

Matching between TMD and twist-3 factorizations in the transversely polarized hyperon production

Shinsuke Yoshida

(South China Normal University)



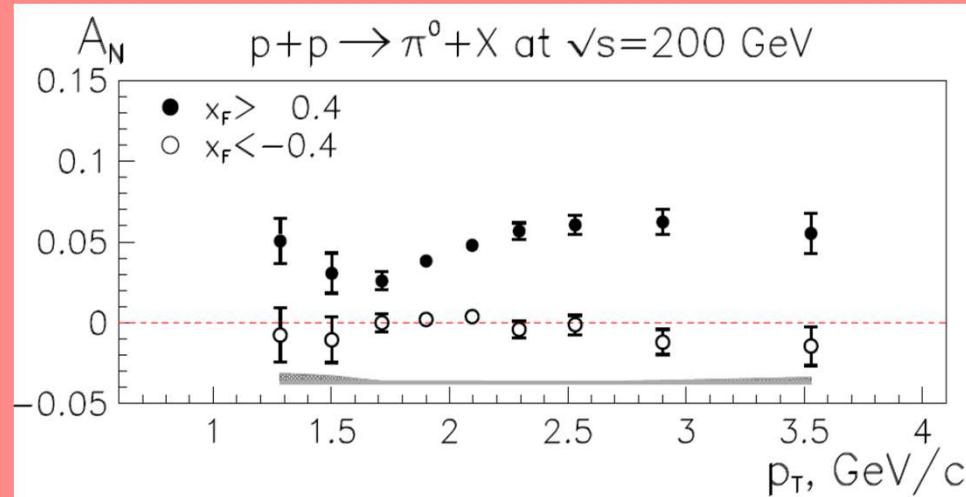
in collaboration with: Riku Ikarashi(Niigata University)

Yuji Koike(Niigata University)

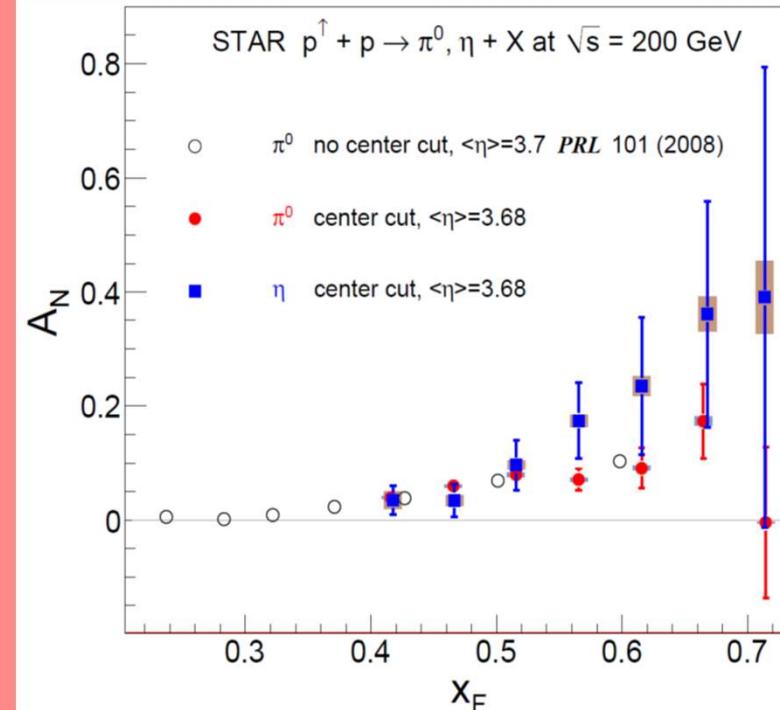
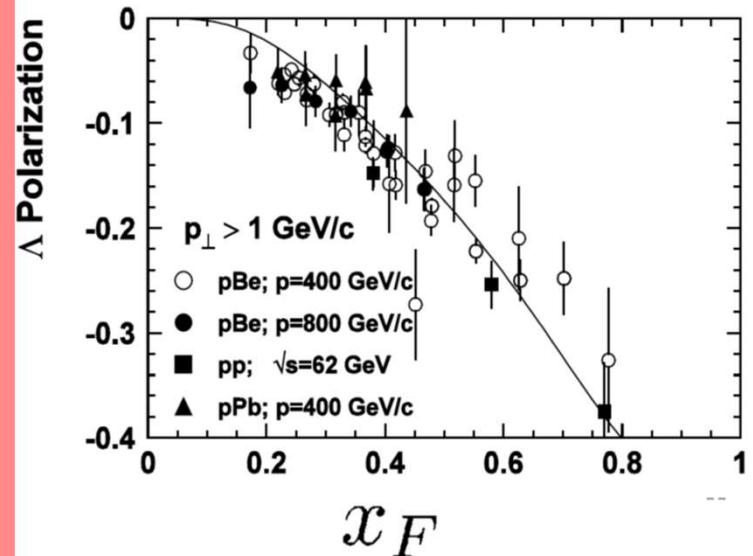
Mysteries in the transverse polarization

Almost half a century mysteries in high-energy QCD physics

$p^\uparrow p \rightarrow \pi X$



$pp \rightarrow \Lambda^\uparrow X$ (in 80's and 90's)



SSA estimated by pQCD

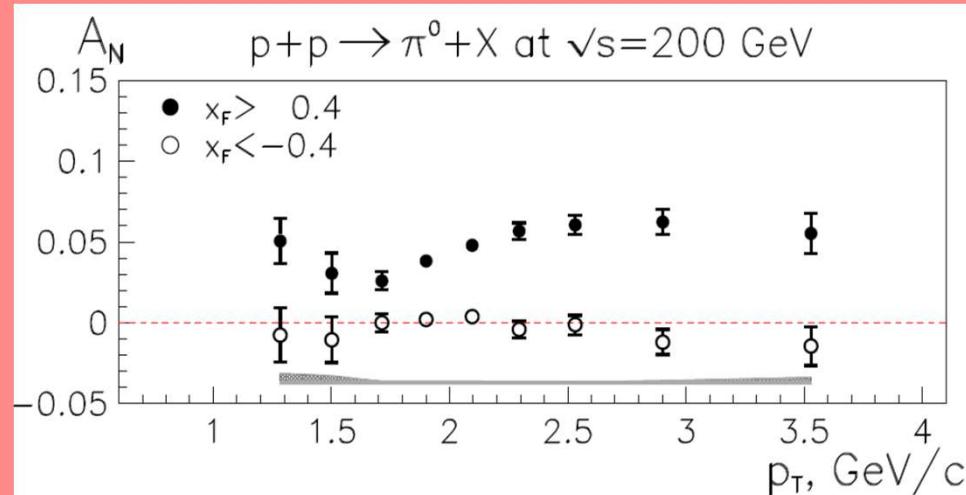
Kane *et al.* ('78)

$$A_N \sim \alpha_s \frac{m_q}{P_T} \sim \text{negligible!}$$

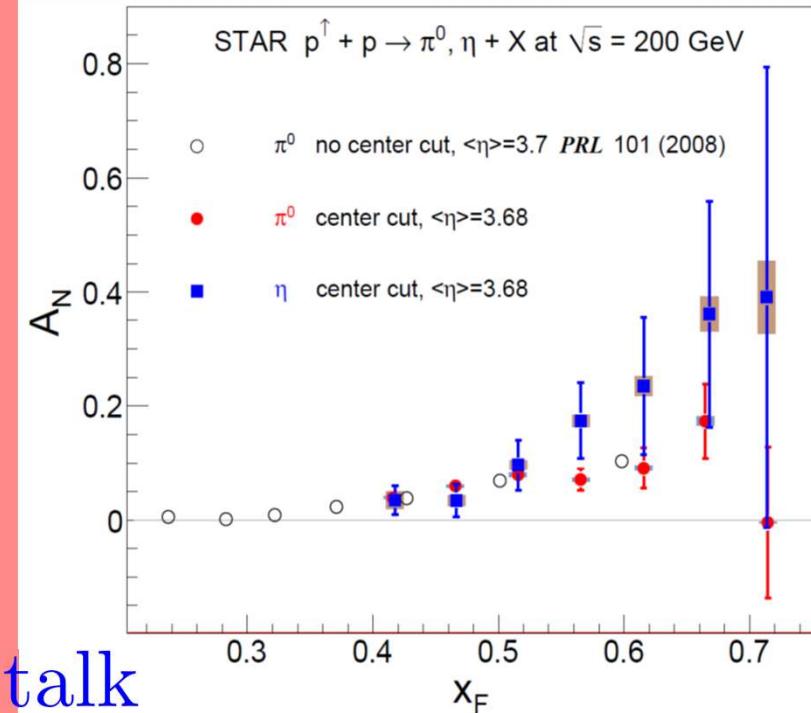
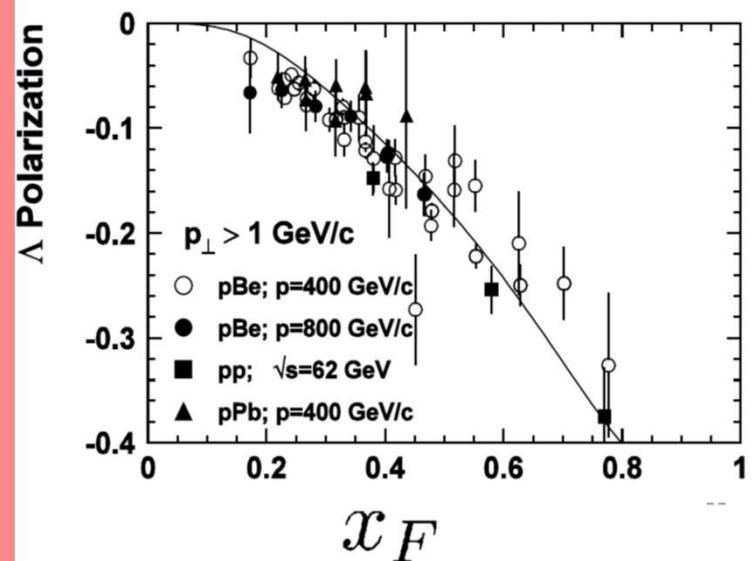
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This talk

SSA estimated by pQCD

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$$A_N \sim \alpha_s \frac{m_q}{P_T} \sim \text{negligible!}$$

pQCD frameworks for the SSA

Transverse Momentum Dependent(TMD) factorization

- Applicable in small $P_T(Q \gg P_T \geq \Lambda_{QCD})$ region cf. Sivers function $f_{1T}^\perp(x, \textcolor{red}{k}_\perp)$
- Nonperturbative functions depend on the transverse momentum of partons

advantage: TMD functions have definite physical interpretation

disadvantage: limited applicable processes $\times pp \rightarrow \pi X$

collinear factorization

- Applicable in large $P_T(P_T \gg \Lambda_{QCD})$ region
- twist-3 multiparton correlation inside hadrons causes the large SSA

advantage: Applicable to many processes such as $pp \rightarrow \pi X$

disadvantage: Physical interpretation of the twist-3 functions is unclear

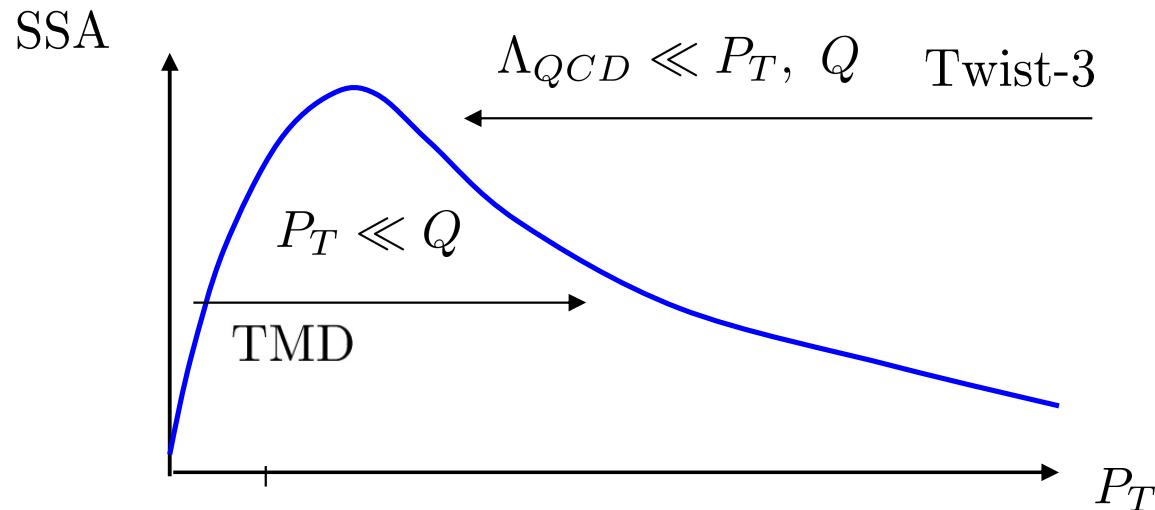
There is overlapping region $\Lambda_{QCD} \ll P_T \ll Q$

There should not be a quantitative ambiguity in this region
if both are the first principle calculations

Relationship between TMD and twist-3

TMD and the collinear twist-3 gives equivalent results in intermediate P_T region in the pion production X. Ji, J.-W. Qiu, W. Vogelsang and F. Yuan, PRD73(2006) PLB638(2006)

F. Yuan and J. Zhou, PRL 103 (2009)



QS-function

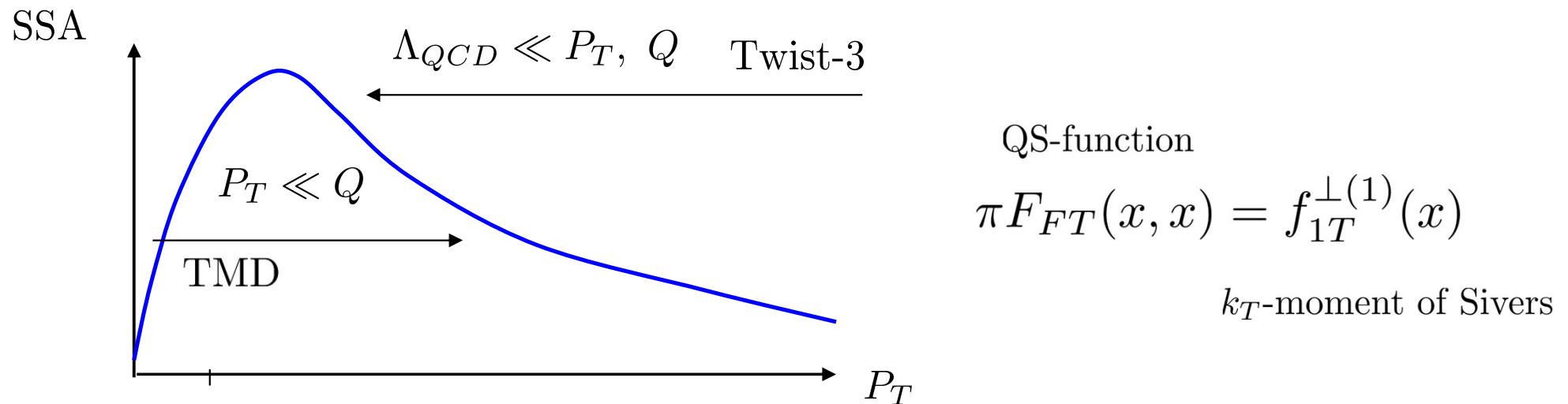
$$\pi F_{FT}(x, x) = f_{1T}^{\perp(1)}(x)$$

k_T -moment of Sivers

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The Latest parameterizations rely on the relations between TMD and the collinear

L. Gamberg *et al.* [Jefferson Lab Angular Momentum (JAM) collaboration], Phys. Rev. D106 (2022)

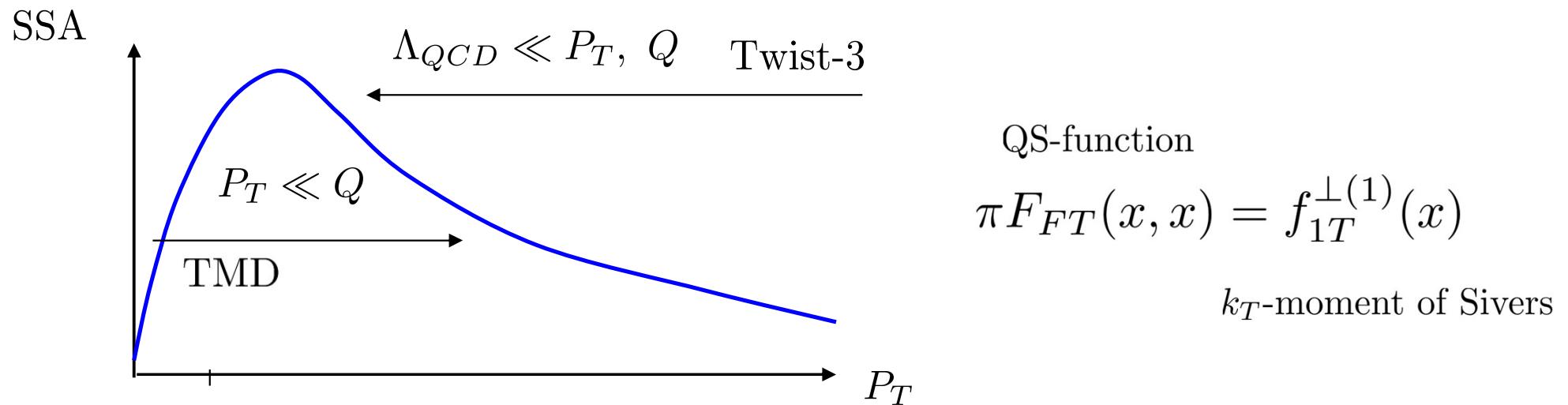
$$f(x) = \int d^2 \vec{k}_T f(x, \vec{k}_T^2) \quad (f = f_1 \text{ or } h_1), \quad \pi F_{FT}(x, x) = \int d^2 \vec{k}_T \frac{\vec{k}_T^2}{2M^2} f_{1T}^{\perp}(x, \vec{k}_T^2) \equiv f_{1T}^{\perp(1)}(x),$$

$$H_1^{\perp(1)}(z) = z^2 \int d^2 \vec{p}_T \frac{\vec{p}_T^2}{2M_h^2} H_1^{\perp}(z, z^2 \vec{p}_T^2).$$

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$$f(x) = \int d^2 \vec{k}_T f(x, \vec{k}_T^2) \quad (f = f_1 \text{ or } h_1), \quad \pi F_{FT}(x, x) = \int d^2 \vec{k}_T \frac{\vec{k}_T^2}{2M^2} f_{1T}^{\perp}(x, \vec{k}_T^2) \equiv f_{1T}^{\perp(1)}(x),$$

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How about the hyperon polarization ?

The distribution part has been done already. J. Zhou, F. Yuan and Z. -T. Liang, PRD78 (2008)

Operator identity in QFT

Two operators are identical

$$\hat{O}_1 = \hat{O}_2$$

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They give the same result under the same operation

$$\langle 0 | T[\hat{O}_1 \hat{A}^\mu \hat{A}^\nu] | 0 \rangle = \langle 0 | T[\hat{O}_2 \hat{A}^\mu \hat{A}^\nu] | 0 \rangle$$

Operator identity in QFT

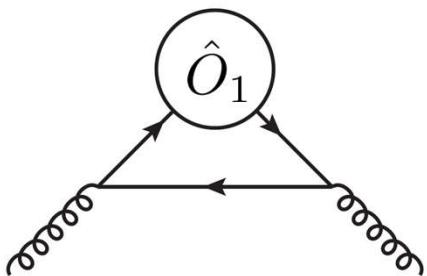
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cf. $\hat{O}_1 = \partial_\mu (\bar{\psi} \gamma^5 \gamma^\mu \psi)$

$$\langle 0 | T[\hat{O}_1 \hat{A}^\mu \hat{A}^\nu] | 0 \rangle \Big|_{g^2} = \text{Diagram} = N_f \frac{g^2}{32\pi^2} \left[4 \frac{1}{p^2} \frac{1}{(p')^2} \epsilon^{\mu\nu\rho\sigma} p_\rho p'_\sigma \right]$$


Operator identity in QFT

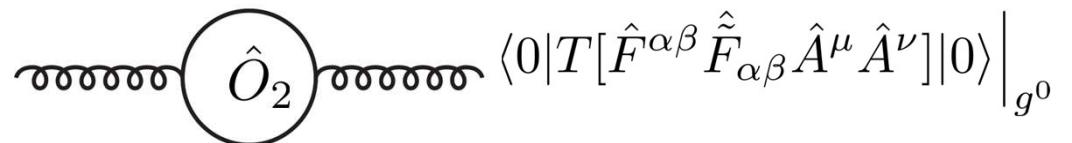
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$$\langle 0 | T[\hat{O}_1 \hat{A}^\mu \hat{A}^\nu] | 0 \rangle \Big|_{g^2} =$$

A triangular loop representing the operator \hat{O}_1 . The top vertex is a circle labeled \hat{O}_1 , with two arrows pointing towards it from the sides. The bottom vertices are wavy lines representing external fields.

$$= N_f \frac{g^2}{32\pi^2} \left[4 \frac{1}{p^2} \frac{1}{(p')^2} \epsilon^{\mu\nu\rho\sigma} p_\rho p'_\sigma \right]$$

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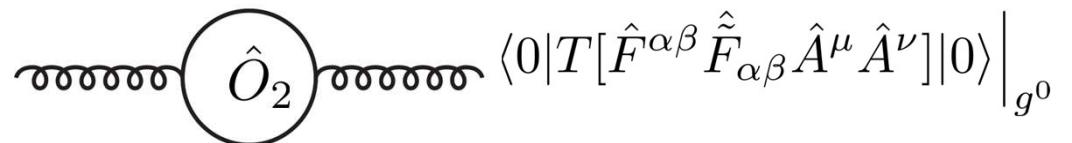
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$$\langle 0 | T[\hat{O}_1 \hat{A}^\mu \hat{A}^\nu] | 0 \rangle \Big|_{g^2} = \begin{array}{c} \text{Feynman diagram for } \hat{O}_1 \\ \text{A triangle loop with a wavy line entering from the left and exiting to the right. The top vertex is a circle labeled } \hat{O}_1. \end{array} = N_f \frac{g^2}{32\pi^2} \left[4 \frac{1}{p^2} \frac{1}{(p')^2} \epsilon^{\mu\nu\rho\sigma} p_\rho p'_\sigma \right]$$

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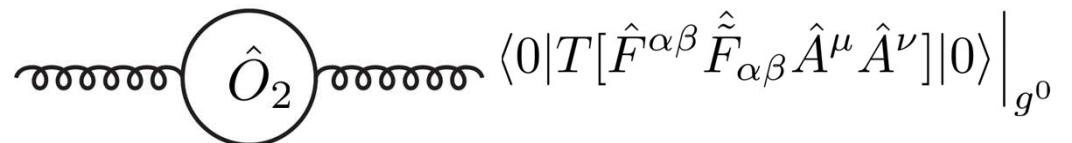
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$$\partial_\mu (\bar{\psi} \gamma^5 \gamma^\mu \psi) = N_f \frac{g^2}{32\pi^2} \hat{F}^{\alpha\beta} \hat{F}_{\alpha\beta}$$

Operator identity in QFT

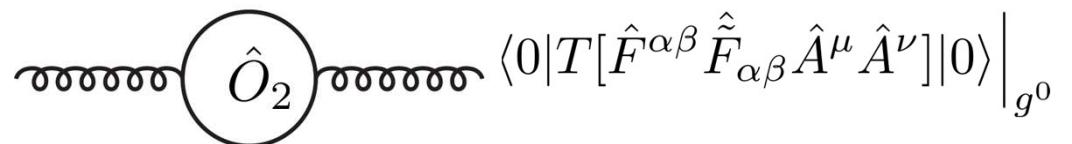
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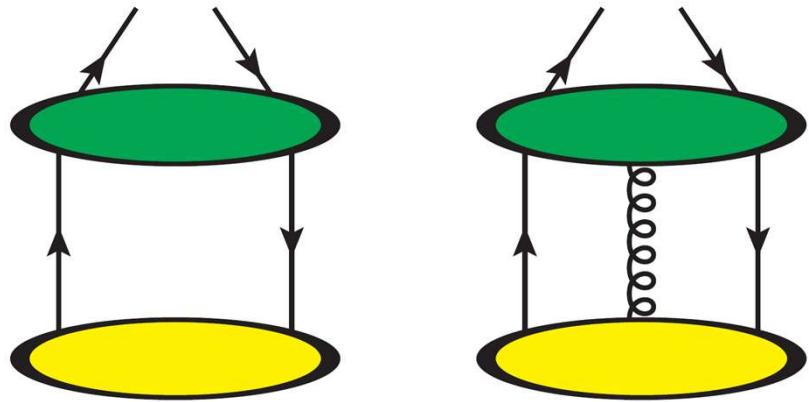
$$\partial_\mu (\hat{\bar{\psi}} \gamma^5 \gamma^\mu \hat{\psi}) = N_f \frac{g^2}{32\pi^2} \hat{F}^{\alpha\beta} \hat{\tilde{F}}_{\alpha\beta}$$

We choose \hat{O}_1 as the TMD operator and calculate perturbative corrections to it

$$D_{1T}^\perp(z_f, p_\perp) = \frac{M_h}{z_f} \frac{1}{(-4\epsilon^{p_\perp P_h w S_\perp})} \sum_X \int \frac{d\xi^+}{2\pi} \int \frac{d^2\xi_\perp}{(2\pi)^2} e^{-ip^- \xi^+ \frac{1}{z}} e^{i\vec{p}_\perp \cdot \xi_\perp} \langle 0 | \mathcal{L}_0 \psi(0) | P_h X \rangle \langle P_h X | \psi(\xi^+, \xi_\perp) \mathcal{L}_\xi^\dagger | 0 \rangle$$

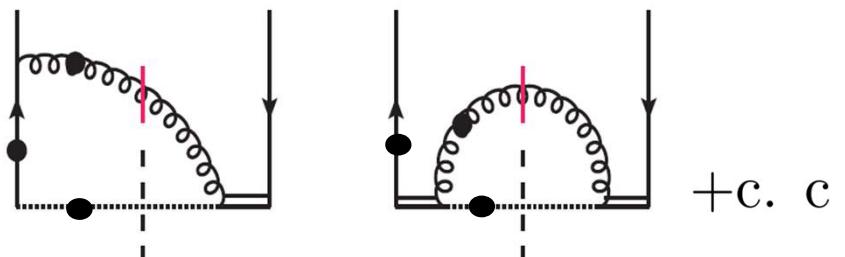
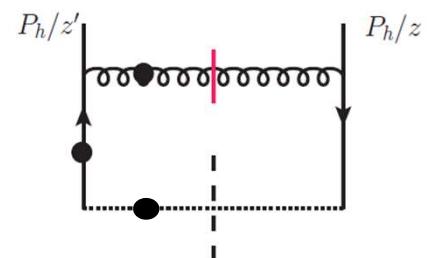
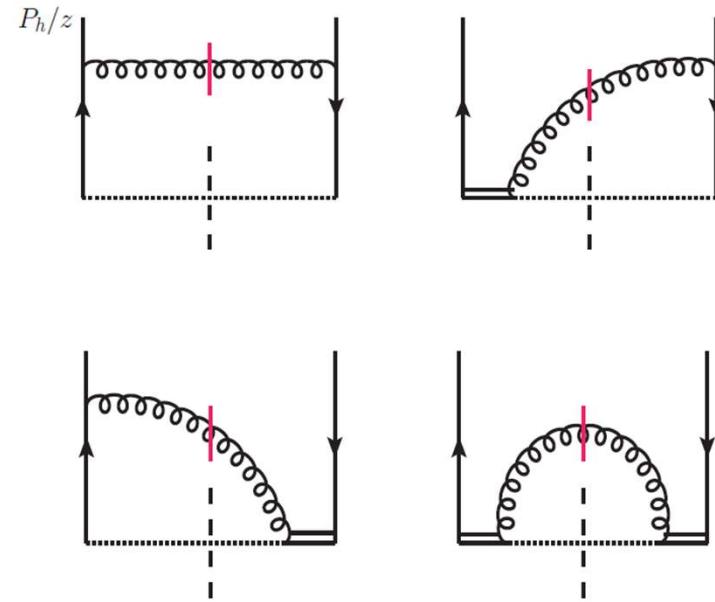
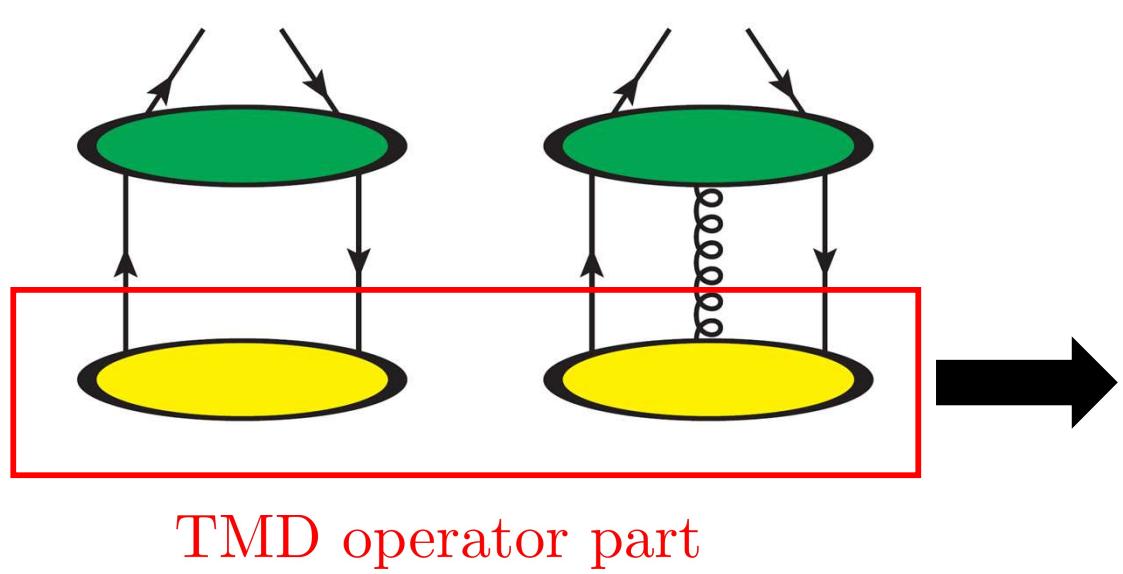
Collinear expansion

Y. Koike, A. Metz, D. Pitonyak, K. Yabe and SY, PRD95 (2017)



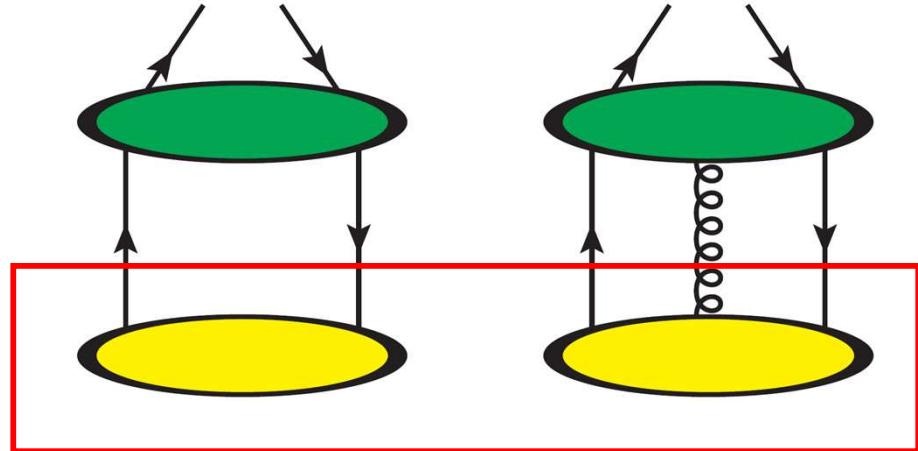
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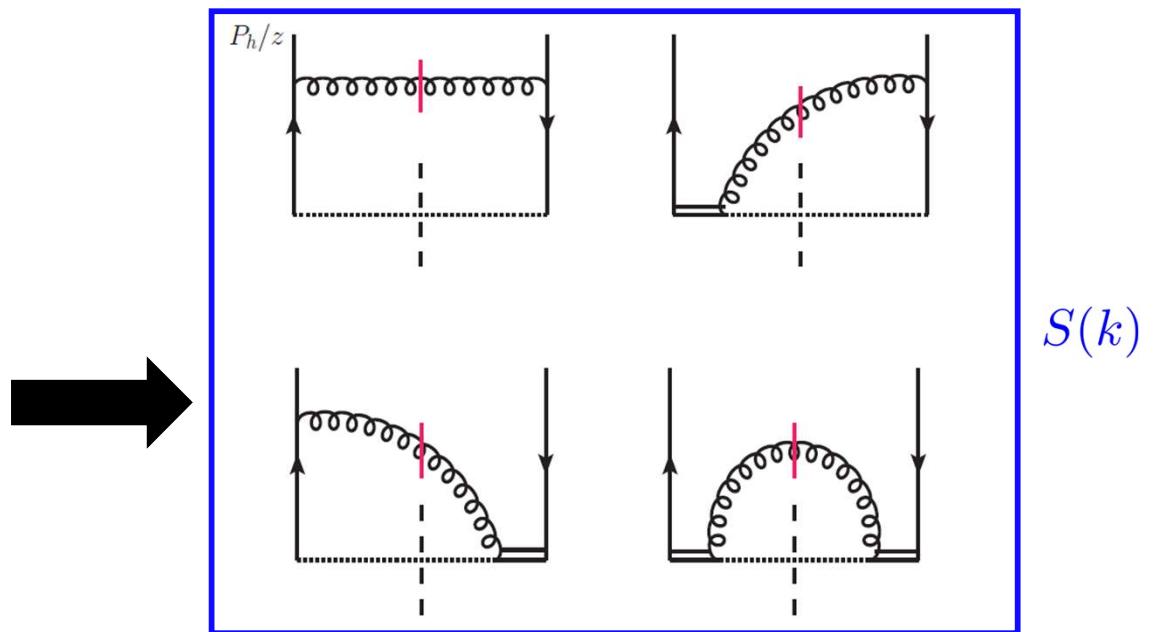


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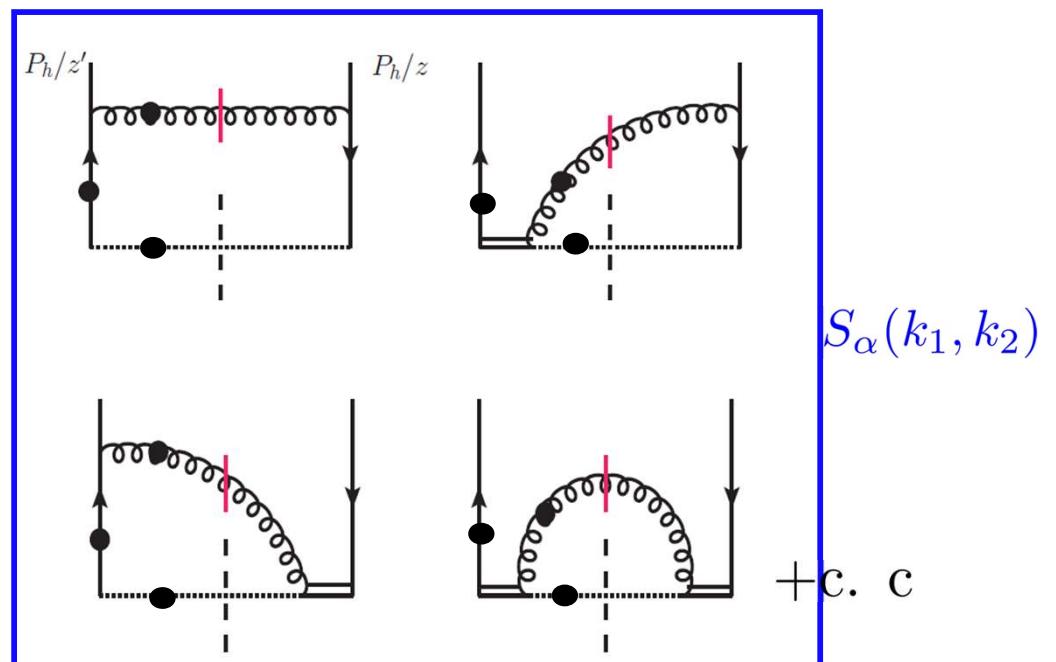


TMD operator part



Ward-Takahashi identity

$$(k_2 - k_1)^\alpha S_\alpha(k_1, k_2) = S(k_2) - S(k_1)$$



+ c. c

$$D_{1T}^\perp(z_f, z_f^2 p_\perp^2) \Big|_{\text{quark-gluon}}$$

$$\begin{aligned} &= \frac{M_h}{z_f} \frac{1}{(-4\epsilon^{p_\perp P_h w S_\perp})} \left[\int \frac{dz}{z^2} \left[\text{Tr}[\Delta(z)H(\frac{P_h}{z})] - i\text{Tr}[\Omega^\alpha{}_\beta \Delta_\partial^\beta(z) \frac{\partial}{\partial k^\alpha} H(k) \Big|_{k=\frac{P_h}{z}}] \right. \right. \\ &\quad \left. \left. + 2\text{Re} \left[i \int \frac{dz'}{z'^2} \text{Tr}[\Omega^\alpha{}_\beta \Delta_F^\beta(z', z) \left(\frac{1}{z} - \frac{1}{z'} \right) H_{L\alpha}(\frac{P_h}{z'}, \frac{P_h}{z})] \right] \right], \end{aligned}$$

operator part is written in this form

$$D_{1T}^\perp(z_f, z_f^2 p_\perp^2) \Big|_{\text{quark-gluon}} \quad \text{intrinsic} \quad \text{kinematical}$$

$$= \frac{M_h}{z_f} \frac{1}{(-4\epsilon^{p_\perp P_h w S_\perp})} \boxed{\int \frac{dz}{z^2} \left[\text{Tr}[\Delta(z) H(\frac{P_h}{z})] - i \text{Tr}[\Omega^\alpha{}_\beta \Delta_\partial^\beta(z) \frac{\partial}{\partial k^\alpha} H(k) \Big|_{k=\frac{P_h}{z}}] \right.} \\ \left. + 2 \text{Re} \left[i \int \frac{dz'}{z'^2} \text{Tr}[\Omega^\alpha{}_\beta \Delta_F^\beta(z', z) \left(\frac{1}{z} - \frac{1}{z'} \right) H_{L\alpha}(\frac{P_h}{z'}, \frac{P_h}{z})] \right], \right.$$

dynamical operator part is written in this form

$$\Delta_{ij}(z) = \frac{1}{N} \sum_X \int \frac{d\lambda}{2\pi} e^{-i\frac{\lambda}{z}} \langle 0 | [\infty w, 0] \psi_i(0) | h(P_h, S_\perp) X \rangle \langle h(P_h, S_\perp) X | \bar{\psi}_j(\lambda w) [\lambda w, \infty w] | 0 \rangle$$

$$= \dots + M_h \epsilon^{\alpha S_\perp w P_h} (\gamma_\alpha)_{ij} \frac{D_T(z)}{z} + \dots$$

$$\Delta_{\partial ij}^\alpha(z) = \frac{1}{N} \sum_X \int \frac{d\lambda}{2\pi} e^{-i\frac{\lambda}{z}} \langle 0 | [\infty w, 0] \psi_i(0) | h(P_h, S_\perp) X \rangle \langle h(P_h, S_\perp) X | \bar{\psi}_j(\lambda w) [\lambda w, \infty w] | 0 \rangle \overleftrightarrow{\partial}^\alpha$$

$$= -i M_h \epsilon^{\alpha S_\perp w P_h} (P_h)_{ij} \frac{D_{1T}^{\perp(1)}(z)}{z} + \dots,$$

$$\Delta_{Fij}^\alpha(z, z_1)$$

$$= \frac{1}{N} \sum_X \int \frac{d\lambda}{2\pi} \int \frac{d\mu}{2\pi} e^{-i\frac{\lambda}{z}} e^{-i\mu(\frac{1}{z} - \frac{1}{z_1})} \langle 0 | \psi_i(0) | h(P_h, S_\perp) X \rangle \langle h(P_h, S_\perp) X | \bar{\psi}_j(\lambda w) g F^{\alpha w}(\mu w) | 0 \rangle$$

$$= M_h \epsilon^{\alpha S_\perp w P_h} (P_h)_{ij} \frac{\widehat{D}_{FT}^*(z, z_1)}{z} - i M_h S_\perp^\alpha (\gamma_5 P_h)_{ij} \frac{\widehat{G}_{FT}^*(z, z_1)}{z} + \dots,$$

Result in TMD side

$$D_{1T}^\perp(z_h, P_{h\perp}^2) \Big|_{\text{quark-gluon}} = \frac{\alpha_s}{2\pi^2} (2M_h^2) \frac{z_f^2}{[P_{h\perp}^2]^2} \left[A + C_F D_{1T}^{\perp(1)}(z_h) \left(\ln \frac{\hat{\zeta}^2}{z_f^2 q_T^2} \right) \right]$$

$$\begin{aligned} A &= \int \frac{dz}{z} \left[\frac{D_T(z)}{z} \left(-C_F(1+2\hat{z}) - \frac{1}{2N_c} \frac{1+\hat{z}^2}{\hat{z}} \right) + \frac{d}{d(1/z)} \left(\frac{D_{1T}^{\perp(1)}}{z} \right) \left(-\frac{1}{2N_c} \frac{1+\hat{z}^2}{\hat{z}} \right) + D_{1T}^{\perp(1)} \left(C_F \frac{\hat{z}(1+\hat{z})}{(1-\hat{z})_+} \right) \right. \\ &\quad + \int d(\frac{1}{z'}) \frac{1/z}{1/z - 1/z'} \left\{ \Im \widehat{D}_{FT}(z, z') \left[\frac{1}{1/z'} \left(\frac{1}{2N_c} \frac{2-\hat{z}}{\hat{z}} \right) + \frac{1}{1/z' - 1/z_f} [C_F + \frac{1}{2N_c}] (1+\hat{z}) \right] \right. \\ &\quad \left. \left. + \Im \widehat{G}_{FT}(z, z') \left[\frac{1}{1/z'} \left(\frac{1}{2N_c} \right) - \frac{1}{1/z' - 1/z_f} [C_F + \frac{1}{2N_c}] (1-\hat{z}) \right] \right\} \right] \end{aligned}$$

Nonperturbative part is written only in terms of the collinear twist-3 functions

$$D_T, D_{1T}^{\perp(1)}, \widehat{D}_{FT}, \widehat{G}_{FT}$$

Result in TMD side

$$D_{1T}^\perp(z_h, P_{h\perp}^2) \Big|_{\text{quark-gluon}} = \frac{\alpha_s}{2\pi^2} (2M_h^2) \frac{z_f^2}{[P_{h\perp}^2]^2} \left[A + C_F D_{1T}^{\perp(1)}(z_h) \left(\ln \frac{\hat{\zeta}^2}{z_f^2 q_T^2} \right) \right]$$

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Nonperturbative part is written only in terms of the collinear twist-3 functions

- TMD cross section in $ep \rightarrow \Lambda^\uparrow X$ $D_T, D_{1T}^{\perp(1)}, \widehat{D}_{FT}, \widehat{G}_{FT}$

$$\begin{aligned} \frac{d^6\sigma}{dx_{bj} dQ^2 dz_f dq_T^2 d\phi d\chi} &= -\sigma_0 \sin \Phi_s z_f^2 \int d^2 k_\perp d^2 p_\perp d^2 \lambda_\perp \delta^2(k_\perp + p_\perp + \lambda_\perp - q_T) \frac{p_\perp^2}{q_T M_h} f(x_{bj}, k_\perp^2) D_{1T}^\perp(z_f, z_f^2 p_\perp^2) S^{-1}(\lambda_\perp^2) \\ &\quad \vec{p}_\perp, \vec{\lambda}_\perp \ll \vec{k}_\perp \simeq \vec{q}_T, \quad \vec{k}_\perp, \vec{\lambda}_\perp \ll \vec{p}_\perp \simeq \vec{q}_T, \quad \vec{k}_\perp, \vec{p}_\perp \ll \vec{\lambda}_\perp \simeq \vec{q}_T \\ &\sim -2\sigma_0 \sin \Phi_s \frac{1}{q_T} M_h \left(f(x_{bj}, k_\perp^2) D_{1T}^{\perp(1)}(z_f) + \frac{P_{h\perp}^2}{2M_h^2} f(x_{bj}) D_{1T}^\perp(z_f, P_{h\perp}^2) + f(x_{bj}) D_{1T}^{\perp(1)}(z_f) S^{-1}(q_T^2) \right), \end{aligned}$$

Result in TMD side

$$D_{1T}^\perp(z_h, P_{h\perp}^2) \Big|_{\text{quark-gluon}} = \frac{\alpha_s}{2\pi^2} (2M_h^2) \frac{z_f^2}{[P_{h\perp}^2]^2} \left[A + C_F D_{1T}^{\perp(1)}(z_h) \left(\ln \frac{\hat{\zeta}^2}{z_f^2 q_T^2} \right) \right]$$

$$\begin{aligned} A &= \int \frac{dz}{z} \left[\frac{D_T(z)}{z} \left(-C_F(1+2\hat{z}) - \frac{1}{2N_c} \frac{1+\hat{z}^2}{\hat{z}} \right) + \frac{d}{d(1/z)} \left(\frac{D_{1T}^{\perp(1)}}{z} \right) \left(-\frac{1}{2N_c} \frac{1+\hat{z}^2}{\hat{z}} \right) + D_{1T}^{\perp(1)} \left(C_F \frac{\hat{z}(1+\hat{z})}{(1-\hat{z})_+} \right) \right. \\ &\quad + \int d(\frac{1}{z'}) \frac{1/z}{1/z - 1/z'} \left\{ \Im \widehat{D}_{FT}(z, z') \left[\frac{1}{1/z'} \left(\frac{1}{2N_c} \frac{2-\hat{z}}{\hat{z}} \right) + \frac{1}{1/z' - 1/z_f} [C_F + \frac{1}{2N_c}] (1+\hat{z}) \right] \right. \\ &\quad \left. \left. + \Im \widehat{G}_{FT}(z, z') \left[\frac{1}{1/z'} \left(\frac{1}{2N_c} \right) - \frac{1}{1/z' - 1/z_f} [C_F + \frac{1}{2N_c}] (1-\hat{z}) \right] \right\} \right] \end{aligned}$$

Nonperturbative part is written only in terms of the collinear twist-3 functions

- TMD cross section in $ep \rightarrow \Lambda^\uparrow X$ $D_T, D_{1T}^{\perp(1)}, \widehat{D}_{FT}, \widehat{G}_{FT}$

$$\begin{aligned} \frac{d^6\sigma}{dx_{bj} dQ^2 dz_f dq_T^2 d\phi d\chi} &= -\sigma_0 \sin \Phi_s z_f^2 \int d^2 k_\perp d^2 p_\perp d^2 \lambda_\perp \delta^2(k_\perp + p_\perp + \lambda_\perp - q_T) \frac{p_\perp^2}{q_T M_h} f(x_{bj}, k_\perp^2) D_{1T}^\perp(z_f, z_f^2 p_\perp^2) S^{-1}(\lambda_\perp^2) \\ &\quad \vec{p}_\perp, \vec{\lambda}_\perp \ll \vec{k}_\perp \simeq \vec{q}_T, \quad \vec{k}_\perp, \vec{\lambda}_\perp \ll \vec{p}_\perp \simeq \vec{q}_T, \quad \vec{k}_\perp, \vec{p}_\perp \ll \vec{\lambda}_\perp \simeq \vec{q}_T \\ &\sim -2\sigma_0 \sin \Phi_s \frac{1}{q_T} M_h \left(f(x_{bj}, k_\perp^2) D_{1T}^{\perp(1)}(z_f) + \frac{P_{h\perp}^2}{2M_h^2} f(x_{bj}) D_{1T}^\perp(z_f, P_{h\perp}^2) + f(x_{bj}) D_{1T}^{\perp(1)}(z_f) S^{-1}(q_T^2) \right), \\ &= -2\sigma_0 \sin \Phi_s \left(\frac{1}{q_T^3} \right) \frac{\alpha_s}{2\pi^2} M_h \left\{ f(x_{bj}) \left[A \right] + C_F D_{1T}^{\perp(1)}(z_f) \int \frac{dx}{x} \left(f(x) \frac{1+\hat{x}^2}{(1-\hat{x})_+} + \frac{1}{2} G(x) [\hat{x}^2 + (1-\hat{x})^2] \right) \right. \\ &\quad \left. + 2C_F f(x_{bj}) D_{1T}^{\perp(1)}(z_f) \ln \frac{Q^2}{q_T^2} \right\} \end{aligned}$$

only in terms of the collinear functions

Result in TMD side

$$D_{1T}^\perp(z_h, P_{h\perp}^2) \Big|_{\text{quark-gluon}} = \frac{\alpha_s}{2\pi^2} (2M_h^2) \frac{z_f^2}{[P_{h\perp}^2]^2} \left[A + C_F D_{1T}^{\perp(1)}(z_h) \left(\ln \frac{\hat{\zeta}^2}{z_f^2 q_T^2} \right) \right]$$

$$\begin{aligned} A &= \int \frac{dz}{z} \left[\frac{D_T(z)}{z} \left(-C_F(1+2\hat{z}) - \frac{1}{2N_c} \frac{1+\hat{z}^2}{\hat{z}} \right) + \frac{d}{d(1/z)} \left(\frac{D_{1T}^{\perp(1)}}{z} \right) \left(-\frac{1}{2N_c} \frac{1+\hat{z}^2}{\hat{z}} \right) + D_{1T}^{\perp(1)} \left(C_F \frac{\hat{z}(1+\hat{z})}{(1-\hat{z})_+} \right) \right. \\ &\quad + \int d(\frac{1}{z'}) \frac{1/z}{1/z - 1/z'} \left\{ \Im \widehat{D}_{FT}(z, z') \left[\frac{1}{1/z'} \left(\frac{1}{2N_c} \frac{2-\hat{z}}{\hat{z}} \right) + \frac{1}{1/z' - 1/z_f} [C_F + \frac{1}{2N_c}] (1+\hat{z}) \right] \right. \\ &\quad \left. \left. + \Im \widehat{G}_{FT}(z, z') \left[\frac{1}{1/z'} \left(\frac{1}{2N_c} \right) - \frac{1}{1/z' - 1/z_f} [C_F + \frac{1}{2N_c}] (1-\hat{z}) \right] \right\} \right] \end{aligned}$$

Nonperturbative part is written only in terms of the collinear twist-3 functions

- TMD cross section in $ep \rightarrow \Lambda^\uparrow X$ $D_T, D_{1T}^{\perp(1)}, \widehat{D}_{FT}, \widehat{G}_{FT}$

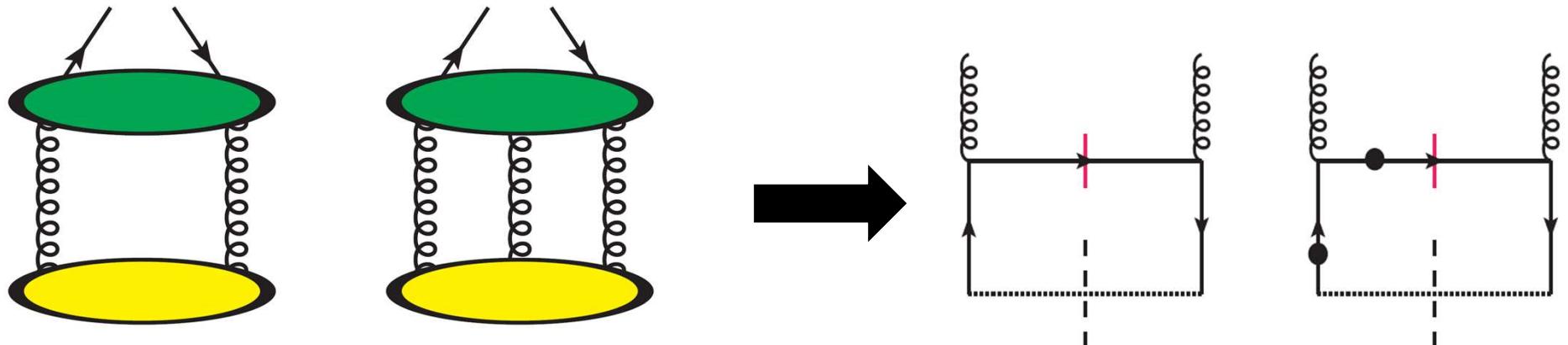
$$\begin{aligned} \frac{d^6\sigma}{dx_{bj} dQ^2 dz_f dq_T^2 d\phi d\chi} &= -\sigma_0 \sin \Phi_s z_f^2 \int d^2 k_\perp d^2 p_\perp d^2 \lambda_\perp \delta^2(k_\perp + p_\perp + \lambda_\perp - q_T) \frac{p_\perp^2}{q_T M_h} f(x_{bj}, k_\perp^2) D_{1T}^\perp(z_f, z_f^2 p_\perp^2) S^{-1}(\lambda_\perp^2) \\ &\quad \vec{p}_\perp, \vec{\lambda}_\perp \ll \vec{k}_\perp \simeq \vec{q}_T, \quad \vec{k}_\perp, \vec{\lambda}_\perp \ll \vec{p}_\perp \simeq \vec{q}_T, \quad \vec{k}_\perp, \vec{p}_\perp \ll \vec{\lambda}_\perp \simeq \vec{q}_T \\ &\sim -2\sigma_0 \sin \Phi_s \frac{1}{q_T} M_h \left(f(x_{bj}, k_\perp^2) D_{1T}^{\perp(1)}(z_f) + \frac{P_{h\perp}^2}{2M_h^2} f(x_{bj}) D_{1T}^\perp(z_f, P_{h\perp}^2) + f(x_{bj}) D_{1T}^{\perp(1)}(z_f) S^{-1}(q_T^2) \right), \\ &= -2\sigma_0 \sin \Phi_s \left(\frac{1}{q_T^3} \right) \frac{\alpha_s}{2\pi^2} M_h \left\{ f(x_{bj}) \left[A \right] + C_F D_{1T}^{\perp(1)}(z_f) \int \frac{dx}{x} \left(f(x) \frac{1+\hat{x}^2}{(1-\hat{x})_+} + \frac{1}{2} G(x) [\hat{x}^2 + (1-\hat{x})^2] \right) \right. \\ &\quad \left. + 2C_F f(x_{bj}) D_{1T}^{\perp(1)}(z_f) \ln \frac{Q^2}{q_T^2} \right\} \end{aligned}$$

only in terms of the collinear functions

What is the relation with the collinear twist-3 result ?

- gluon mixing

R. Ikarashi, Y. Koike, K. Yabe and SY, PRD106 (2022)



$$D_{1T}^\perp(z_f, p_\perp'^2) \Big|_{\text{3-gluon}} = -\frac{\alpha_s}{2\pi^2}(M_h^2) \frac{z_f^2}{[p_\perp^2]^2} [B]$$

$$\begin{aligned}
 B = & \int dz z \left[\hat{G}_T^{(1)}(z) \sigma_1 + \Delta \hat{H}_T^{(1)}(z) \sigma_2 \right. \\
 & + \int d(\frac{1}{z'}) \frac{1}{1/z - 1/z'} \Im \left(\hat{N}_1(\frac{1}{z'}, \frac{1}{z}) \sigma_{N1} + \hat{N}_2(\frac{1}{z'}, \frac{1}{z}) \sigma_{N2} + \hat{N}_2(\frac{1}{z} - \frac{1}{z'}, \frac{1}{z}) \sigma_{N3} \right) \\
 & + \int d(\frac{1}{z'}) \frac{1/z}{(1/z - 1/z')^2} \Im \left(\hat{N}_1(\frac{1}{z'}, \frac{1}{z}) \sigma_{DN1} + \hat{N}_2(\frac{1}{z'}, \frac{1}{z}) \sigma_{DN1} + \hat{N}_2(\frac{1}{z} - \frac{1}{z'}, \frac{1}{z}) \sigma_{DN3} \right) \\
 & + \int d(\frac{1}{z'}) \frac{1}{1/z - 1/z'} \Im \left(\hat{O}_1(\frac{1}{z'}, \frac{1}{z}) \sigma_{O1} + \hat{O}_2(\frac{1}{z'}, \frac{1}{z}) \sigma_{O2} + \hat{O}_2(\frac{1}{z} - \frac{1}{z'}, \frac{1}{z}) \sigma_{O2} \right) \\
 & + \int d(\frac{1}{z'}) \frac{1/z}{(1/z - 1/z')^2} \Im \left(\hat{O}_1(\frac{1}{z'}, \frac{1}{z}) \sigma_{DN1} + \hat{O}_2(\frac{1}{z'}, \frac{1}{z}) \sigma_{DN1} + \hat{O}_2(\frac{1}{z} - \frac{1}{z'}, \frac{1}{z}) (-\sigma_{DN3}) \right) \\
 & + \frac{1}{C_F} \int d(\frac{1}{z'}) \left\{ \Im \hat{D}_{FT}(\frac{1}{z'}, \frac{1}{z'} - \frac{1}{z}) \left(-\frac{1}{2} \sigma_{N1} + \frac{1/z}{1/z - 1/z'} \sigma_{DF2} + \frac{1/z}{1/z' - 1/z_f} \sigma_{DF4} \right) \right. \\
 & \left. + \Im \hat{G}_{FT}(\frac{1}{z'}, \frac{1}{z'} - \frac{1}{z}) \left(\frac{1/z}{1/z - 1/z'} \sigma_{GF2} + \frac{1/z}{1/z' - 1/z_f} \sigma_{GF4} \right) \right\}
 \end{aligned}$$

Collinear twist-3 side

Y. Koike, K. Takada, S. Usui, K. Yabe and SY, PRD105 (2022)

$$\begin{aligned}
& \frac{d^6 \Delta \sigma^{\text{tw3-frag}}}{dx_{bj} dQ^2 dz_f dq_T^2 d\phi d\chi} \\
&= \frac{\alpha_{em}^2 \alpha_s (-M_h)}{16\pi^2 x_{bj}^2 S_{ep}^2 Q^2} \sum_{k=1}^9 A_k (\phi - \chi) \mathcal{S}_k \int \frac{dx}{x} f_1(x) \int \frac{dz}{z} \delta \left(\frac{q_T^2}{Q^2} - \left(1 - \frac{1}{\hat{x}} \right) \left(1 - \frac{1}{\hat{z}} \right) \right) \times \left[\frac{D_T(z)}{z} \hat{\sigma}_T^k - \left\{ \frac{d}{d(1/z)} \frac{D_{1T}^{\perp(1)}(z)}{z} \right\} \hat{\sigma}_{\perp D}^k - D_{1T}^{\perp(1)}(z) \hat{\sigma}_{\perp}^k \right. \\
&\quad + \int \frac{dz'}{z'^2} P \left(\frac{1}{1/z - 1/z'} \right) \left\{ \Im \widehat{D}_{FT}(z, z') \left[\frac{z'}{z} \hat{\sigma}_{DF1}^k + \frac{1}{z} \left(\frac{1}{z'} - \frac{1}{(1 - q_T^2/Q^2) z_f} \right)^{-1} \hat{\sigma}_{DF2}^k \right] \right. \\
&\quad \left. \left. + \Im \widehat{G}_{FT}(z, z') \left[\frac{z'}{z} \hat{\sigma}_{GF1}^k + \frac{1}{z} \left(\frac{1}{z'} - \frac{1}{(1 - q_T^2/Q^2) z_f} \right)^{-1} \hat{\sigma}_{GF2}^k \right] \right\} \right].
\end{aligned}$$

Collinear twist-3 side

Y. Koike, K. Takada, S. Usui, K. Yabe and SY, PRD105 (2022)

$$\begin{aligned}
 & \frac{d^6 \Delta \sigma^{\text{tw3-frag}}}{dx_{bj} dQ^2 dz_f dq_T^2 d\phi d\chi} \\
 &= \frac{\alpha_{em}^2 \alpha_s (-M_h)}{16\pi^2 x_{bj}^2 S_{ep}^2 Q^2} \sum_{k=1}^9 A_k (\phi - \chi) \mathcal{S}_k \int \frac{dx}{x} f_1(x) \int \frac{dz}{z} \delta \left(\frac{q_T^2}{Q^2} - \left(1 - \frac{1}{\hat{x}} \right) \left(1 - \frac{1}{\hat{z}} \right) \right) \\
 & \quad \times \left[\frac{D_T(z)}{z} \hat{\sigma}_T^k - \left\{ \frac{d}{d(1/z)} \frac{D_{1T}^{\perp(1)}(z)}{z} \right\} \hat{\sigma}_{\perp D}^k - D_{1T}^{\perp(1)}(z) \hat{\sigma}_{\perp}^k \right. \\
 & \quad + \int \frac{dz'}{z'^2} P \left(\frac{1}{1/z - 1/z'} \right) \left\{ \Im \widehat{D}_{FT}(z, z') \left[\frac{z'}{z} \hat{\sigma}_{DF1}^k + \frac{1}{z} \left(\frac{1}{z'} - \frac{1}{(1 - q_T^2/Q^2) z_f} \right)^{-1} \hat{\sigma}_{DF2}^k \right] \right. \\
 & \quad \left. \left. + \Im \widehat{G}_{FT}(z, z') \left[\frac{z'}{z} \hat{\sigma}_{GF1}^k + \frac{1}{z} \left(\frac{1}{z'} - \frac{1}{(1 - q_T^2/Q^2) z_f} \right)^{-1} \hat{\sigma}_{GF2}^k \right] \right\} \right].
 \end{aligned}$$

84 hard cross sections

$$\begin{aligned}
 & q\gamma \rightarrow q \\
 & q\gamma \rightarrow g
 \end{aligned}$$

Collinear twist-3 side

Y. Koike, K. Takada, S. Usui, K. Yabe and SY, PRD105 (2022)

$$\begin{aligned}
 & \frac{d^6 \Delta \sigma^{\text{tw3-frag}}}{dx_{bj} dQ^2 dz_f dq_T^2 d\phi d\chi} \\
 &= \frac{\alpha_{em}^2 \alpha_s (-M_h)}{16\pi^2 x_{bj}^2 S_{ep}^2 Q^2} \sum_{k=1}^9 A_k (\phi - \chi) \mathcal{S}_k \int \frac{dx}{x} f_1(x) \int \frac{dz}{z} \delta \left(\frac{q_T^2}{Q^2} - \left(1 - \frac{1}{\hat{x}} \right) \left(1 - \frac{1}{\hat{z}} \right) \right) \\
 & \quad \times \left[\frac{D_T(z)}{z} \hat{\sigma}_T^k - \left\{ \frac{d}{d(1/z)} \frac{D_{1T}^{\perp(1)}(z)}{z} \right\} \hat{\sigma}_{\perp D}^k - D_{1T}^{\perp(1)}(z) \hat{\sigma}_{\perp}^k \right. \\
 & \quad + \int \frac{dz'}{z'^2} P \left(\frac{1}{1/z - 1/z'} \right) \left\{ \Im \widehat{D}_{FT}(z, z') \left[\frac{z'}{z} \hat{\sigma}_{DF1}^k + \frac{1}{z} \left(\frac{1}{z'} - \frac{1}{(1 - q_T^2/Q^2) z_f} \right)^{-1} \hat{\sigma}_{DF2}^k \right] \right. \\
 & \quad \left. \left. + \Im \widehat{G}_{FT}(z, z') \left[\frac{z'}{z} \hat{\sigma}_{GF1}^k + \frac{1}{z} \left(\frac{1}{z'} - \frac{1}{(1 - q_T^2/Q^2) z_f} \right)^{-1} \hat{\sigma}_{GF2}^k \right] \right\} \right].
 \end{aligned}$$

84 hard cross sections

$$\begin{aligned}
 & q\gamma \rightarrow q \\
 & q\gamma \rightarrow g
 \end{aligned}$$

$$\hat{\sigma}_1^1 = 4 \frac{Q^2}{q_T^3} (1 - \hat{z}) \frac{1 - 9\hat{x} + 14\hat{x}^2 - 6\hat{x}^3 + 2\hat{z}^2(1 - 6\hat{x} + 6\hat{x}^2) + \hat{z}(-4 + 23\hat{x} - 30\hat{x}^2 + 12\hat{x}^3)}{\hat{x}\hat{z}(1 - \hat{x} - \hat{z})}$$

Collinear twist-3 side

Y. Koike, K. Takada, S. Usui, K. Yabe and SY, PRD105 (2022)

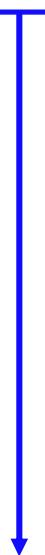
$$\frac{d^6 \Delta \sigma^{\text{tw3-frag}}}{dx_{bj} dQ^2 dz_f dq_T^2 d\phi d\chi}$$

$$= \frac{\alpha_{em}^2 \alpha_s (-M_h)}{16\pi^2 x_{bj}^2 S_{ep}^2 Q^2} \sum_{k=1}^9 A_k (\phi - \chi) \mathcal{S}_k \int \frac{dx}{x} f_1(x) \int \frac{dz}{z} \delta \left(\frac{q_T^2}{Q^2} - \left(1 - \frac{1}{\hat{x}}\right) \left(1 - \frac{1}{\hat{z}}\right) \right)$$

—————

$$\begin{aligned} & \times \left[\frac{D_T(z)}{z} \hat{\sigma}_T^k - \left\{ \frac{d}{d(1/z)} \frac{D_{1T}^{\perp(1)}(z)}{z} \right\} \hat{\sigma}_{\perp D}^k - D_{1T}^{\perp(1)}(z) \hat{\sigma}_{\perp}^k \right. \\ & + \int \frac{dz'}{z'^2} P \left(\frac{1}{1/z - 1/z'} \right) \left\{ \Im \widehat{D}_{FT}(z, z') \left[\frac{z'}{z} \hat{\sigma}_{DF1}^k + \frac{1}{z} \left(\frac{1}{z'} - \frac{1}{(1 - q_T^2/Q^2) z_f} \right)^{-1} \hat{\sigma}_{DF2}^k \right] \right. \\ & \left. \left. + \Im \widehat{G}_{FT}(z, z') \left[\frac{z'}{z} \hat{\sigma}_{GF1}^k + \frac{1}{z} \left(\frac{1}{z'} - \frac{1}{(1 - q_T^2/Q^2) z_f} \right)^{-1} \hat{\sigma}_{GF2}^k \right] \right\} \right]. \end{aligned}$$

$$q_T^2 \ll Q^2$$



84 hard cross sections

$$\begin{aligned} q\gamma &\rightarrow q \\ q\gamma &\rightarrow g \end{aligned}$$

X. Ji, J.-W. Qiu, W. Vogelsang and F. Yuan, PRD73(2006)
PLB638(2006)

$$\hat{x} \hat{z} \left[\frac{1}{(1 - \hat{z})_+} \delta(1 - \hat{x}) + \frac{1}{(1 - \hat{x})_+} \delta(1 - \hat{z}) + \delta(1 - \hat{x}) \delta(1 - \hat{z}) \ln \frac{Q^2}{q_T^2} \right]$$

$$\hat{\sigma}_1^1 = 4 \frac{Q^2}{q_T^3} (1 - \hat{z}) \frac{1 - 9\hat{x} + 14\hat{x}^2 - 6\hat{x}^3 + 2\hat{z}^2(1 - 6\hat{x} + 6\hat{x}^2) + \hat{z}(-4 + 23\hat{x} - 30\hat{x}^2 + 12\hat{x}^3)}{\hat{x} \hat{z} (1 - \hat{x} - \hat{z})}$$

Collinear twist-3 side

Y. Koike, K. Takada, S. Usui, K. Yabe and SY, PRD105 (2022)

$$\frac{d^6 \Delta \sigma^{\text{tw3-frag}}}{dx_{bj} dQ^2 dz_f dq_T^2 d\phi d\chi}$$

$$= \frac{\alpha_{em}^2 \alpha_s (-M_h)}{16\pi^2 x_{bj}^2 S_{ep}^2 Q^2} \sum_{k=1}^9 A_k (\phi - \chi) \mathcal{S}_k \int \frac{dx}{x} f_1(x) \int \frac{dz}{z} \delta \left(\frac{q_T^2}{Q^2} - \left(1 - \frac{1}{\hat{x}}\right) \left(1 - \frac{1}{\hat{z}}\right) \right)$$

—————

$$\begin{aligned} & \times \left[\frac{D_T(z)}{z} \hat{\sigma}_T^k - \left\{ \frac{d}{d(1/z)} \frac{D_{1T}^{\perp(1)}(z)}{z} \right\} \hat{\sigma}_{\perp D}^k - D_{1T}^{\perp(1)}(z) \hat{\sigma}_{\perp}^k \right. \\ & + \int \frac{dz'}{z'^2} P \left(\frac{1}{1/z - 1/z'} \right) \left\{ \Im \widehat{D}_{FT}(z, z') \left[\frac{z'}{z} \hat{\sigma}_{DF1}^k + \frac{1}{z} \left(\frac{1}{z'} - \frac{1}{(1 - q_T^2/Q^2) z_f} \right)^{-1} \hat{\sigma}_{DF2}^k \right] \right. \\ & \left. \left. + \Im \widehat{G}_{FT}(z, z') \left[\frac{z'}{z} \hat{\sigma}_{GF1}^k + \frac{1}{z} \left(\frac{1}{z'} - \frac{1}{(1 - q_T^2/Q^2) z_f} \right)^{-1} \hat{\sigma}_{GF2}^k \right] \right\} \right]. \end{aligned}$$

$$q_T^2 \ll Q^2$$

84 hard cross sections

$$q\gamma \rightarrow q$$

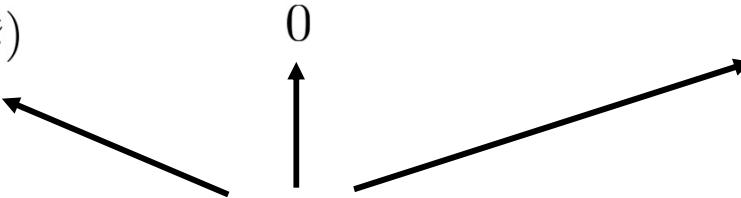
$$q\gamma \rightarrow g$$

X. Ji, J.-W. Qiu, W. Vogelsang and F. Yuan, PRD73(2006)
PLB638(2006)

$$\hat{x}\hat{z} \left[\frac{1}{(1-\hat{z})_+} \delta(1-\hat{x}) + \frac{1}{(1-\hat{x})_+} \delta(1-\hat{z}) + \delta(1-\hat{x})\delta(1-\hat{z}) \ln \frac{Q^2}{q_T^2} \right]$$

$$\hat{\sigma}_1^1 = -4C_F \frac{Q^2}{q_T^3} (1 + 2\hat{z})$$

0



$$\hat{\sigma}_1^1 = 4 \frac{Q^2}{q_T^3} (1 - \hat{z}) \frac{1 - 9\hat{x} + 14\hat{x}^2 - 6\hat{x}^3 + 2\hat{z}^2(1 - 6\hat{x} + 6\hat{x}^2) + \hat{z}(-4 + 23\hat{x} - 30\hat{x}^2 + 12\hat{x}^3)}{\hat{x}\hat{z}(1 - \hat{x} - \hat{z})}$$

$$\frac{d^6\Delta\sigma^{\rm tw3-frag}}{dx_{bj}dQ^2dz_f dq_T^2 d\phi d\chi}$$

$$= -2\sigma_0 \sin \Phi_s \Big(\frac{1}{q_T^3} \Big) \frac{\alpha_s}{2\pi^2} M_h \Big\{ f(x_{bj}) \Big[A \Big] + C_F D_{1T}^{\perp(1)}(z_f) \int \frac{dx}{x} \Big(f(x) \frac{1+\hat{x}^2}{(1-\hat{x})_+} + \frac{1}{2} G(x) [\hat{x}^2 + (1-\hat{x})^2] \Big) \\ + 2C_F f(x_{bj}) D_{1T}^{\perp(1)}(z_f) \ln \frac{Q^2}{q_T^2} \Big\}$$

$$A~=~\int \frac{dz}{z}\Bigg[\frac{D_T(z)}{z}\Big(-C_F(1+2\hat{z})-\frac{1}{2N_c}\frac{1+\hat{z}^2}{\hat{z}}\Big)+\frac{d}{d(1/z)}\Big(\frac{D_{1T}^{\perp(1)}}{z}\Big)\Big(-\frac{1}{2N_c}\frac{1+\hat{z}^2}{\hat{z}}\Big)+D_{1T}^{\perp(1)}\Big(C_F\frac{\hat{z}(1+\hat{z})}{(1-\hat{z})_+}\Big)\\+\int d(\frac{1}{z'})\frac{1/z}{1/z-1/z'}\Big\{\Im\widehat{D}_{FT}(z,z')\Big[\frac{1}{1/z'}\Big(\frac{1}{2N_c}\frac{2-\hat{z}}{\hat{z}}\Big)+\frac{1}{1/z'-1/z_f}[C_F+\frac{1}{2N_c}](1+\hat{z})\Big]\\+\Im\widehat{G}_{FT}(z,z')\Big[\frac{1}{1/z'}\Big(\frac{1}{2N_c}\Big)-\frac{1}{1/z'-1/z_f}[C_F+\frac{1}{2N_c}](1-\hat{z})\Big]\Big\}\Bigg]$$

Consistent with the TMD result !

- twist-3 gluon fragmentation effect

R. Ikarashi, Y. Koike, K. Yabe and SY, PRD105 (2022)

$$g\gamma \rightarrow q$$

$$\begin{aligned}
& \frac{d^6\sigma}{dx_{bj}dQ^2dz_f dq_T^2 d\phi d\chi} \\
&= \frac{\alpha_{em}^2 \alpha_s M_h}{16\pi^2 x_{bj}^2 S_{ep}^2 Q^2} \sum_k \mathcal{A}_k(\phi - \chi) \mathcal{S}_k \int_{x_{min}}^1 \frac{dx}{x} \int_{z_{min}}^1 \frac{dz}{z} z^2 f_1(x) \delta\left(\frac{q_T^2}{Q^2} - \left(1 - \frac{1}{\hat{x}}\right) \left(1 - \frac{1}{\hat{z}}\right)\right) \\
&\times \left\{ \widehat{G}_T^{(1)}(z) \hat{\sigma}_G^k + \Delta \widehat{H}_T^{(1)}(z) \hat{\sigma}_H^k \right. \\
&+ \int d\left(\frac{1}{z'}\right) \left[\frac{1}{1/z - 1/z'} \Im \left(\widehat{N}_1 \left(\frac{1}{z'}, \frac{1}{z} \right) \hat{\sigma}_{N1}^k + \widehat{N}_2 \left(\frac{1}{z'}, \frac{1}{z} \right) \hat{\sigma}_{N2}^k + \widehat{N}_2 \left(\frac{1}{z} - \frac{1}{z'}, \frac{1}{z} \right) \hat{\sigma}_{N3}^k \right) \right. \\
&+ \frac{1}{z} \left(\frac{1}{1/z - 1/z'} \right)^2 \Im \left(\widehat{N}_1 \left(\frac{1}{z'}, \frac{1}{z} \right) \hat{\sigma}_{DN1}^k + \widehat{N}_2 \left(\frac{1}{z'}, \frac{1}{z} \right) \hat{\sigma}_{DN2}^k + \widehat{N}_2 \left(\frac{1}{z} - \frac{1}{z'}, \frac{1}{z} \right) \hat{\sigma}_{DN3}^k \right) \\
&+ \frac{1}{1/z - 1/z'} \Im \left(\widehat{O}_1 \left(\frac{1}{z'}, \frac{1}{z} \right) \hat{\sigma}_{O1}^k + \widehat{O}_2 \left(\frac{1}{z'}, \frac{1}{z} \right) \hat{\sigma}_{O2}^k + \widehat{O}_2 \left(\frac{1}{z} - \frac{1}{z'}, \frac{1}{z} \right) \hat{\sigma}_{O3}^k \right) \\
&+ \frac{1}{z} \left(\frac{1}{1/z - 1/z'} \right)^2 \Im \left(\widehat{O}_1 \left(\frac{1}{z'}, \frac{1}{z} \right) \hat{\sigma}_{DO1}^k + \widehat{O}_2 \left(\frac{1}{z'}, \frac{1}{z} \right) \hat{\sigma}_{DO2}^k + \widehat{O}_2 \left(\frac{1}{z} - \frac{1}{z'}, \frac{1}{z} \right) \hat{\sigma}_{DO3}^k \right) \Big] \longrightarrow \\
&+ \int d\left(\frac{1}{z'}\right) \frac{2}{C_F} \left[\Im \widetilde{D}_{FT} \left(\frac{1}{z'}, \frac{1}{z'} - \frac{1}{z} \right) \left(\hat{\sigma}_{DF1}^k + \frac{1}{z} \frac{1}{1/z - 1/z'} \hat{\sigma}_{DF2}^k + \frac{z'}{z} \hat{\sigma}_{DF3}^k \right. \right. \\
&+ \frac{1}{1 - (1 - q_T^2/Q^2)z_f/z'} \hat{\sigma}_{DF4}^k + \frac{1}{1 - (1 - q_T^2/Q^2)z_f(1/z - 1/z')} \hat{\sigma}_{DF5}^k \Big) \\
&+ \Im \widetilde{G}_{FT} \left(\frac{1}{z'}, \frac{1}{z'} - \frac{1}{z} \right) \left(\frac{1}{z} \frac{1}{1/z - 1/z'} \hat{\sigma}_{GF2}^k + \frac{z'}{z} \hat{\sigma}_{GF3}^k \right. \\
&+ \frac{1}{1 - (1 - q_T^2/Q^2)z_f/z'} \hat{\sigma}_{GF4}^k + \frac{1}{1 - (1 - q_T^2/Q^2)z_f(1/z - 1/z')} \hat{\sigma}_{GF5}^k \Big) \Big] \Big\}
\end{aligned}$$

Consistent with the gluon-mixing term in TMD !

Summary

- The twist-3 cross section for the SSA in Λ^\dagger production was completed very recent at the lowest order
- We calculated the one-loop corrections to the TMD operator and found the relation with the collinear twist-3 functions
- Using this relation, we confirmed that the TMD and the collinear twist-3 give the consistent results in $\Lambda_{QCD} \ll P_T \ll Q$
- The consistency will play a role in future phenomenological studies as that in the pion production has done

Back up

pQCD frameworks for the SSA

Two possible ways are known as extensional pQCD frameworks

$$\text{Conventional parton model} \quad \int \frac{d\lambda}{2\pi} e^{i\lambda x} \langle p | \bar{\psi}_j(0) \psi(\lambda n) | p \rangle = \underbrace{i(\gamma^5 \sigma^{\alpha\beta} P_\beta)_{ij}}_{\text{Chiral-odd}} \times \underbrace{S_{\perp\alpha}}_{\text{The parent proton is transversely polarized}} \times h_1(x)$$

Transverse Momentum Dependent(TMD) factorization

$$\int \frac{d\lambda}{2\pi} \int \frac{d^2\xi_\perp}{2\pi} e^{i\lambda x} e^{-i\vec{\xi}_\perp \cdot \vec{k}_T} \langle p | \bar{\psi}_j(0) \psi(\lambda n, x_\perp) | p \rangle = \frac{1}{2} f(x, \vec{k}_T) (\not{p})_{ij} + \underbrace{\frac{1}{2M_N} (\gamma_\alpha)_{ji} \epsilon^{\alpha\nu\rho\sigma} P_\nu \vec{k}_T \not{p}_\rho}_{\text{Chiral-even!}} S_{\perp\sigma} f_{1T}^\perp(x, \vec{k}_T) \cdots$$

- Applicable in small $P_T(Q \gg P_T \geq \Lambda_{QCD})$ region $\frac{m_q}{P_T} \rightarrow \frac{\Lambda_{QCD}}{P_T}$

Twist-3 mechanism in the collinear factorization

$$\int \frac{d\lambda}{2\pi} \int \frac{d\mu}{2\pi} e^{i\lambda x} e^{i\mu(x_2 - x_1)} \langle p | \bar{\psi}_j(0) g F^{\alpha+}(\mu n) \psi(\lambda n) | p \rangle = -\frac{M_N}{2} \underbrace{\epsilon^{\alpha\beta-+}(\not{p})_{ij}}_{\text{Chiral-even!}} \underbrace{S_{\perp\beta}}_{\text{Chiral-even!}} F_{FT}(x_1, x_2) + \cdots$$

- Applicable in large $P_T(P_T \gg \Lambda_{QCD})$ region

There is overlapping region $\Lambda_{QCD} \ll P_T \ll Q$

There should not be a quantitative ambiguity in this region if both are the first principle calculations

$$\begin{aligned} \frac{d^6\sigma}{dx_{bj}dQ^2dz_f dq_T^2 d\phi d\chi} &= -\sigma_0 \sin \Phi_s z_f^2 \int d^2k_\perp d^2p_\perp d^2\lambda_\perp \delta^2(k_\perp + p_\perp + \lambda_\perp - q_T) \frac{p_\perp^2}{q_T M_h} f(x_{bj}, k_\perp^2) D_{1T}^\perp(z_f, z_f^2 p_\perp^2) S^{-1}(\lambda_\perp^2) \\ &\quad \vec{p}_\perp, \vec{\lambda}_\perp \ll \vec{k}_\perp \simeq \vec{q}_T, \quad \vec{k}_\perp, \vec{\lambda}_\perp \ll \vec{p}_\perp \simeq \vec{q}_T, \quad \vec{k}_\perp, \vec{p}_\perp \ll \vec{\lambda}_\perp \simeq \vec{q}_T \\ &\sim -2\sigma_0 \sin \Phi_s \frac{1}{q_T} M_h \left(f(x_{bj}, k_\perp^2) D_{1T}^{\perp(1)}(z_f) + \frac{P_{h\perp}^2}{2M_h^2} f(x_{bj}) D_{1T}^\perp(z_f, P_{h\perp}^2) + f(x_{bj}) D_{1T}^{\perp(1)}(z_f) S^{-1}(q_T^2) \right), \end{aligned}$$

$$\begin{aligned} \text{X. Ji, J.-W. Qiu, W. Vogelsang and F. Yuan, PRD73(2006)} \\ \text{PLB638(2006)} \end{aligned}$$

$$f(x_{bj}, q_T^2) = \frac{\alpha_s}{2\pi^2} \frac{1}{q_T^2} C_F \int \frac{dx}{x} f(x) \left[\frac{1 + \hat{x}^2}{(1 - \hat{x})_+} + \delta(1 - \hat{x}) \left(\ln \frac{x_{bj}^2 \zeta^2}{q_T^2} - 1 \right) \right]$$

$$S^{-1}(q_T^2) = -\frac{\alpha_s}{2\pi^2} \frac{1}{q_T^2} C_F (\ln \rho^2 - 1), \quad \ln \frac{x_{bj}^2 \zeta^2 \hat{\zeta}^2}{z_f^2 \rho^2} = \ln Q^4$$

Wilson line

F. Yuan and J. Zhou, PRL 103 (2009)

$$D_{1T}^\perp(z_f,p_\perp) = \frac{M_h}{z_f} \frac{1}{(-4\epsilon^{p_\perp P_h w S_\perp})} \sum_X \int \frac{d\xi^+}{2\pi} \int \frac{d^2\xi_\perp}{(2\pi)^2} e^{-ip^-\xi^+\frac{1}{z}} e^{i\vec{p}_\perp \cdot \xi_\perp} \langle 0 | \mathcal{L}_0 \psi(0) | P_h X \rangle \langle P_h X | \psi(\xi^+, \xi_\perp) \mathcal{L}_\xi^\dagger | 0 \rangle$$

$$\mathcal{L}_\xi = \exp\left(-ig \int_0^\infty d\lambda v \cdot A(\lambda v + \xi)\right)$$