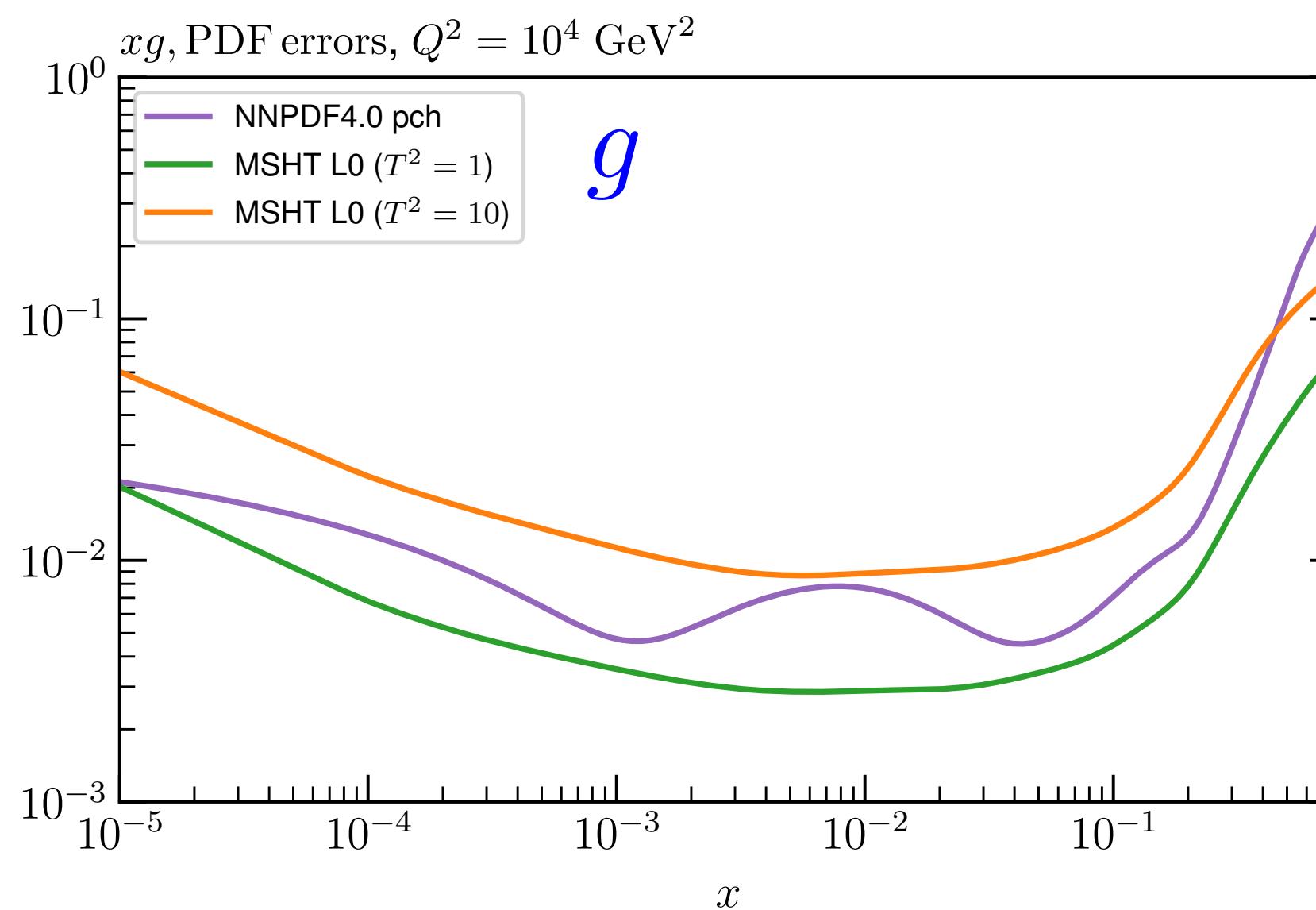
★ Gluon (singlet at intermediate x): $\sigma(\text{MSHT}, T^2 = 1) \lesssim \sigma(\text{NNPDF}) \lesssim \sigma(\text{MSHT}, T^2 = 10)$

MSHT, $T^2 = 1$

MSHT, $T^2 = 10$

NNPDF4.0 pch

- With rather similar overall trends with x .
- Exception at high x where NNPDF uncertainty becomes larger.



Full fit: comparison

- Can also consider result of fit to real data entering NNPDF4.0 fit. To restate: exactly **same data and theory**, with **only difference** from PDF input **parameterisation**.
- Will in addition consider case where positivity is imposed at PDF (and cross section) level, as in NNPDF fit, for $x_i \in \{5 \cdot 10^{-7}, 0.9\}$. Not something that is done in MSHT fits!

	NNPDF4.0 pch	MSHT fit	MSHT fit (w positivity)
$\chi^2_{t_0}$	5928.3 (1.282)	5736.7 (1.240)	5837.8 (1.262)
$\Delta\chi^2_{t_0} :$	$-191.6 (-0.04)$	$-90.5 (-0.02)$	

→ Fit quality with **MSHT** parameterisation is **significantly better** than result of central NNPDF set.

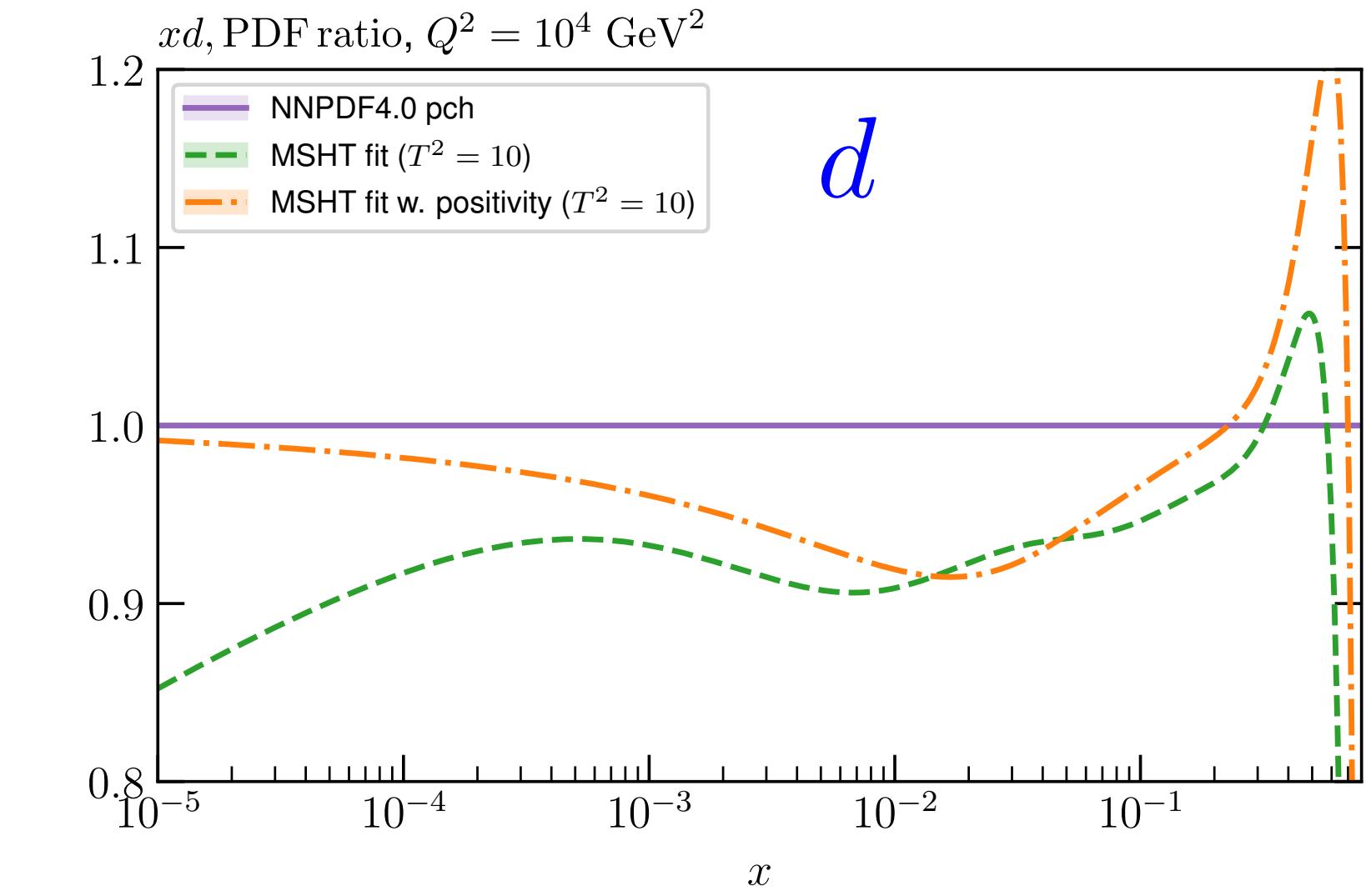
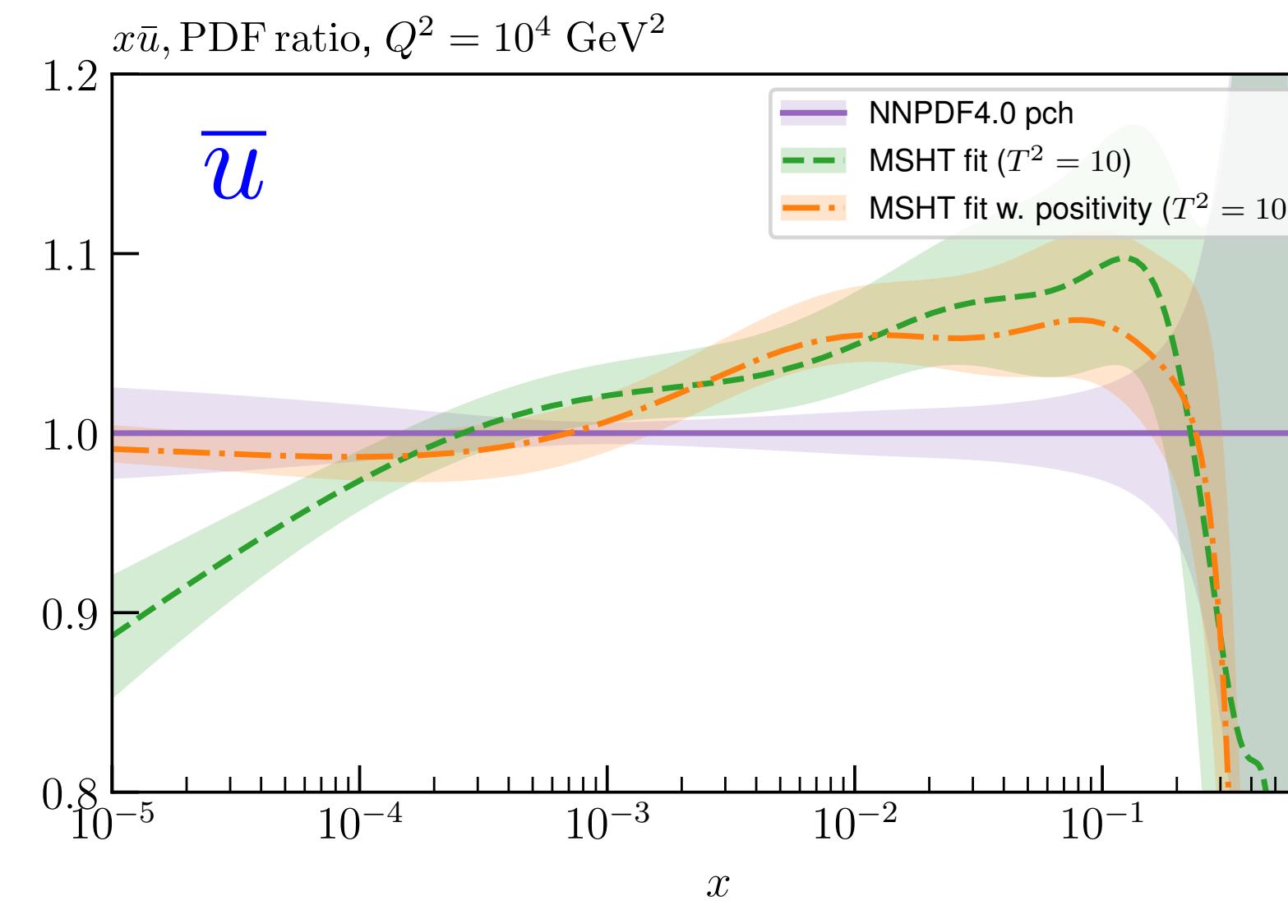
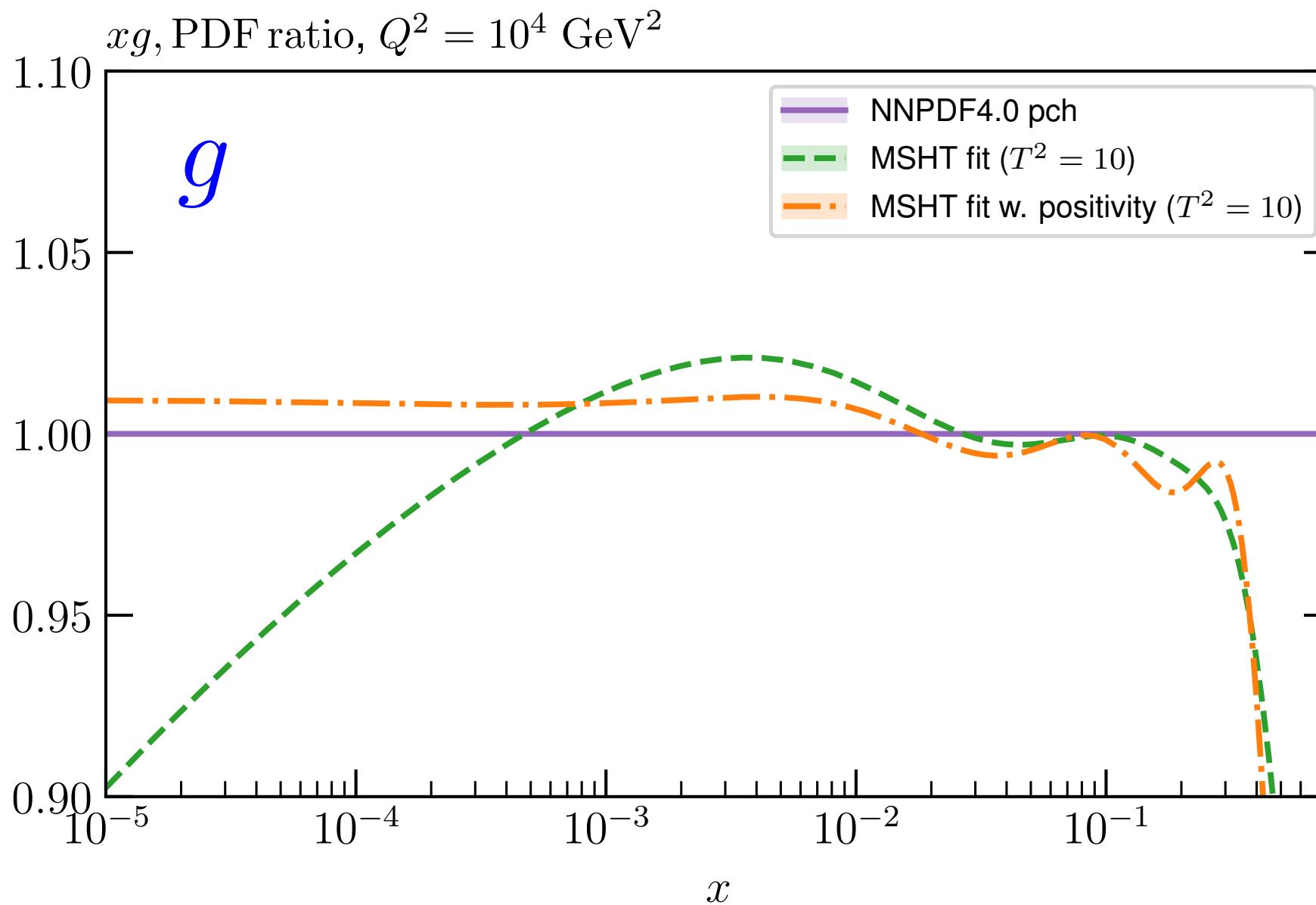
- **Positivity** clearly plays significant role - completely dominated by low x gluon. Indeed removing it from pure NNPDF fit gives $O(100)$ improvement. But **not the only difference** here. **Integrability checked and not issue: Backup**
- Improvement spread across fixed target, HERA DIS (without positivity) and LHC DY data.
- Do not expect central replica $\chi^2_{\text{rep},0}$ to be absolute minimum of χ^2 but difference too large for this. Overfitting seems unlikely given fixed parameterisation, though not impossible?

PDFs

- Comparing PDFs, see clear effect of positivity at low x : driven by known fact that default result prefers gluon to be negative at low x, Q^2 . Trend also seen if positivity removed from pure NNPDF fit.
- Imposing **positivity** gives much better agreement at low x , but clear **difference** in flavour decomposition **remains**.
- Notable that qualitatively this follows trend of using flavour rather than evolution basis (though difference larger and not identical).

See Backup

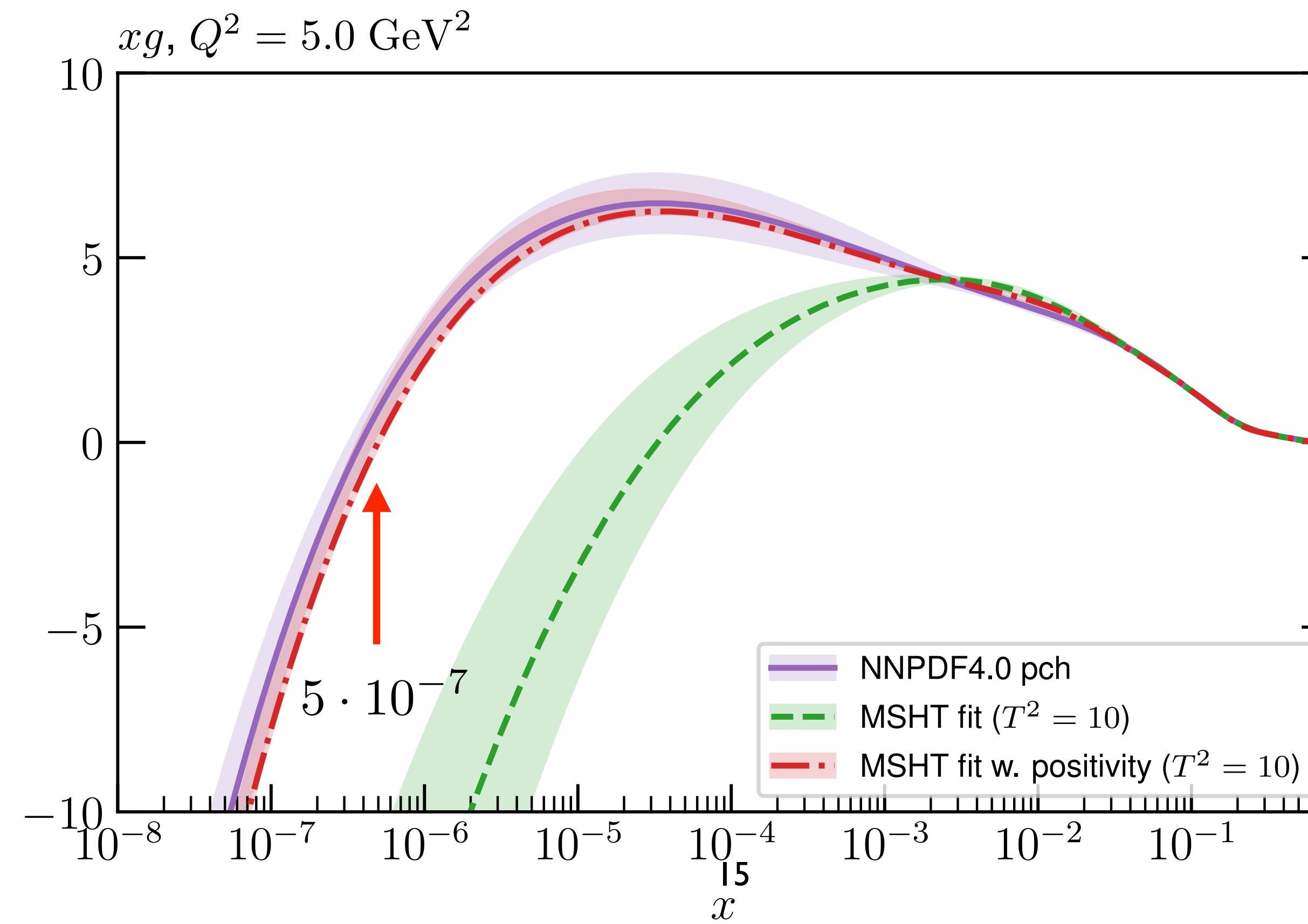
NNPDF4.0pch MSHT (pos) $T^2 = 10$ MSHT, $T^2 = 10$



Positivity?

- General arguments for imposing strict positivity on PDFs outside of current data region rely on perturbative stability. Not clear for (very) low gluon - sensitivity to resummation etc.
- More importantly - **all cases** are actually **negative** at low x ! Notable that the NNPDF gluon still prefers to be as negative as possible, i.e. just below the minimum x_i value where positivity imposed.

Driving fit in undesirable way?



Comparison to NNPDF uncertainties

- Can do same comparison of MSHT vs. NNPDF PDF uncertainties but now in global fit. Completely like-for-like. Results very similar to closure test comparison:

★ Quark flavour decomposition: $\sigma(\text{NNPDF}) \sim \sigma(\text{MSHT}, T^2 = 1)$

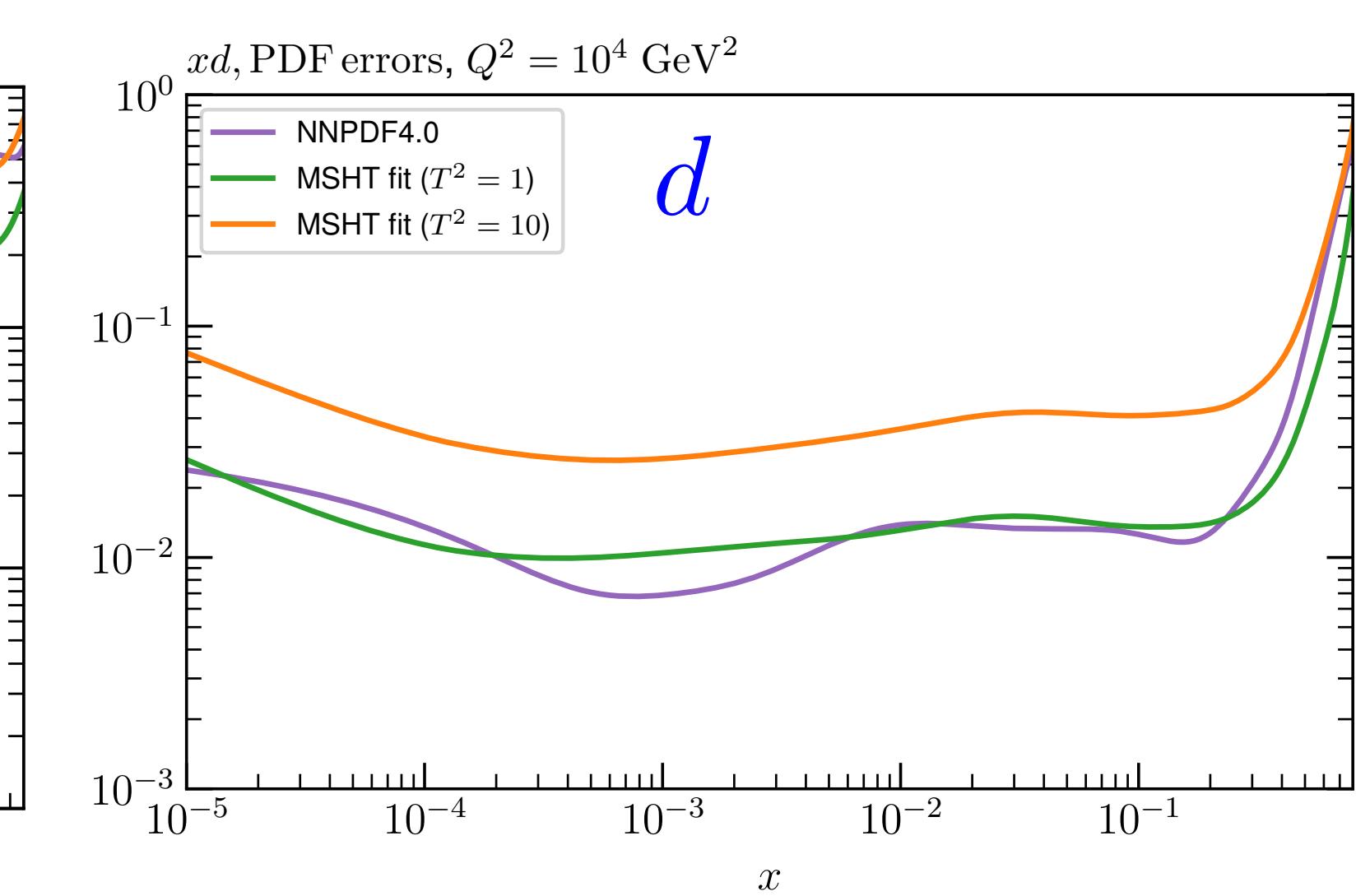
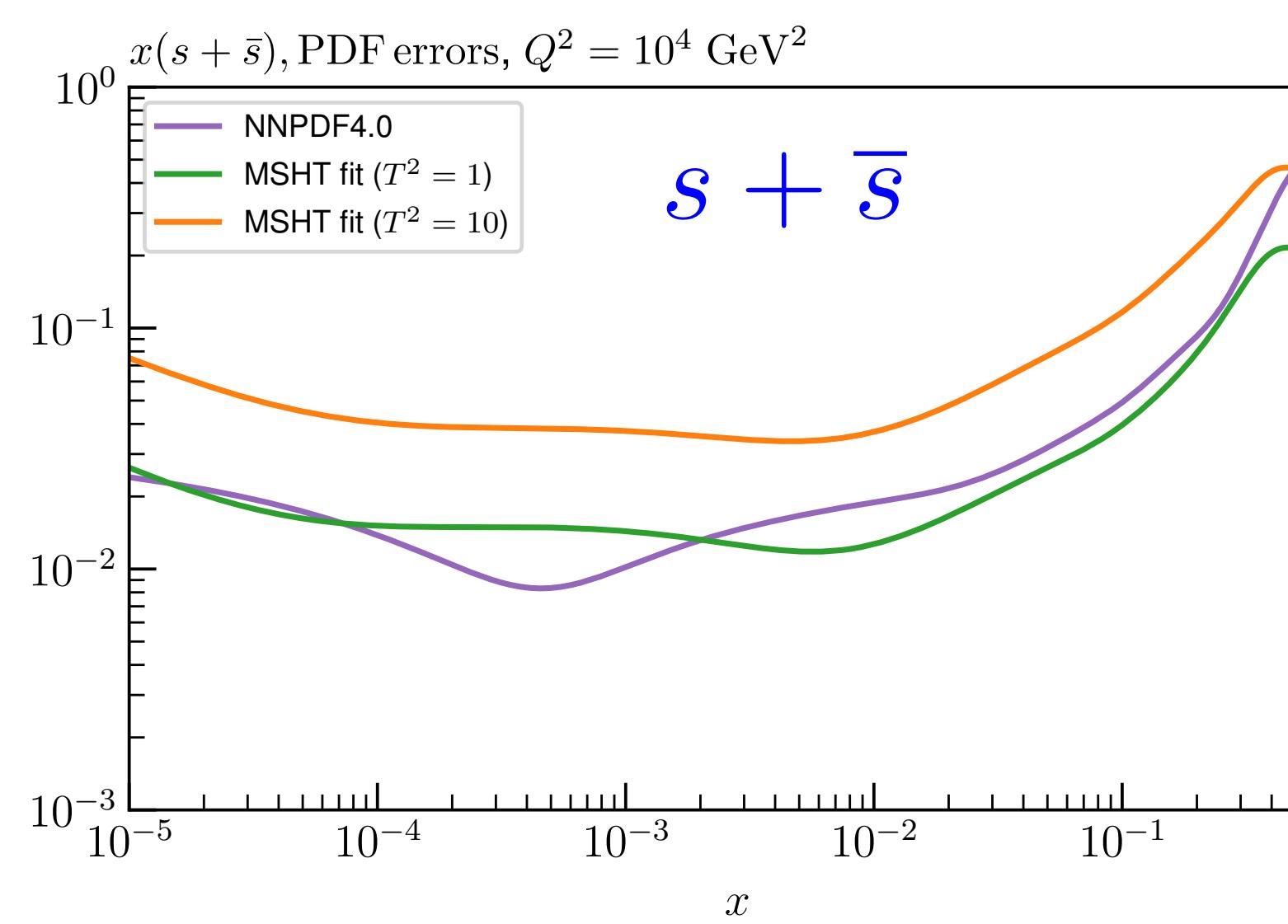
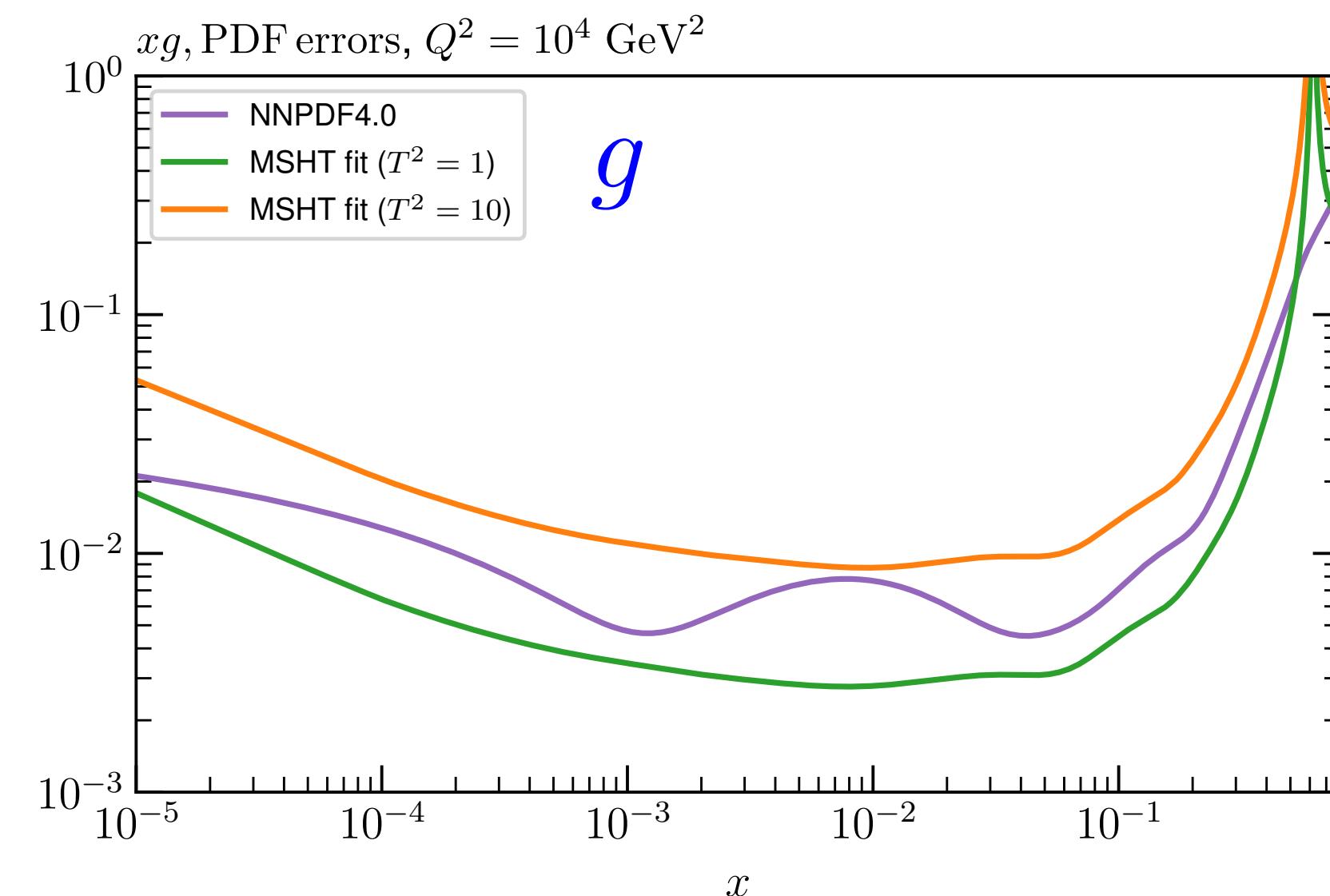
★ Gluon (singlet at intermediate x): $\sigma(\text{MSHT}, T^2 = 1) \lesssim \sigma(\text{NNPDF}) \lesssim \sigma(\text{MSHT}, T^2 = 10)$

MSHT, $T^2 = 1$

MSHT, $T^2 = 10$

NNPDF4.0pch

- With rather similar overall trends with x .
- Exception at high x where NNPDF uncertainty can become larger.

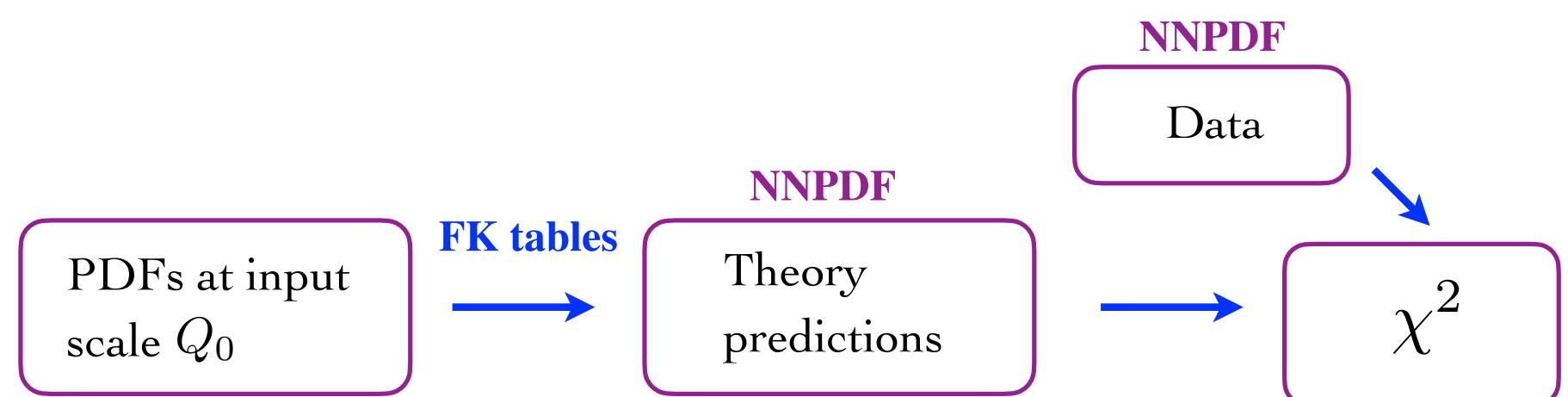


Full fit with fitted charm

- Readily extend previous study to include **fitted charm**. NNPDF theory inputs change accordingly, while PDFs parametrised at $Q_0 = 1.65 \text{ GeV} (> m_c)$ rather than 1 GeV, and parameterise charm:

$$xc_+(x, Q_0) = A_{c+} x^{\delta_{c+}} (1-x)^{\eta_{c+}} \left(1 + \sum_{i=1}^6 a_{c,i} T_i(y(x)) \right)$$

$$xc_-(x, Q_0) = 0$$



- Find:

	NNPDF4.0	MSHT fit	MSHT fit (w positivity)
$\chi^2_{t_0}$	5692.1 (1.233)	5645.2 (1.222)	5651.0 (1.224)
$\Delta\chi^2_{t_0} :$	<u>-46.9 (0.011)</u>	<u>-41.1 (0.009)</u>	

→ Fit quality with **MSHT** parameterisation again **better** than result of central NNPDF set, albeit by less than in perturbative charm case.

- Improvement spread across fixed target and LHC DY data.
- Role of **positivity** now **marginal** (confirmed with direct NNPDF fit).

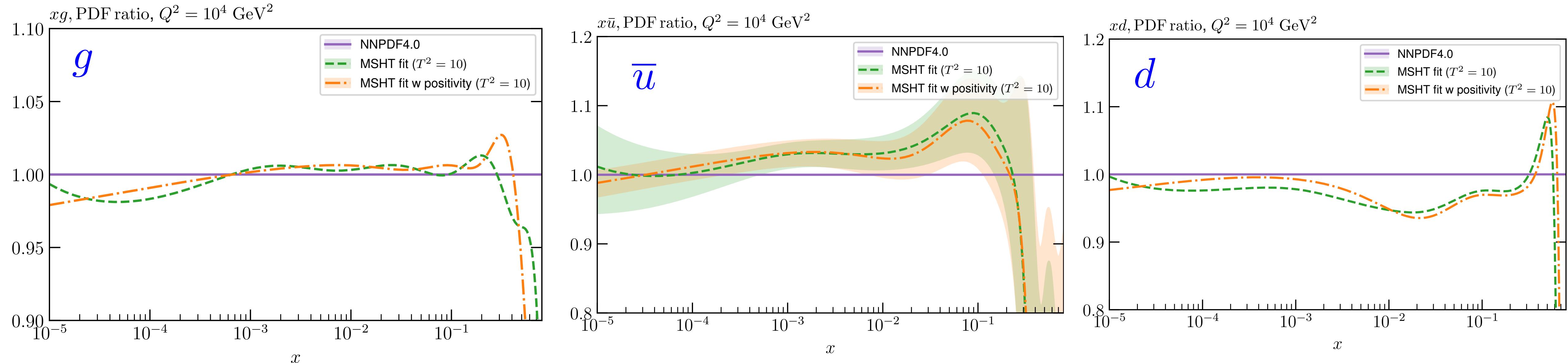
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PDFs

- Confirm that role of **positivity** now **marginal**. Also seen if positivity removed from pure NNPDF fit.
- Difference less than in p. charm case, but **clear difference** in **flavour decomposition** remains.
- Again qualitatively this follows trend of using flavour rather than evolution basis (though difference larger and not identical).
- Show $T^2 = 10$ for concreteness and fact that there is reasonable (not perfect) agreement within these not relevant factor, given fit quality is better for MSHT sets.

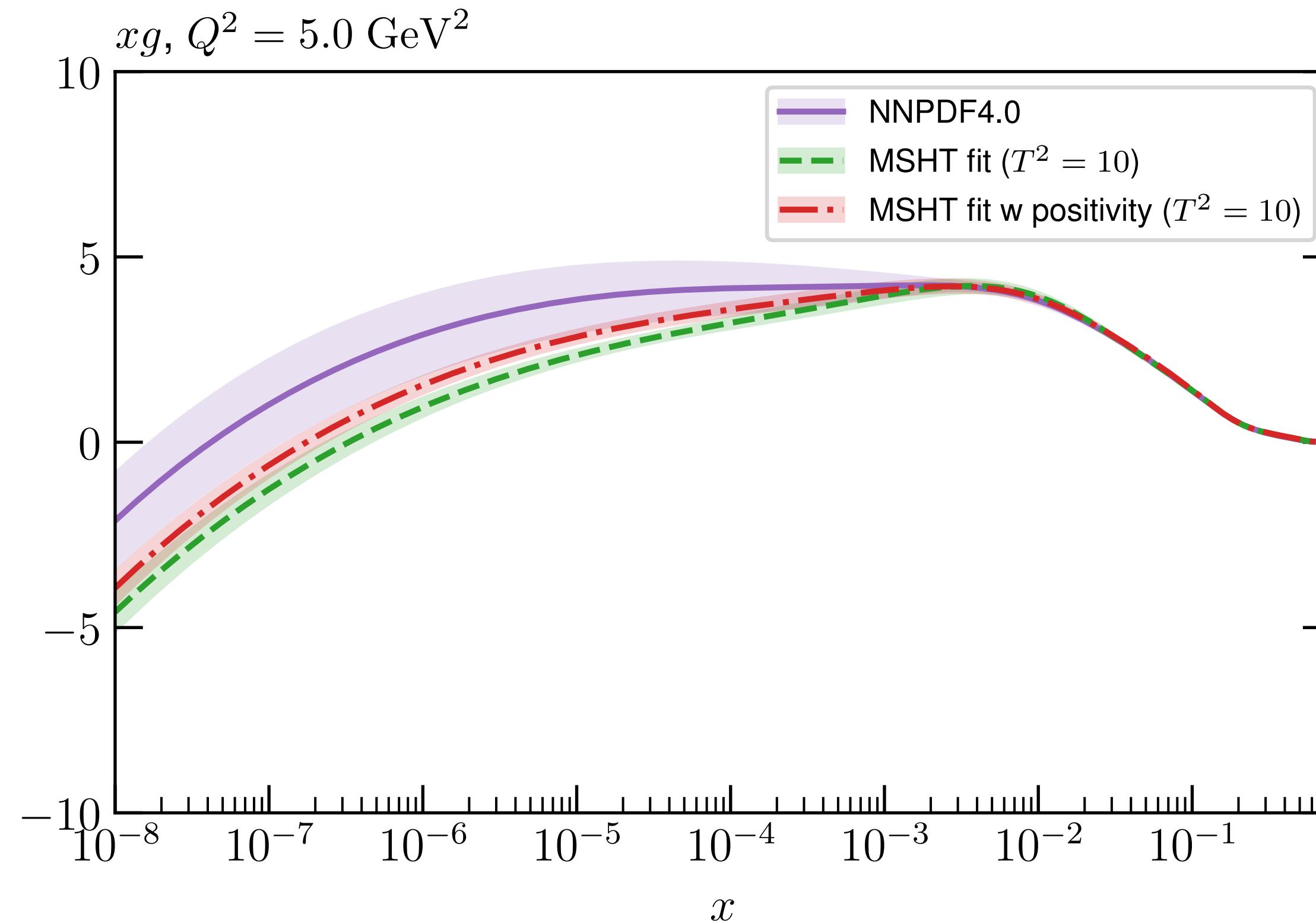
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NNPDF4.0pch MSHT (pos) $T^2 = 10$ MSHT, $T^2 = 10$



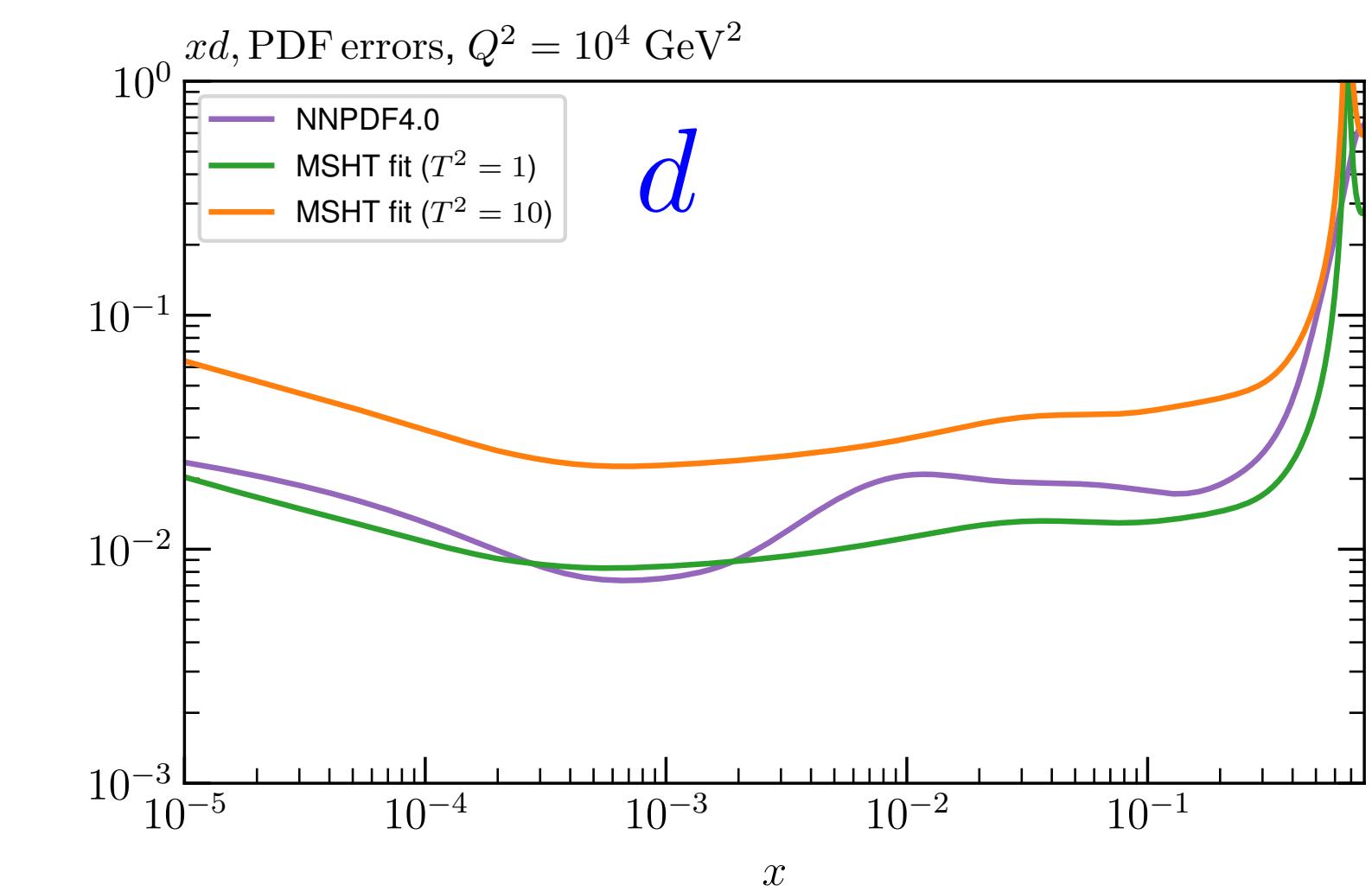
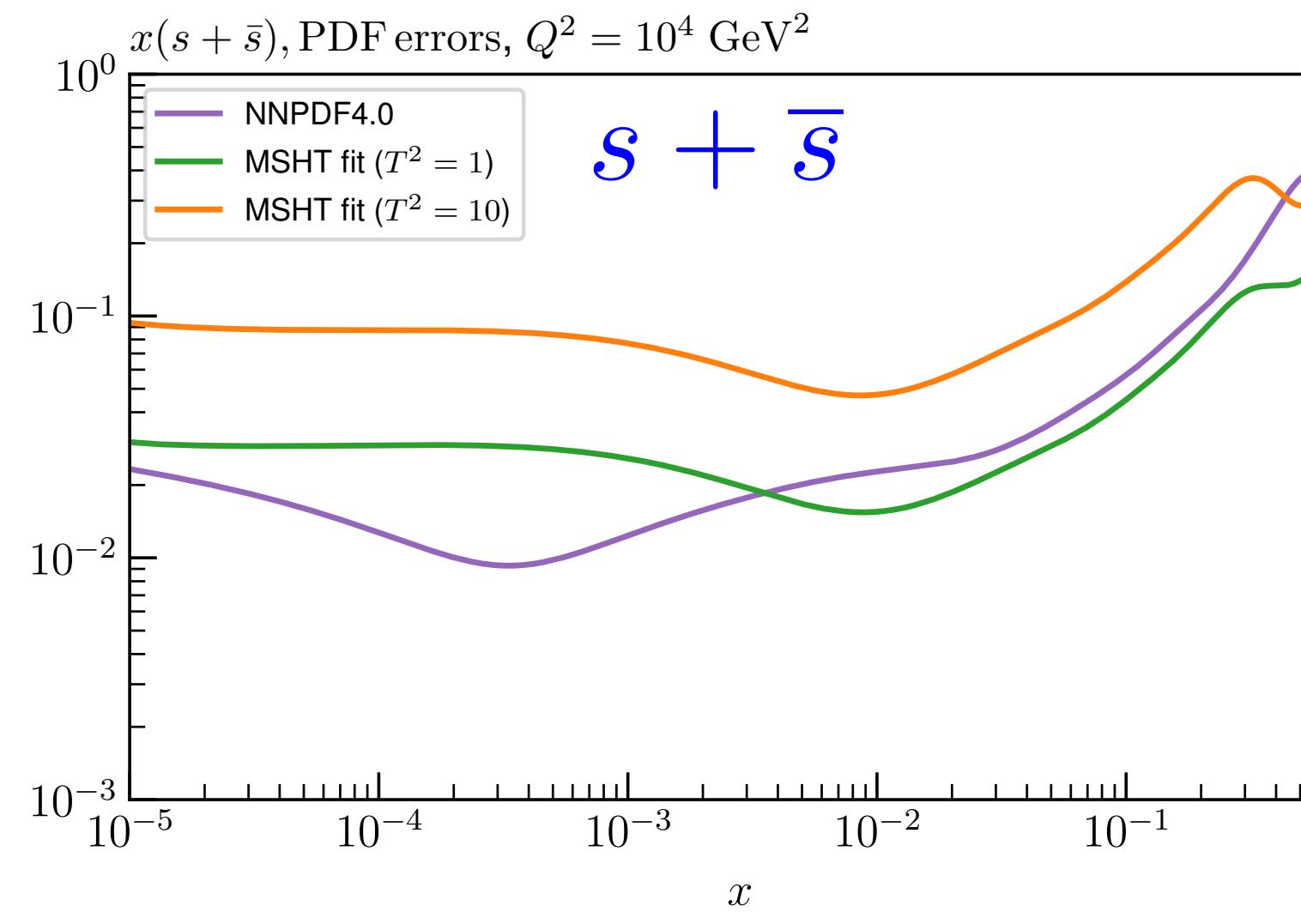
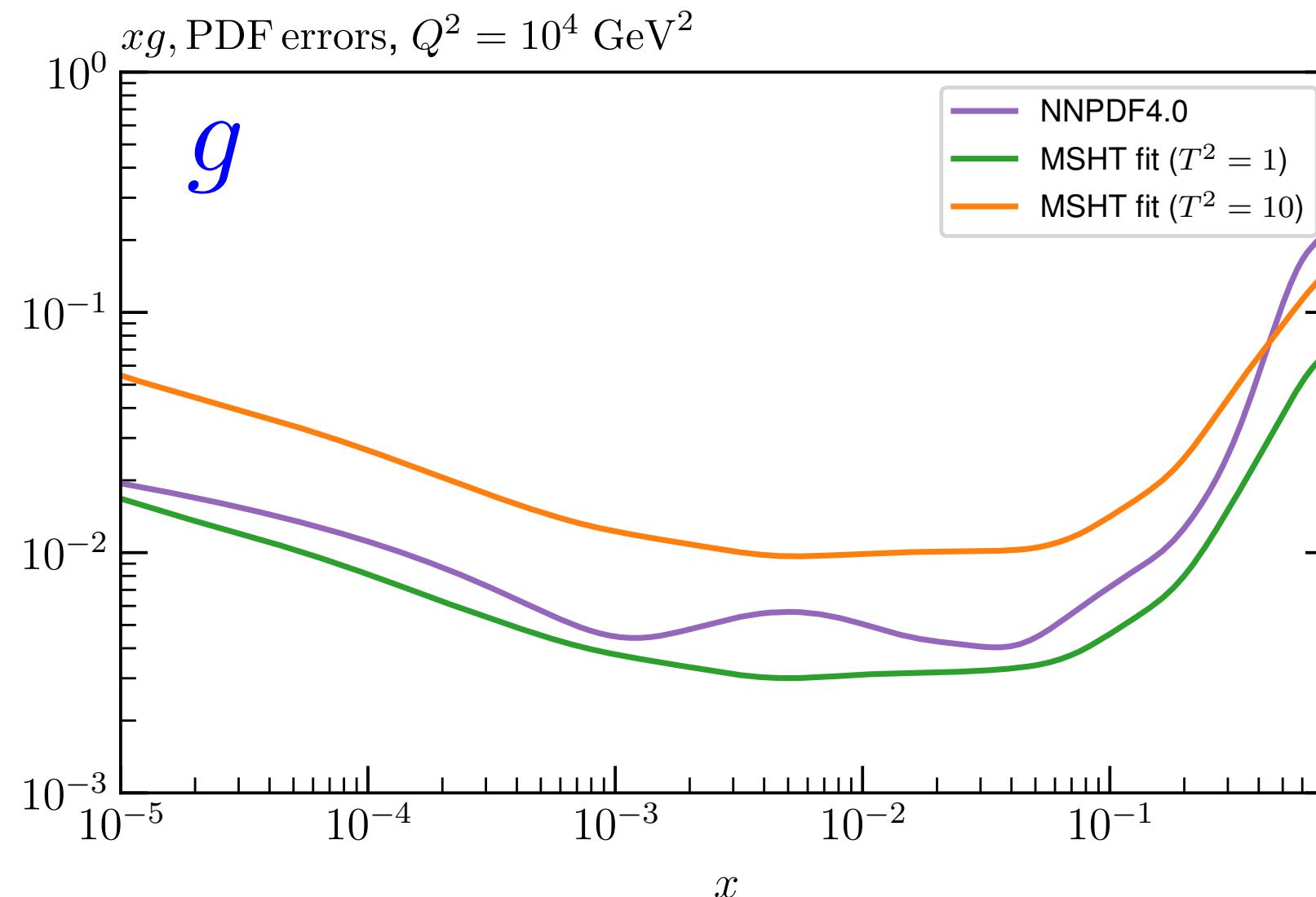
Positivity?

- With fitted charm tendency for preferred gluon to be negative reduced and pushed to lower x : why positivity requirement plays minor role - also seen in direct NNPDF fit.
- Not clear why this is (under investigation) but given intrinsic charm is expected to be high x phenomena might be concern?



Comparison to NNPDF uncertainties

- Can do same comparison of MSHT vs. NNPDF PDF uncertainties but now with fitted charm. Again completely like-for-like. Results very similar to closure test comparison:
 - ★ Quark flavour decomposition: $\sigma(\text{NNPDF}) \sim \sigma(\text{MSHT}, T^2 = 1)$
 - ★ Gluon (singlet at intermediate x): $\sigma(\text{MSHT}, T^2 = 1) \lesssim \sigma(\text{NNPDF}) \lesssim \sigma(\text{MSHT}, T^2 = 10)$
- MSHT, $T^2 = 1$
 MSHT, $T^2 = 10$
 NNPDF4.0
- With rather similar overall trends with x .
 - Exception at high x where NNPDF uncertainty can become larger.
 - Some trend for gluon to be a little closer to $T^2 = 1$ case.

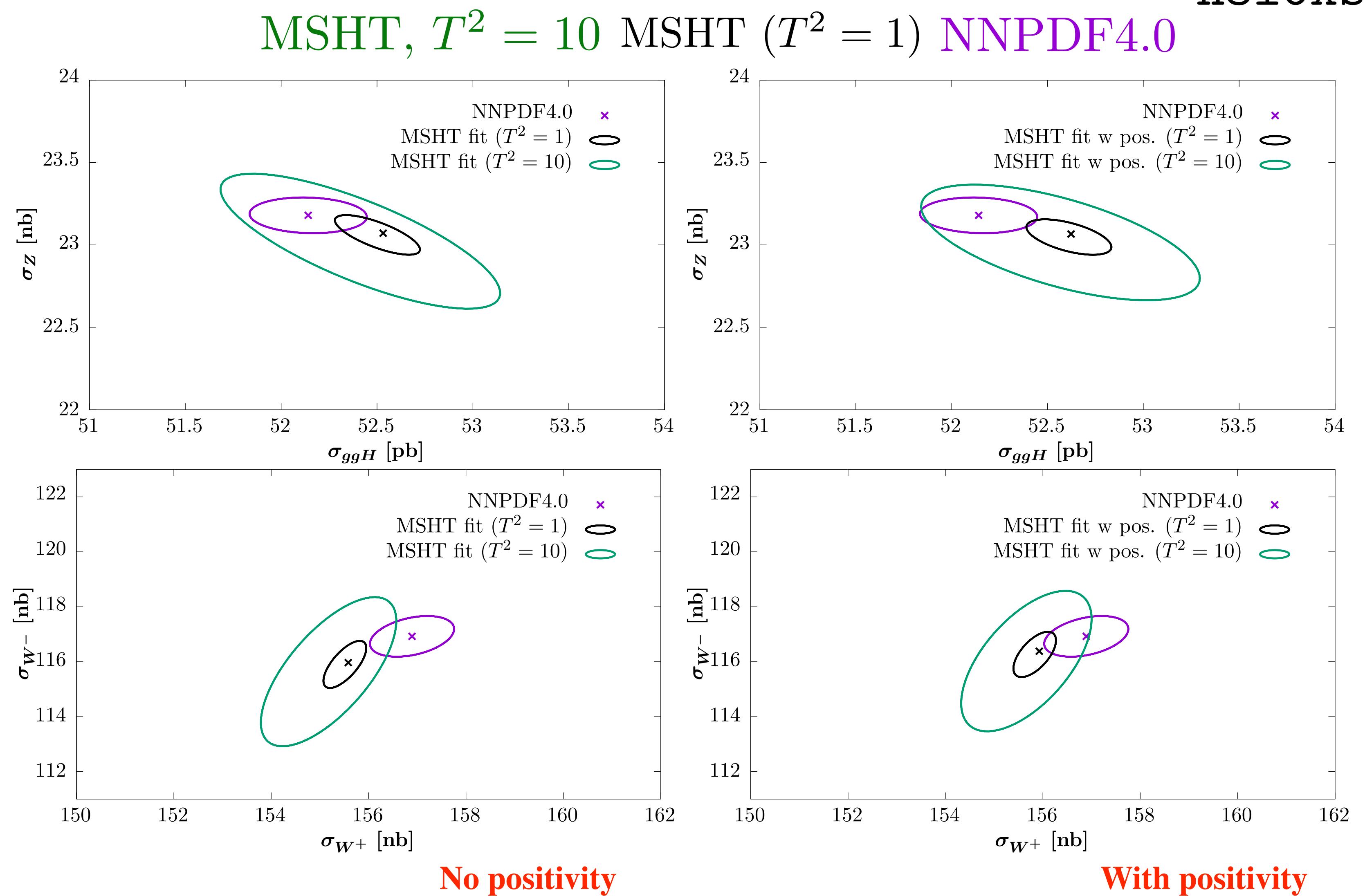


Cross Sections

- Consider ggH, and W, Z cross sections (14 TeV) in fitted charm case:

- ★ NNPDF uncertainties ~ MSHT ($T^2 = 1$) but **significantly smaller** than $T^2 = 10$
- ★ NNPDF and MSHT fit basically consistent within $T^2 = 1$ uncertainties but not relevant factor given fit qualities.

p. charm: Backup



Summary

- In this talk I have presented:
 - ★ First global **closure test** of fixed parameterisation (MSHT) approach: is parameterisation flexible enough to give faithful description of global pseudodata?
 - ◆ Yes: no issue in passing (unfluctuated) global closure test.
 - ★ First completely direct **comparison** between fixed parameterisation (MSHT) and NN approaches. How do these compare in full global fit?
 - ◆ At level of errors $\sigma(\text{NNPDF}) \sim \sigma(\text{MSHT}, T^2 = 1)$ in general with some exceptions
 - gluon larger though less than $T^2 = 10$ (MSHT20 default).
 - ◆ At level of PDFs, surprisingly find fit quality is **lower** in **MSHT** fixed parameterisation case, and outside of NNPDF uncertainties. Reason for this is currently unclear. Positivity clearly important in p. charm case, but not only source (PDF basis?).

Summary

- Returning to the original possibilities (focus on MSHT as only considered here). We need to - and can
 - work out which is true:
 1. NNPDF4.0 uncertainty not conservative enough (too small).
 2. MSHT uncertainty too conservative (too large).
 3. MSHT fit less accurate, due to parameterisation inflexibility, and hence enlarged errors needed (less precise).

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 1. NNPDF4.0 uncertainty not conservative enough (too small).
 2. MSHT uncertainty too conservative (too large).
 3. ~~MSHT fit less accurate, due to parameterisation inflexibility, and hence enlarged errors needed (less precise).~~
- Successful closure test + comparison to NNPDF4.0 global fit suggests 3 is not dominant issue (at least in data region) for MSHT. Can be issue for less flexible ones and for MSHT in extrapolation regions.

Summary

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- Successful closure test + comparison to NNPDF4.0 global fit suggests 3 is not dominant issue (at least in data region) for MSHT. Can be issue for less flexible ones and for MSHT in extrapolation regions.
- First direct comparison to NNPDF4.0 global fit finds that this gives inherently different (smaller) uncertainties than MSHT fixed parameterisation, keeping everything else equal.

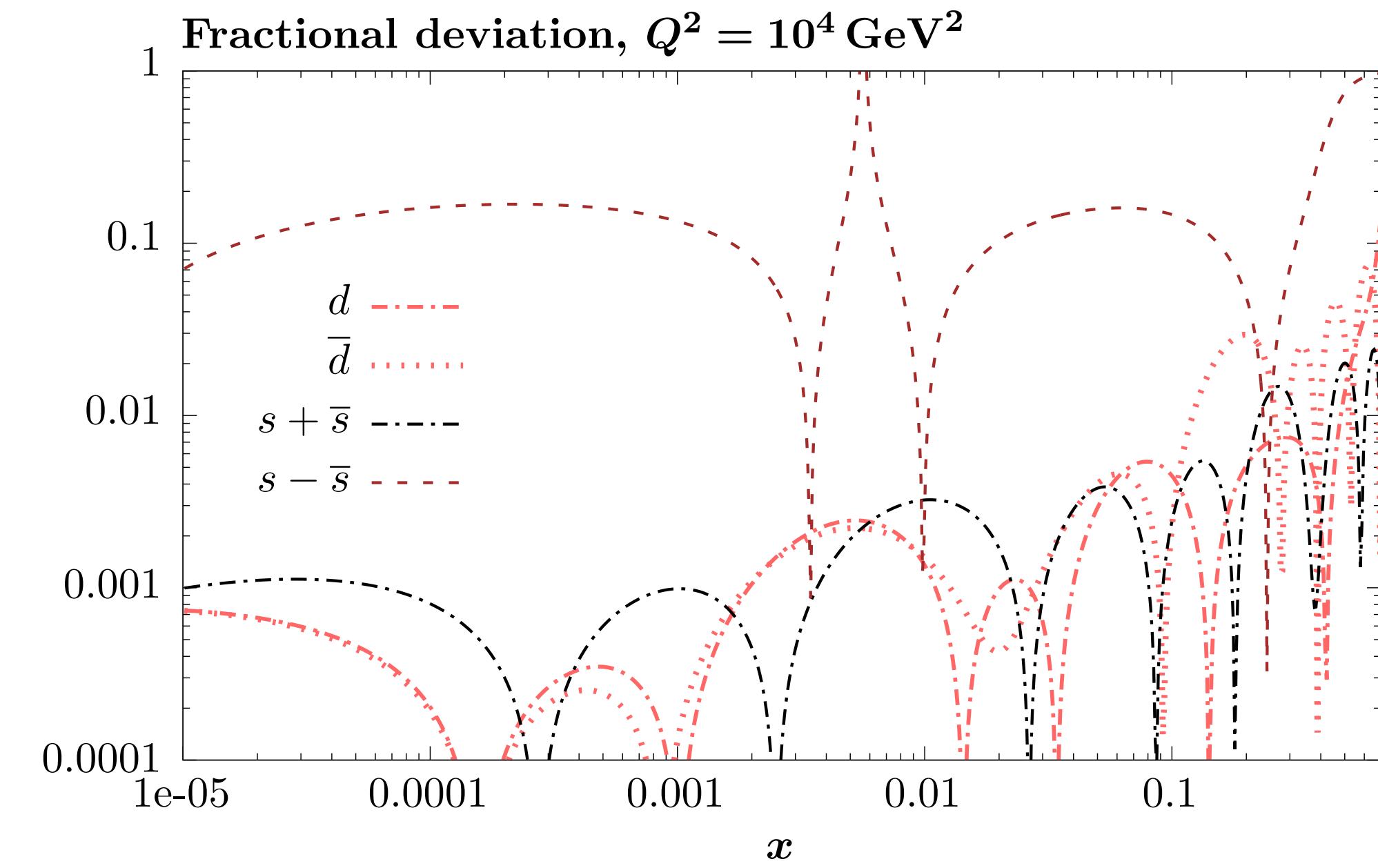
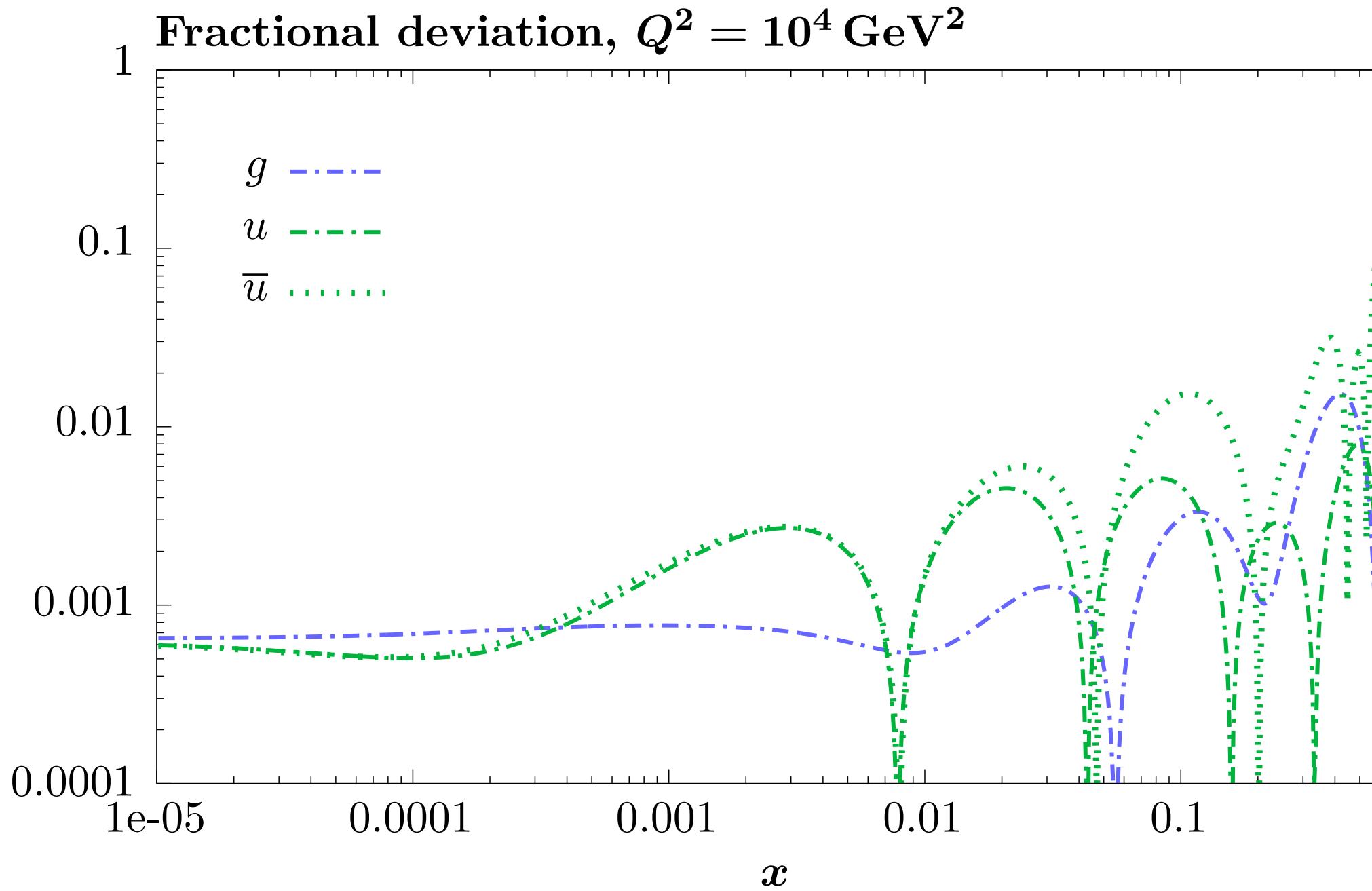
Summary

- Returning to the original possibilities (focus on MSHT as only considered here). We need to - and can
 - work out which is true:
 1. NNPDF4.0 uncertainty not conservative enough (too small).
 2. MSHT uncertainty too conservative (too large).
 3. ~~MSHT fit less accurate, due to parameterisation inflexibility, and hence enlarged errors needed (less precise).~~
- Successful closure test + comparison to NNPDF4.0 global fit suggests 3 is not dominant issue (at least in data region) for MSHT. Can be issue for less flexible ones and for MSHT in extrapolation regions.
- First direct comparison to NNPDF4.0 global fit finds that this gives inherently different (smaller) uncertainties than MSHT fixed parameterisation, keeping everything else equal.
- Put together, this implies that **either 1 or 2 is true** (or both). This study has not addressed which, though question of tolerance discussed elsewhere, but either way suggests more work needed.
- **Future steps:** extend to L1 closure, look again at question of tolerance in closure test framework...

Backup

Closure test - warm up

- Before considering global closure test, consider related question. Given PDF-level pseudodata, how closely can MSHT parameterisation match it Basic point: for LHC precision aim for sub-1% agreement.
- 500 PDF points logarithmically in $x \in \{10^{-5}, 0.99\}$ scattered by 1% uncertainty, for $u_V, d_V, S, s_+, s_-, g, \bar{d}/\bar{u}$
- Take NNPDF4.0 (p. charm) as input and plot fractional deviation. Find this is $\ll 0.01$ for most of the x region. Biggest deviations at high x and for s_- (as expected - MSHT parameterisation limited at moment).



- Encouraging, but rather artificial - really want to see how deviation compares in data region of global fit.

MSHT parameterisation

$$u_V(x, Q_0^2) = A_u(1-x)^{\eta_u} x^{\delta_u} \left(1 + \sum_{i=1}^6 a_{u,i} T_i(y(x)) \right)$$

$$d_V(x, Q_0^2) = A_d(1-x)^{\eta_d} x^{\delta_d} \left(1 + \sum_{i=1}^6 a_{d,i} T_i(y(x)) \right)$$

$$S(x, Q_0^2) = A_S(1-x)^{\eta_S} x^{\delta_S} \left(1 + \sum_{i=1}^6 a_{S,i} T_i(y(x)) \right)$$

$$s_+(x, Q_0^2) = A_{s+}(1-x)^{\eta_{s+}} x^{\delta_S} \left(1 + \sum_{i=1}^6 a_{s+,i} T_i(y(x)) \right)$$

$$g(x, Q_0^2) = A_g(1-x)^{\eta_g} x^{\delta_g} \left(1 + \sum_{i=1}^4 a_{g,i} T_i(y(x)) \right) + A_{g-}(1-x)^{\eta_{g-}} x^{\delta_{g-}}$$

$$s_-(x, Q_0^2) = A_{s-}(1-x)^{\eta_{s-}} (1-x/x_0) x^{\delta_{s-}}$$

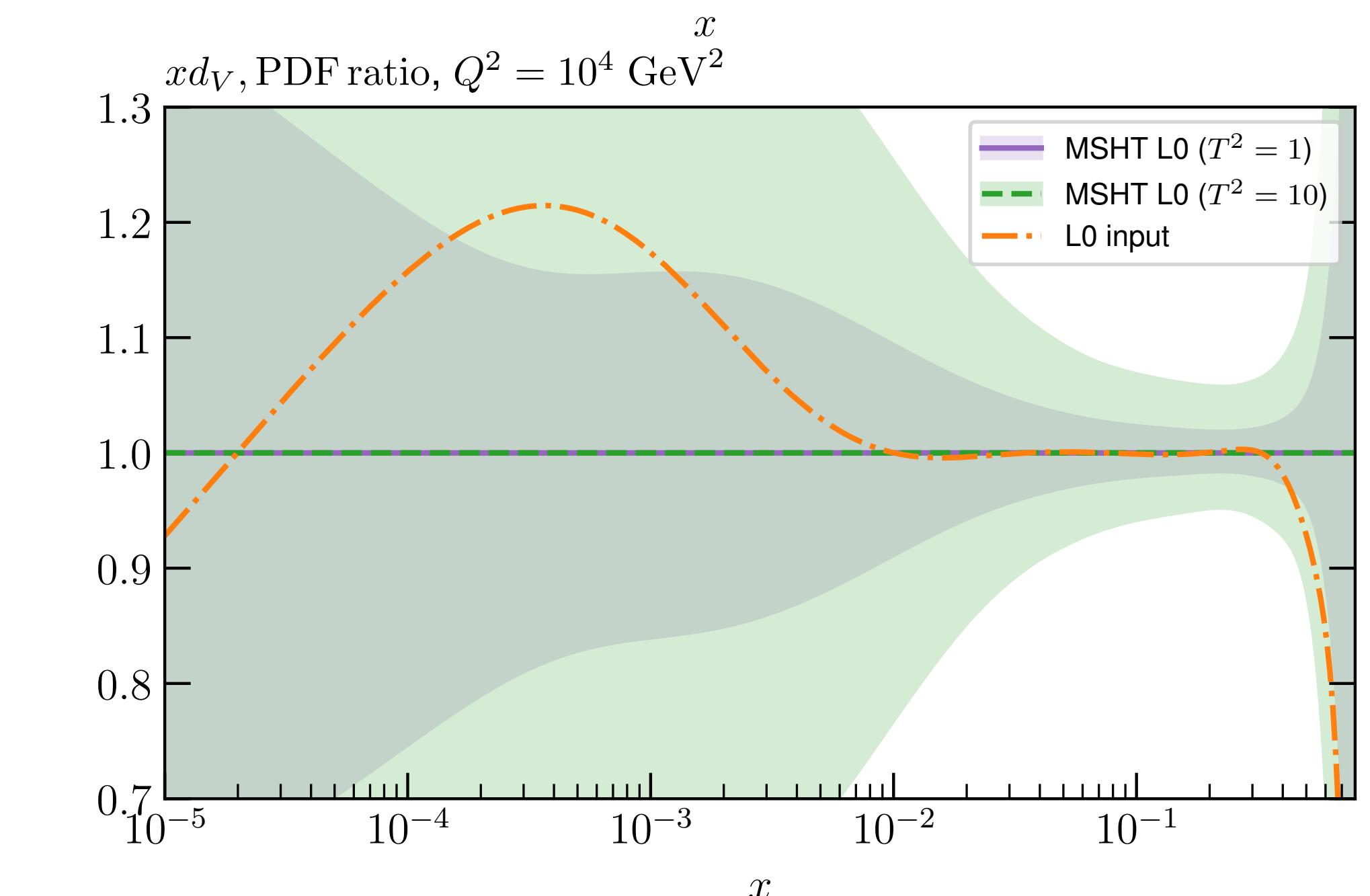
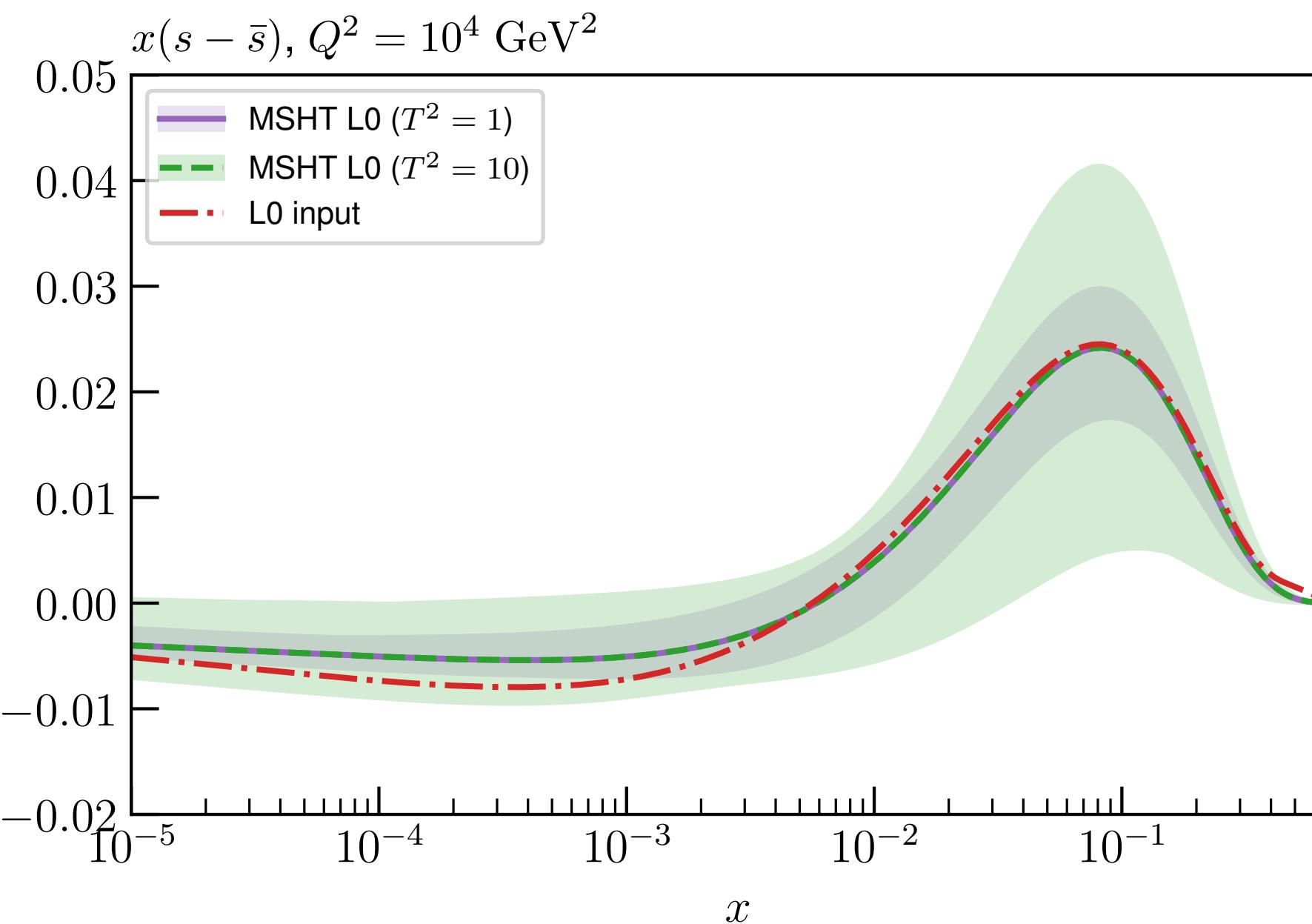
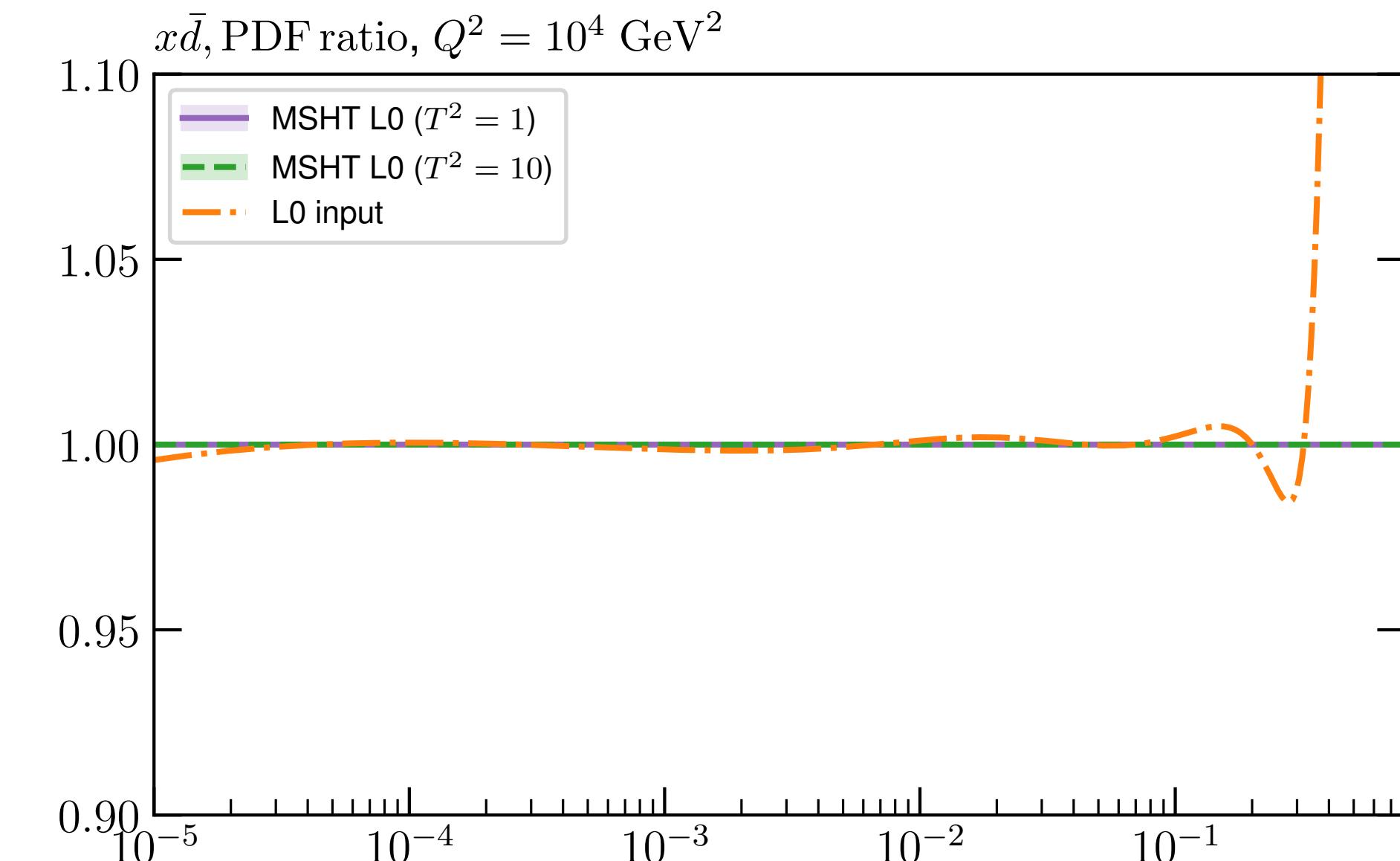
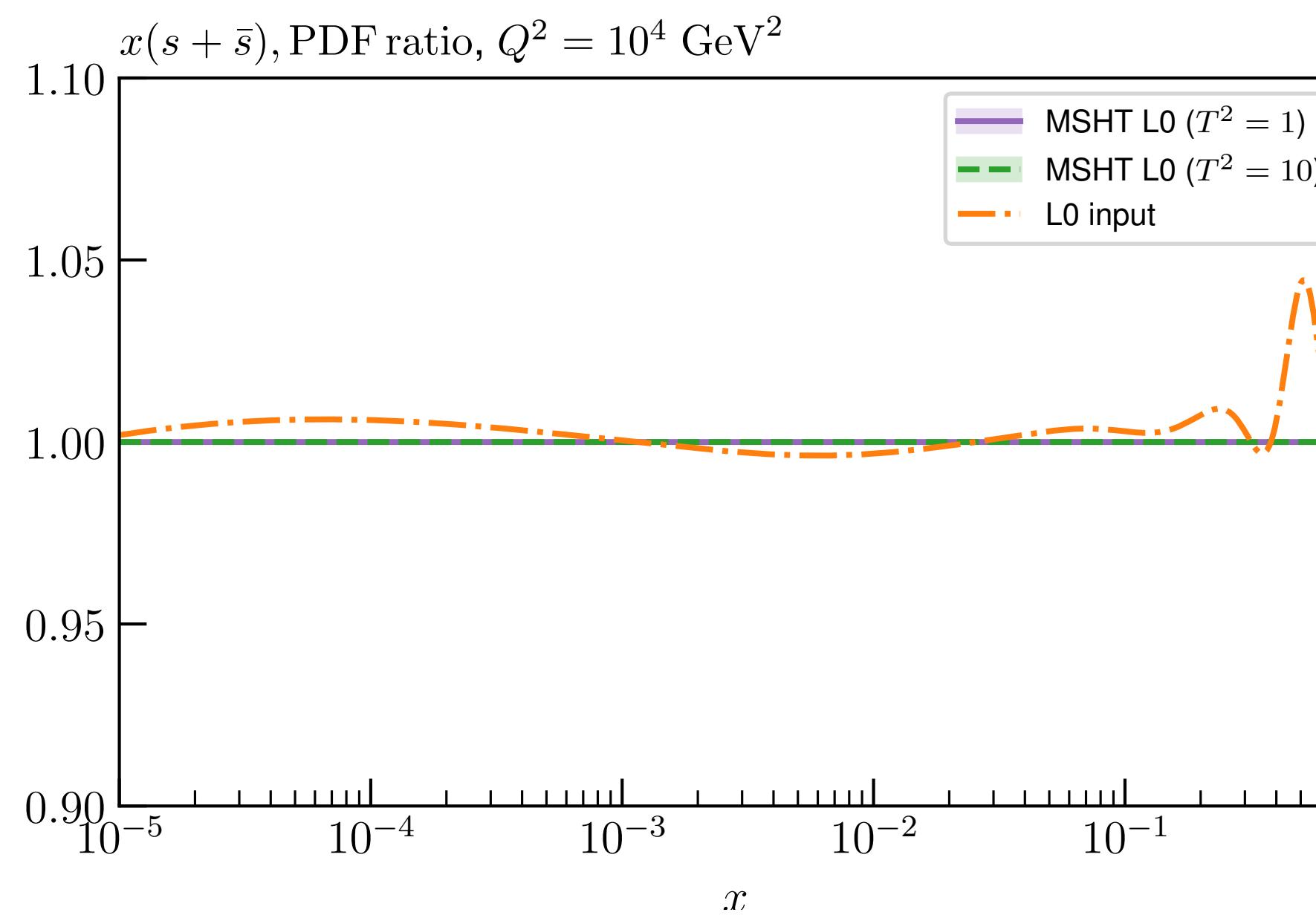
$$(\bar{d}/\bar{u})(x, Q_0^2) = A_\rho(1-x)^{\eta_\rho} \left(1 + \sum_{i=1}^6 a_{\rho,i} T_i(y(x)) \right)$$

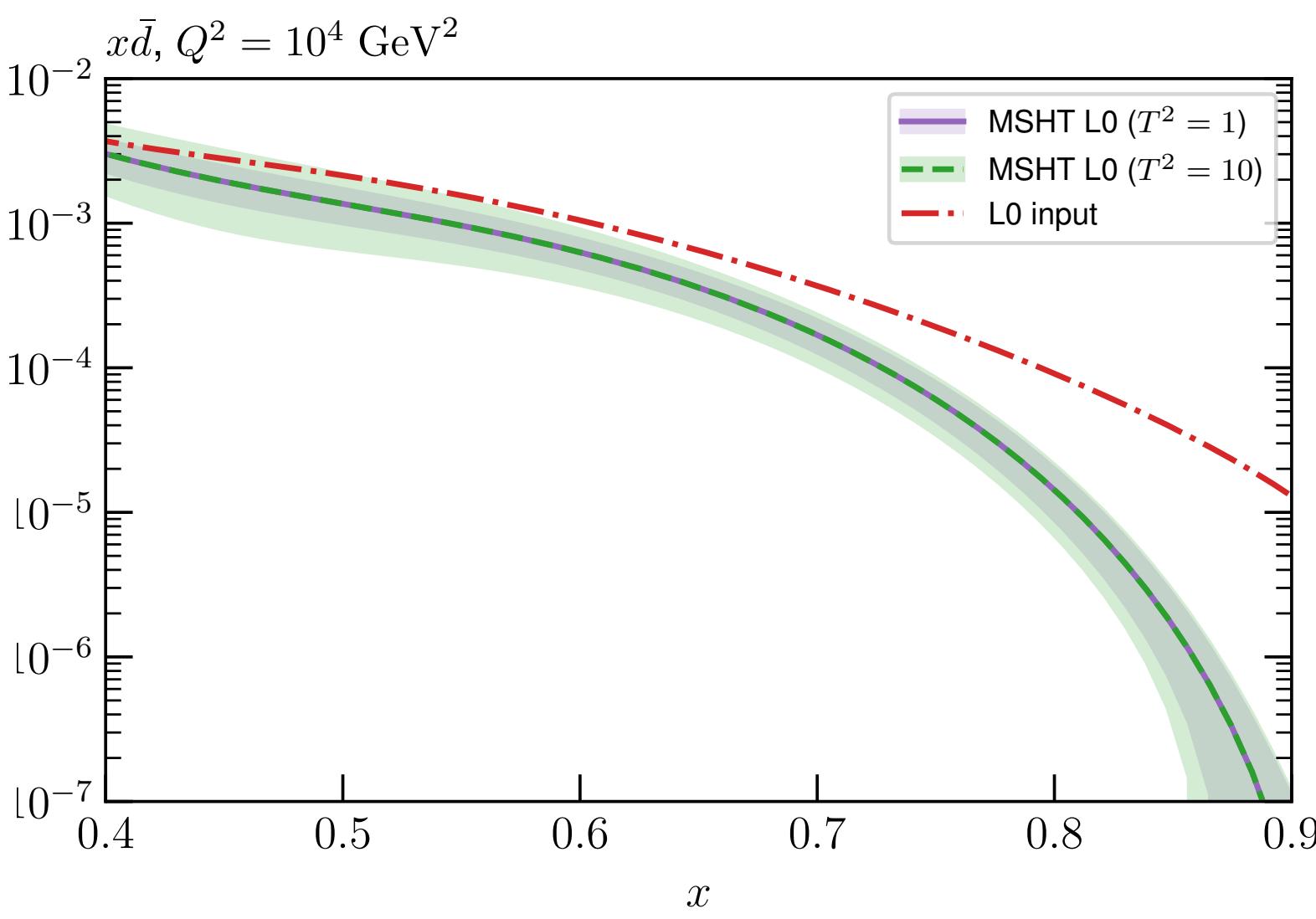
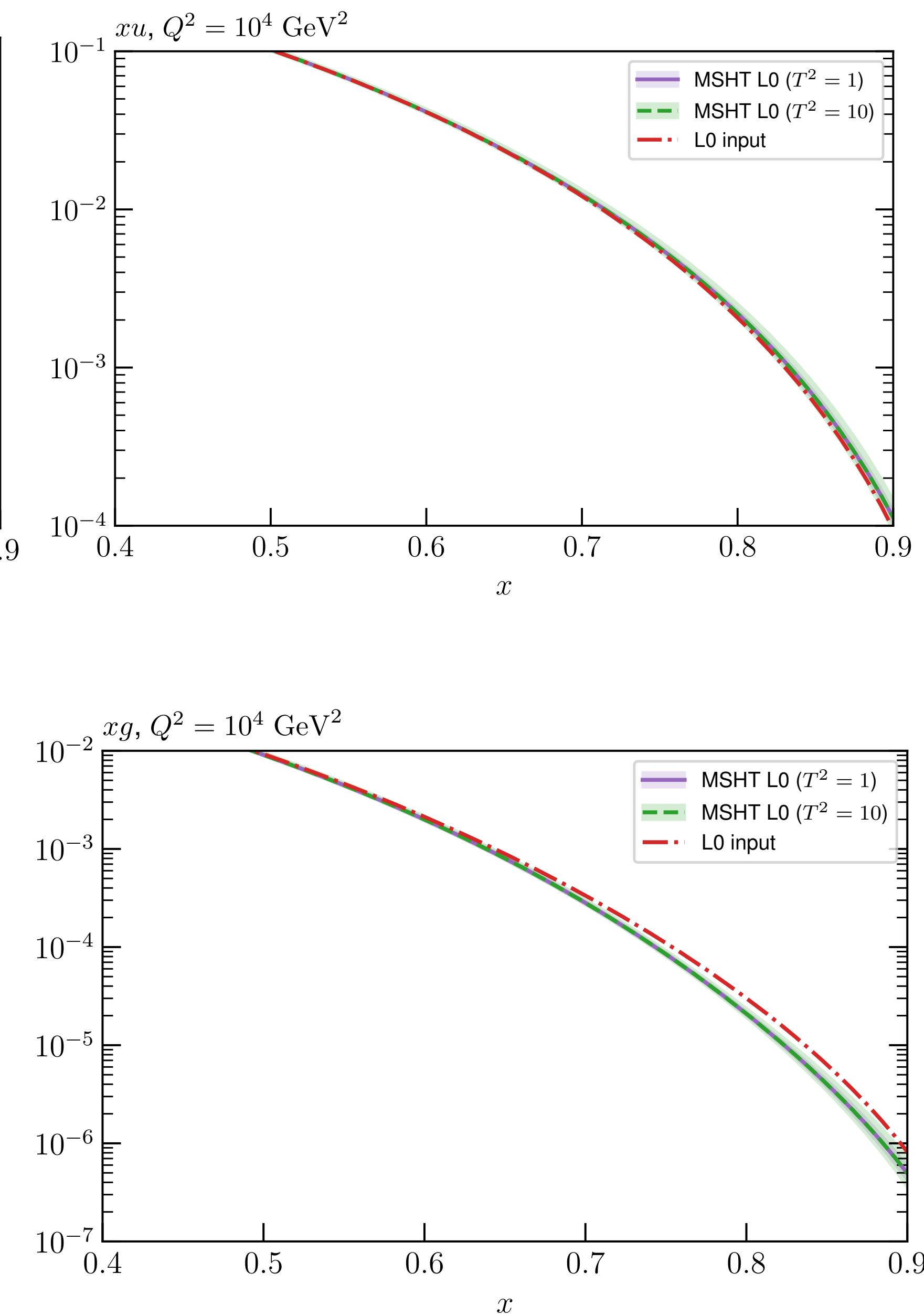
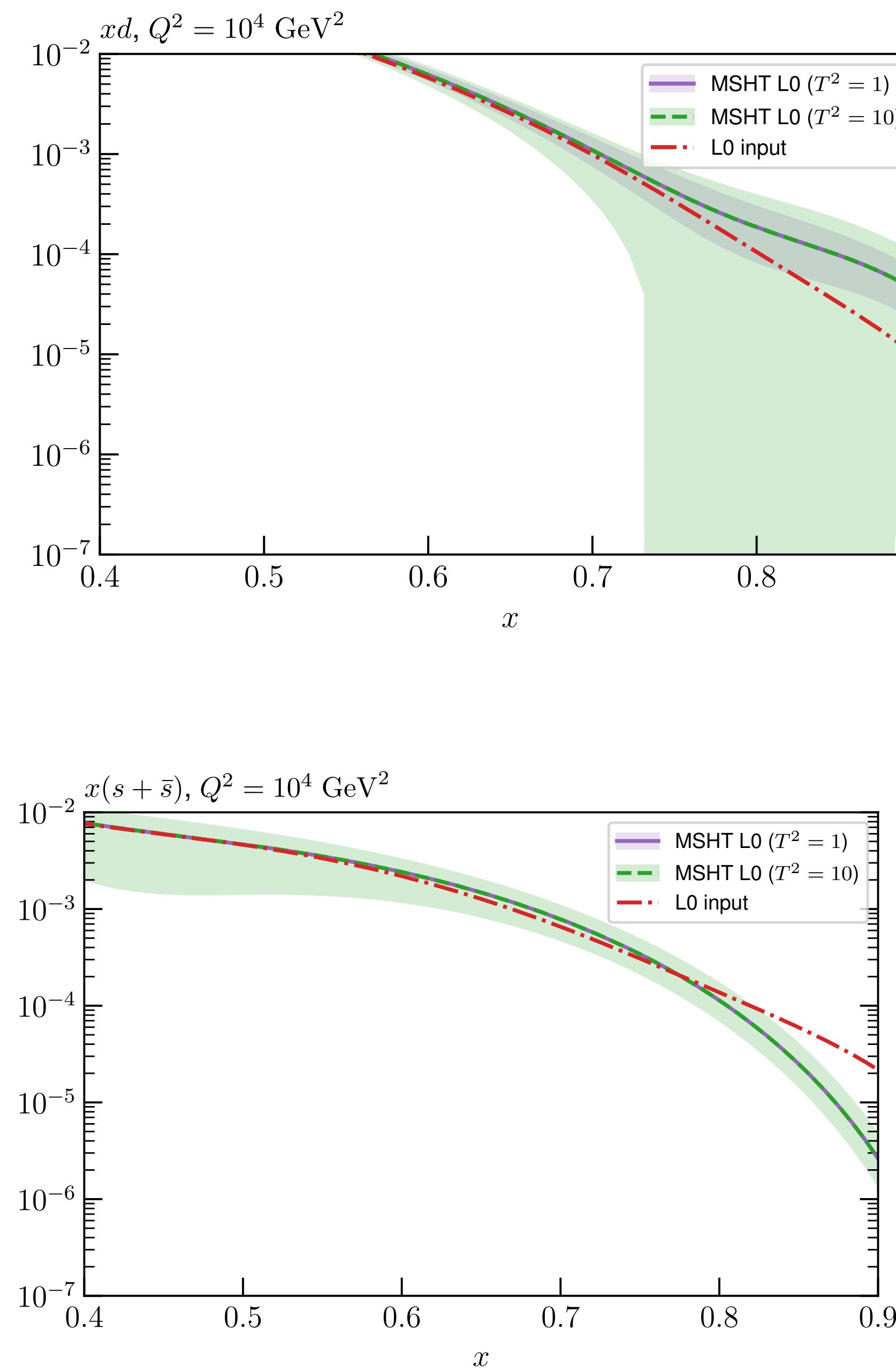
$$y(x) = 1 - 2\sqrt{x}$$

$$S(x) = 2(\bar{u}(x) + \bar{d}(x)) + s(x) + \bar{s}(x)$$

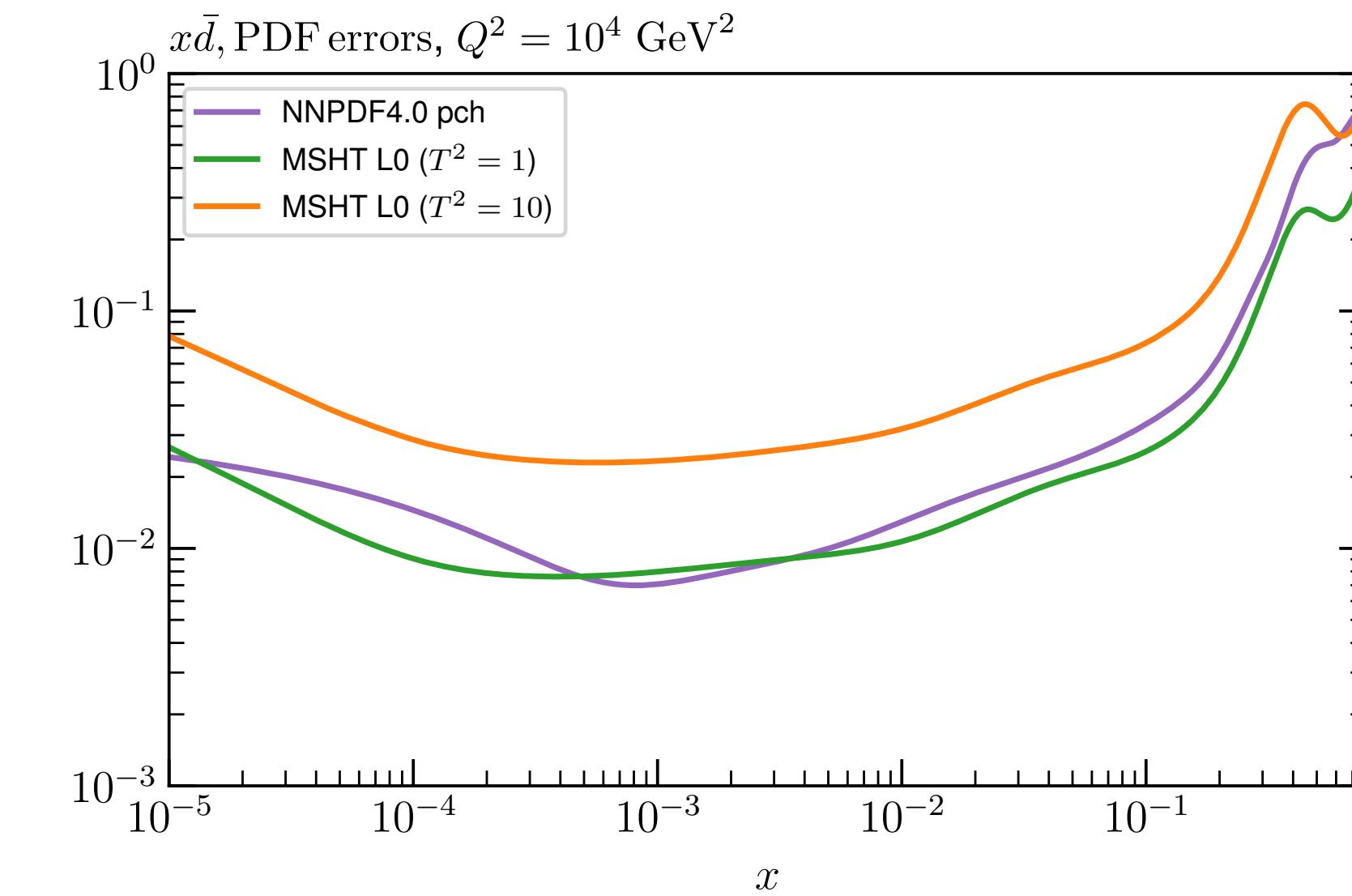
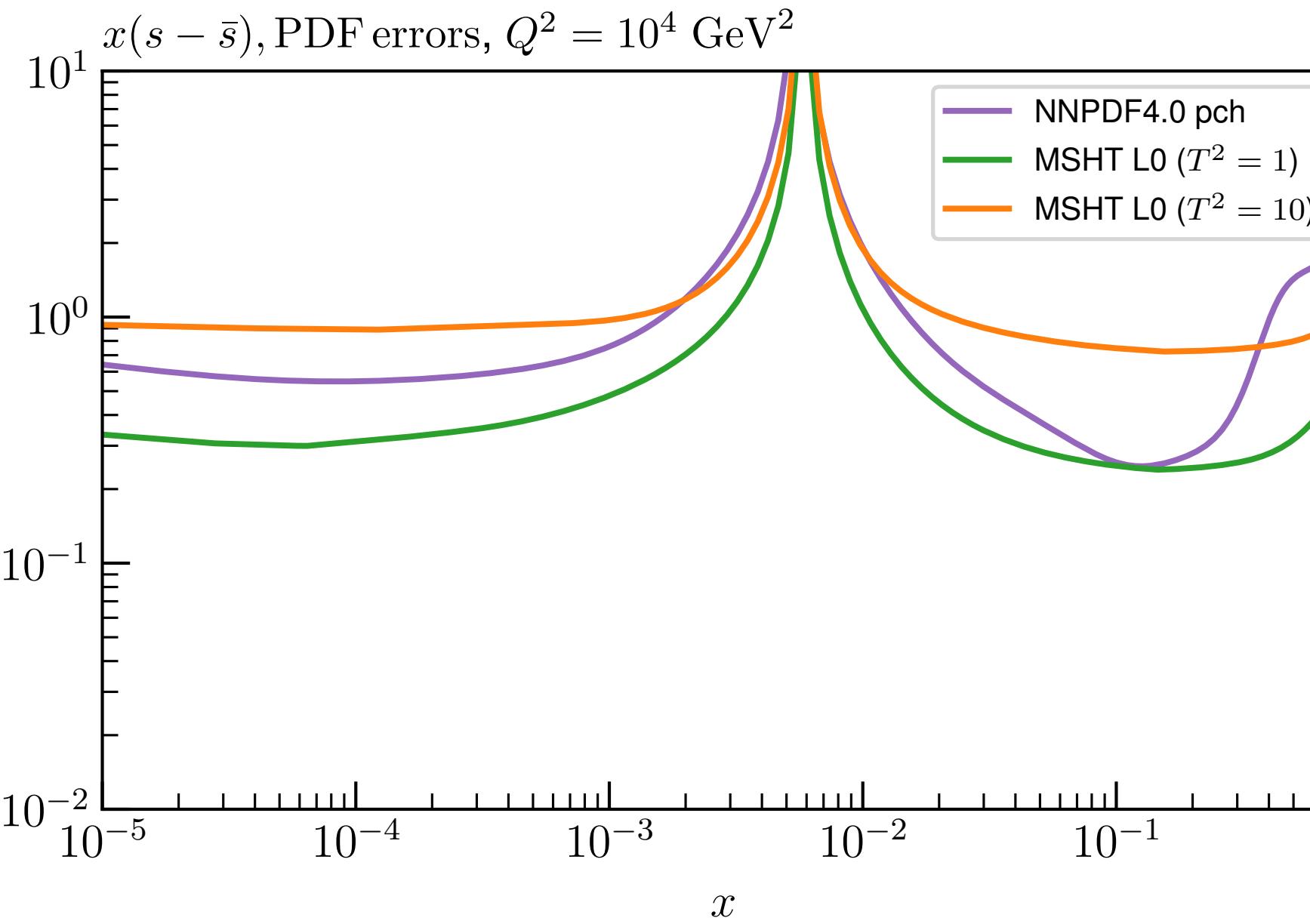
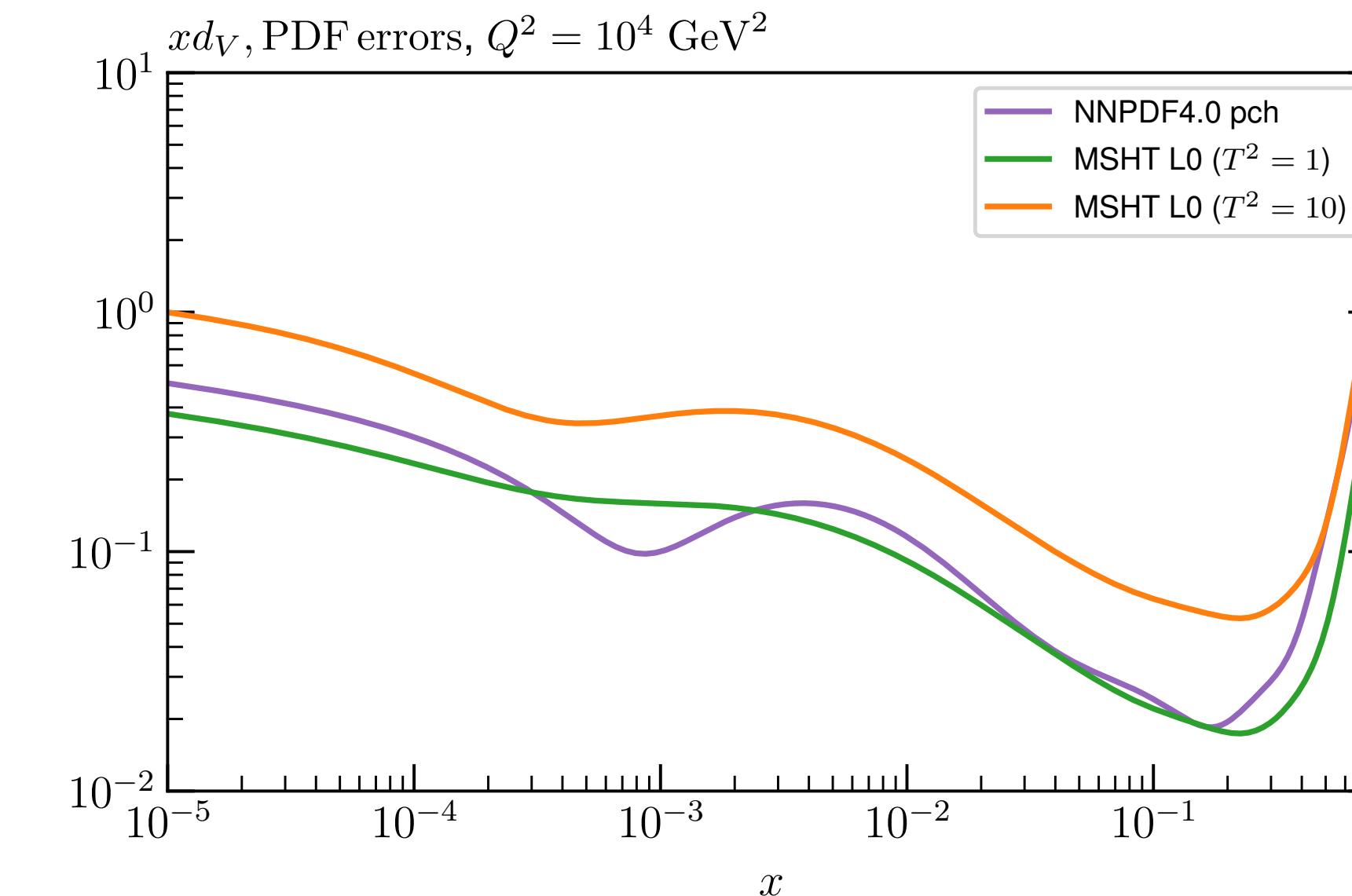
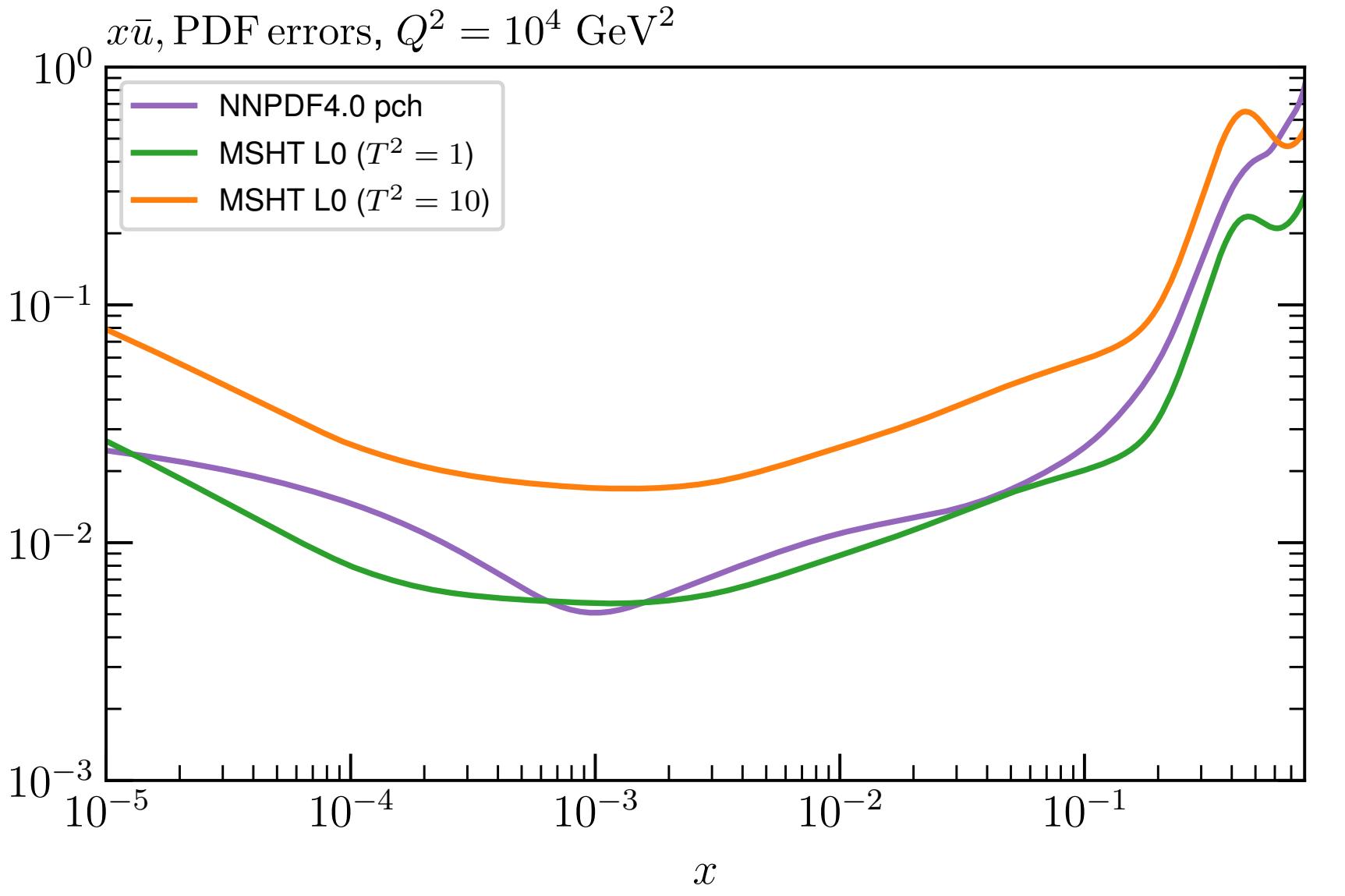
T_i : Chebyshev Polynomials

Global Closure PDFs





Comparison to NNPDF errors - Closure



Positivity

- We take:

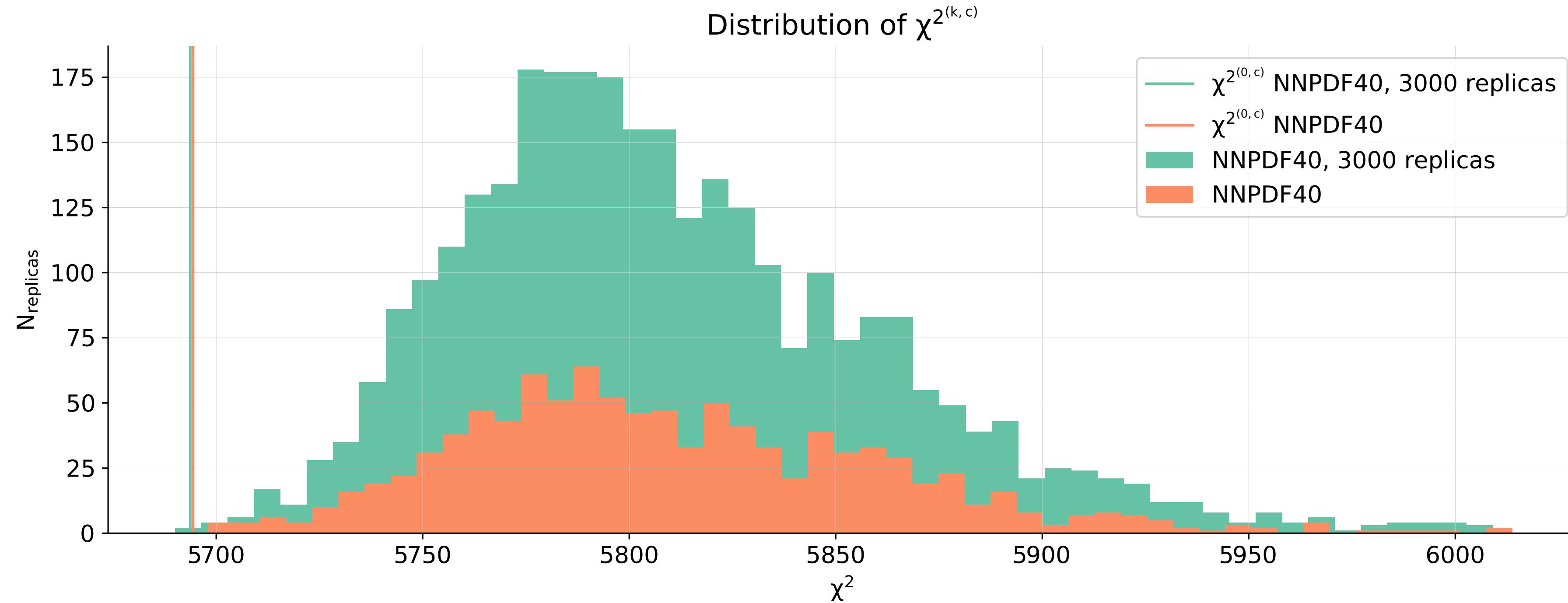
$$\chi_{\text{tot}}^2 \rightarrow \chi_{\text{tot}}^2 + \sum_{k=1}^8 \Lambda_k \sum_{i=1}^{n_i} \text{Elu}_\alpha \left(-\tilde{f}_k(x_i, Q^2) \right), \quad \text{with} \quad \Lambda_k = 10^3 \\ x_i \in \{5 \cdot 10^{-7}, 0.9\}$$

$$\text{Elu}_\alpha(t) = \begin{cases} t & \text{if } t > 0 \\ \alpha(e^t - 1) & \text{if } t < 0 \end{cases},$$

See NNPDF,
arXiv:2109.02653

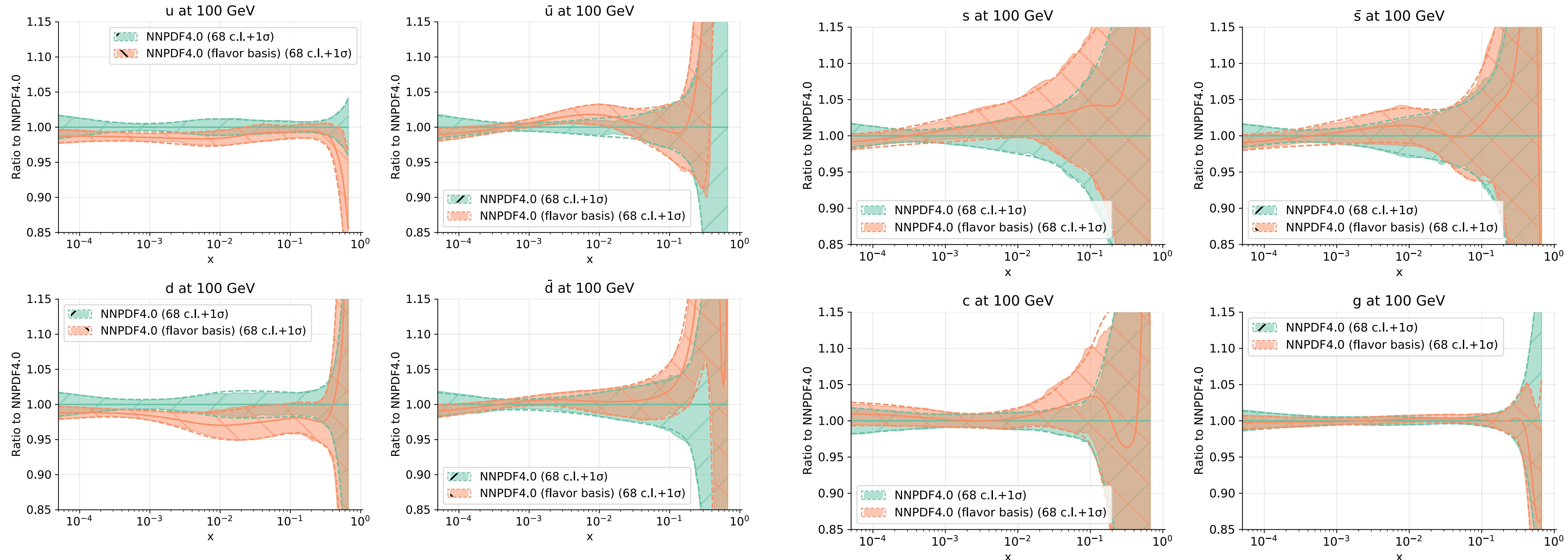
- And similarly for cross section constraints.

NNPDF chi² spread

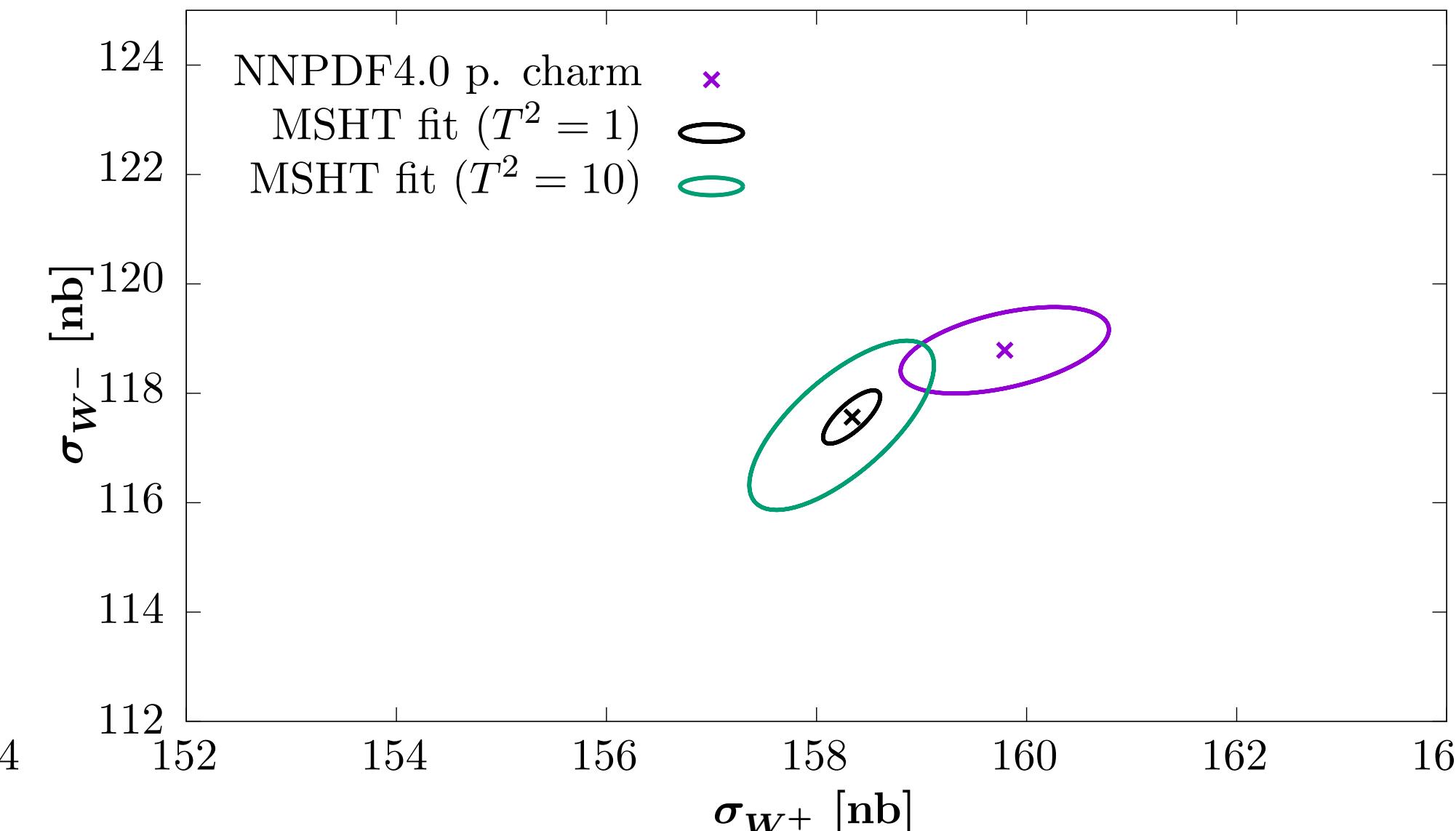
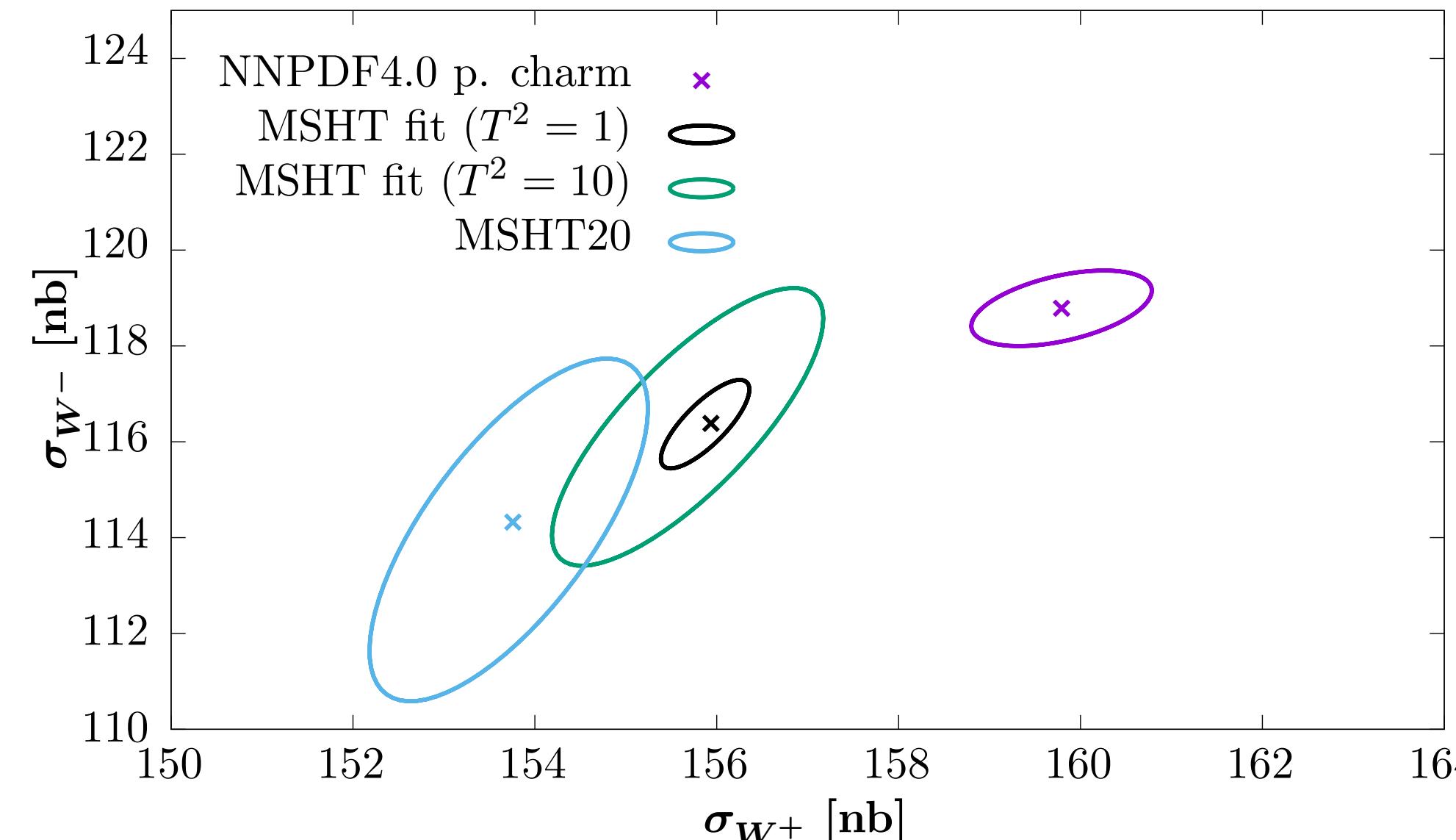
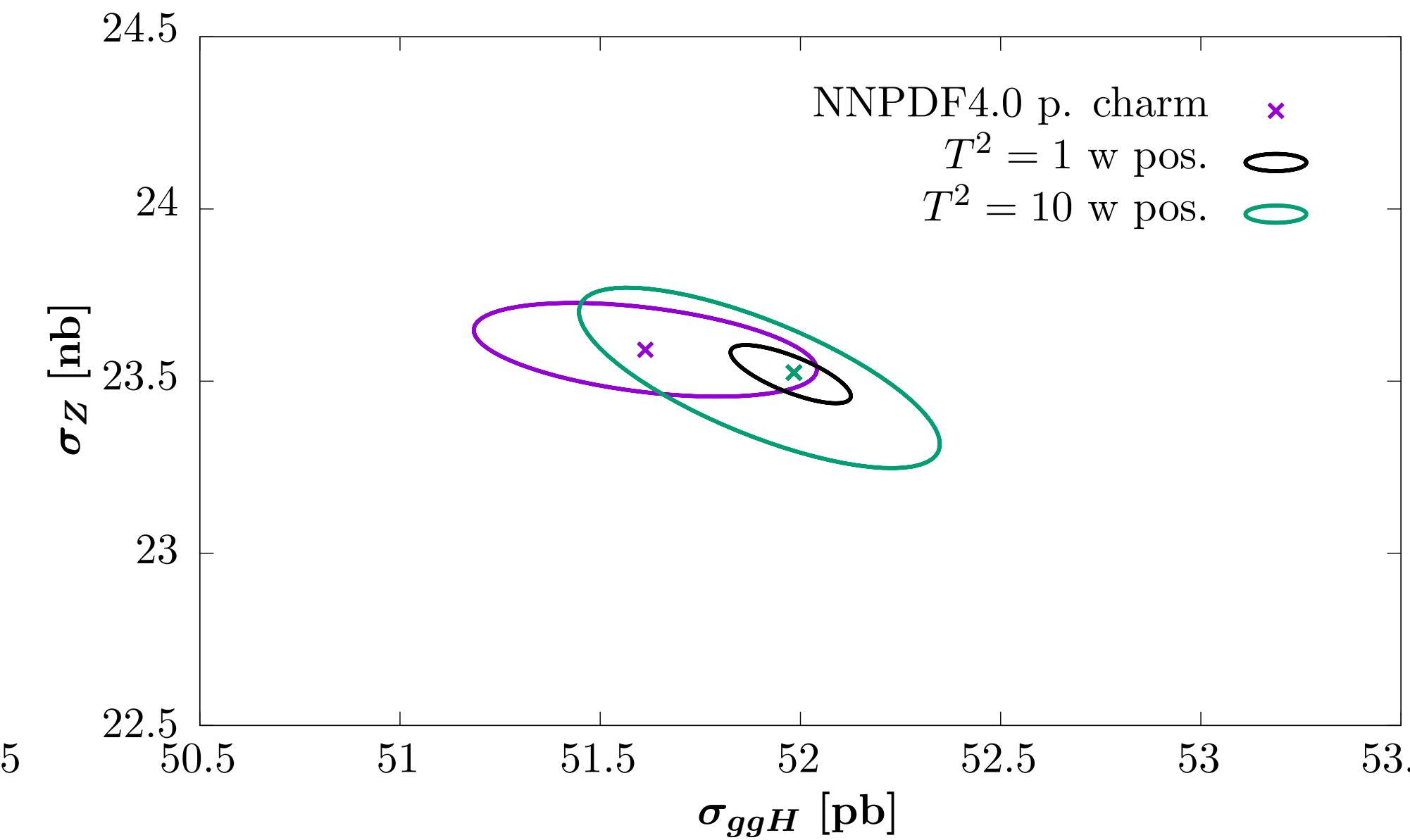
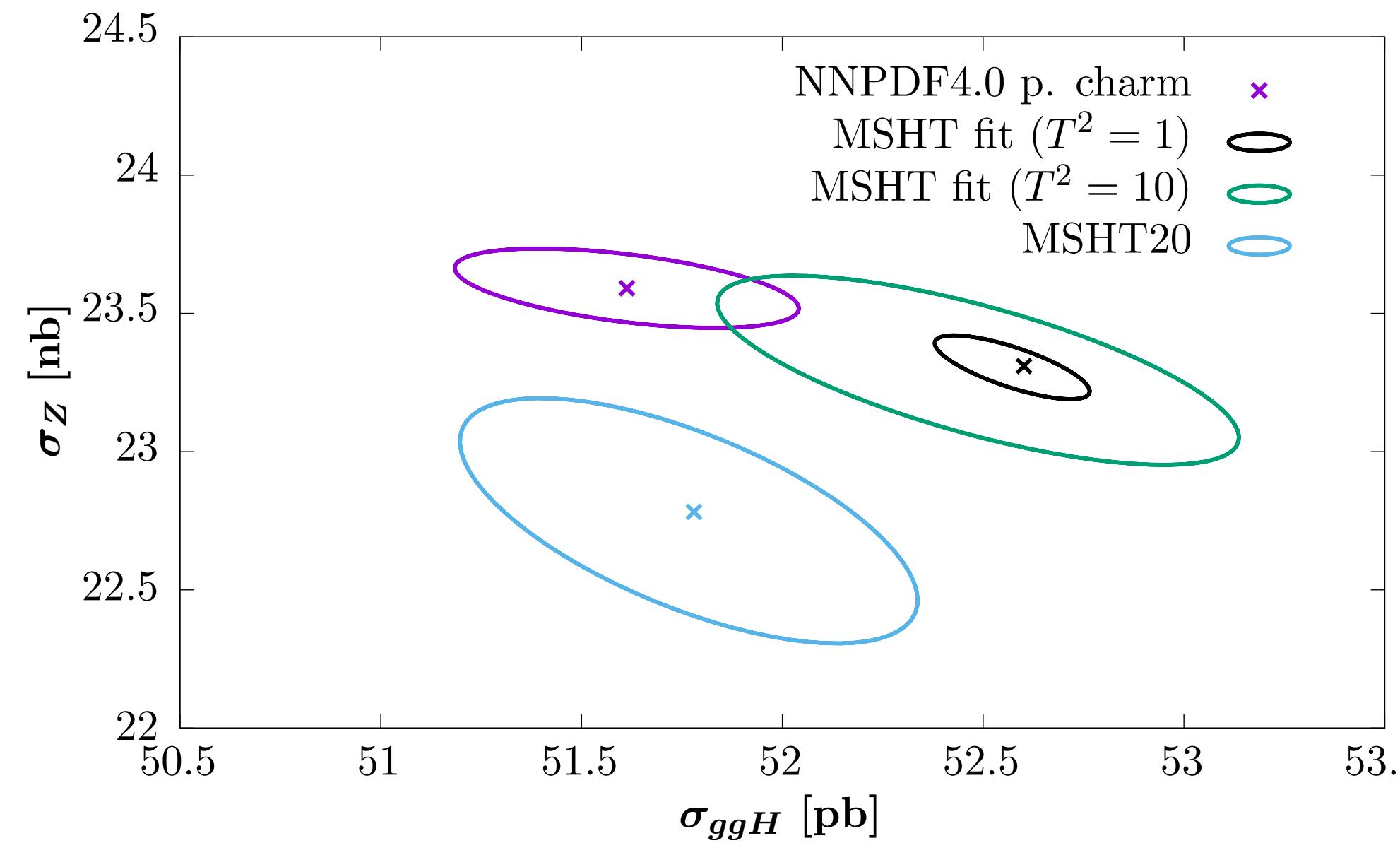


- In general do not expect central replica $\chi_{\text{rep},0}^2$ to be absolute minimum of χ^2 in NN approach due to overfitting regularization and statistical nature of replica ensemble. However former seems ruled out here (fixed parameterization) and latter should only give a handful of points lower.

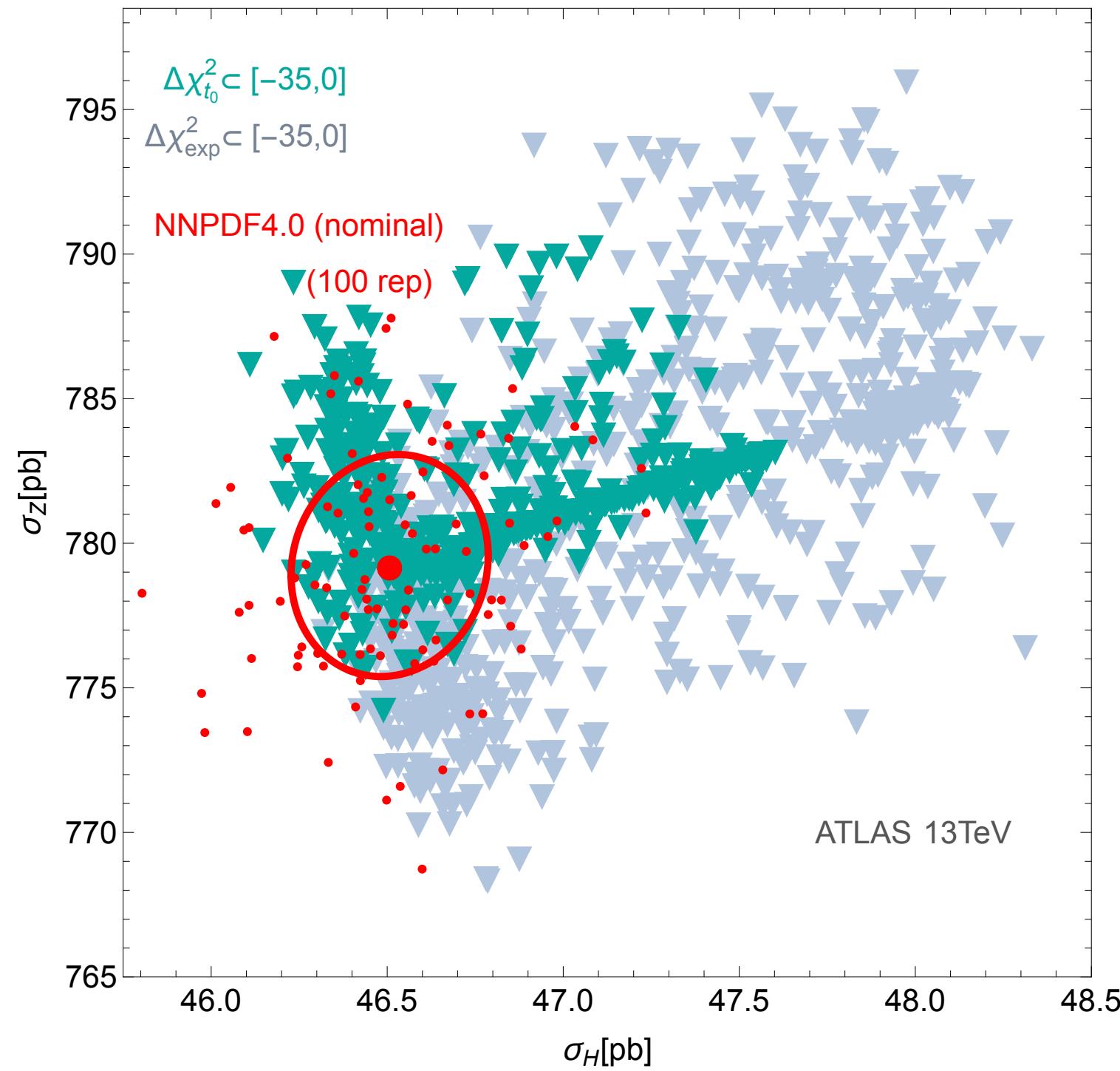
NNPDF flavour basis



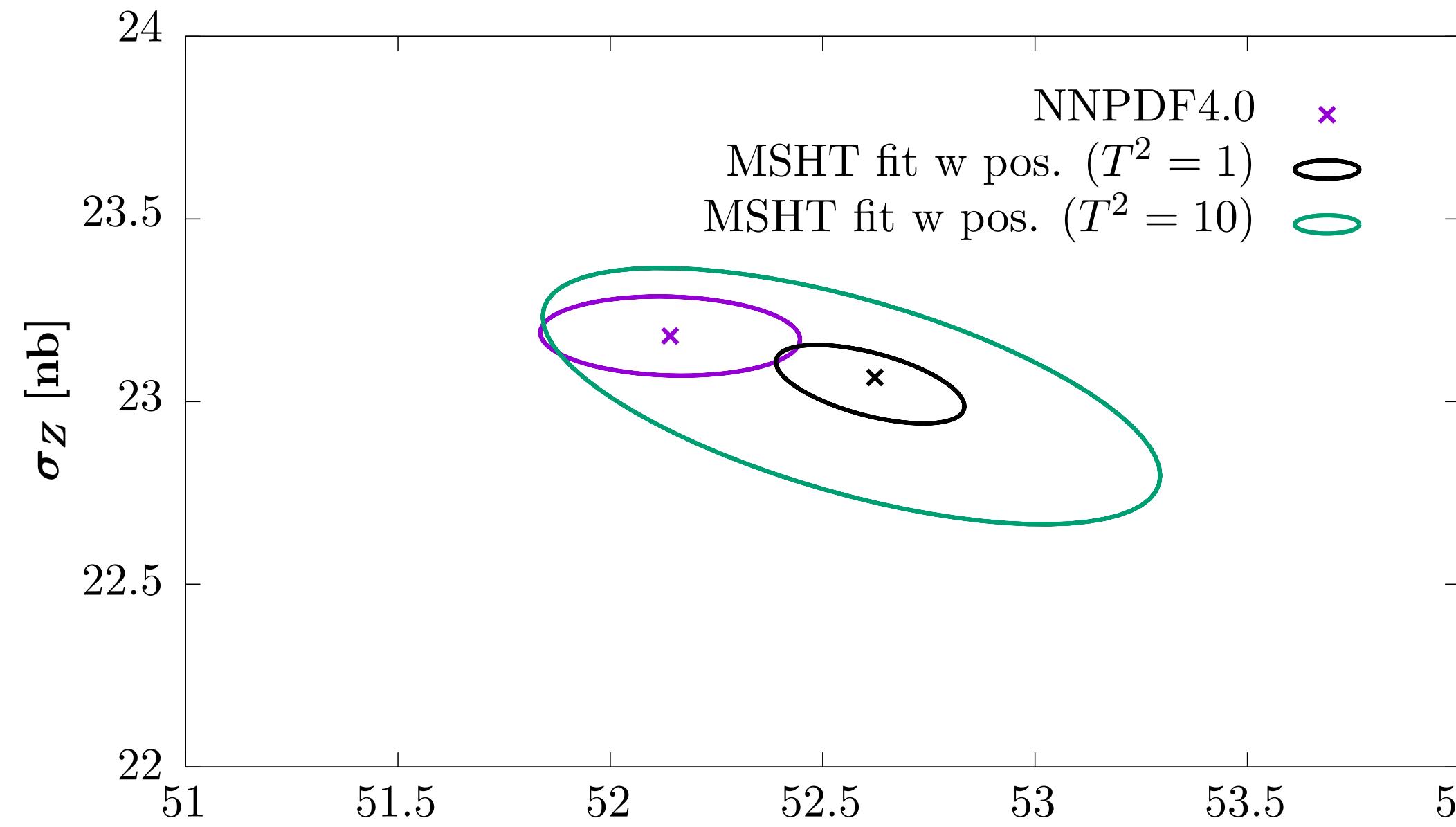
Cross Sections - p. charm



Comparison to hopscotch



A. Courtoy et al., arXiv: 2205.10444

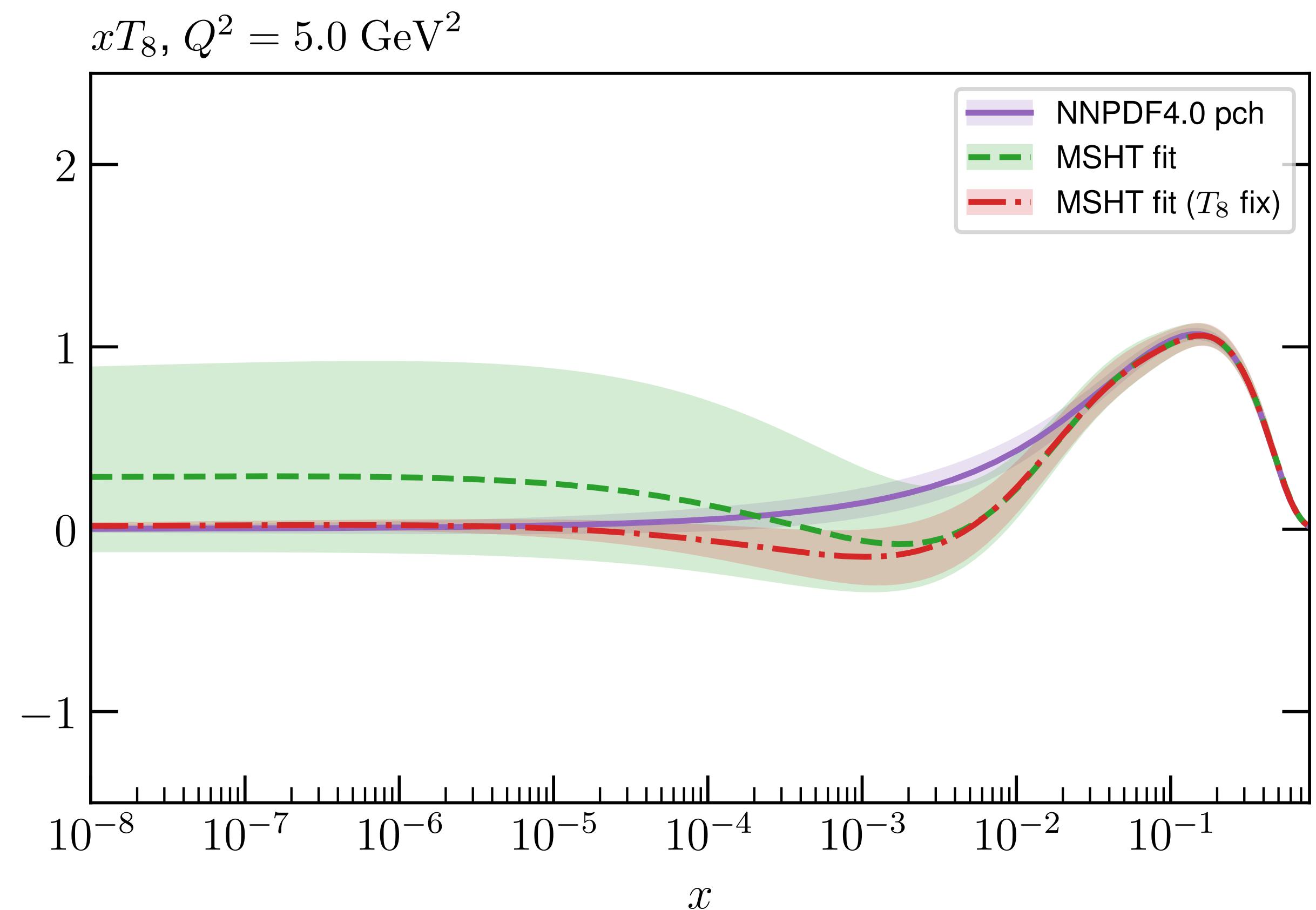


Integrability

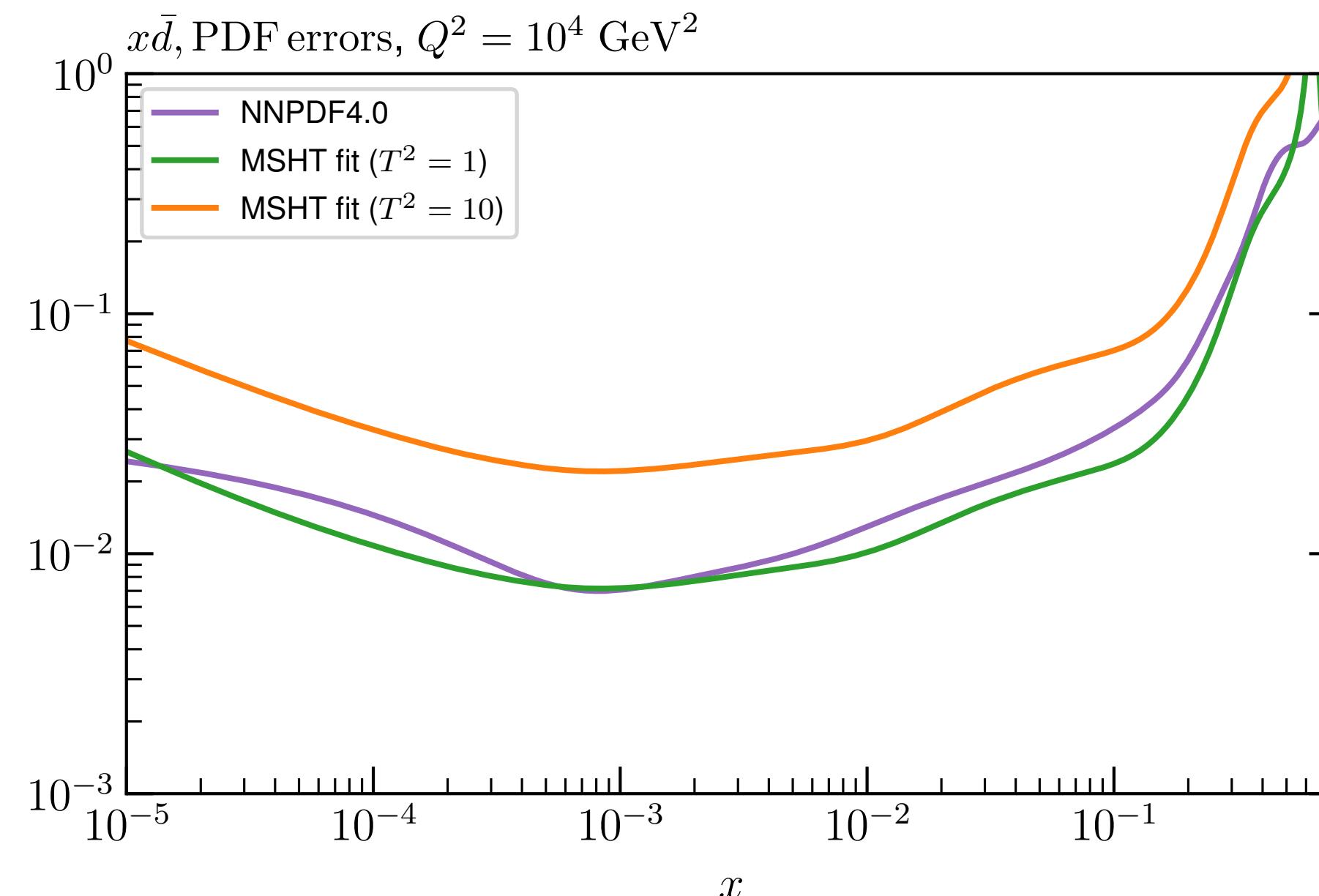
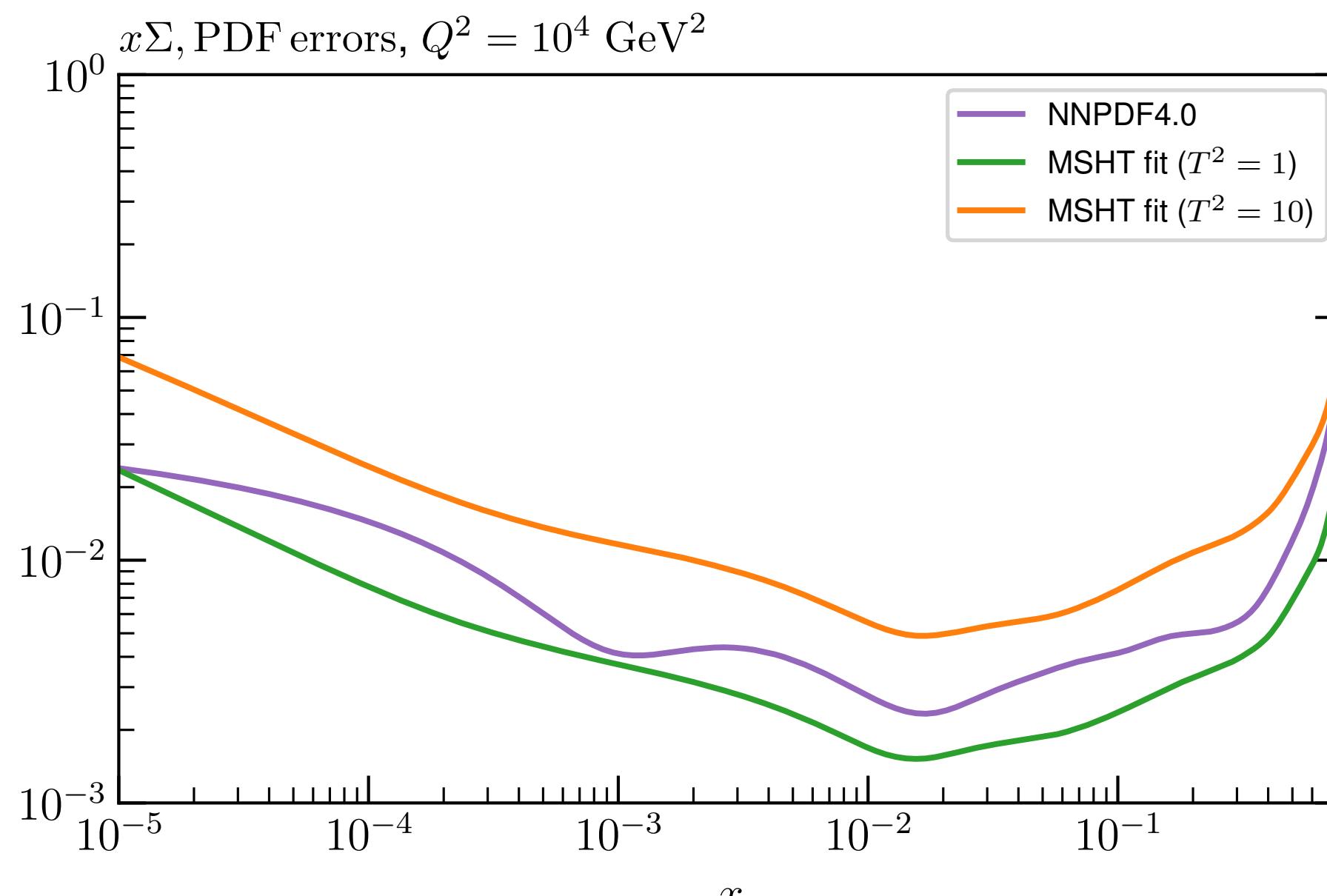
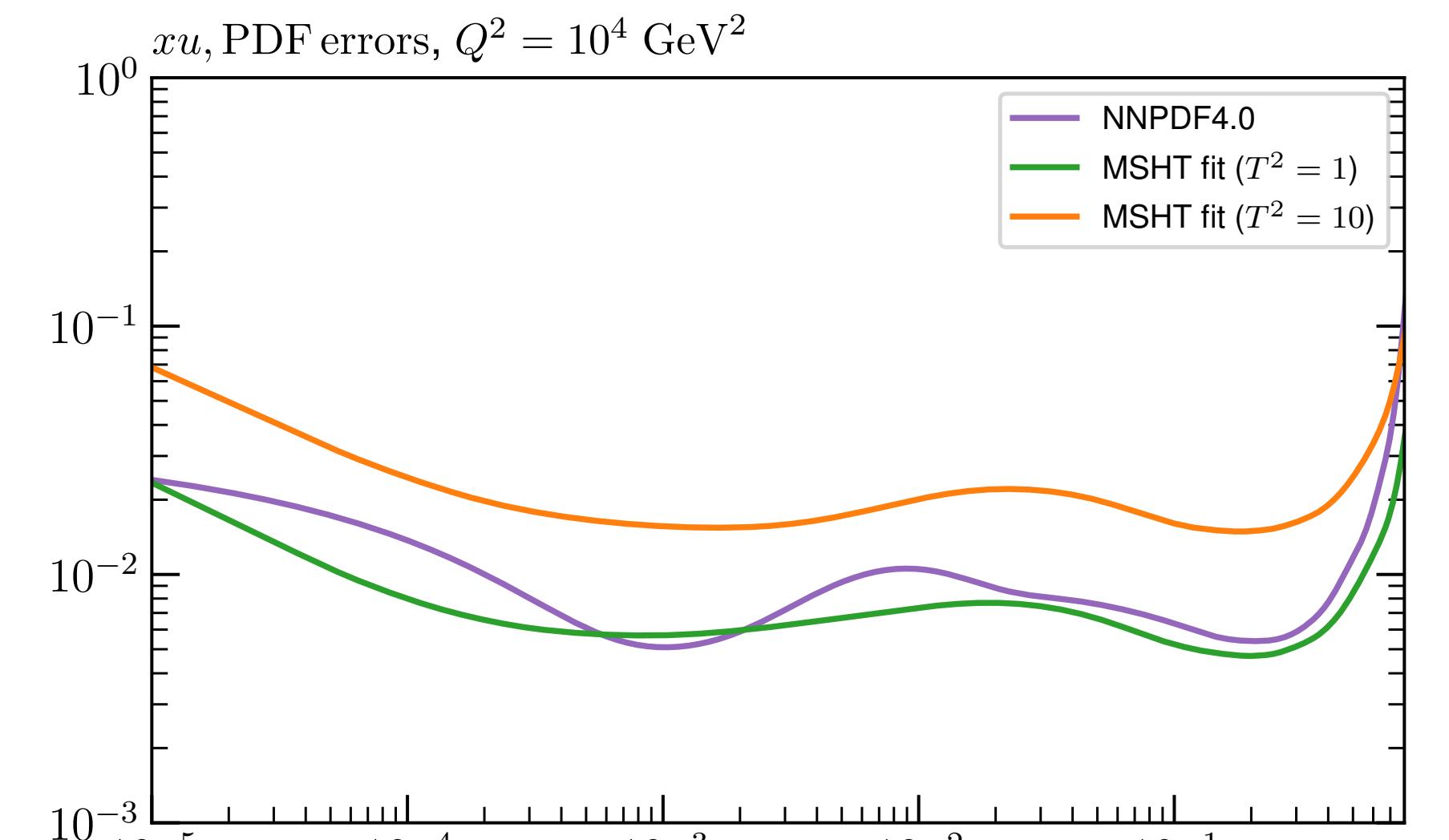
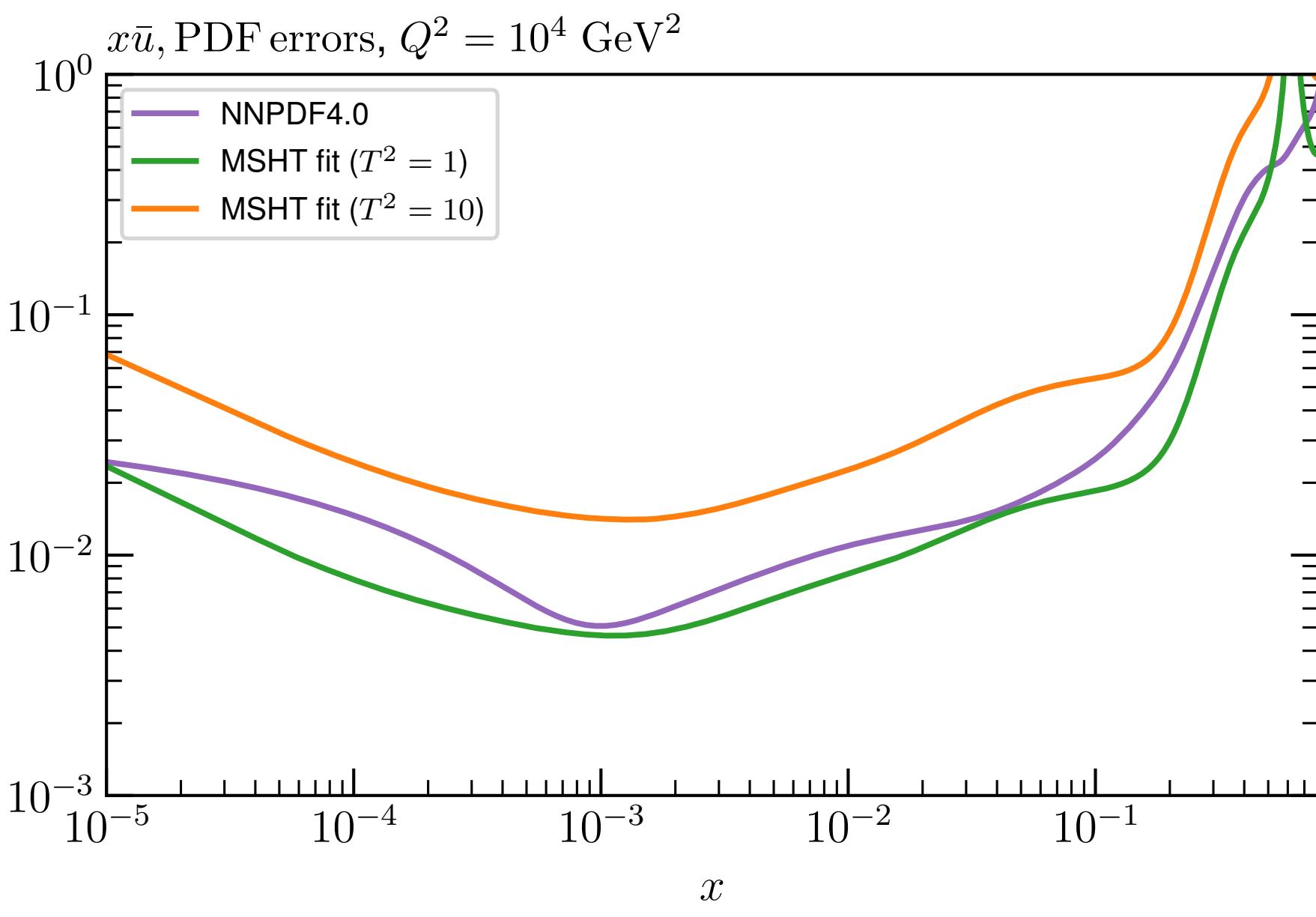
$$\lim_{x \rightarrow 0} x f_k(x, Q) = 0, \quad \forall Q, \quad f_k = T_3, T_8.$$

$$\chi_{\text{tot}}^2 \rightarrow \chi_{\text{tot}}^2 + \sum_k \Lambda_k^{(\text{int})} \sum_{i=1}^{n_i} \left[x f_k \left(x_{\text{int}}^{(i)}, Q_i^2 \right) \right]^2, \quad x_{\text{int}}^{(i)} = 10^{-9}, 10^{-8}, 10^{-7}. \quad \Lambda_k^{(\text{int})} = 100$$

- Biggest deviation from this for MSHT is for the T_8 combination.
- In MSHT fixed parameterisation can impose that this vanishes at low x by simply fixing strangeness normalization.
- Gives \sim same fit quality for p. charm, and ~ 3 points worse for fitted.



Uncertainties - p. charm



Uncertainties - fitted charm

