The Sivers asymmetry of vector meson production in SIDIS process







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2. Framework

- TMD evolution
- The FFs of ρ^0 meson
- The calculation of Sivers asymmetry

3. Summary

Outline



Quark			
	L	Т	
		h_1^\perp () - ()	
	$g_1 \longrightarrow - \bigcirc \rightarrow$	h_{1L}^{\perp} $\triangleright \bullet $	
)	g_{1T}^{\perp}	$\begin{array}{ccc} h_1 & & & \bullet \\ h_1^{\perp} & & \bullet \end{array}$	U: Unpolari L: Longitud
		$T^{\mu}1T$ (?) – (?)	T: Transver

$$\begin{aligned} \frac{d\sigma^{h}}{dxdyd\phi_{S}dzd\phi_{h}d\mathbf{P}_{h\perp}^{2}} &= \frac{\alpha^{2}}{xyQ^{2}}\frac{y^{2}}{2(1-\varepsilon)}\left(1+\frac{\gamma^{2}}{2x}\right) \\ \times \left\{ \begin{bmatrix} F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)}\cos\phi_{h}F_{UU}^{\cos\phi_{h}} + \varepsilon\cos\left(2\phi\right) \\ + \lambda_{l} \left[\sqrt{2\varepsilon(1-\varepsilon)}\sin\phi_{h}F_{LU}^{\sin\phi_{h}}\right] \\ + S_{L} \left[\sqrt{2\varepsilon(1+\varepsilon)}\sin\phi_{h}F_{UL}^{\sin\phi_{h}} + \varepsilon\sin\left(2\phi_{h}\right)F_{UL}^{\sin2\phi_{h}}\right] \\ + S_{L} \lambda_{l} \left[\sqrt{1-\varepsilon^{2}}F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)}\cos\phi_{h}F_{LL}^{\cos\phi_{h}}\right] \\ + S_{T} \left[\sin\left(\phi_{h} - \phi_{S}\right)\left(F_{UT,T}^{\sin(\phi_{h} - \phi_{S})} + \varepsilon F_{UT,L}^{\sin(\phi_{h} - \phi_{S})}\right) \\ + \varepsilon\sin\left(\phi_{h} + \phi_{S}\right)F_{UT}^{\sin(\phi_{h} + \phi_{S})} + \varepsilon\sin\left(3\phi_{h} - \phi_{S}\right)F_{UT}^{\sin(3\phi_{h}} \\ + \sqrt{2\varepsilon(1+\varepsilon)}\sin\phi_{S}F_{UT}^{\sin\phi_{S}} + \sqrt{2\varepsilon(1+\varepsilon)}\sin\left(2\phi_{h} - \phi_{S}\right) \\ + S_{T} \lambda_{l} \left[\sqrt{1-\varepsilon^{2}}\cos\left(\phi_{h} - \phi_{S}\right)F_{LT}^{\cos(\phi_{h} - \phi_{S})} + \sqrt{2\varepsilon(1-\varepsilon)}\cos\phi_{h}\right] \right\} \end{aligned}$$

 $F_{X Y, Z} \propto PDF \otimes FF X$: Lepton polarization *Y*: Target polarization *Z*: γ^* polarization

 $\left. \phi_h \right) F_{UU}^{\cos 2\phi_h}
ight|$

Transversely single spin asymmetry (TSSA):

$$A_{UT} = \frac{\sigma\left(\vec{S}_{T}\right) - \sigma\left(-\vec{S}_{T}\right)}{\sigma\left(\vec{S}_{T}\right) + \sigma\left(-\vec{S}_{T}\right)}$$



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$$\frac{d\sigma^{h}}{dxdyd\phi_{S}dzd\phi_{h}d\mathbf{P}_{h\perp}^{2}} = \frac{\alpha^{2}}{xyQ^{2}}\frac{y^{2}}{2(1-\varepsilon)}\left(1+\frac{\gamma^{2}}{2x}\right)$$

$$\times \left\{ \begin{bmatrix} F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)}\cos\phi_{h}F_{UU}^{\cos\phi_{h}} + \varepsilon\cos\left(2\varphi\right) \\ + \lambda_{l} \left[\sqrt{2\varepsilon(1-\varepsilon)}\sin\phi_{h}F_{LU}^{\sin\phi_{h}}\right] \\ + S_{L} \left[\sqrt{2\varepsilon(1+\varepsilon)}\sin\phi_{h}F_{UL}^{\sin\phi_{h}} + \varepsilon\sin\left(2\phi_{h}\right)F_{UL}^{\sin2\phi_{h}}\right] \\ + S_{L} \lambda_{l} \left[\sqrt{1-\varepsilon^{2}}F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)}\cos\phi_{h}F_{LL}^{\cos\phi_{h}}\right] \\ + S_{T} \left[\sin\left(\phi_{h} - \phi_{S}\right)\left(F_{UT,T}^{\sin(\phi_{h} - \phi_{S})} + \varepsilon F_{UT,L}^{\sin(\phi_{h} - \phi_{S})}\right) \\ + \varepsilon\sin\left(\phi_{h} + \phi_{S}\right)F_{UT}^{\sin(\phi_{h} + \phi_{S})} + \varepsilon\sin\left(3\phi_{h} - \phi_{S}\right)F_{UT}^{\sin(3\phi_{h}} \\ + \sqrt{2\varepsilon(1+\varepsilon)}\sin\phi_{S}F_{UT}^{\sin\phi_{S}} + \sqrt{2\varepsilon(1+\varepsilon)}\sin\left(2\phi_{h} - \phi_{S}\right) \\ + S_{T} \lambda_{l} \left[\sqrt{1-\varepsilon^{2}}\cos\left(\phi_{h} - \phi_{S}\right)F_{LT}^{\cos(\phi_{h} - \phi_{S})} + \sqrt{2\varepsilon(1-\varepsilon)}\cos\varepsilon\left(2\phi_{h} - \phi_{S}\right)F_{LT}^{\cos(2\phi_{h} - \phi_{S})}\right]\right\}$$

 $F_{X Y, Z} \propto PDF \otimes FF X$: Lepton polarization *Y*: Target polarization *Z*: γ^* polarization

 $\left. \phi_h \right) F_{UU}^{\cos 2\phi_h} \right|$













$$\frac{d\sigma^{h}}{dxdyd\phi_{S}dzd\phi_{h}d\mathbf{P}_{h\perp}^{2}} = \frac{\alpha^{2}}{xyQ^{2}}\frac{y^{2}}{2(1-\epsilon)}\left(1+\frac{\gamma^{2}}{2x}\right)$$

$$\times \left\{ \begin{bmatrix} F_{UU,T} + \epsilon F_{UU,L} + \sqrt{2\epsilon(1+\epsilon)}\cos\phi_{h}F_{UU}^{\cos\phi_{h}} + \epsilon\cos\left(2\phi_{h}\right) \\ + \lambda_{l} \left[\sqrt{2\epsilon(1-\epsilon)}\sin\phi_{h}F_{LU}^{\sin\phi_{h}}\right] \\ + S_{L} \left[\sqrt{2\epsilon(1+\epsilon)}\sin\phi_{h}F_{UL}^{\sin\phi_{h}} + \epsilon\sin\left(2\phi_{h}\right)F_{UL}^{\sin2\phi_{h}}\right] \\ + S_{L} \lambda_{l} \left[\sqrt{1-\epsilon^{2}}F_{LL} + \sqrt{2\epsilon(1-\epsilon)}\cos\phi_{h}F_{LL}^{\cos\phi_{h}}\right] \\ + S_{T} \left[\sin\left(\phi_{h} - \phi_{S}\right)\left(F_{UT,T}^{\sin(\phi_{h} - \phi_{S})} + \epsilon F_{UT,L}^{\sin(\phi_{h} - \phi_{S})}\right) \\ + \epsilon\sin\left(\phi_{h} + \phi_{S}\right)F_{UT}^{\sin(\phi_{h} + \phi_{S})} + \epsilon\sin\left(3\phi_{h} - \phi_{S}\right)F_{UT}^{\sin(3\phi_{h}} \\ + \sqrt{2\epsilon(1+\epsilon)}\sin\phi_{S}F_{UT}^{\sin\phi_{S}} + \sqrt{2\epsilon(1+\epsilon)}\sin\left(2\phi_{h} - \phi_{S}\right) \\ + S_{T} \lambda_{l} \left[\sqrt{1-\epsilon^{2}}\cos\left(\phi_{h} - \phi_{S}\right)F_{LT}^{\cos(\phi_{h} - \phi_{S})} + \sqrt{2\epsilon(1-\epsilon)}\cos\phi_{h}F_{LT}^{\cos(\phi_{h} - \phi_{S})}\right] \right\}$$





Collins and Sivers transverse-spin asymmetries in inclusive muoproduction of ρ^0 mesons

The COMPASS Collaboration

Abstract

The production of vector mesons in deep inelastic scattering is an interesting yet scarsely explored channel to study the transverse spin structure of the nucleon and the related phenomena. The COMPASS collaboration has performed the first measurement of the Collins and Sivers asymmetries for inclusively produced ρ^0 mesons. The analysis is based on the data set collected in deep inelastic scattering in 2010 using a 160 GeV/c μ^+ beam impinging on a transversely polarized NH₃ target. The ρ^0 mesons are selected from oppositely charged hadron pairs, and the asymmetries are extracted as a function of the Bjorken-x variable, the transverse momentum of the pair and the fraction of the energy z carried by the pair. Indications for positive Collins and Sivers asymmetries are observed.

The COMPASS, C. Alice, A. Amoroso, V. Andrieux et al., Phys.Lett.B 843 (2023) 137950



The Sivers asymmetry of ρ^0 meson in SIDIS process has been measured firstly.



 $A_{UT}^{sin(\varphi_{hh}^{}-\varphi_{S}^{})}$







TMD factorization and evolution:

$$F_{\text{Sivers}}^{\alpha} (x_B, z_h, P_{h\perp}, Q) = \int \frac{d^2b}{(2\pi)^2} e^{i\vec{P}_{h\perp}\cdot\vec{b}/z_h} \widetilde{F}_{\text{Sivers}}^{\alpha} (x_B, z_h, b, Q) + Y_{\text{Sivers}}^{\alpha} (x_B, z_h, P_{h\perp}, Q)$$

$$F_{UU}(x_B, z_h, P_{h\perp}, Q) = \int \frac{d^2b}{(2\pi)^2} e^{i\vec{P}_{h\perp} \cdot \vec{b}/z_h} \widetilde{F}_{UU}(x_B, z_h, b, Q) + Y_{UU}(x_B, z_h, P_{h\perp}, Q)$$

We focus on the region $P_{h\perp} \ll Q$, where the TMD factorization approximatively applies.

Framework

$$A_{UT}^{\sin(\phi_h - \phi_S)} = \frac{F_{UT}^{\sin(\phi_h - \phi_S)}}{F_{UU}}$$
$$\sin(\phi_h - \phi_S) F_{UT}^{\sin(\phi_h - \phi_S)} = \varepsilon_{\perp \alpha \beta} S_{\perp}^{\alpha} F_{S}^{\beta}$$

Dominates in $P_{h\perp} \ll Q$

Dominates in $P_{h\perp} \gtrsim Q$



TMD Factorization $\widetilde{F}_{UU}(x_B, z_h, b, Q) = H(\mu, Q) \sum_{\sigma} e_q^2 \widetilde{f}_{1,q/p}(x_B, b, \mu, \zeta_1) \widetilde{D}_{1,h/q}(z_h, b, \mu, \zeta_2)$ $\widetilde{F}_{\text{Sivers}}^{\alpha}(x_B, z_h, b, Q) = H(\mu, Q) \sum_{\widetilde{r}} e_q^2 \left(-iMb^{\alpha}\right) \widetilde{f}_{1T,q/p}^{\perp}(x_B, b, \mu, \zeta_1) \widetilde{D}_{1,h/q}(z_h, b, \mu, \zeta_2)$

$$f_{1,q/p}\left(x,k_{\perp}^{2},\mu,\zeta\right) = \frac{1}{2\pi} \int \mathrm{d}b \ bJ_{0}\left(bk_{\perp}\right) \widetilde{f}_{1,q/p}\left(x,b,\mu,\zeta\right)$$
$$D_{1,h/q}\left(z,p_{\perp}^{2},\mu,\zeta\right) = \frac{1}{2\pi} \int \mathrm{d}b \ bJ_{0}\left(b\frac{p_{\perp}}{z}\right) \widetilde{D}_{1h/q}\left(z,b,\mu,\zeta\right)$$
$$\frac{k_{\perp}}{M} f_{1T}^{\perp}\left(x,k_{\perp}^{2},\mu,\zeta\right) = \int_{0}^{\infty} \frac{dbb^{2}M}{2\pi} J_{1}\left(bk_{\perp}\right) \widetilde{f}_{1T}^{\perp}\left(x,b,\mu,\zeta\right)$$

 μ : the renormalization scale $\mu^2 = Q^2$ $\zeta_1 \zeta_2 = Q^4$ ζ : the rapidity scale



symmetrical choice

TMD evolution obtain the TMD functions with any scale.

TMDs evolution equations: CS equation:

$$\zeta \frac{d\tilde{F}(x,b;\mu,\zeta)}{d\zeta} = -\mathcal{D}(b,\mu)\tilde{F}(x,b;\mu,\zeta) \qquad \mu^2 \frac{d\tilde{F}(x,b,\mu,\zeta)}{d\mu^2} = \frac{\gamma_F(\mu,\zeta)}{2}\tilde{F}(x,b,\mu,\zeta)$$

 \mathcal{D} : the rapidity anomalous dimension

The solution: $\widetilde{F}(x,b;\mu_f,\zeta_f) = R[b;(\mu_f,\zeta_f) -$

 $R[b; (\mu_i, \zeta_i) \to (\mu_f, \zeta_f)] = \exp$

$$\zeta \frac{\mathrm{d}}{\mathrm{d}\zeta} \gamma_F(\mu,\zeta) = -\mu \frac{\mathrm{d}}{\mathrm{d}\mu} \mathcal{D}(\mu,\zeta)$$

 \bar{F} stands for any TMD function.

RG equation:

 γ_F : the TMD anomalous dimension

$$\rightarrow (\mu_i, \zeta_i)] \widetilde{F}(x, b; \mu_i, \zeta_i) \\ \left[\int_P \left(\gamma_F(\mu, \zeta) \frac{d\mu}{\mu} - \mathcal{D}(\mu, b) \frac{d\zeta}{\zeta} \right) \right]$$

The finite-order perturbative calculation destroys the path independence property.



TMD evolution:

$$\frac{\partial \ln \tilde{F}(x, b, \mu, \zeta)}{\partial \ln \mu^2} = \frac{\gamma_F(\mu, \zeta)}{2}$$
$$\frac{\partial \ln \tilde{F}(x, b; \mu, \zeta)}{\partial \ln \zeta} = -\mathcal{D}(b, \mu)$$

$$\vec{E} = \left(\frac{\gamma_F(\mu, \zeta)}{2}, -\mathcal{D}(\mu, \zeta)\right)$$

In the ζ -prescription, the initial scales μ and ζ belong to a null-evolution line, that is expressed as $(\mu, \zeta_{\mu}(b))$



 $\widetilde{F}(x,b;$

Ignazio Scimemi, Alexey Vladimirov, JHEP 06 (2020) 137

$$\begin{split} \zeta_{\mu}(\mu,b) = & \zeta_{\mu}^{\text{pert}}(\mu,b) e^{-b^2/B_{\text{NP}}^2} \\ & + \zeta_{\mu}^{\text{exact}}(\mu,b) \left(1 - e^{-b^2/B_{\text{N}}^2}\right) \\ \mathcal{D}(\mu,b) = & \mathcal{D}_{\text{resum}}(\mu,b_*) + d_{\text{NP}}(b) \\ & b_* = \frac{b}{\sqrt{1 + \frac{b^2}{B_{NP}^2}}} \end{split}$$

$$(Q, Q^2) = \exp\left[\int_P \left(\gamma_F(\mu, \zeta) \frac{d\mu}{\mu} - \mathcal{D}(\mu, b) \frac{d\zeta}{\zeta}\right)\right] \widetilde{F}(x, b)$$
$$= \left(\frac{Q^2}{\zeta_Q(b)}\right)^{-\mathcal{D}(b, Q)} \widetilde{F}(x, b)$$



Within the TMD factorization and TMD evolution, the Sivers asymmetry: $\sin(\phi)$

$$A_{UT}^{\sin(\phi_h - \phi_S)} = \frac{F_{UT}^{\sin(\phi_h - \phi_S)}}{F_{UU}}$$

$$F_{UT}^{\sin(\phi_h - \phi_s)}\left(x_B, z_h, P_{h\perp}, Q\right) = -H\left(\mu, Q\right) M \sum_q e_q^2 \int_0^\infty \frac{\mathrm{d}b}{2\pi} b^2 J_1\left(\frac{bP_{h\perp}}{z_h}\right) \left(\frac{Q^2}{\zeta_Q(b)}\right)^{-2\mathcal{D}(b,Q)}$$
$$\times \tilde{f}_{1T,q/p}^{\perp}\left(x_B, b\right) \tilde{D}_{1,h/q}\left(z_h, b\right)$$

$$F_{UU}(x_B, z_h, P_{h\perp}, Q) = H(\mu, Q) \sum_q e_q^2 \int_0^\infty \frac{\mathrm{d}b}{2\pi} b J_0\left(\frac{bP_{h\perp}}{z_h}\right) \left(\frac{Q^2}{\zeta_Q(b)}\right)^{-2\mathcal{D}(b,Q)}$$
$$\times \tilde{f}_{1,q/p}(x_B, b) \tilde{D}_{1,h/q}(z_h, b)$$

 $-\phi_S)$



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Chunhua Zeng, Tianbo Liu, Peng Sun, Yuxiang Zhao, Phys.Rev.D 106 (2022) 9,094039

Sivers function parameterization: BPV20

Marcin Bury, Alexei Prokudin, Alexey Vladimirov, JHEP 05 (2021) 151





Unpolarized FFs of ρ^0 meson

The unpolarized FFs of ρ^0 meson:

No appropriate parametrization for FFs of ρ^0 meson.

 $F^{\rho^0}\left(z,Q^2\right) = \frac{1}{\sigma} \frac{d\sigma\left(z,Q^2\right)}{\sigma}$



Approach : perform a global fit of Pythia's ρ^0 meson data.

$$\frac{\left(e^+e^- \to \rho^0 X\right)}{dz} = \frac{1}{N_{tot}} \frac{\Delta N \left(e^+e^- \to \rho^0 X\right)}{\Delta z}$$

Q = 110GeV Charge conjugate symmetry, Isospinit symmetry $\sum_{k=0.9}^{6.9} \sum_{k=0}^{6.9} \sum_{k=0}^{6.9$ beta_g 9.951 ± 8.561 ^{0.2} $D_{\rho^0/g}(z,\mu_0) = N_g \times z^{\alpha_g} \times (1-z)^{\beta_g}$ $\mu_0^2 = 1.2 GeV^2$ Only take the gluon, $u, \bar{u}, d, \bar{d}, s, \bar{s}$ into account







Unpolarized FFs of ρ^0 meson

Result:



$$egin{aligned} \chi^2/d.\,o.\,f. &= 749.4/743 = 1.0086 \ function & N \ D_{
ho^0/u} & 0.4224 \pm 0.0013 & -0.6118 \ D_{
ho^0/s} & 0.3346 \pm 0.0044 & -0.777 \ D_{
ho^0/g} & 129.038 \pm 21.1586 & 3.2235 \end{aligned}$$



The TMD FFs

Obtain TMD FFs by collinear FFs:

$$D_{1,h/f}\left(z,b\right) = \frac{1}{z^2} \sum_{f'} \int_{z}^{1} \frac{\mathrm{d}y}{y} y^2 \mathbb{C}_{f \to f'}\left(y,b,\mu_{\mathrm{OPE}}^{\mathrm{FF}}\right) D_{1,h/f'}\left(\frac{z}{y},\mu_{\mathrm{OPE}}^{\mathrm{FF}}\right) D_{\mathrm{NP}}\left(z,b\right)$$



Numerical Result

The Sivers asymmetry of ρ^0 meson



Sivers function parameterization: **ZLSZ**

Sivers function parameterization: **BPV20**







Sivers asymmetry at the EIC's kinematics



0.015 $A_{UT}^{\sin(\phi_h - \phi_S)}$ 0.010 BPV20 0.01y = 0.20.005 $z_h = 0.48$ $P_{h\perp} = 0.4 \text{GeV}$ 0.000 0.0050 $A_{UT}^{\sin(\phi_h-\phi_S)}$ 0.0025 ZLSZ 0.0000 0.000 0.10.20.0 x_B



Sivers asymmetry at the EicC's kinematics

EicC $\sqrt{s} = 16.7 \text{GeV}$

BPV20



ZLSZ

Numerical Result

Sivers asymmetry of *K** mesons.

EIC = 100 GeV



ZLSZ



Numerical Result

Sivers asymmetry of *K** mesons.

BPV20

ZLSZ

EIC = 100 GeV





Sivers asymmetry of *K** mesons.

EicC $\sqrt{s} = 16.7 \text{GeV}$



 K^{*-}

Sivers asymmetry of *K** mesons.

BPV20

ZLSZ



Numerical Result

Comparing the Sivers asymmetry between EIC and EicC

BPV20







- 1. The Sivers function extracted from pion's and kaon's data can be well matched with ρ^0 's data. The universality of Sivers function.
- 2. There is a large difference between ZLSZ parameterization and BPV20 parameterization. EIC and EicC can provide a test of the Sivers function's universality and constrain the extraction of the Sivers function.

Thanks!