Partonic collinear structure by quantum computing

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State-of-art determination of collinear PDFs

Experimental data based global analysis



Talk by Lucian Harland-Lang on Monday

Equal time lattice calculation



Talk by Huey-Wen Lin on Friday



A new method for first principle calculation



Sign problem in Monte Carlo sampling of distribution and real time evolution

$$\langle \mathcal{O} \rangle = \frac{\int D\phi(x) e^{-S_E[\phi]} \mathcal{O}}{\int D\phi(x) e^{-S_E[\phi]}}$$

quantum computing: each gate is the result of time evolution

$$U(t) |\psi\rangle = e^{-it(H_0 + H_1 + \dots)} |\psi\rangle$$



Simulate hadron partonic structure on quantum computer Operator definition of quark PDF $f_{q/p}(x) = \int_{-\infty}^{\infty} \frac{dy^{-}}{2\pi} e^{ixp^{+}y^{-}} \langle p \,| \,\bar{\psi}(0) \frac{\gamma^{+}}{2} \mathcal{W}(0, y^{-}) \psi(y^{-}) \,| \, p \rangle$ $y^{-} = (y_0 - y_3)/\sqrt{2}$

PDFs are extremely challenge to simulate directly in Euclidean lattice calculation, due to multidimensional oscillating integral.

QC can naturally simulate real-time dynamics.

We are far from QCD Quantum Supremacy, start from a toy model for proof of concept study







- Talk by Christian Bauer on Friday



Simulate hadron partonic structure on quantum computer A toy model - 1+1D NJL (Gross, Neveu, 1974), no gauge field $\mathcal{L} = \bar{\psi}_{\alpha} (i\gamma^{\mu}\partial_{\mu} - m_{\alpha})\psi_{\alpha} + g(\bar{\psi}_{\alpha}\psi_{\alpha})^2$

$$f(x) = \int dz^- e^{-ixM_h z^-} \langle h | \bar{\psi}(z^-) \gamma^+ \psi(0) | h \rangle =$$

• Challenges in quantum computing $|h\rangle$ Map QFT to qubits+gates system

- Prepare the external hadronic state $|h\rangle$
- Evaluate the real-time dynamical correlation function
- Measurement of final observable



• Quantum field to qubits+gates $\mathscr{L} = \bar{\psi}(i\partial - m)\psi + g(\bar{\psi}\psi)^2$

• Discretization: staggered fermion, put different fermion components, flavors on different sites

$$\psi = \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix} \to \begin{pmatrix} \phi_{2n} \\ \phi_{2n+1} \end{pmatrix}$$

Jordan-Wigner transformation

$$\phi_n = \prod_{i < n} Z_i (X + iY)_n$$

• Discretized PDF:

$$f(x) \to \sum_{i,j} \sum_{z} \frac{1}{4\pi} e^{-ixM_{h}z} \langle h | e^{iHz} \phi_{-2z+i}^{\dagger} e^{-ixM_{h}z} \langle h | e^{iHz} \phi_{-2z+i}^{\dagger$$



$${}^{iHz}\phi_j |\,h
angle$$

 $\left[X_n Y_{n+1} - Y_n X_{n+1}\right]$



Hadron state preparation - VQE

- Hadron states are the eigenstates of the Hamiltonian with certain quantum numbers.
- 2103.08505 + ...
- Prepare the state by variational quantum eigensolver (VQE) VQE is a hybrid method involves both classical and quantum computers



show its power in quantum chemistry



Hadron state preparation - VQE

- I. Construct a trial hadronic state $|\psi_{lk}\rangle$, and a symmetry-preserving unitary operator $U(\theta)$
- II. The k-th state with quantum number l $|\psi_{lk}(\theta)\rangle = U(\theta) |\psi_{lk}\rangle$
- III. Optimization for hadronic state, minimize the cost function (PRL 113, 020505)

$$E_{l}(\theta) = \sum_{i=1}^{\kappa} w_{li} \langle \psi_{li}(\theta) | H | \psi_{li}(\theta) \rangle$$

IV. $|h\rangle = U(\theta^*) |\psi_{lk}\rangle$, θ^* is the optimized parameter set

Li et al (QuNu), PRD (letter, 2022)



Step II is carried out on quantum computer, all the others are computed on classical machine







Hadron state preparation

• Construct $U(\theta)$: quantum alternating operator ansatz (QAOA)

- I. Divide the hamiltonian, each term inherits the symmetries of H, $H = H_1 + H_2 + H_3 + H_4$
- II. $U(\theta)$ consists p layers, each layer evolve H_i with time duration θ_{ij} , $U(\theta) \equiv \prod_{i=1}^{P} \prod_{j=1}^{n} \exp(i\theta_{ij}H_j)$ $i=1 \ j=1$
- III. Prepare the input reference states for QAOA $|\psi_{\Omega,1}\rangle = |010101...01\rangle$ Naive vacuum $|\psi_{\Omega,2}\rangle =$

 $+...+|0101,...,10\rangle$ "quark pair" excitation





Simulate hadron partonic structure on quantum computer Evaluate the real-time dynamical correlation function $S_{mn}(t) = \langle h | e^{iHt} \Xi_m^3 \sigma_m^i e^{-iHt} \Xi_n^3 \sigma_n^J | h \rangle$

PDFs can be written as a sum of such correlation functions

Measure the observable with one auxiliary qubit

Measure the ancillary qubit on X(Y) basis to get the real (imaginary) part of $S_{mn}(t)$

- $|\alpha\rangle_a |0\rangle_b \to \frac{\sqrt{2}}{2} |\alpha\rangle_a (|0\rangle_b + |1\rangle_b) \to |\phi\rangle \equiv$
- $\langle \phi | I_a \otimes X_b | \phi \rangle = \frac{1}{2} + Re(\langle \alpha | \hat{O} | \alpha \rangle)$ $\langle \phi | I_a \otimes Y_b | \phi \rangle = \frac{1}{2} Im(\langle \alpha | \hat{O} | \alpha \rangle)$

 $S_{mn}(t) = \langle h | \hat{O}_{mn}(t) | h \rangle \hat{O}_{mn}(t)$



$$\equiv \frac{\sqrt{2}}{2} (|\alpha\rangle_a |0\rangle_b + \hat{O} |\alpha\rangle_a |1\rangle_b)$$

$$P_{mn}(t) = e^{iHt} \Xi_m^3 \sigma_m^i e^{-iHt} \Xi_n^3 \sigma_n^j$$



Num

erical results from quantum computing						
Measurement of hadron mass $M_h = \langle h H h \rangle - \langle \Omega H \Omega \rangle$						
g	0.2	0.4	0.6	0.8	1.0	N - 12
$M_{h,\mathrm{QC}}a$	1.002	1.810	2.674	3.534	4.352	
$M_{h,{ m NUM}}a$	1.001	1.801	2.659	3.509	4.342	ma = 0.2

- Considering the current limitations of using real quantum devices, the results are generated using a classical simulation of the quantum circuit
- Measure the mass of the lowest-lying *ud*-like hadron in NJL model with 2 flavors, QAOA has good accuracy
- For small quark mass, the dominant contribution comes from the interaction rather than the quark masses
- For ma = 0.8, the quark masses are dominant



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Numerical results from quantum computing

quark PDF of the lowest-lying zero-charge hadron

quark PDF in position space

ma = 0.8 N = 18 $n_f = 1$

- The real part is consistent with 0 $f_q(x) = f_{\bar{q}}(x) = -f_q(-x)$
 - The bound state behavior

computing ro-charge hadron



 \mathcal{Z}

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Numerical results from quantum computing Li et al (QuNu), PRD (letter, 2022) quark PDF of the lowest-lying zero-charge hadron

- Good agreement between quantum computing and numerical diagonalization
- The non-vanishing contributions in the x > 1 are partly due to the finite volume effect
- We observe the expected peak around x = 0.5 and qualitative agreement with pion PDFs





Quantum computing for LCDA

to find a pair of valence $q\bar{q}$ inside a meson

LCDA is an essential ingredient in exclusive high-energy QCD processes, e.g. form factor in the process $\gamma^* \gamma \rightarrow \pi^0$

$$F(Q^2) = f_{\pi} \int_0^1 dx \, T_H(x, Q^2; \mu) \phi_{\pi}(x; q) dx$$

$$\phi(x) = \frac{1}{f} \int dz e^{-i(x-1)n \cdot Pz} \langle \Omega | \bar{\psi}(z) \rangle$$

The current knowledge on LCDA is limited, mainly on models and lattice calculations

First try using quantum computing

LCDA - light cone distribution amplitude, describes the probability amplitude









LCDA on quantum computer

Quantum circuit



Li et al (QuNu), SCPMA (2023)



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LCDA on quantum computer Numerical results



- hadron mass
- Converges to asymptotic result in weak coupling limit



peak gets narrower with decreasing coupling constant or increasing



Summary and outlook

- First attempt for direct computing of "PDFs"
 - 1. Use NJL model as a proof of concept study
 - 2. parton distribution function and light cone distribution amplitude
- The field is still at early age of development, many more need to be done
 - 1. gauge field encoding
 - 2. Extend to higher dimensions for TMDs and spin dependent processes
 - 3. quantum noise

Thanks for your attention!



