

# Dimuon production in neutrino-nucleus collisions - the SIDIS approach

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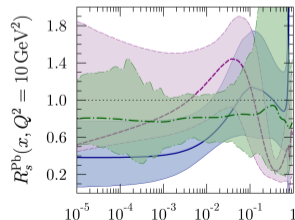
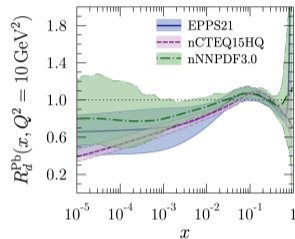
DIS 2024, 9.4.2024, Grenoble



Centre of Excellence  
in Quark Matter

## Strange-quark distribution

- ▶ Still quite poorly known, even for free protons
- ▶ Relevant for
  - ▶  $W$  and  $Z$  production at the LHC
  - ▶ precision determination of electroweak parameters such as  $W$  mass and weak mixing angle
- ▶ Mostly constrained by neutrino-nucleus data
  - ▶ Nuclear target needed for sufficient statistics
  - ▶ Plays an important role in free-proton fits  $\rightarrow$  correlation between proton and nuclear PDFs

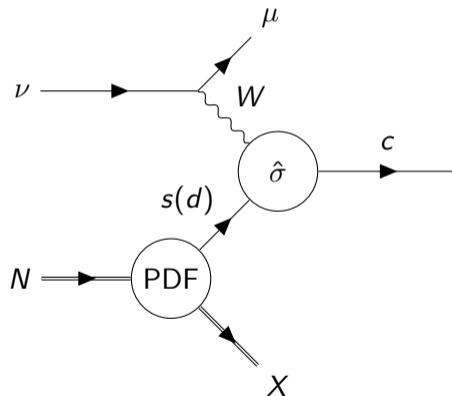


## Dimuon production in $\nu N$ collisions

- ▶ Charm production probes the strange-quark distribution
- ▶ Can't detect charm directly, so dimuons instead
- ▶ Usually computed by assuming factorization:

$$d\sigma(\nu N \rightarrow \mu\mu X) \simeq \mathcal{A} \mathcal{B}_\mu d\sigma(\nu N \rightarrow \mu c X)$$

- ▶ True at LO, but breaks down at NLO
- ▶ Requires external acceptance  $\mathcal{A}$
- ▶ Semileptonic branching fraction  $\mathcal{B}_\mu$  averaged over charm mesons is an effective quantity

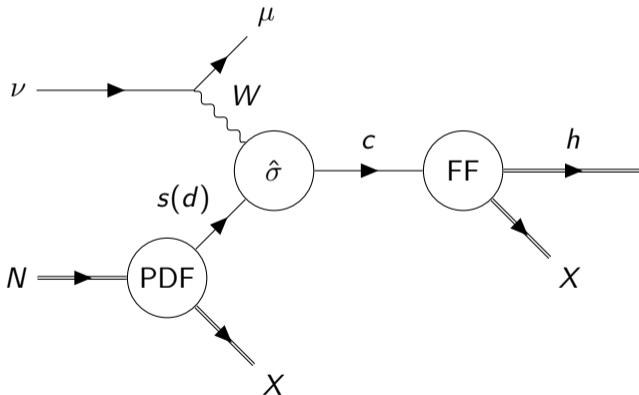


## SIDIS

- ▶ SIDIS: fragmentation of charm quark
- ▶ DGLAP-evolved NLO FFs
  - ▶ kkks08 for  $D^0$  and  $D^+$
  - ▶ bkk05 for  $D_s$  and  $\Lambda_c$
- ▶ Slow rescaling variable

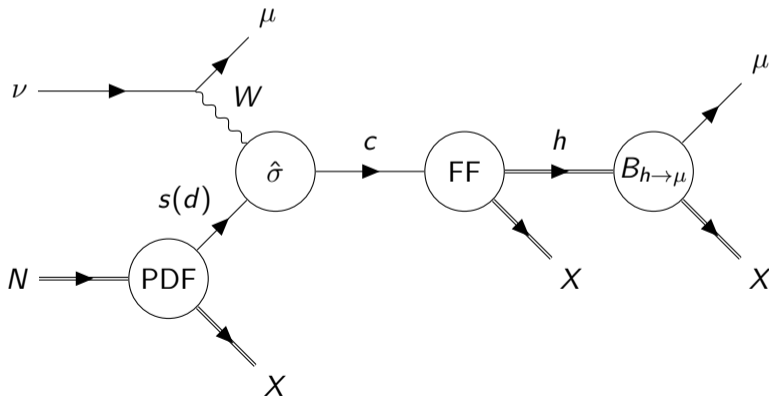
$$\chi = x \left( 1 + \frac{m^2}{Q^2} \right)$$

takes leading quark-mass effects into account



## Dimuon production in $\nu N$ collisions - the SIDIS approach

- ▶ Use SIDIS to produce a charmed hadron  $\rightarrow$  well-known
- ▶ Decay of charmed hadron to a muon  $\rightarrow$  how?



## General formalism

- ▶ Assuming production of hadron  $h$  and its decay factorizes, the cross section can be written as

$$\sigma(\nu N \rightarrow \mu\mu X) = \int d\sigma(\nu N \rightarrow \mu h X) \frac{\Gamma_{h \rightarrow \mu}}{\Gamma_{\text{tot}}}$$

- ▶ Write the decay width as [ $w = (P_\mu \cdot P_h)/m_h^2$ ]

$$\Gamma_{h \rightarrow \mu} = \frac{1}{2m_h} \int \frac{d^3\mathbf{P}_\mu}{E_\mu} d_{h \rightarrow \mu}(w) \implies \frac{d\Gamma_{h \rightarrow \mu}}{d|\mathbf{P}_\mu|} = \frac{\pi}{m_h} \frac{|\mathbf{P}_\mu|^2}{E_\mu} \int d(\cos\theta) d_{h \rightarrow \mu}(w)$$

- ▶ Introduce cut on muon energy

$$\Gamma_{h \rightarrow \mu}(E_h, E_\mu^{\min}) = \frac{\pi}{m_h} \int d\rho \rho E_h^2 \int d(\cos\theta) d_{h \rightarrow \mu}(w) \Big|_{E_\mu = \rho E_h \geq E_\mu^{\min}}$$

## Fitting the decay function

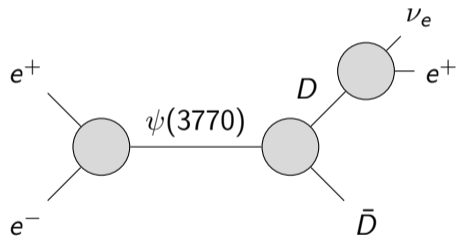
- ▶ Decay function parametrization

$$d_{D \rightarrow \mu}(w) = N w^\alpha (1 - \gamma w)^\beta \theta(0 \leq w \leq 1/\gamma)$$

- ▶ Fit the function

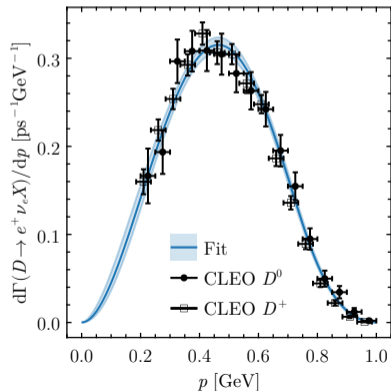
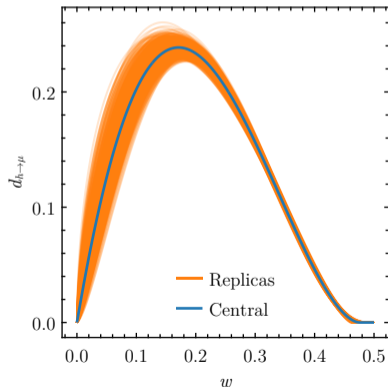
$$\frac{d\Gamma(D)}{d|\mathbf{P}_e|} = \frac{\pi}{m_D} |\mathbf{P}_e| \int d(\cos\theta) d_{D \rightarrow \mu}(w)$$

to CLEO data [PRL 97 (2006) 251801]



## Fit uncertainty

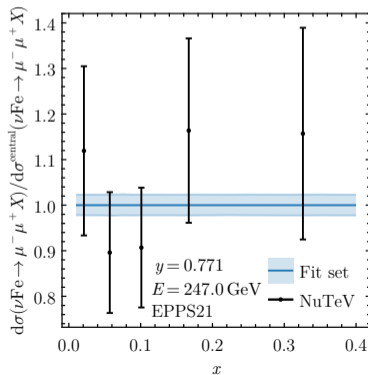
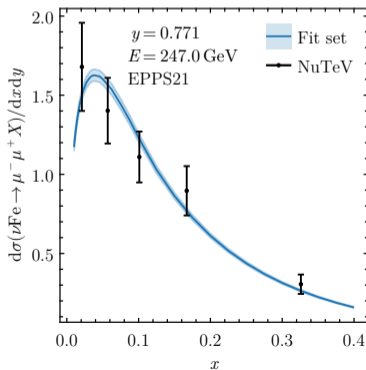
- ▶ Fit parameters  $N$ ,  $\alpha$ ,  $\beta$ , and  $\gamma$  are highly correlated
- ▶ Use 1000 Monte-Carlo replicas to estimate uncertainty





## Application to dimuon production

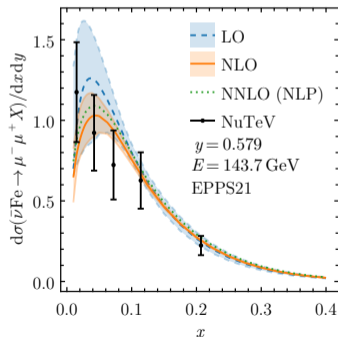
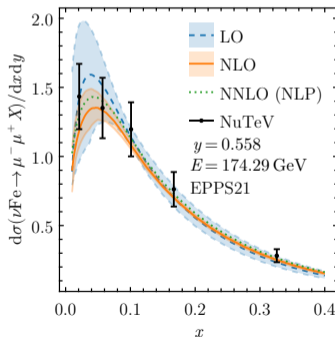
$$\frac{d\sigma(\nu N \rightarrow \mu\mu X)}{dx dy} = \sum_h \int dz \frac{d\sigma(\nu N \rightarrow \mu h X)}{dz dx dy} \frac{1}{\Gamma_{\text{tot}}^h} \Gamma_{h \rightarrow \mu}(E_h = zyE_\nu, E_\mu^{\text{min}})$$



## Scale uncertainty

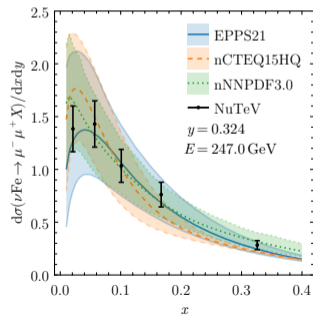
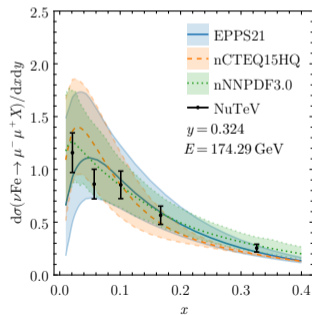
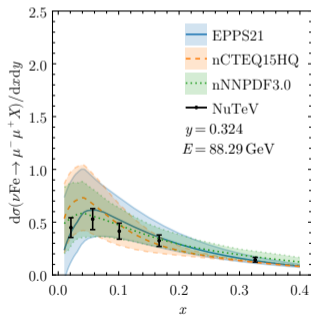
- ▶ Vary renormalization, factorization, and fragmentation scales
- ▶ Approximative quark-to-quark next-to-leading power NNLO contribution

[Abele et al; Phys.Rev.D 104 (2021) 9, 094046]



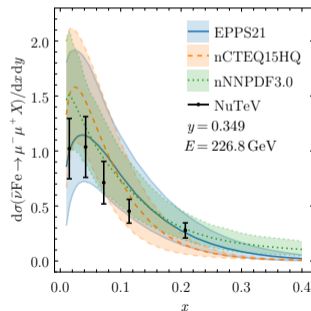
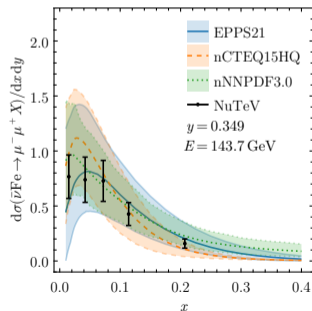
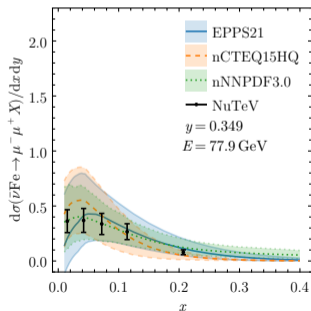
## PDF comparison

- ▶ Good agreement with NuTeV data [PRL 99 (2007) 192001]
- ▶ PDFs agree within uncertainty bands, but overall shapes differ → reflects the shape of the strange-quark distribution



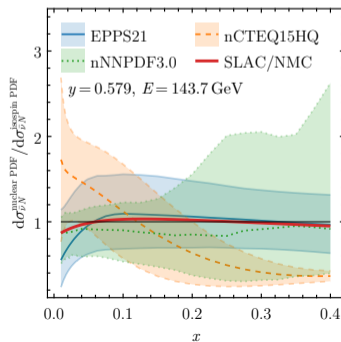
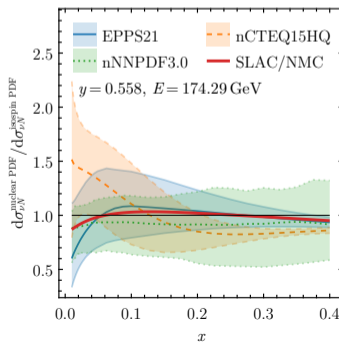
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## Nuclear modification

- ▶ Multiplicative nuclear modification factor  $d\sigma_{\nu N}^{\text{nPDF}} / d\sigma_{\nu N}^{\text{isospin PDF}}$
- ▶ Isospin-corrected PDFs  $(u_A^p, d_A^p) = \frac{Z}{A}(u^p, d^p) + \frac{N}{A}(d^p, u^p)$
- ▶ SLAC/NMC form for reference [Seligman, PhD thesis, Columbia U. 1997]
- ▶ Significant uncertainties from nPDF uncertainties



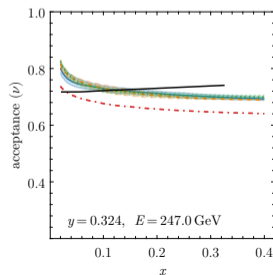
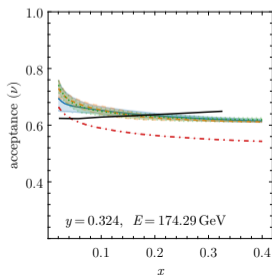
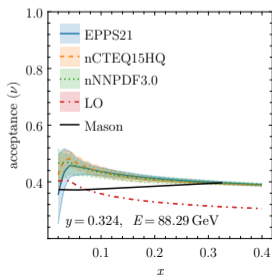
## Acceptance correction

- ▶ Dimuon cross section now has muon energy cut built-in → can compute acceptance from

$$\mathcal{A} = \frac{d\sigma(\nu N \rightarrow \mu\mu X)}{\mathcal{B}_\mu d\sigma(\nu N \rightarrow \mu c X)}$$

- ▶ Compare against existing Monte-Carlo calculation, used e.g. by NNPDF4.0

[D. A. Mason, PhD thesis, Oregon U. 2006]



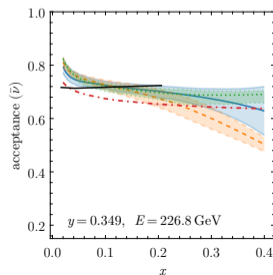
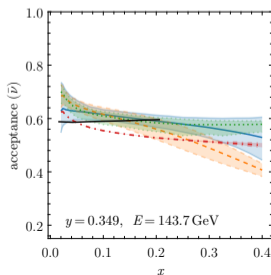
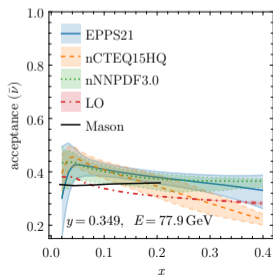
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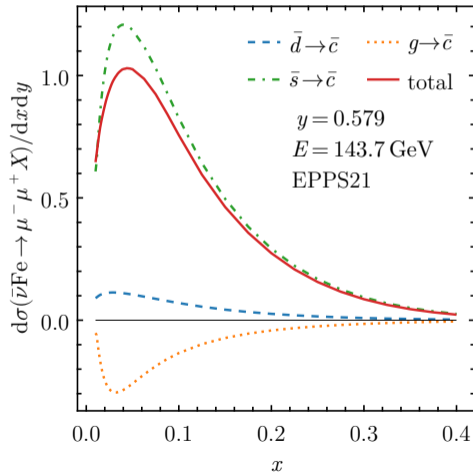
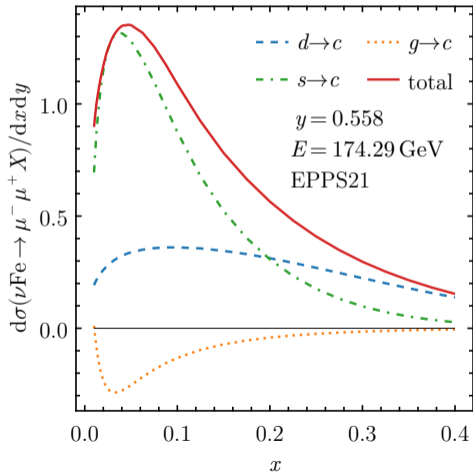
## Conclusions

- ▶ We have computed dimuon production in neutrino-nucleus collisions directly using SIDIS and a decay function
  - ▶ Decay function can be obtained by fitting to independent data
  - ▶ Assumption that dimuon production is proportional to charm production is not necessary
  - ▶ Effective acceptance correction shows systematic differences to a commonly used calculation
- ▶ Good agreement with NuTeV (and CCFR) dimuon data
- ▶ **The SIDIS approach provides a consistent and self-contained calculation of dimuon production without the need for an external acceptance correction**

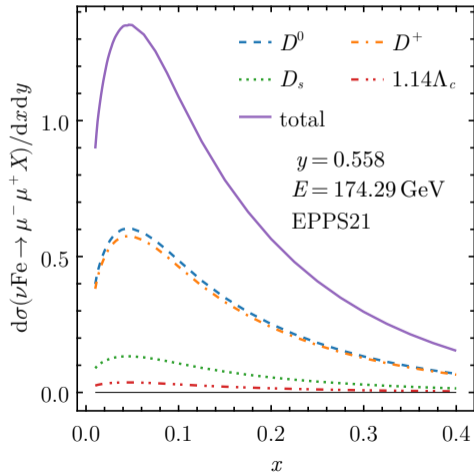




## Channel decomposition



# Fragmentation



## Quark-mass effects

- ▶ Leading quark-mass effects can be included with the slow rescaling variable

$$\chi = x \left( 1 + \frac{m^2}{Q^2} \right)$$

- ▶ The slow rescaling variable captures all mass effects at LO, and leading mass effects at NLO
- ▶ Relevant effects: shrinks the phase space, shifts the argument of PDFs

## Decay formalism (1/2)

- ▶ Assume that the **production** of a hadron  $h$  and its **decay** factorizes in  $AB \rightarrow h \rightarrow CX$

$$\sigma(AB \rightarrow CX) = \frac{1}{2s} \int d(\text{PS})(2\pi)^4 \delta^{(4)}(P_A + P_B - P_C - P_{X_1} - P_{X_2}) \\ \times (\text{production}) \frac{1}{(P_h^2 - m_h^2)^2 + m_h^2 \Gamma_{\text{tot}}^2} (\text{decay})$$

- ▶ Narrow-width approximation

$$\frac{1}{(P_h^2 - m_h^2)^2 + m_h^2 \Gamma_{\text{tot}}^2} \simeq \frac{\pi}{m_h \Gamma_{\text{tot}}} \delta(P_h^2 - m_h^2)$$

## Decay formalism (2/2)

- ▶ Insert unity as

$$1 = \int \frac{dP_h^2}{2\pi} \int \frac{d^3\mathbf{P}_h}{(2\pi)^3 2E_h} (2\pi)^4 \delta^{(4)}(P_h - P_C - P_{X_2})$$

- ▶ Identify production cross section and decay width

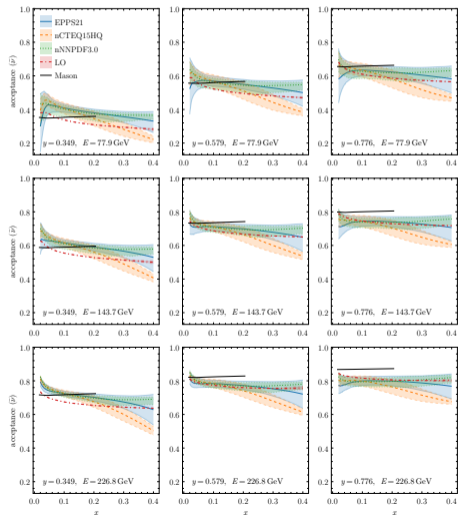
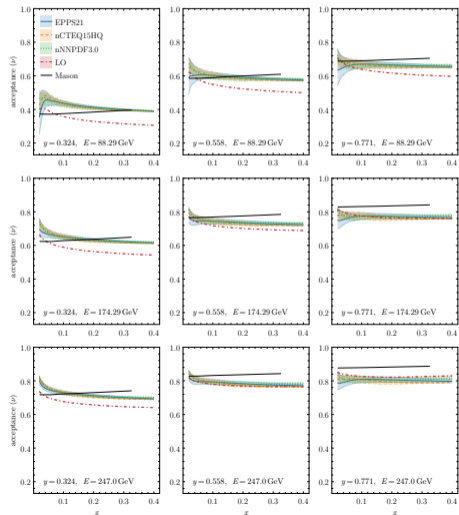
$$d\sigma(AB \rightarrow hX) = \frac{1}{2s} \int d(\text{PS}) (\text{production}) \delta^{(4)}(P_A + P_B - P_{X_1} - P_h - P_C)$$

$$\Gamma_{h \rightarrow C} = \frac{1}{2m_h} \int d(\text{PS}) (\text{decay}) \delta^{(4)}(P_h - P_{X_2} - P_C)$$

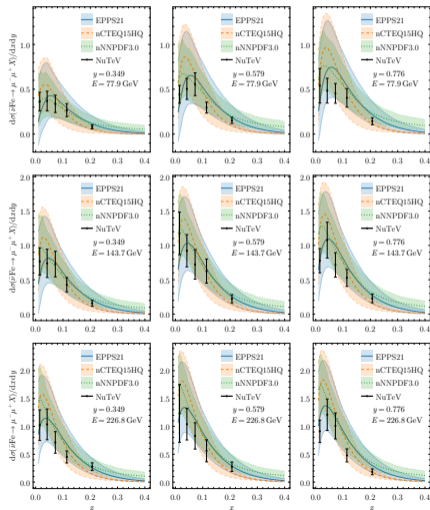
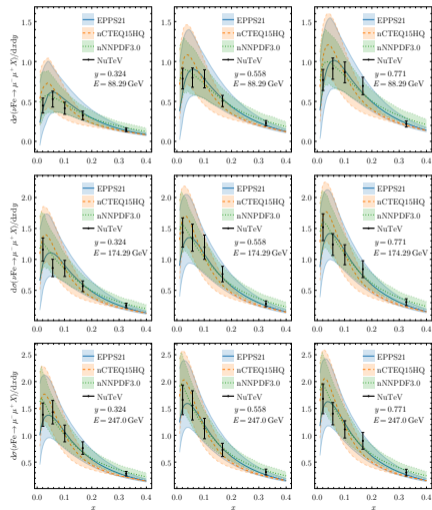
- ▶ With these, the cross section becomes

$$\sigma(AB \rightarrow CX) = \int d(\text{PS}) d\sigma(AB \rightarrow hX) \frac{\Gamma_{h \rightarrow C}}{\Gamma_{\text{tot}}}$$

# Acceptance systematics



# PDF comparison systematics (NuTeV)





# PDF comparison systematics (CCFR)

