



The XXXI International Workshop on Deep Inelastic Scattering and Related Subjects (DIS2024)

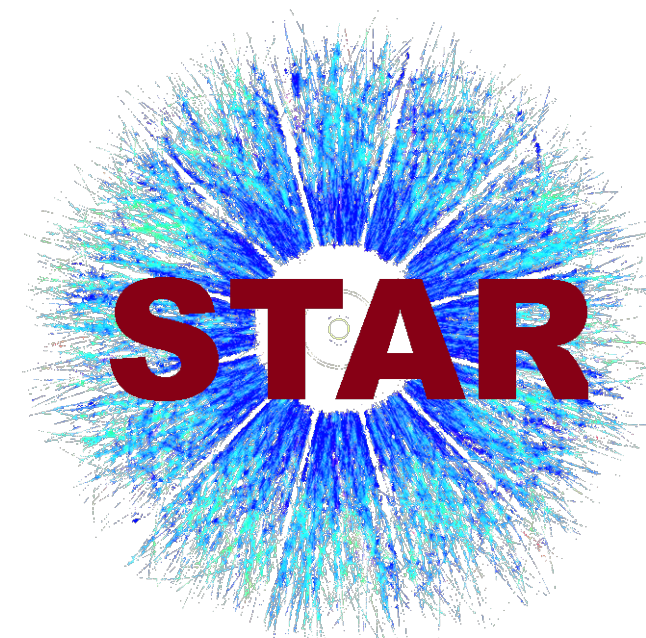
Apr 8–12, 2024

Maison MINATEC, Grenoble, FRANCE

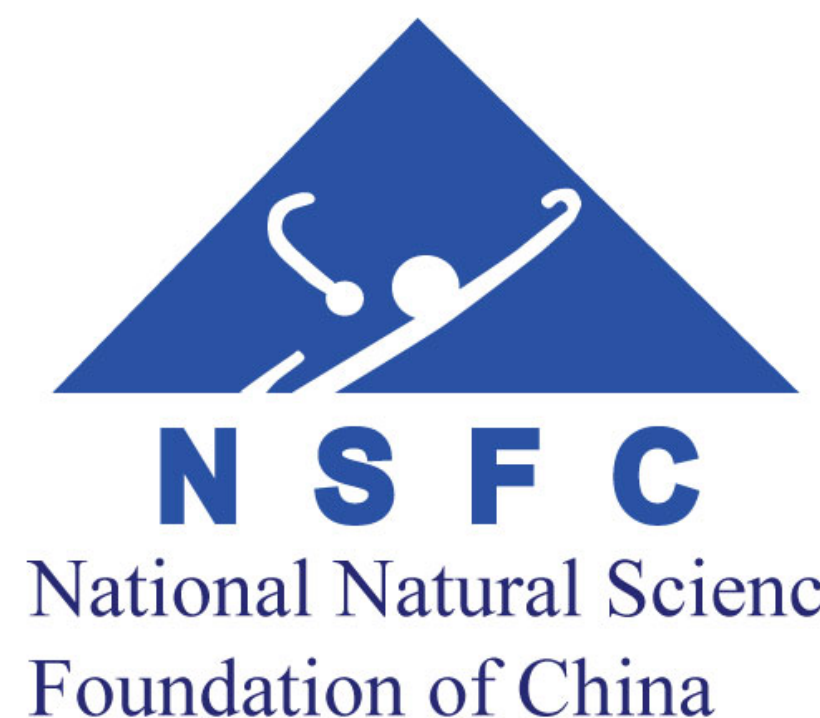
# Probing gluon and strange quark helicity distribution in the proton at STAR

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山东大学  
SHANDONG UNIVERSITY





# Outline

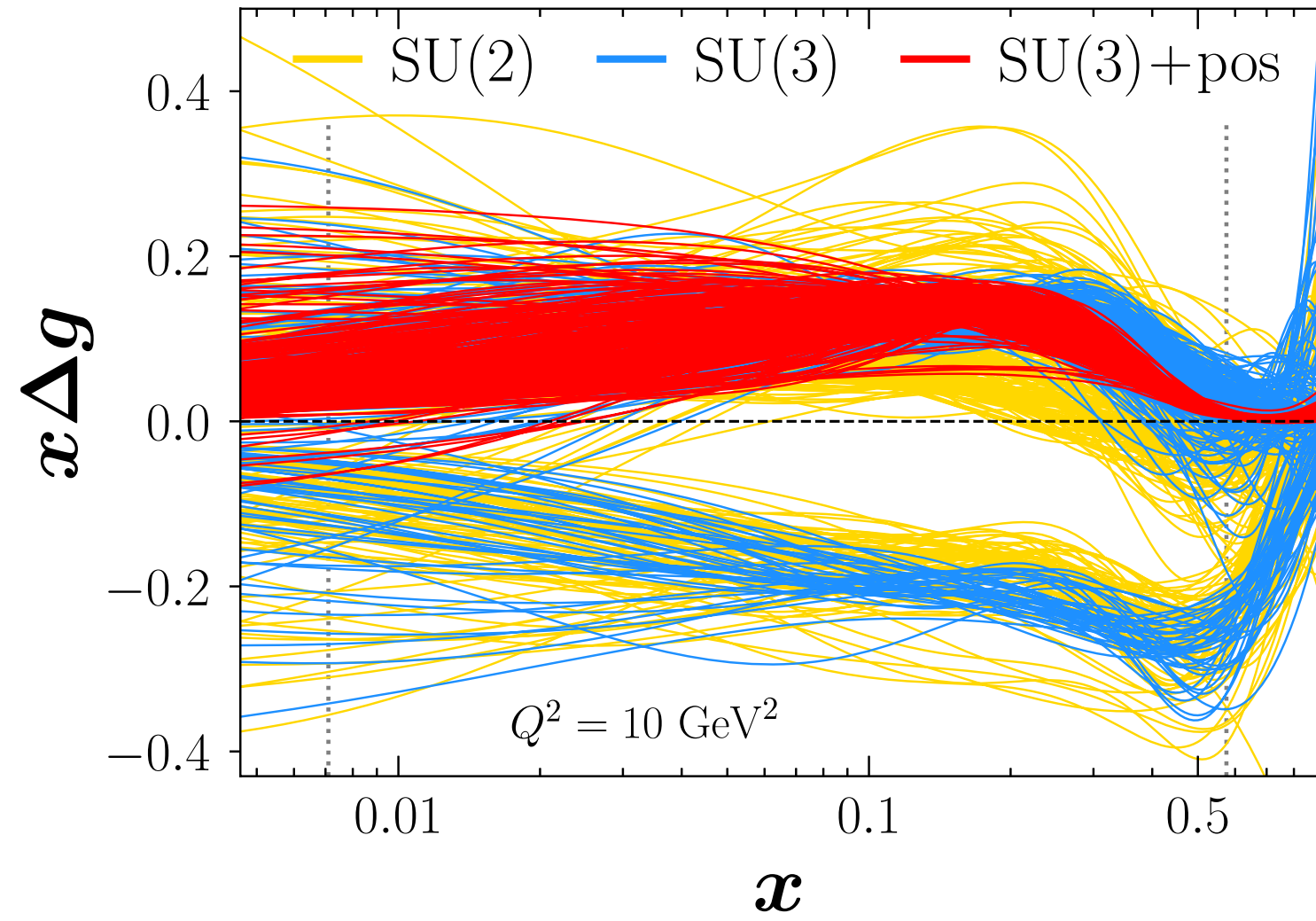
- Motivation
- Introduction to RHIC-STAR
- Longitudinal double spin asymmetry  $A_{LL}$  for  $\pi^\pm$ -tagged jets
- Longitudinal double spin asymmetry  $A_{LL}$  for  $\Lambda$ ,  $\bar{\Lambda}$  and  $K_S^0$
- Longitudinal spin transfer  $D_{LL}$  of  $\Lambda$  and  $\bar{\Lambda}$
- Summary



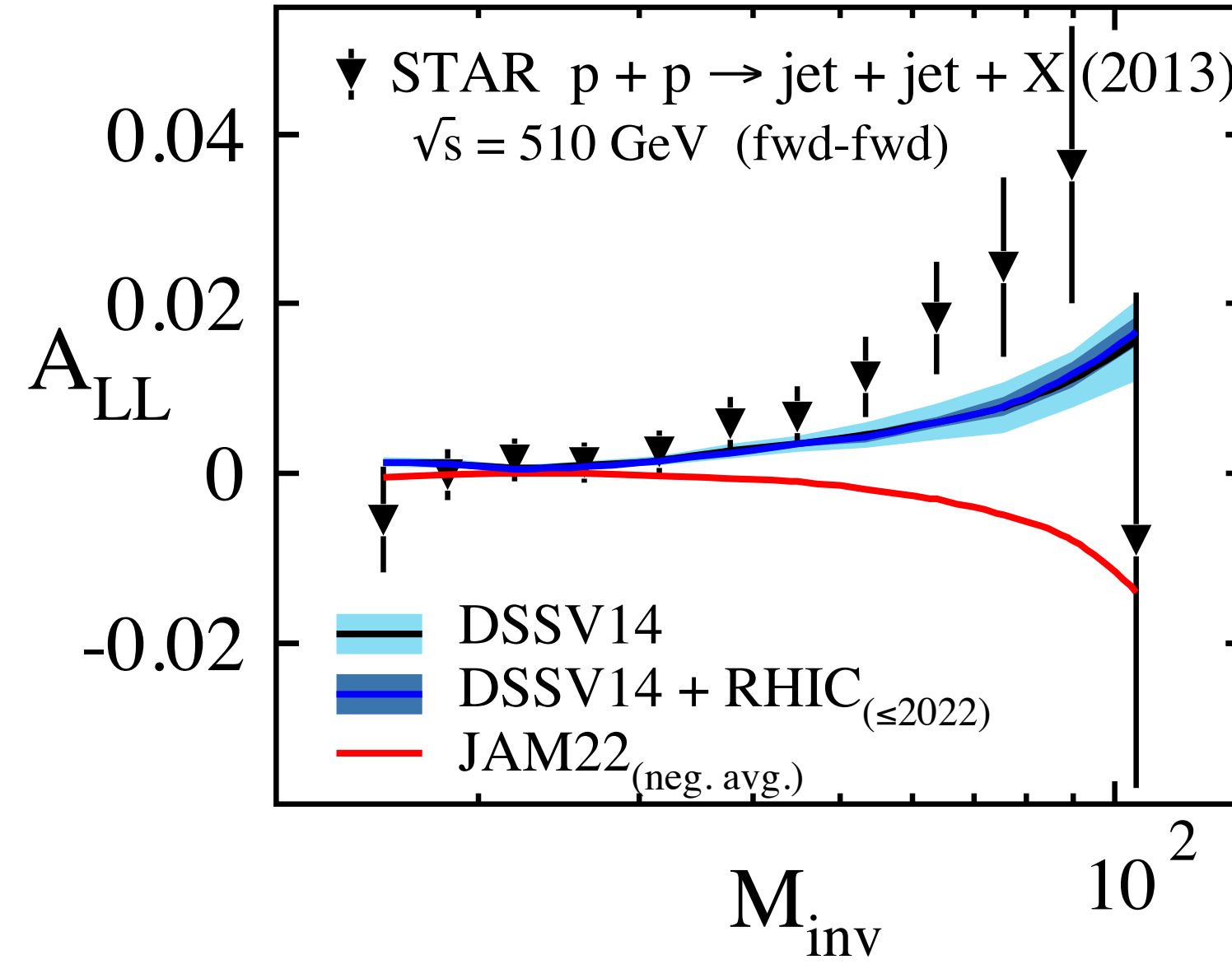
# Gluon helicity: positive or negative?



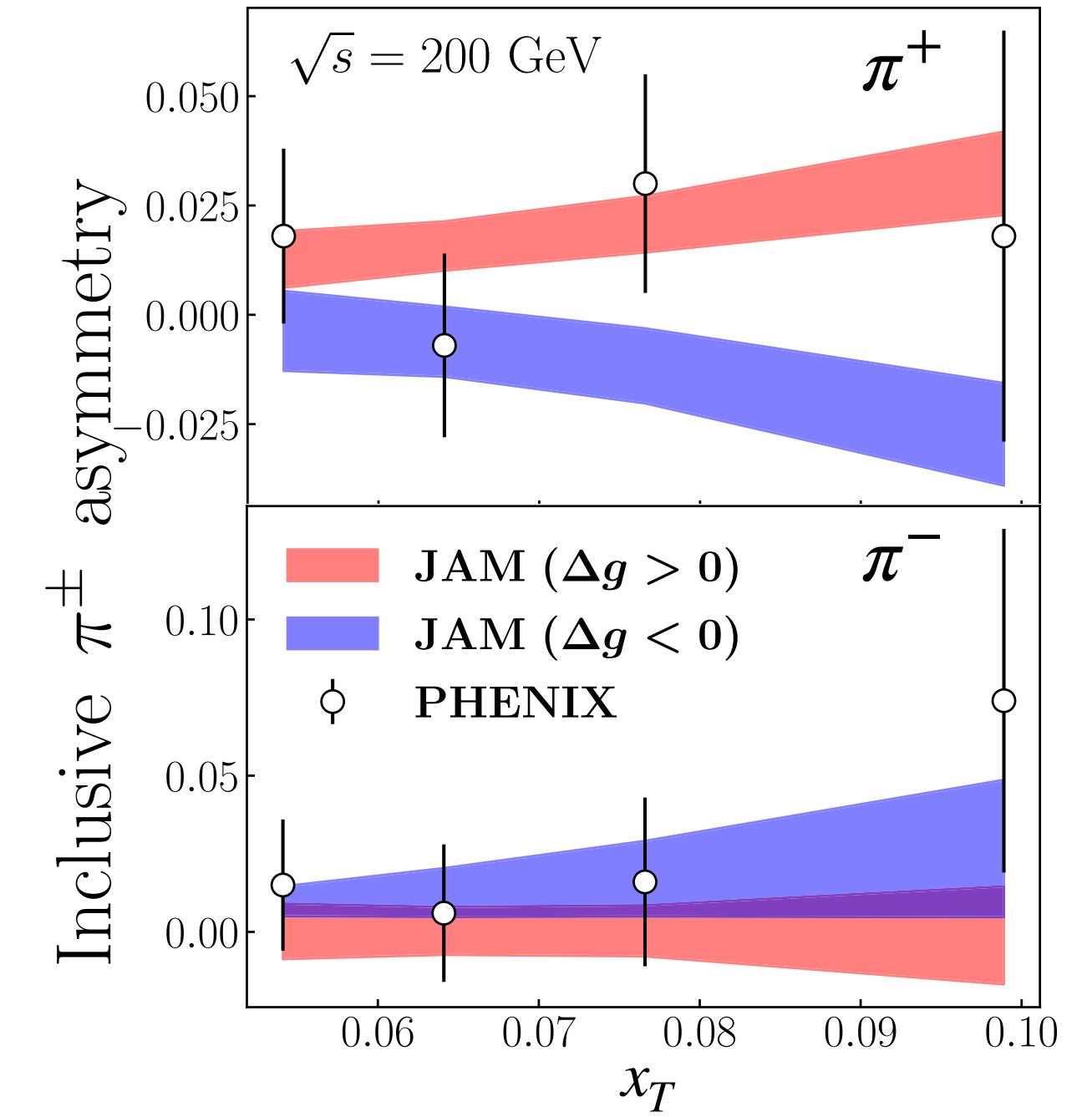
[JAM], Phys. Rev. D **105**, 074022 (2022).



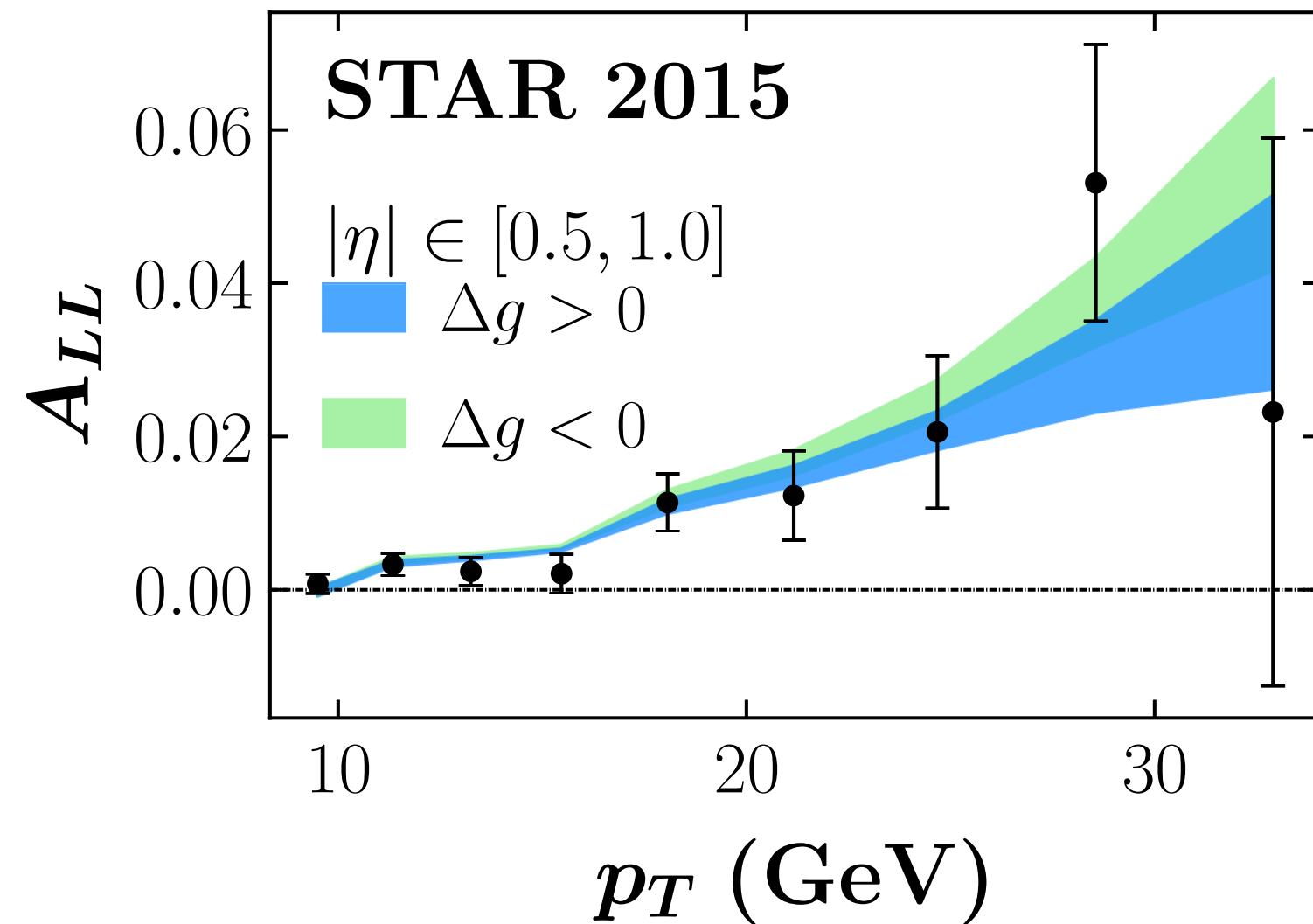
E.-C. Aschenauer et al., arXiv:2302.00605.



[JAM], Phys. Rev. D **107**, 034033 (2023).



[JAM], Phys. Rev. D **105**, 074022 (2022).

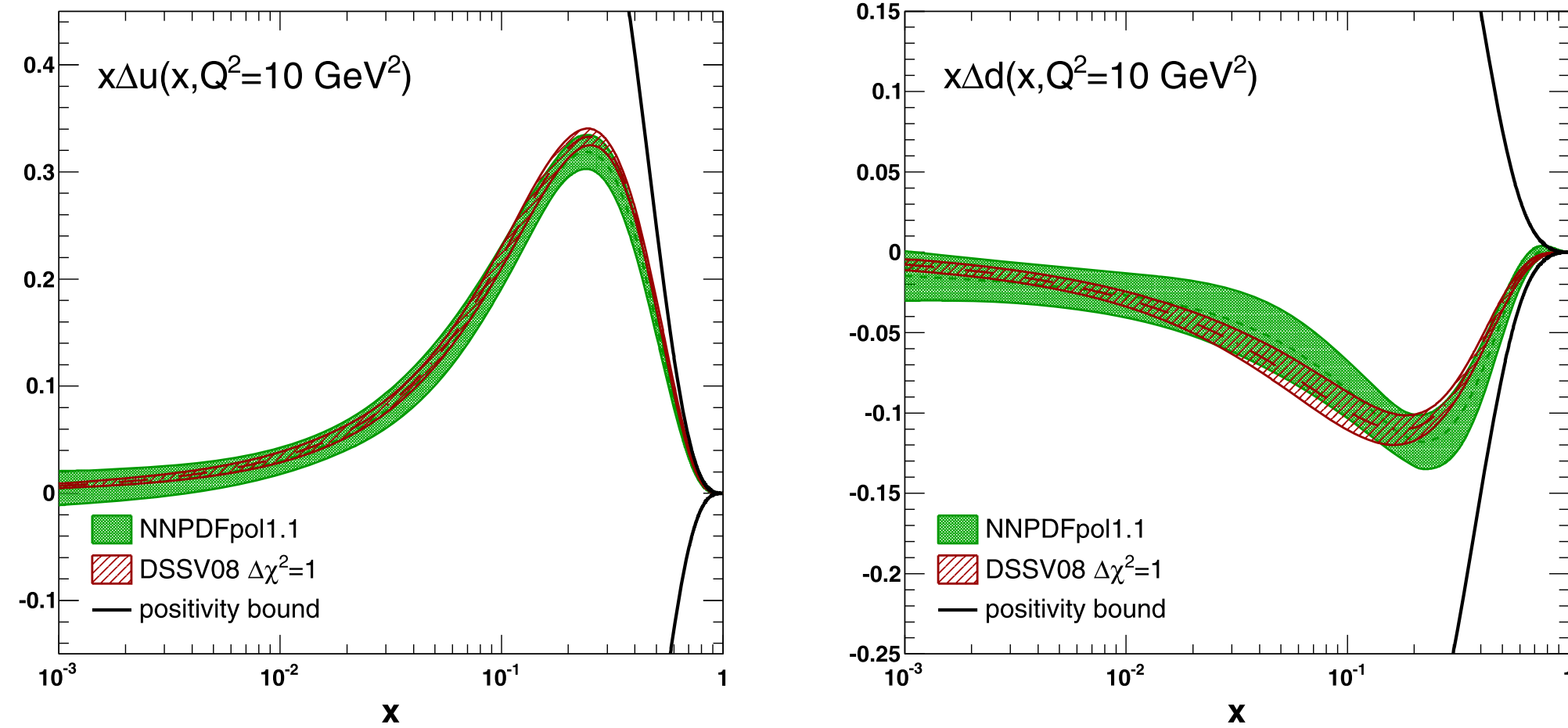


- Inclusive jet  $A_{LL}$  is insensitive to the sign of the  $\Delta g$
- Di-jet  $A_{LL}$  disfavors negative  $\Delta g$  solution
- $\pi^\pm$ -tagged jet  $A_{LL}$  provides an additional way to constrain the sign of  $\Delta g$

# Constraining gluon polarization with $\pi^\pm$ -tagged jet $A_{LL}$

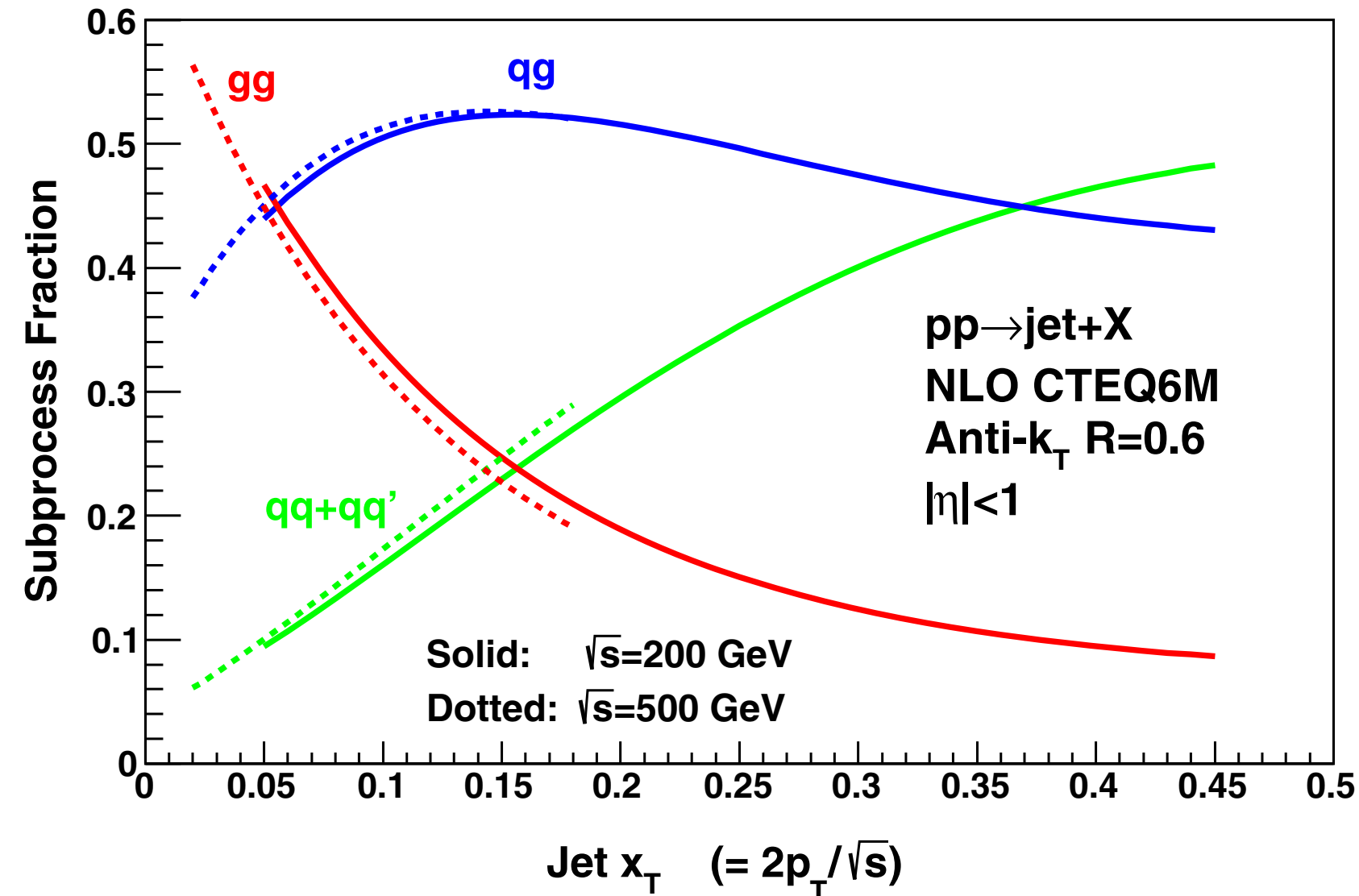


[NNPDF], Nucl. Phys. B **887**, 276 (2014).



$$A_{LL}^{\pi^\pm} = \frac{\sigma_{++} - \sigma_{+-}}{\sigma_{++} + \sigma_{+-}} = \frac{\Sigma \Delta f_i \otimes \Delta f_j \otimes \Delta \hat{\sigma} \otimes D_k^{\pi^\pm}}{\Sigma f_i \otimes f_j \otimes \hat{\sigma} \otimes D_k^{\pi^\pm}}$$

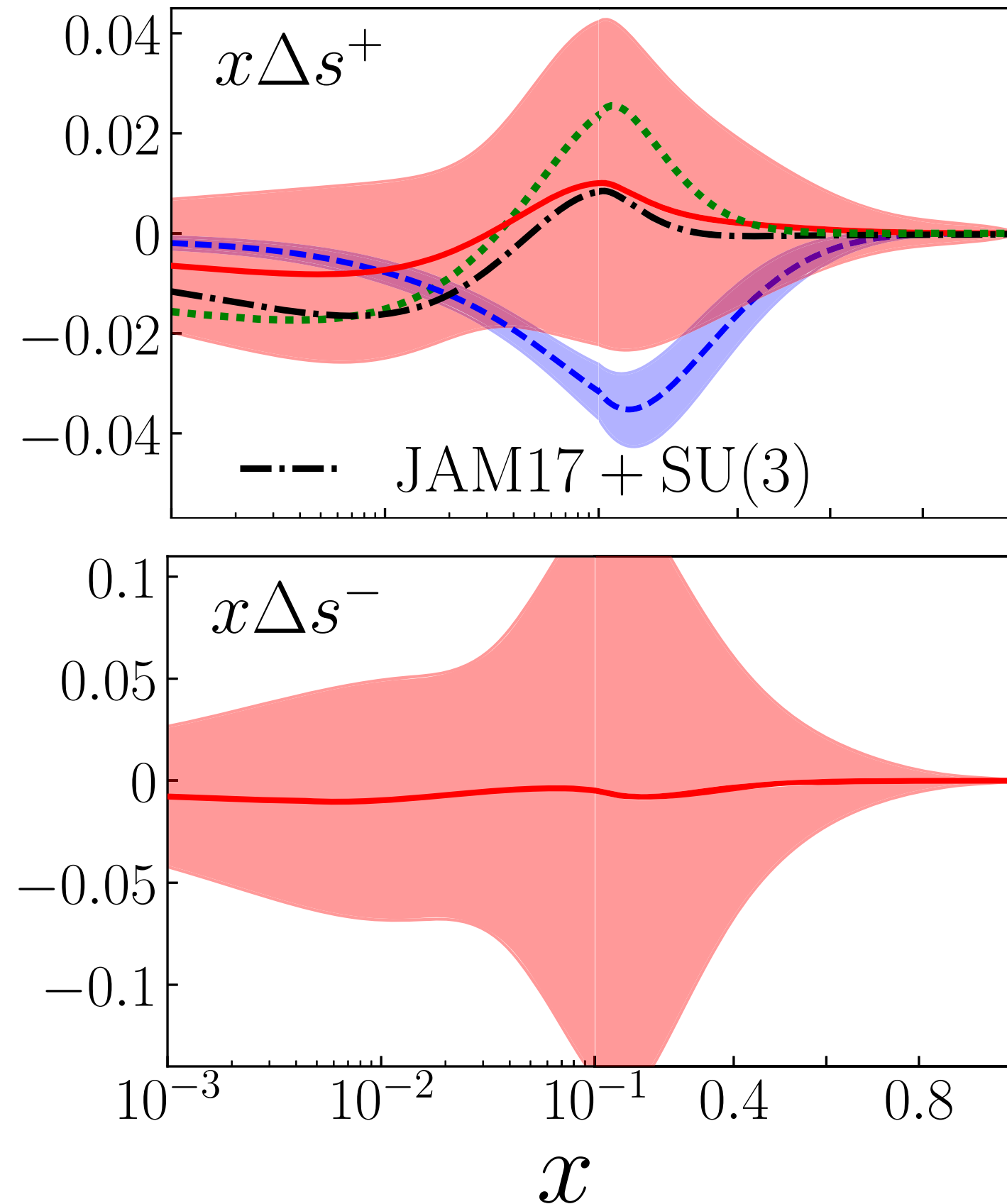
[STAR], Phys. Rev. D **100**, 052005 (2019).



- $\Delta u > 0$  and  $\Delta d < 0$
- $u$ - $g$  and  $d$ - $g$  scatterings are sensitive to the sign of  $\Delta g$
- $u$  quark favors  $\pi^+$ ,  $d$  quark favors  $\pi^-$
- $q$ - $g$  scattering is the dominated process at RHIC energy
- $\Delta g > 0 \rightarrow A_{LL}^{\pi^+} > A_{LL}^{\pi^-}$
- $\Delta g < 0 \rightarrow A_{LL}^{\pi^+} < A_{LL}^{\pi^-}$



# Probing strange quark helicity distribution



[JAM], Phys. Rev. Lett. **119**, 132001 (2017).

- Poor constraints on the (anti-)strange quark helicity distributions  $(\Delta\bar{s}) \Delta s$

- $A_{LL}$  of  $\Lambda, \bar{\Lambda}$  and  $K_S^0$

$$A_{LL}^{\Lambda} = \frac{\sigma_{++} - \sigma_{+-}}{\sigma_{++} + \sigma_{+-}} = \frac{\Sigma\Delta f_i \otimes \Delta f_j \otimes \Delta\hat{\sigma} \otimes D_k^{\Lambda}}{\Sigma f_i \otimes f_j \otimes \hat{\sigma} \otimes D_k^{\Lambda}}$$

- ▶ Valence  $s$  or  $\bar{s}$  inside  $\Lambda, \bar{\Lambda}$  and  $K_S^0$
- ▶  $s$  or  $\bar{s}$  prefers  $\Lambda, \bar{\Lambda}$  and  $K_S^0$  in the fragmentation process
- ▶  $A_{LL}$  of  $\Lambda, \bar{\Lambda}$  and  $K_S^0$  can shed light on the  $\Delta s$  and  $\Delta\bar{s}$

# Probing strange quark helicity distribution

- Longitudinal spin transfer  $D_{LL}$  of  $\Lambda$  and  $\bar{\Lambda}$  in p+p collisions

$$D_{LL}^{\Lambda} \equiv \frac{d\sigma^{p^+p \rightarrow \Lambda^+X} - d\sigma^{p^+p \rightarrow \Lambda^-X}}{d\sigma^{p^+p \rightarrow \Lambda^+X} + d\sigma^{p^+p \rightarrow \Lambda^-X}} = \frac{d\Delta\sigma}{d\sigma}$$

$$d\Delta\sigma \propto \Delta f_a(x_a) f_b(x_b) \Delta\sigma^{ab \rightarrow cd} \Delta D^{\Lambda}(z)$$

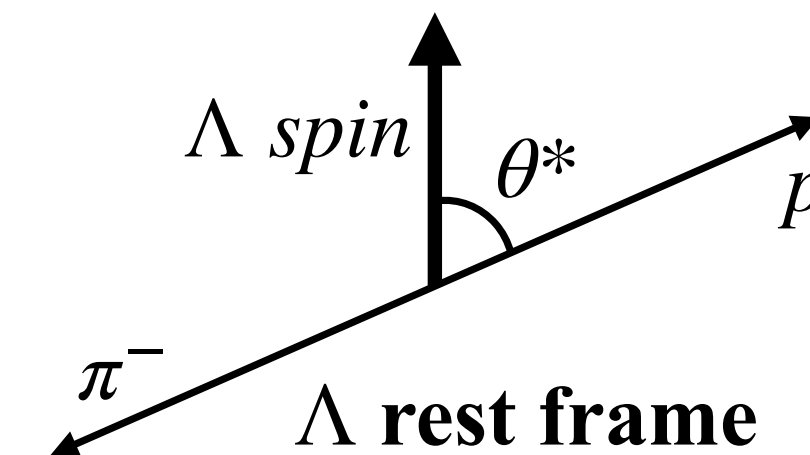
helicity distribution     
 pQCD calculable     
 longitudinally polarized FFs

- Polarization of  $\Lambda$  can be measured via its weak decay

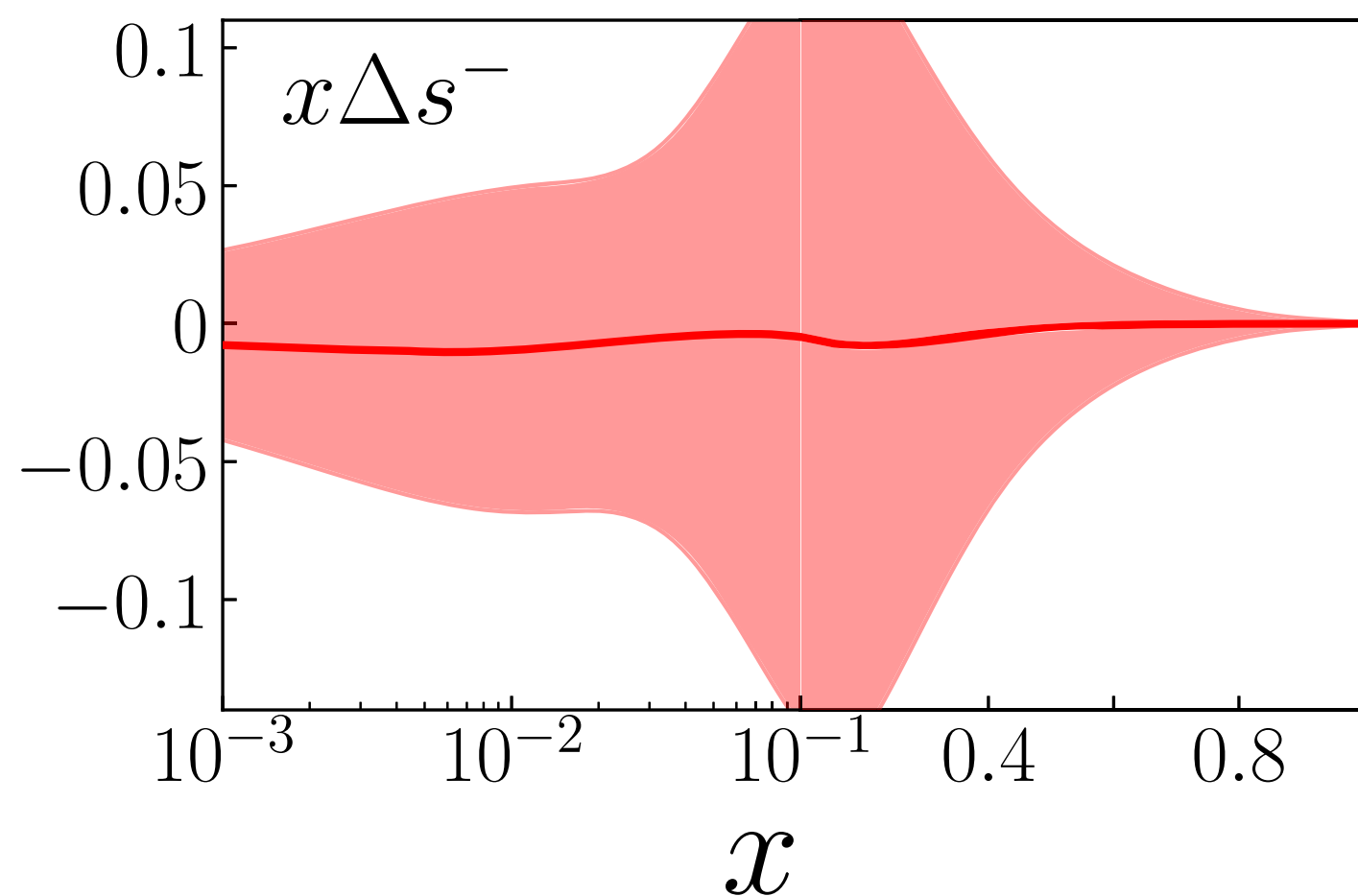
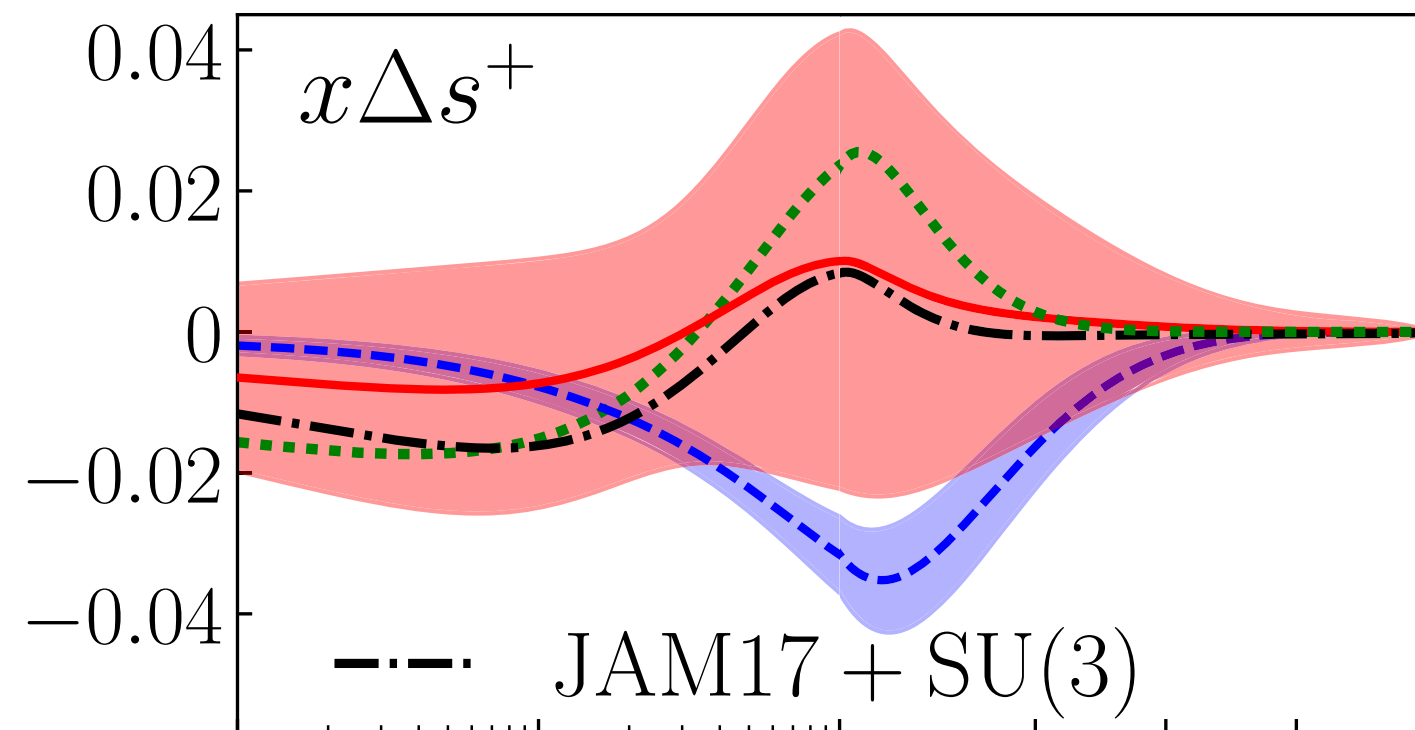
$$dN \sim (1 + \alpha P_{\Lambda} \cos \theta^*) d \cos \theta^*$$

$\alpha$ : weak decay constant

$P_{\Lambda}$ :  $\Lambda$  polarization



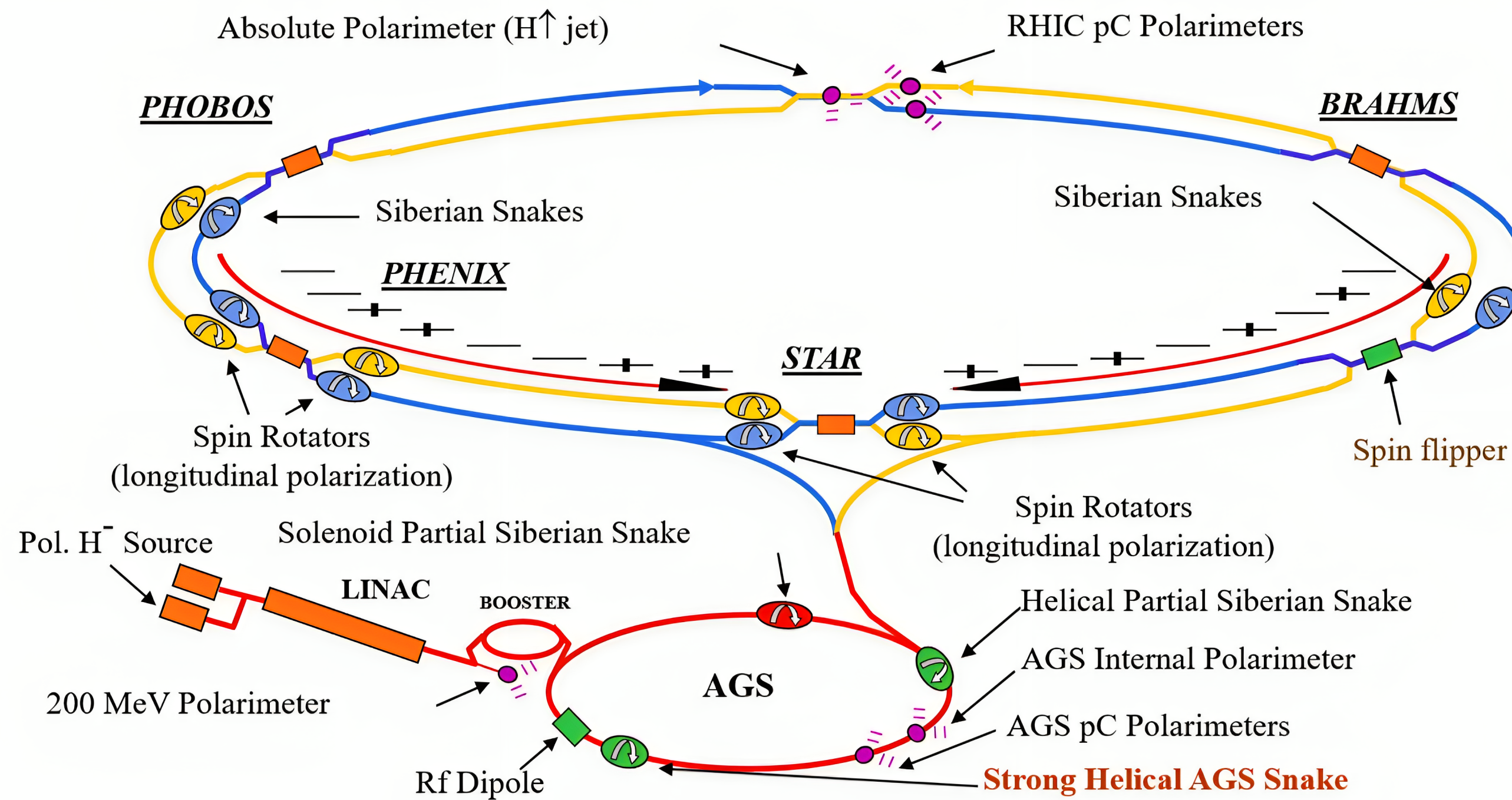
- Large fraction of  $\Lambda$  spin is carried by its  $s$  quark



[JAM], Phys. Rev. Lett. **119**, 132001 (2017).



# The Relativistic Heavy Ion Collider



- The first and only polarized p+p collider in the world
- Collides both longitudinally and transversely polarized proton beams at  $\sqrt{s} = 200$  GeV and 500/510 GeV

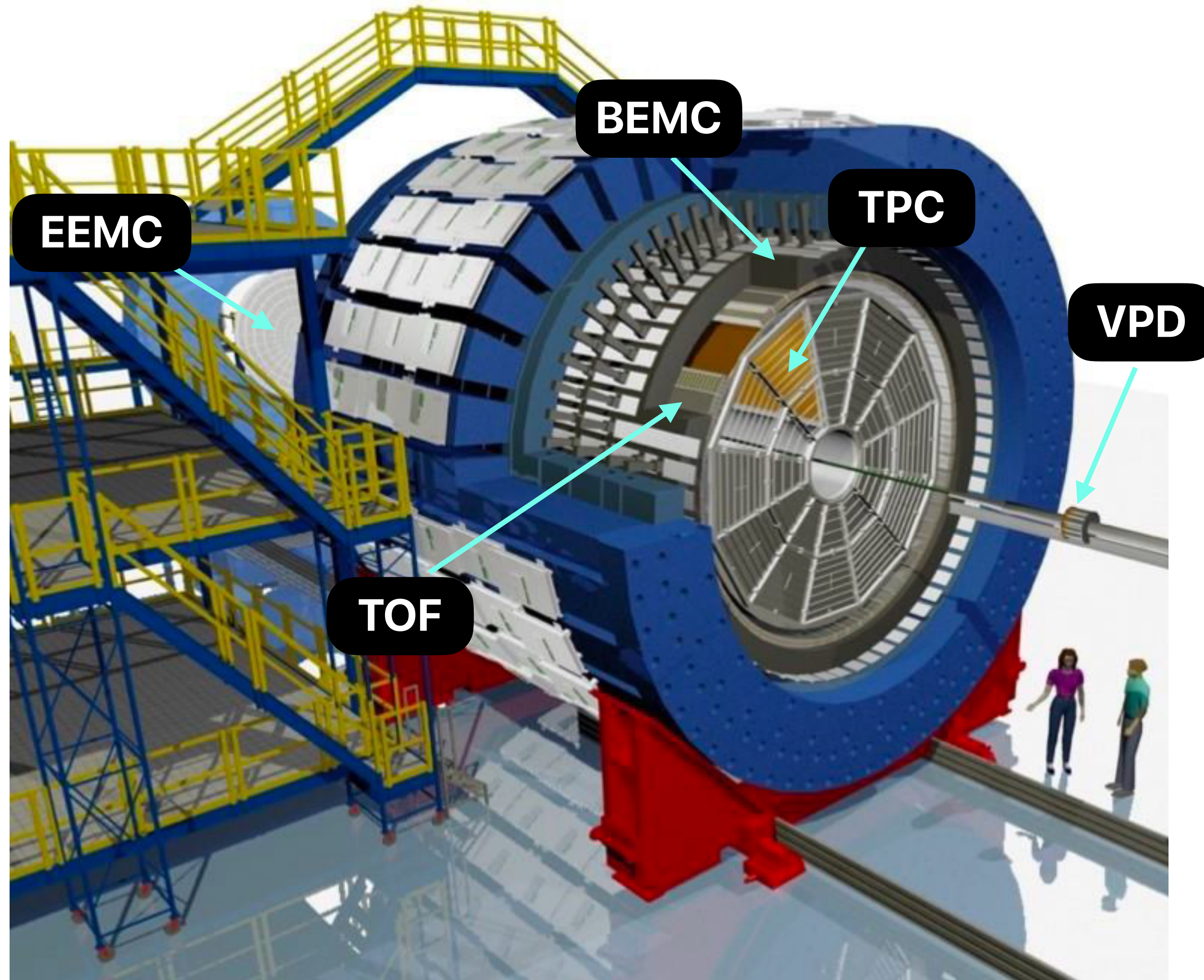
Longitudinally polarized  $p + p$  collision samples taken at STAR

Year	$\sqrt{s}$ (GeV)	$\int L$ (pb <sup>-1</sup> )	$P_{beam}$
2009	200	19	57% / 57%
2015	200	52	52% / 56%
2012	510	82	50% / 53%
2013	510	300	51% / 52%

used in this analysis →



# The Solenoidal Tracker at RHIC



- Time Projection Chamber (TPC)
  - ▶  $|\eta| < 1.3$  and  $0 \leq \phi \leq 2\pi$
  - ▶ Tracking and particle identification (PID)
- Time of Flight detector (TOF)
  - ▶  $|\eta| < 0.9$  and  $0 \leq \phi \leq 2\pi$
  - ▶ Particle identification
- Electromagnetic Calorimeter (EMC)
  - ▶ Barrel EMC (BEMC):  $|\eta| < 1.0$  and  $0 \leq \phi \leq 2\pi$
  - ▶ Endcap EMC (EEMC):  $1.086 < \eta < 2.0$  and  $0 \leq \phi \leq 2\pi$
  - ▶ Reconstruction of photon,  $\pi^0$ , jet ..., and serves as trigger detectors
- Vertex Position Detector (VPD)
  - ▶  $4.24 < |\eta| < 5.1$
  - ▶ Monitor the relative luminosities and determine the primary vertex



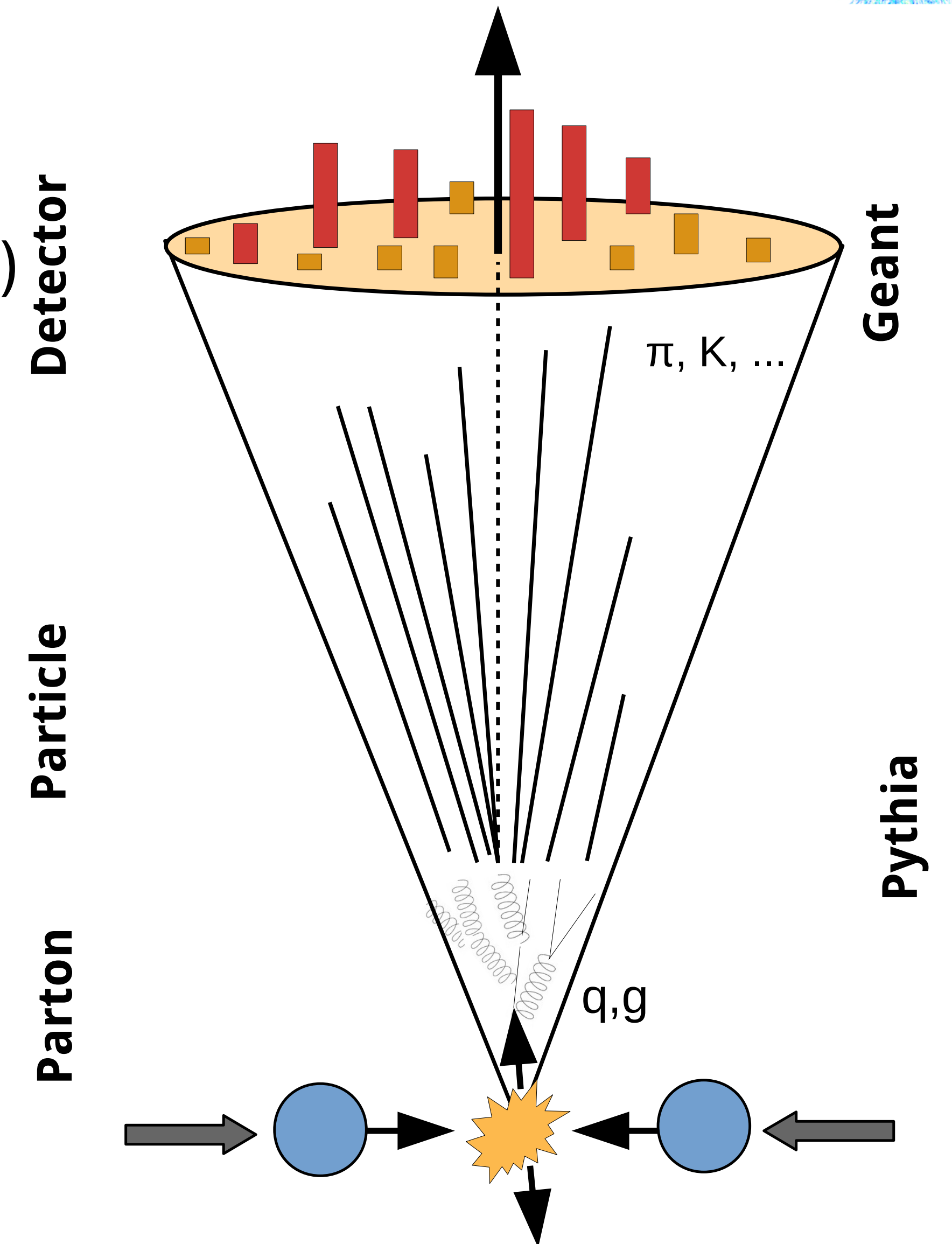
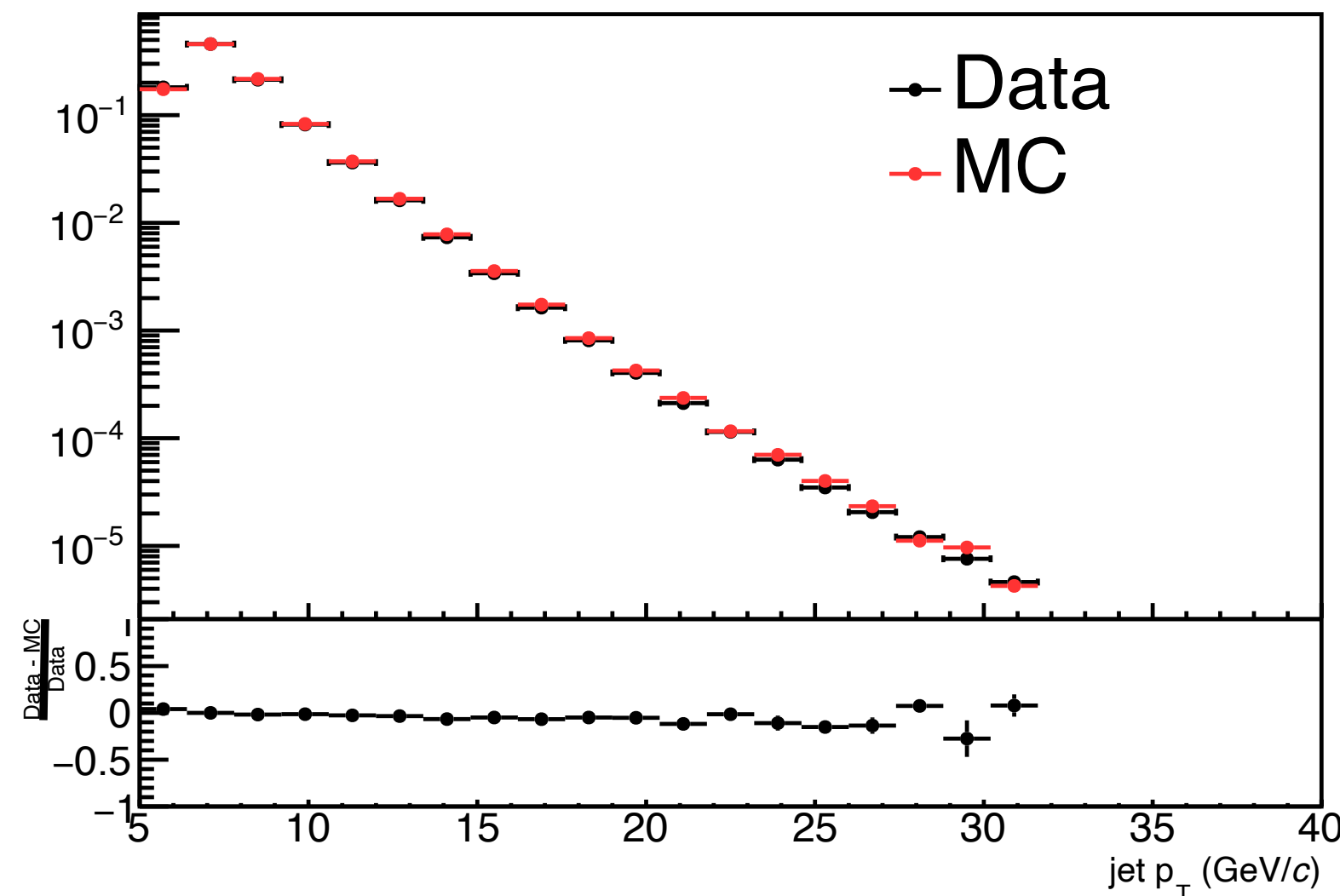


# Part I: Longitudinal double spin asymmetry $A_{LL}$ of $\pi^{\pm}$ -tagged jets

# Jet Reconstruction

- High- $p_T$  jet triggers (JP1 and JP2) are used based on energy deposits in EMC
- Jet reconstruction (TPC tracks + energy deposits in EMC)
  - ▶ Anti- $k_T$  algorithm, with  $R = 0.6$
  - ▶ Simulation: PYTHIA6 + GEANT3 + Zero-bias events
    - Jet was reconstructed at parton, particle and detector level
  - ▶ Jet  $p_T$  was corrected back to particle level
- Jets tagged with  $\pi^\pm$  with  $z > 0.2$  or  $z > 0.3$

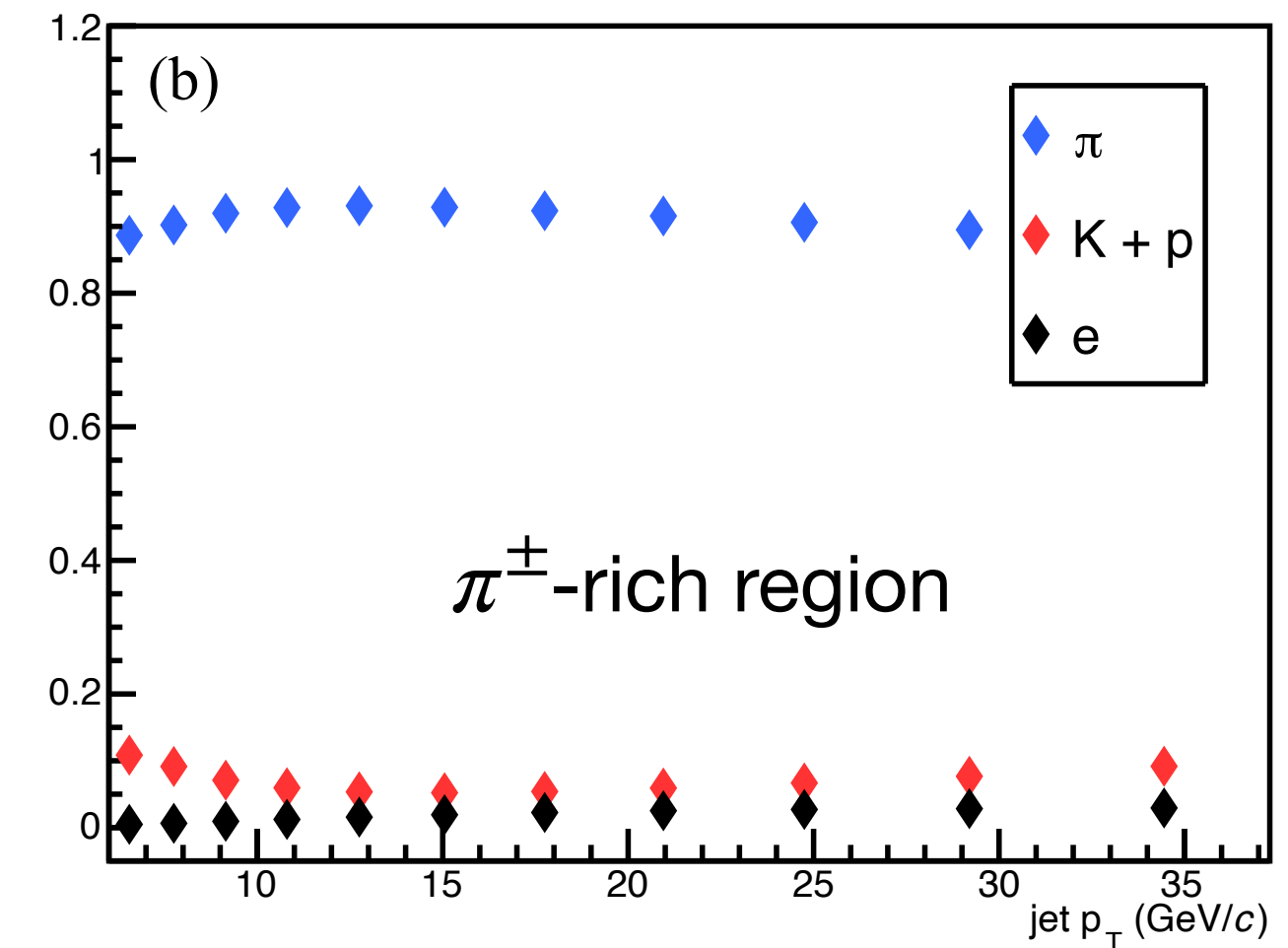
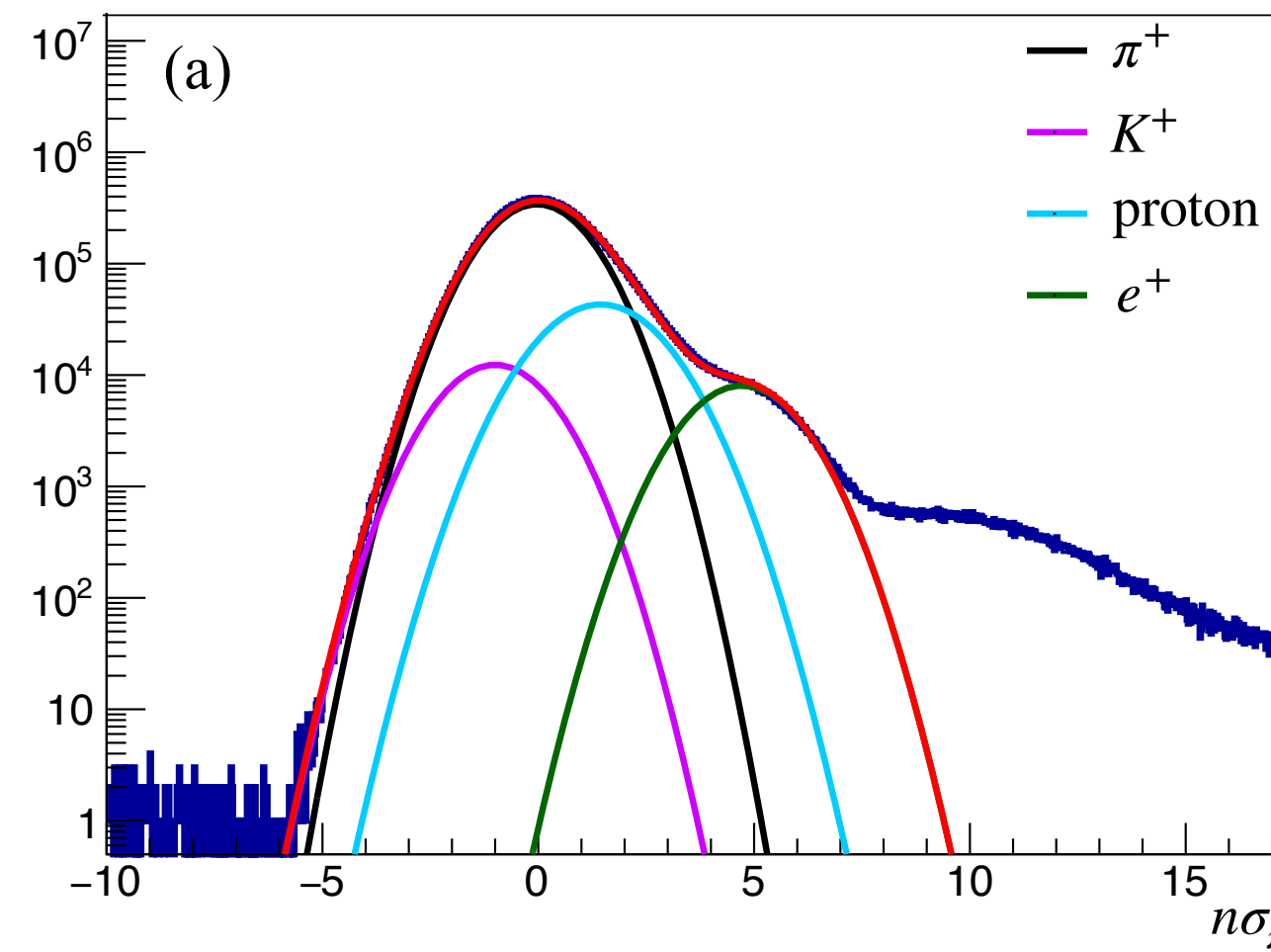
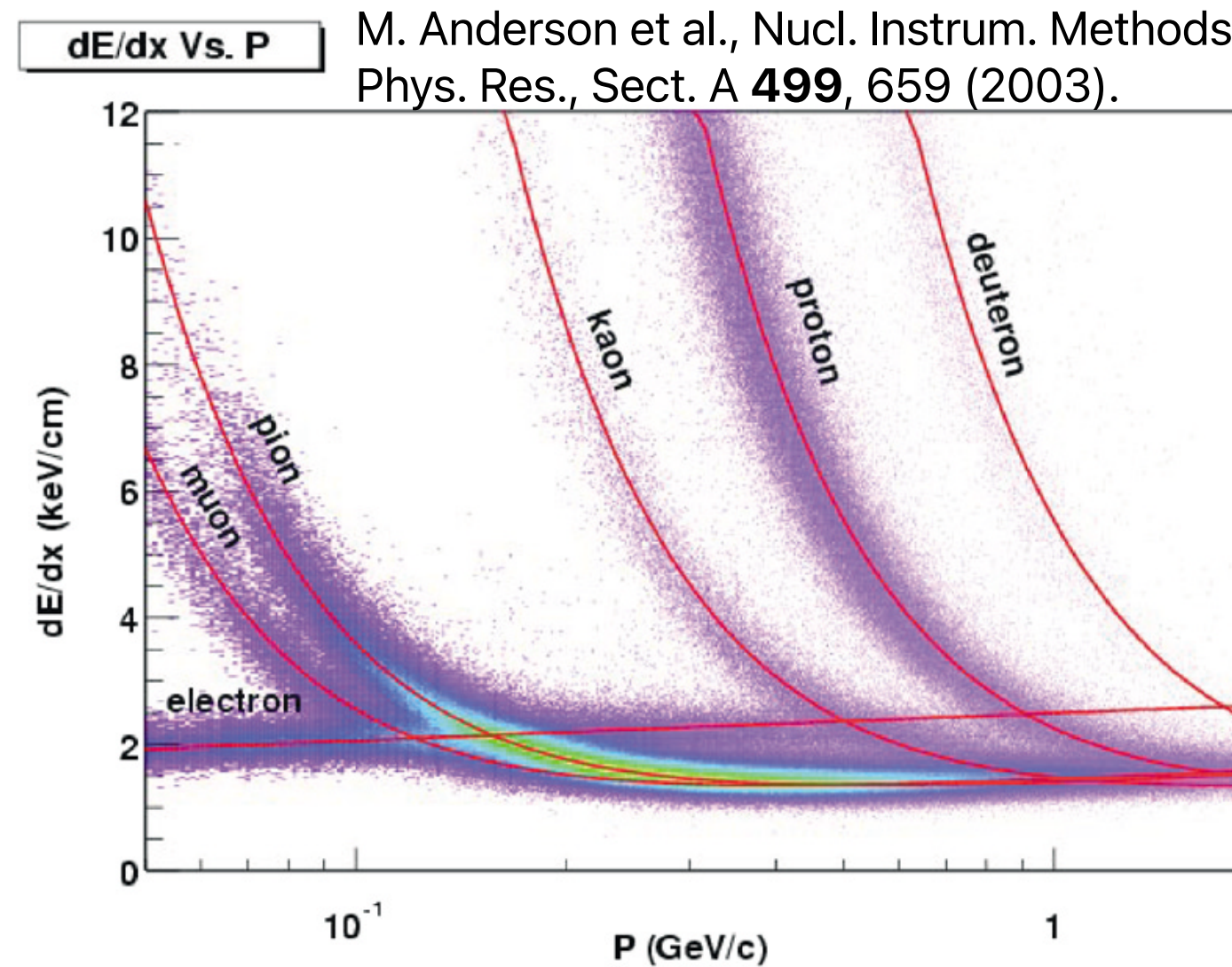
$$z \equiv \frac{\vec{p}_\pi \cdot \vec{p}_{jet}}{|\vec{p}_{jet}|^2}$$





# $\pi^\pm$ PID

- $\pi^\pm$  are identified based on their energy loss inside the TPC  $n\sigma(\pi) = \frac{1}{\sigma_{\text{exp}}} \ln \left( \frac{dE/dx_{\text{obs}}}{dE/dx_{\pi, \text{cal}}} \right)$
- Particle purity is estimated with multi-Gaussian fitting of the  $n\sigma_\pi$  distribution
- 3 particle rich regions ( $\pi^\pm, K^\pm + p(\bar{p}), e^\pm$ )



$$A_{LL} = \frac{1}{P_B P_Y} \frac{(N_{++} + N_{--}) - R_3(N_{+-} + N_{-+})}{(N_{++} + N_{--}) + R_3(N_{+-} + N_{-+})}$$

+ ( - ) denotes the beam helicity

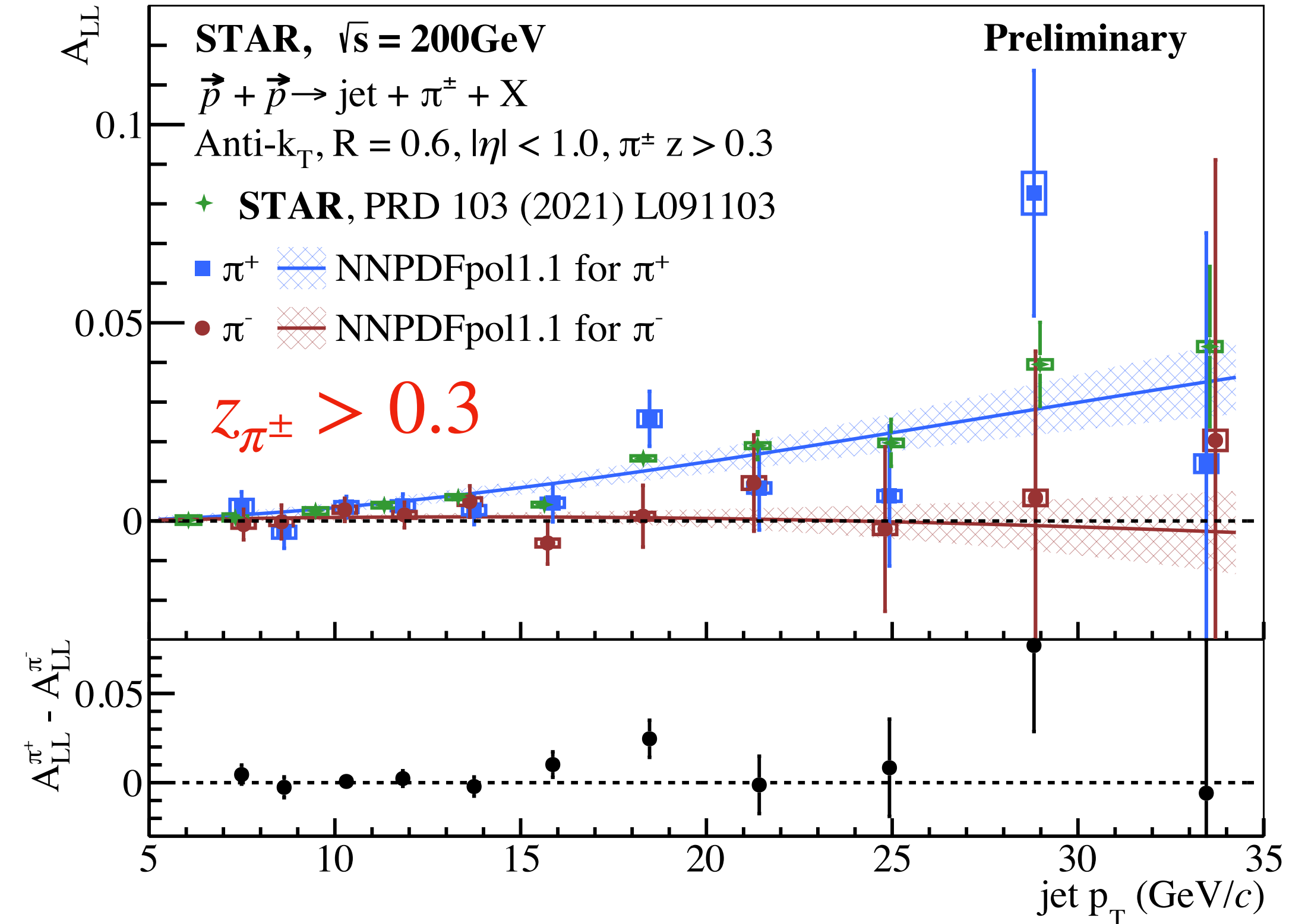
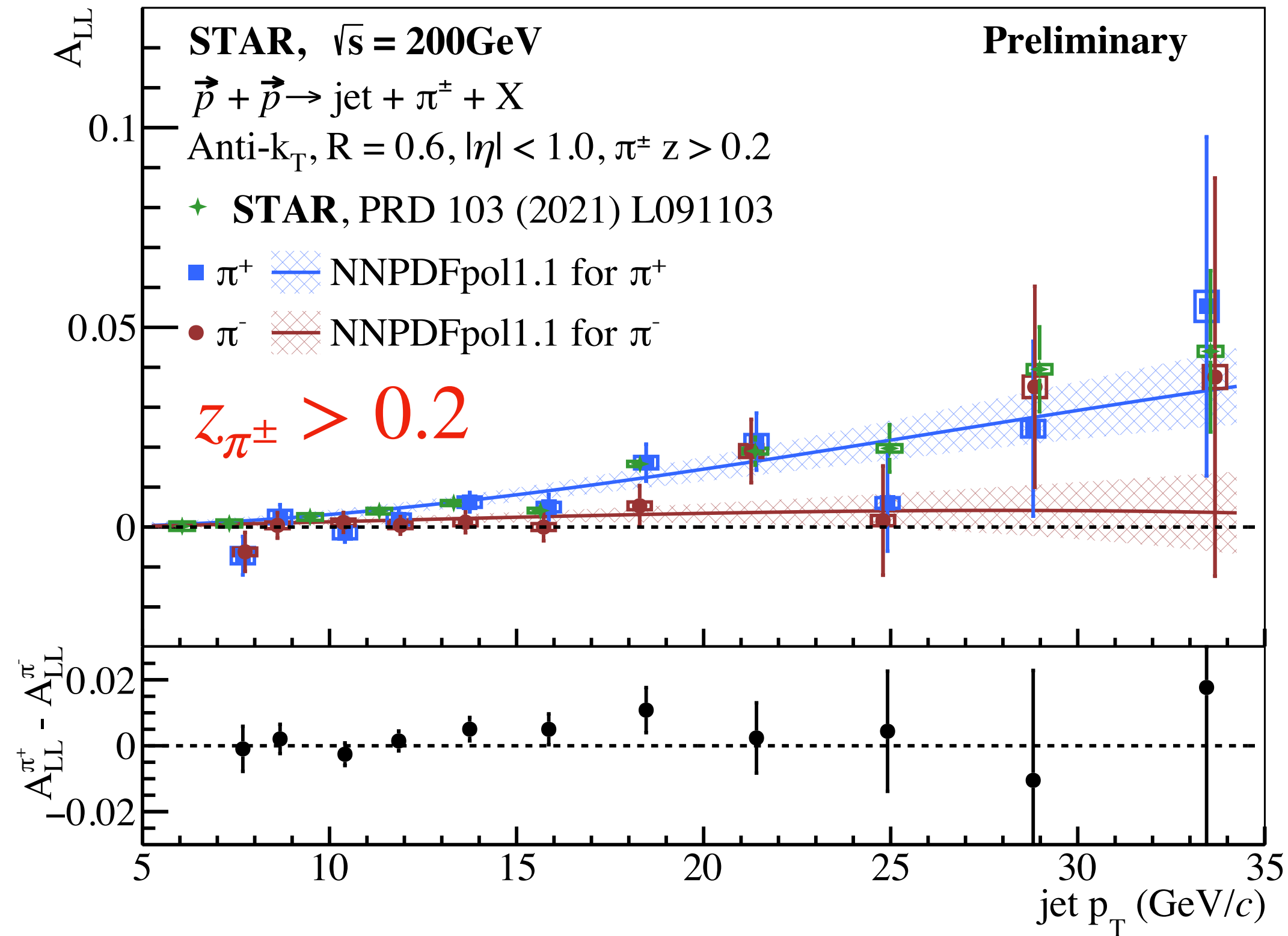
$N_{++}$  etc are the jet yields for different beam helicity configurations

$P_B$  and  $P_Y$  are beam polarizations

$R_3$  is the relative luminosity calculated with the VPD

$$A_{LL}^{raw} = \sum_{\pi, K+p, e} f_i A_{LL}^i \begin{bmatrix} f_{\pi_{rich}}^\pi & f_{\pi_{rich}}^{K+p} & f_{\pi_{rich}}^e \\ f_{K+p_{rich}}^\pi & f_{K+p_{rich}}^{K+p} & f_{K+p_{rich}}^e \\ f_{e_{rich}}^\pi & f_{e_{rich}}^{K+p} & f_{e_{rich}}^e \end{bmatrix} \begin{bmatrix} A_\pi \\ A_{K+p} \\ A_e \end{bmatrix} = \begin{bmatrix} A_{\pi_{rich}}^{raw} \\ A_{K+p_{rich}}^{raw} \\ A_{e_{rich}}^{raw} \end{bmatrix}$$

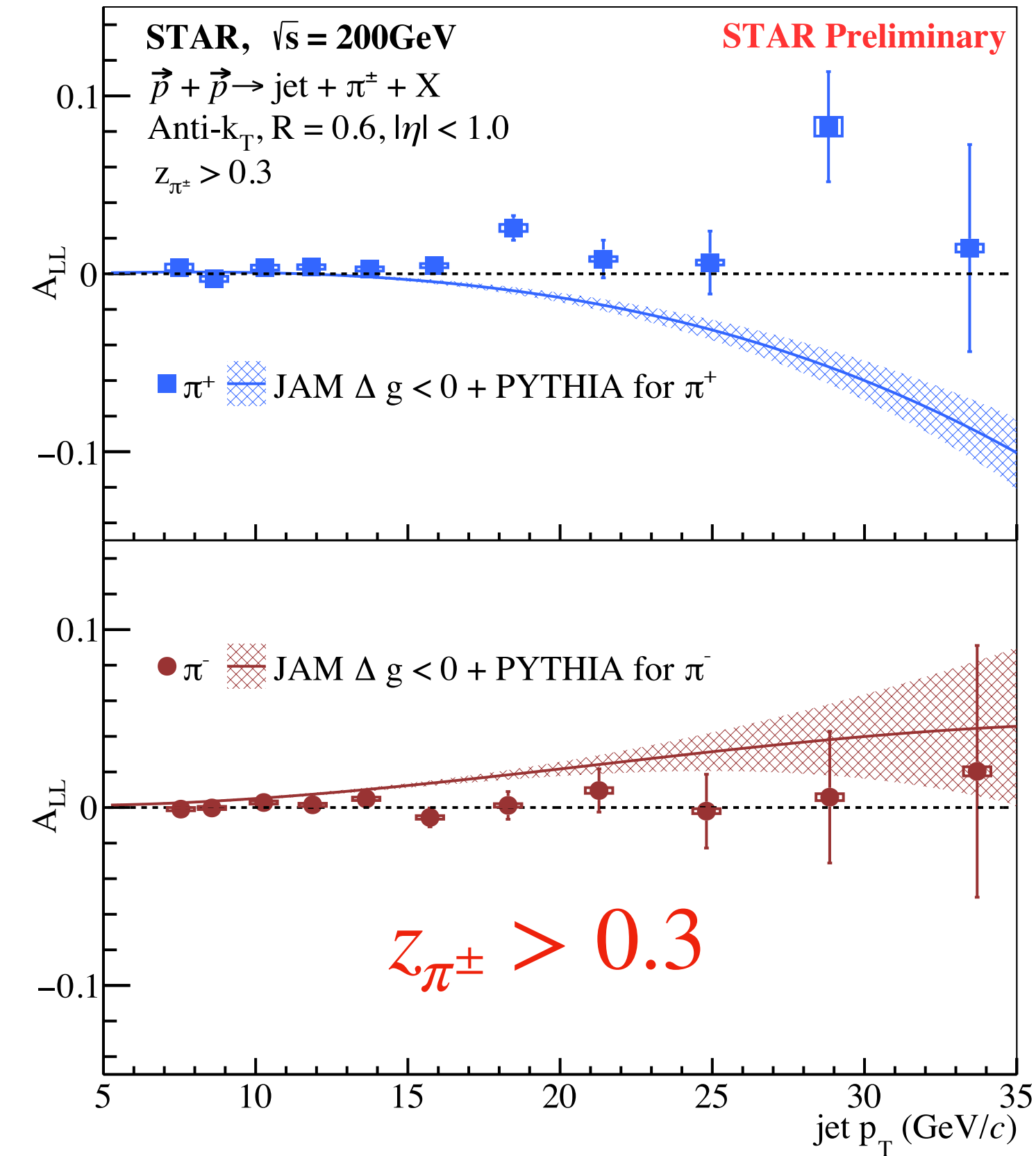
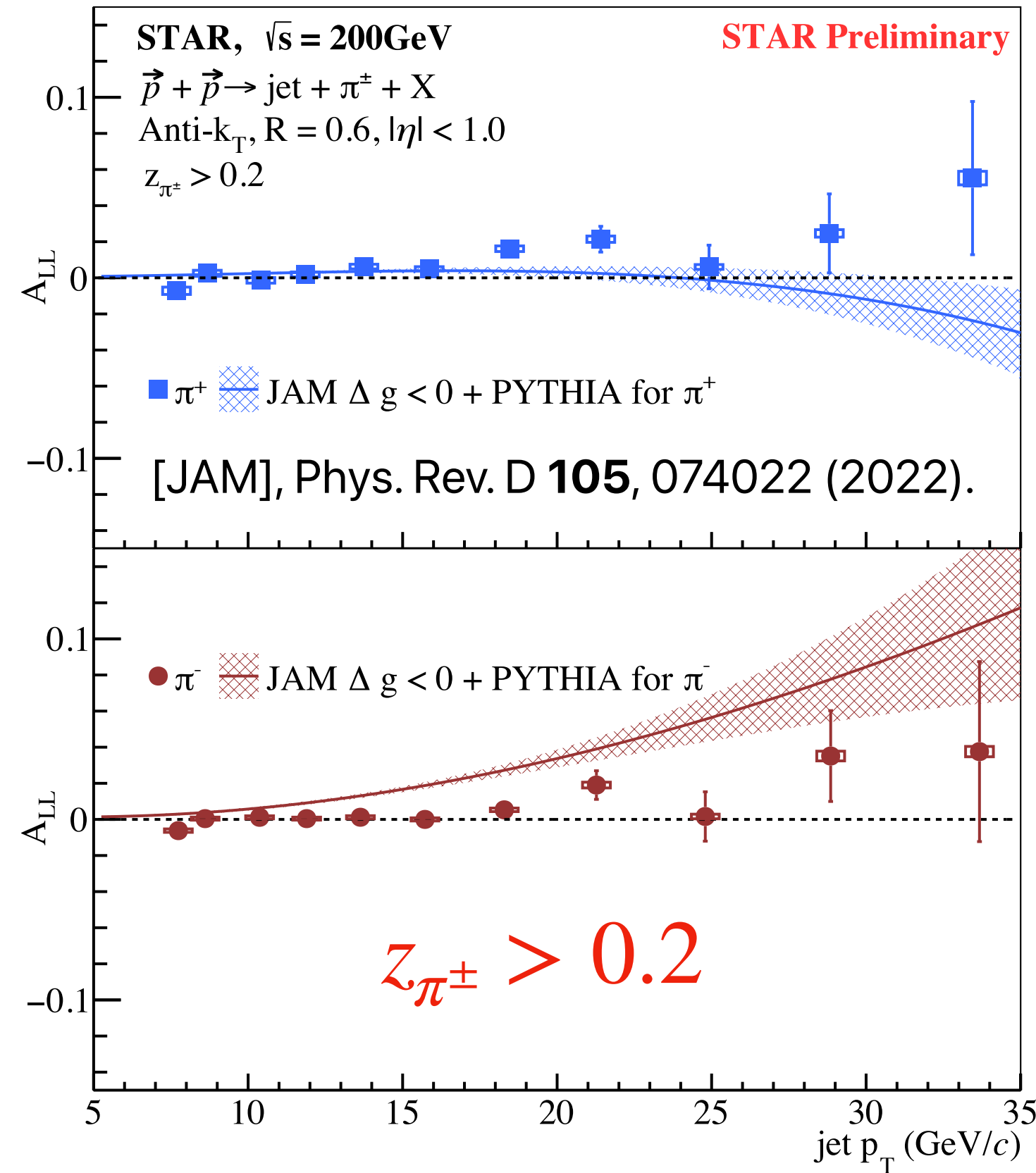
# Results of $A_{LL}$ vs jet $p_T$



- Indication of  $A_{LL}^{\pi^+} > A_{LL}^{\pi^-}$
- NNPDFpol1.1 predicts  $A_{LL}^{\pi^+} > A_{LL}^{\pi^-}$  with positive gluon helicity  $\Delta g$
- The results are close to the predictions



# Results of $A_{LL}$ vs jet $p_T$



- JAM+PYTHIA predicts different trends of  $A_{LL}^\pm$
- The measurements are not consistent with the JAM+PYTHIA prediction with  $\Delta g < 0$

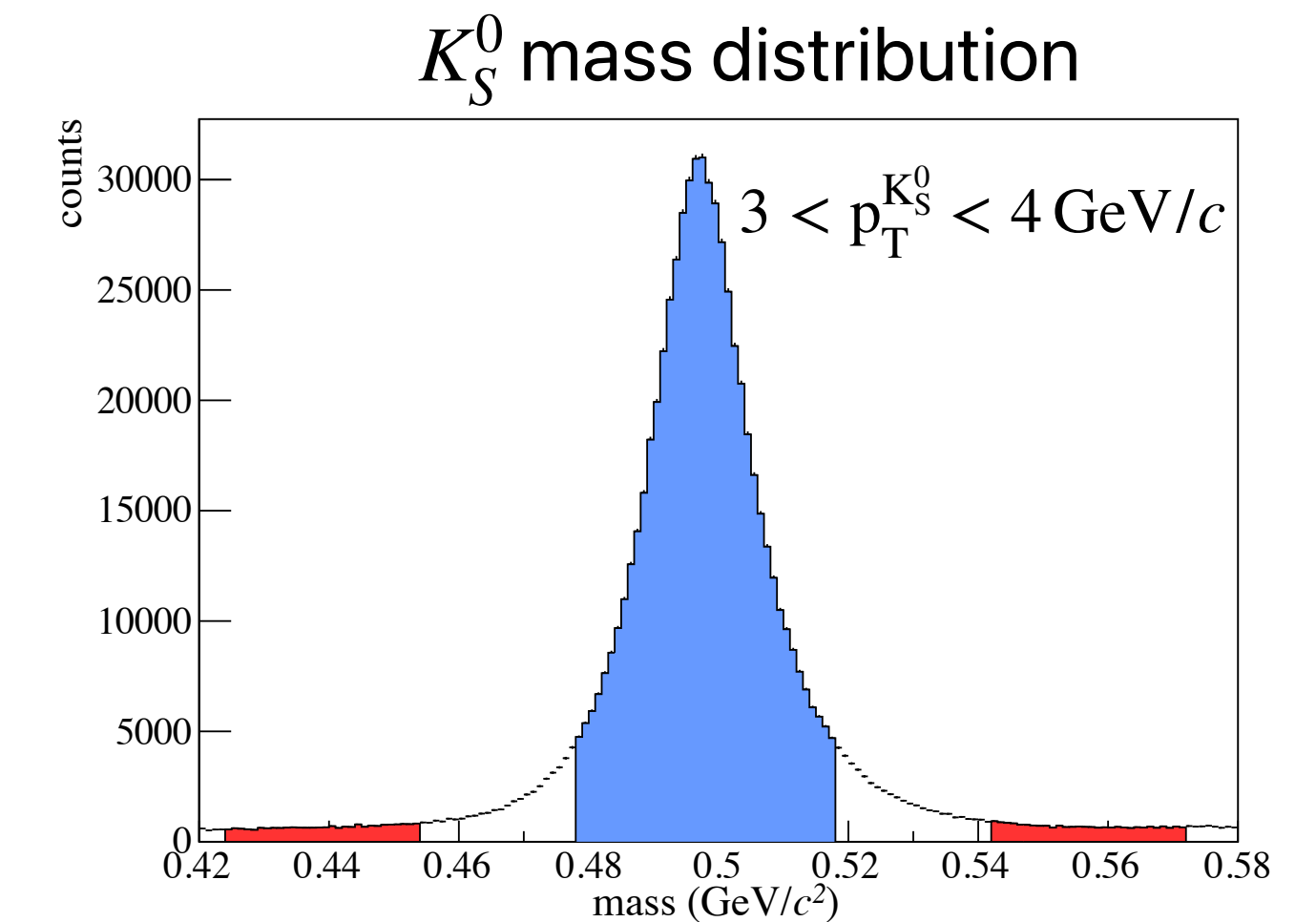
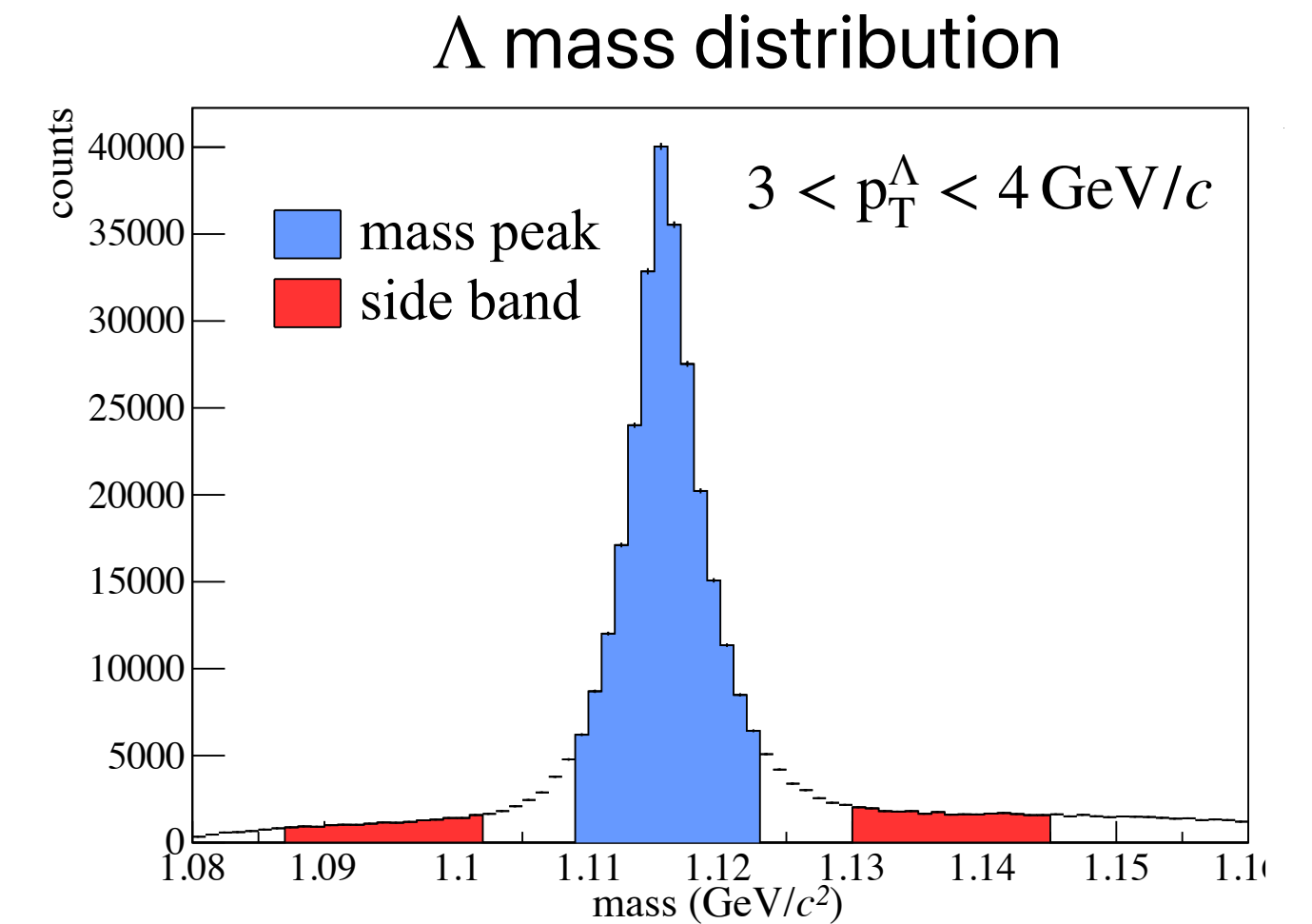


## Part II: Longitudinal double spin asymmetry $A_{LL}$ of $\Lambda$ , $\bar{\Lambda}$ and $K_S^0$

# $\Lambda$ and $K_S^0$ Selection & Jet Reconstruction



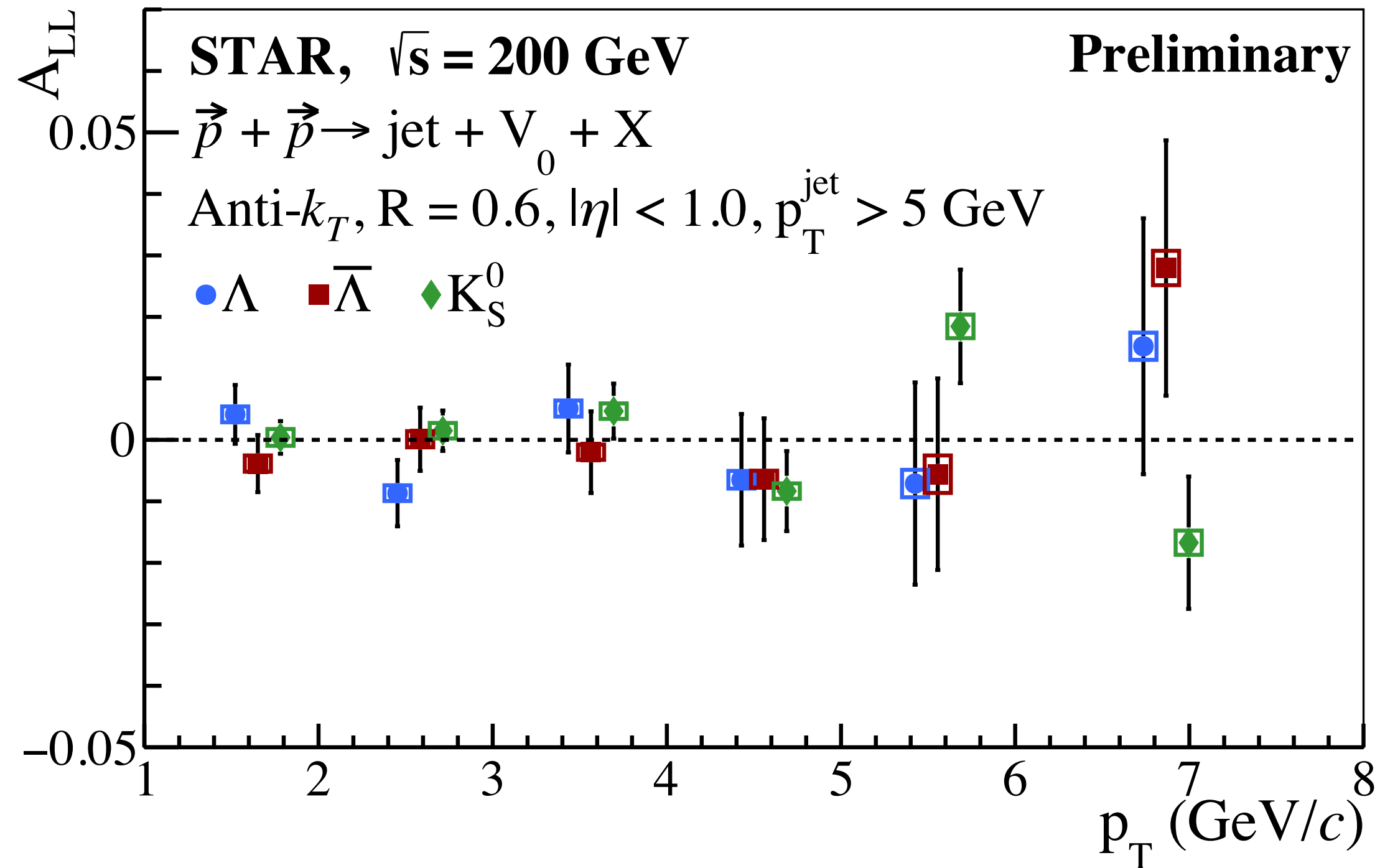
- $\Lambda$  and  $K_S^0$  selection
  - ▶  $\Lambda(\bar{\Lambda}) \rightarrow p(\bar{p}) + \pi^-(\pi^+)$ ,  $K_S^0 \rightarrow \pi^+ + \pi^-$
  - ▶  $p(\bar{p})$  and  $\pi^\pm$  tracks were measured with the TPC
  - ▶ Sets of topological cuts were applied to reduce background
  - ▶ Residual background fraction  $r$  was estimated with side-band method
- Jet reconstruction
  - ▶ Jet reconstructed with anti- $k_T$  algorithm with  $R = 0.6$
  - ▶  $\Lambda$  and  $K_S^0$  candidate as input for jet reconstruction
  - ▶ In-jet  $\Lambda$  and  $K_S^0$  are used to make sure they originate from the hard scattering



**Background subtraction**  $\longrightarrow$  
$$A_{LL} = \frac{A_{LL}^{raw} - rA_{LL}^{bkg}}{1 - r}$$

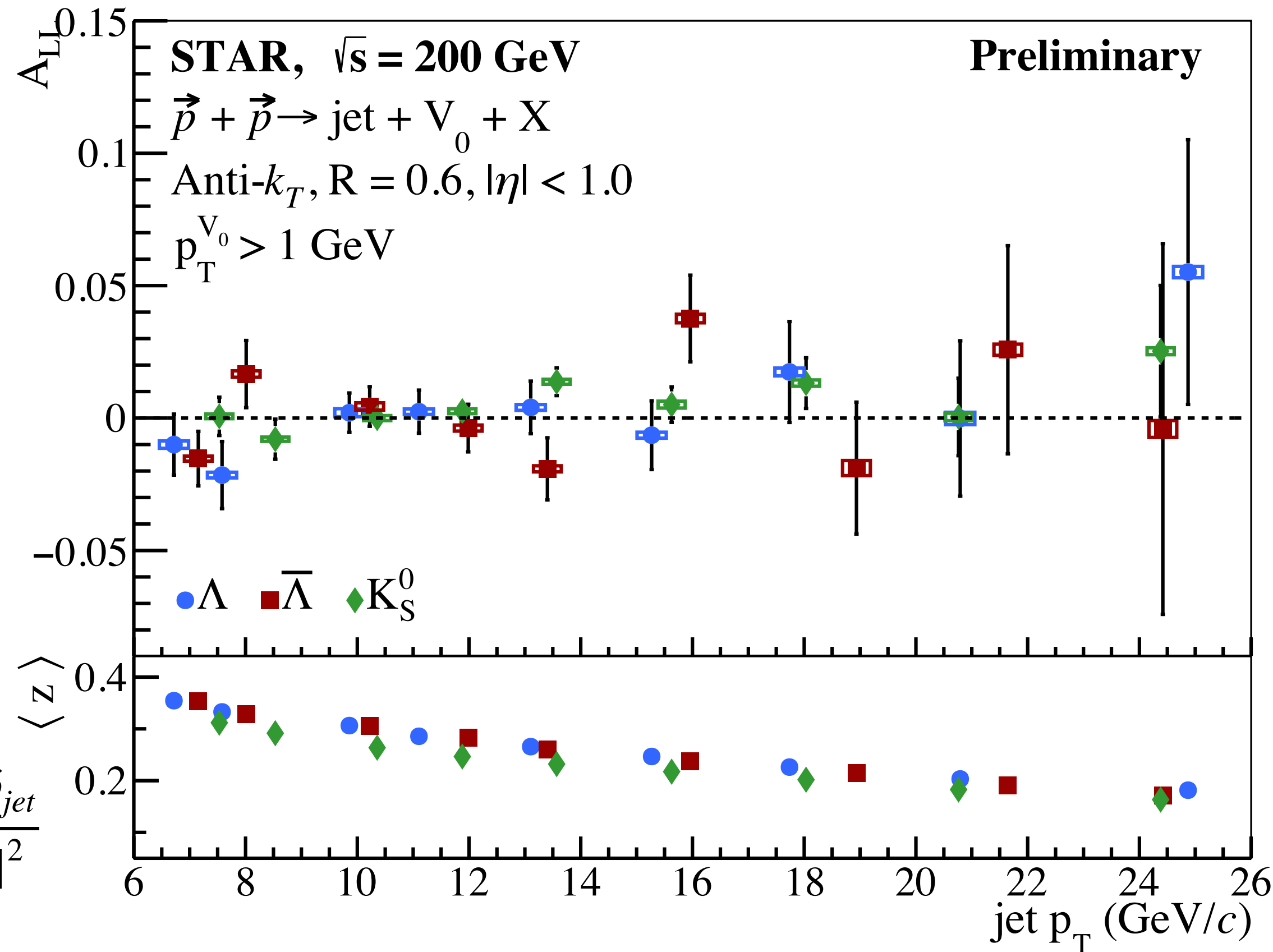


# $A_{LL}$ vs $p_T$ for $\Lambda$ , $\bar{\Lambda}$ and $K_S^0$



- First measurement  $A_{LL}$  vs  $p_T$  for  $\Lambda$ ,  $\bar{\Lambda}$  and  $K_S^0$  in polarized p+p collisions
- The results are independent of particle  $p_T$
- The results are consistent with zero
- Indication of small helicity distributions of  $s$  and  $\bar{s}$

# $\Lambda$ , $\bar{\Lambda}$ and $K_S^0$ tagged jet $A_{LL}$



- A subset of inclusive jets
- No jet  $p_T$  dependence
- Results are consistent with zero
- Provide constraints on strange quark helicity distribution



# Part III: Longitudinal spin transfer $D_{LL}$ of $\Lambda$ and $\bar{\Lambda}$





# $D_{LL}$ in p+p collision

## Prediction of $D_{LL}$ at RHIC energy

- Definition

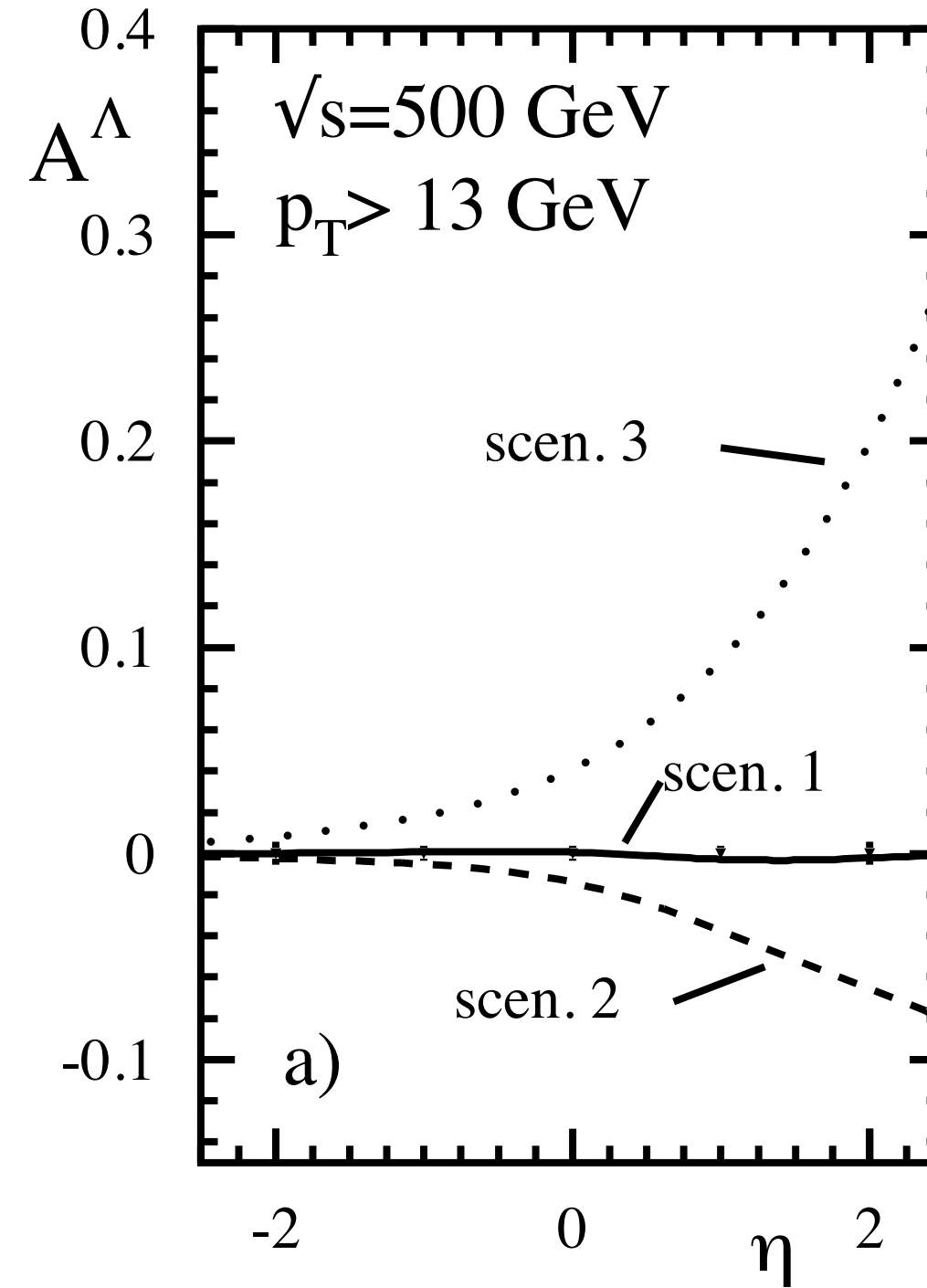
$$D_{LL}^{\Lambda} \equiv \frac{d\sigma^{p^+p \rightarrow \Lambda^+X} - d\sigma^{p^+p \rightarrow \Lambda^-X}}{d\sigma^{p^+p \rightarrow \Lambda^+X} + d\sigma^{p^+p \rightarrow \Lambda^-X}} = \frac{d\Delta\sigma}{d\sigma}$$

$$d\Delta\sigma \propto \Delta f_a(x_a) f_b(x_b) \Delta\sigma^{ab \rightarrow cd} \Delta D^{\Lambda}(z)$$

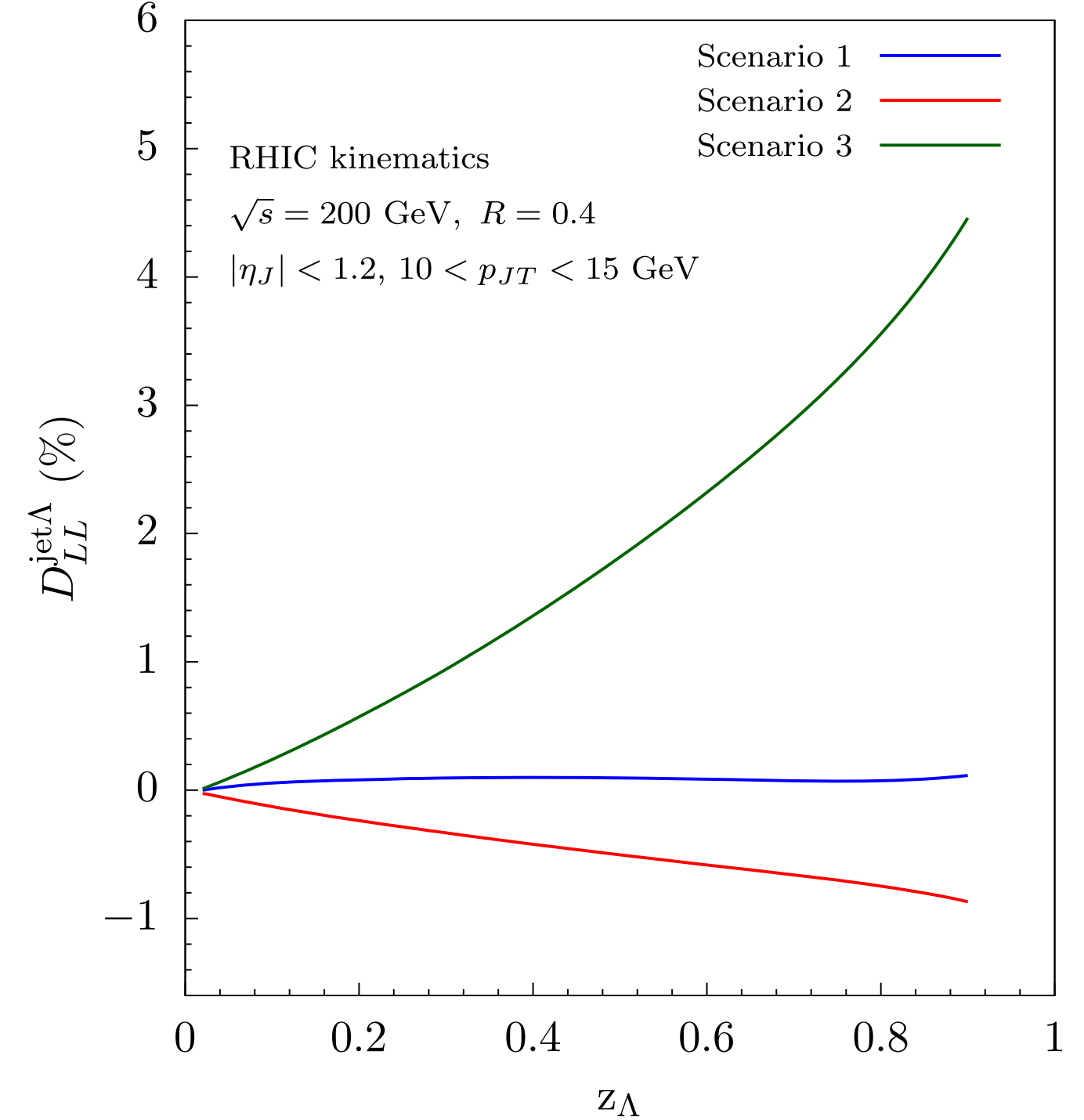
helicity distribution     
 pQCD calculable     
 longitudinally polarized FFs

- $D_{LL}$  can shed light on both polarized fragmentation functions (FFs) and the helicity distributions of  $s(\bar{s})$
- $D_{LL}$  vs  $z$  can provide direct probe to the polarized FFs

D. de Florian, M. Stratmann, and W. Vogelsang, Phys. Rev. Lett. **81**, 4 (1998).



Z.-B. Kang, K. Lee, and F. Zhao, Physics Letters B **809**, 135756 (2020).



**scenario 1:** only s quark can contribute to  $\Lambda$  polarization.

**scenario 2:** u and d quarks have the same contribution to polarized  $\Lambda$  but with an opposite sign from s quark.

**scenario 3:** u, d and s quarks have the same contribution to the polarized  $\Lambda$

# $D_{LL}$ Extraction

- $D_{LL}$  is measured with the asymmetry of  $\Lambda(\bar{\Lambda})$  yields as a function of  $\cos \theta^*$

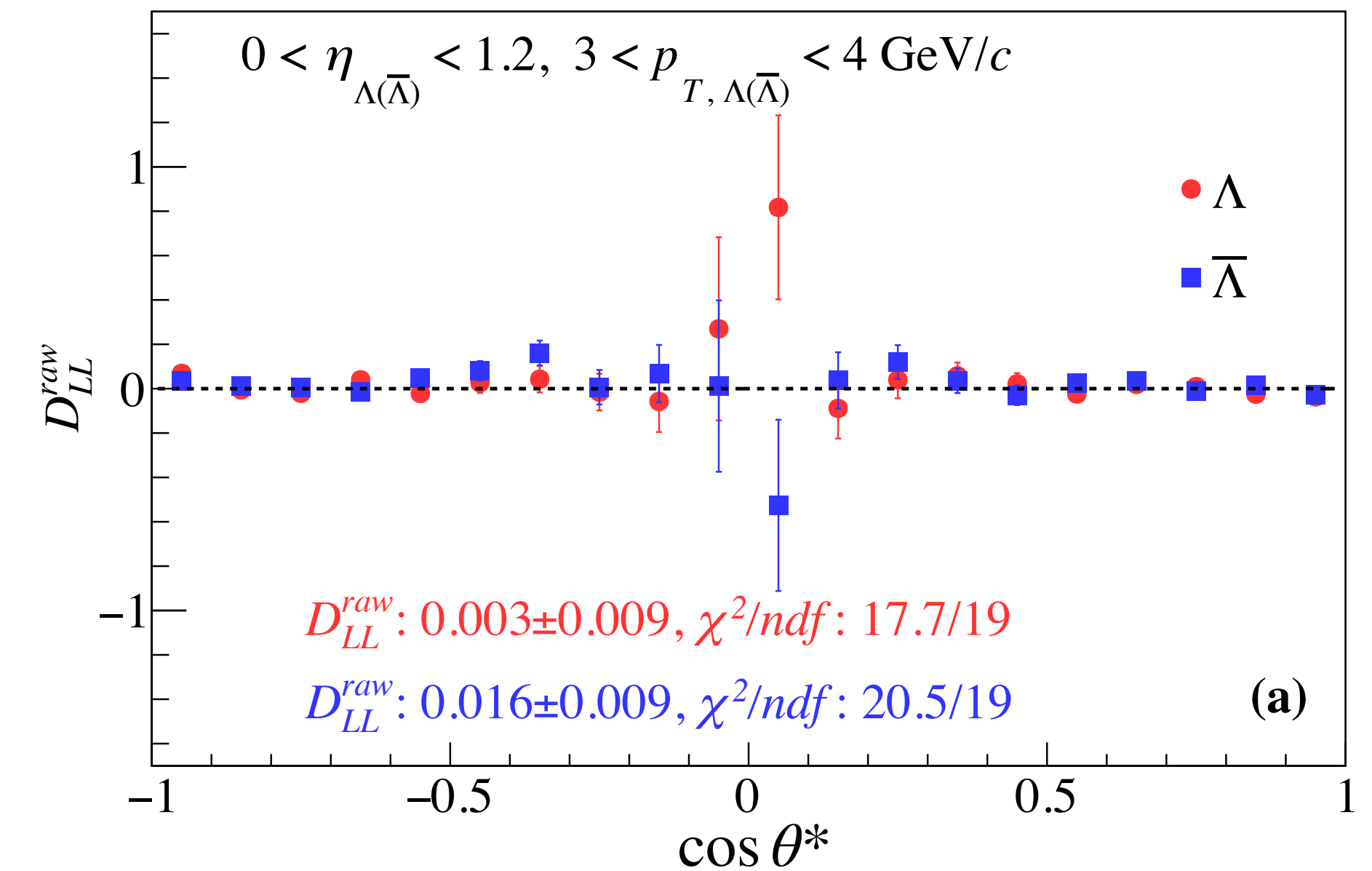
$$D_{LL} = \frac{1}{\alpha P_{beam} \langle \cos \theta^* \rangle} \frac{N^+ - RN^-}{N^+ + RN^-} \quad \text{Acceptance canceled}$$

firstly used in STAR, *Phys. Rev. D* 80, 111102 (2009).

**Background subtraction**

$$D_{LL} = \frac{D_{LL}^{raw} - rD_{LL}^{bkg}}{1 - r}$$

- ▶  $N^{+(-)}$ : the  $\Lambda$  yields with positive (negative) beam helicity
- ▶  $R$ : relative luminosity measured by the VPD
- ▶  $\alpha$ : decay parameter of  $\Lambda$ ,  $\alpha_{\Lambda} = 0.732$ ,  $\alpha_{\Lambda} = -\alpha_{\bar{\Lambda}}$
- ▶  $P_{beam}$ : the beam polarization

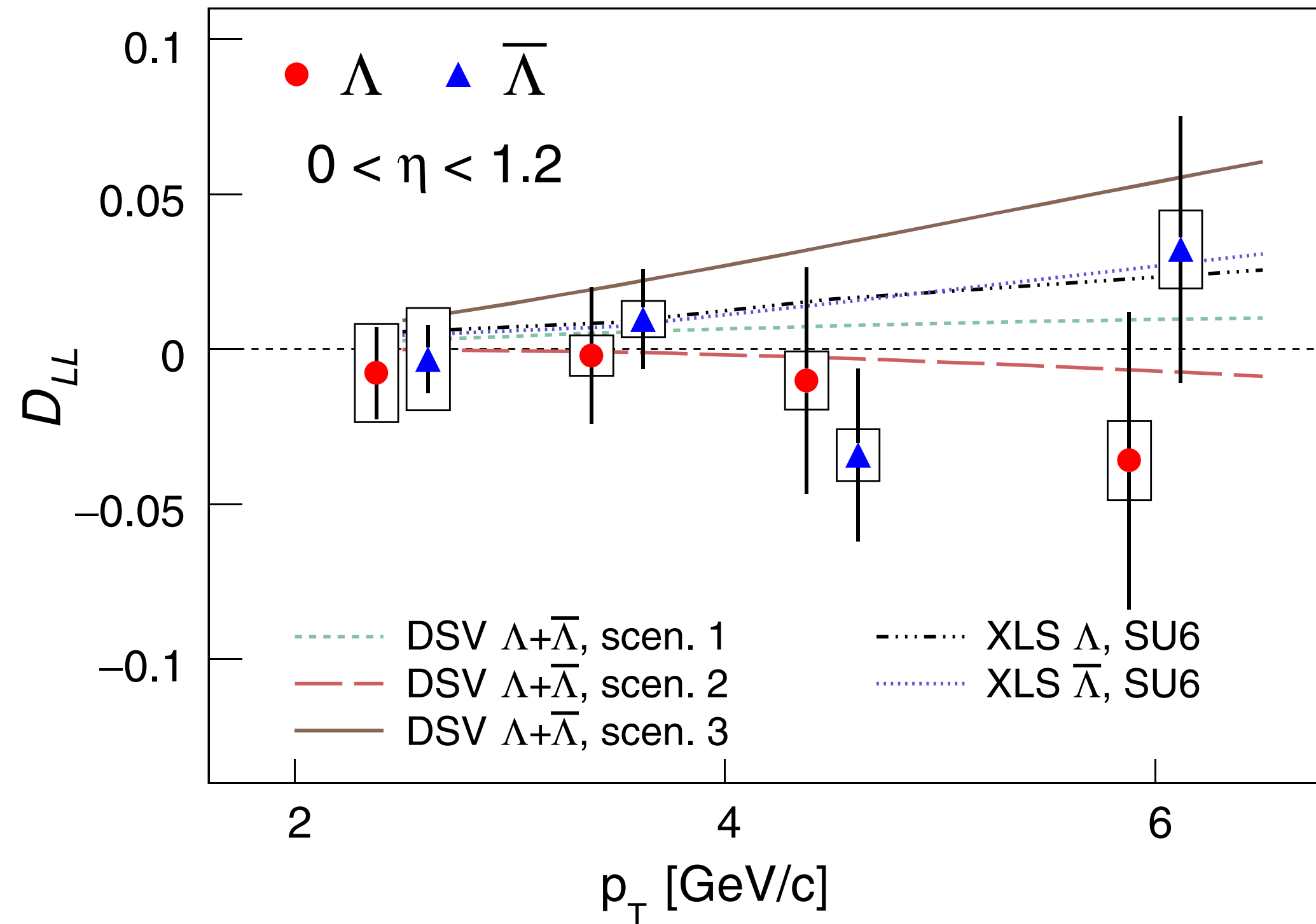


# Previous $D_{LL}$ vs $p_T$ results with STAR 2009 data



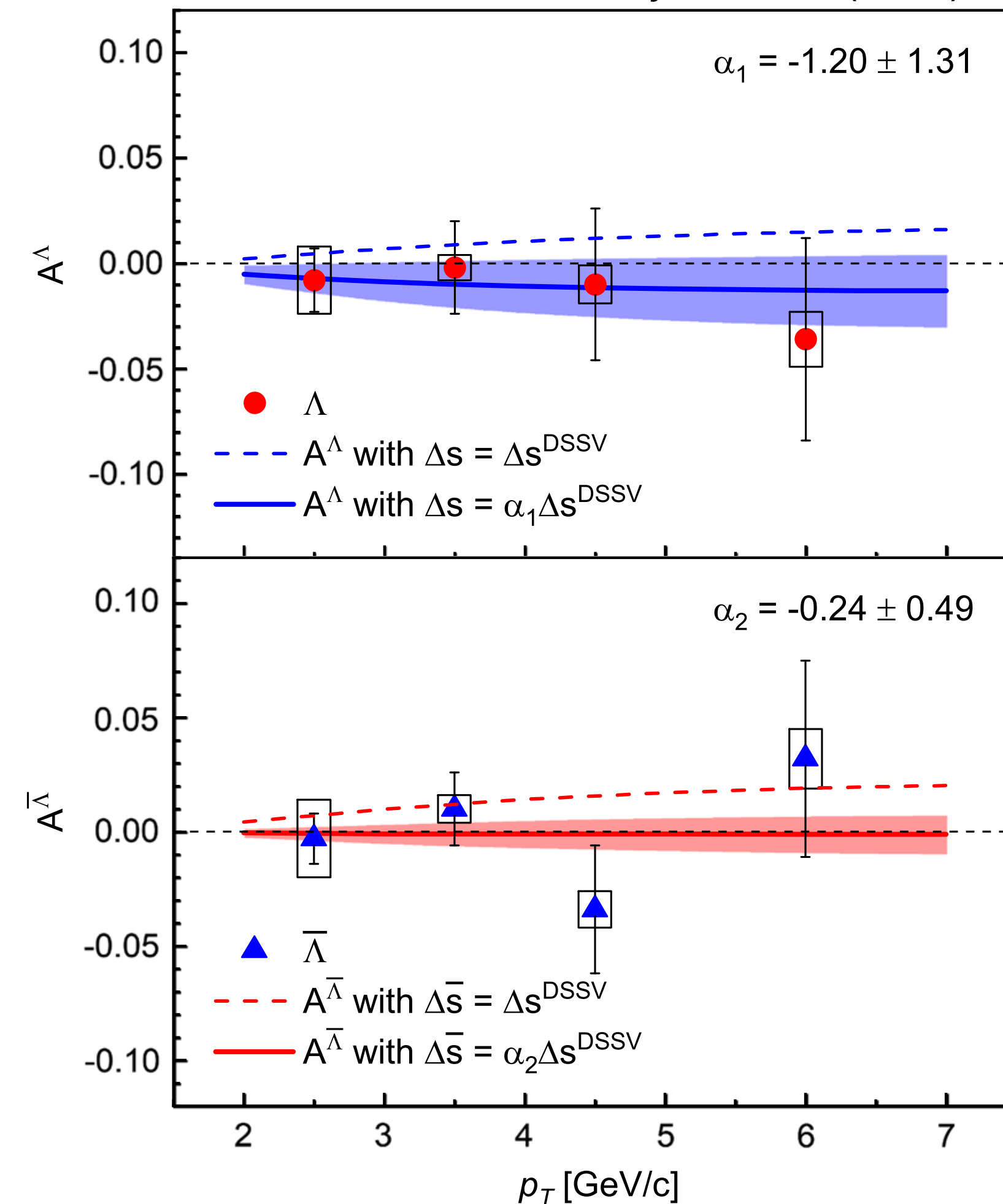
- Theoretical models, when fit to data, provide constraints to (anti)strange quark polarization

STAR, *Phys. Rev. D* **98**, 112009 (2018).



- Statistically limited.
- In agreement with models

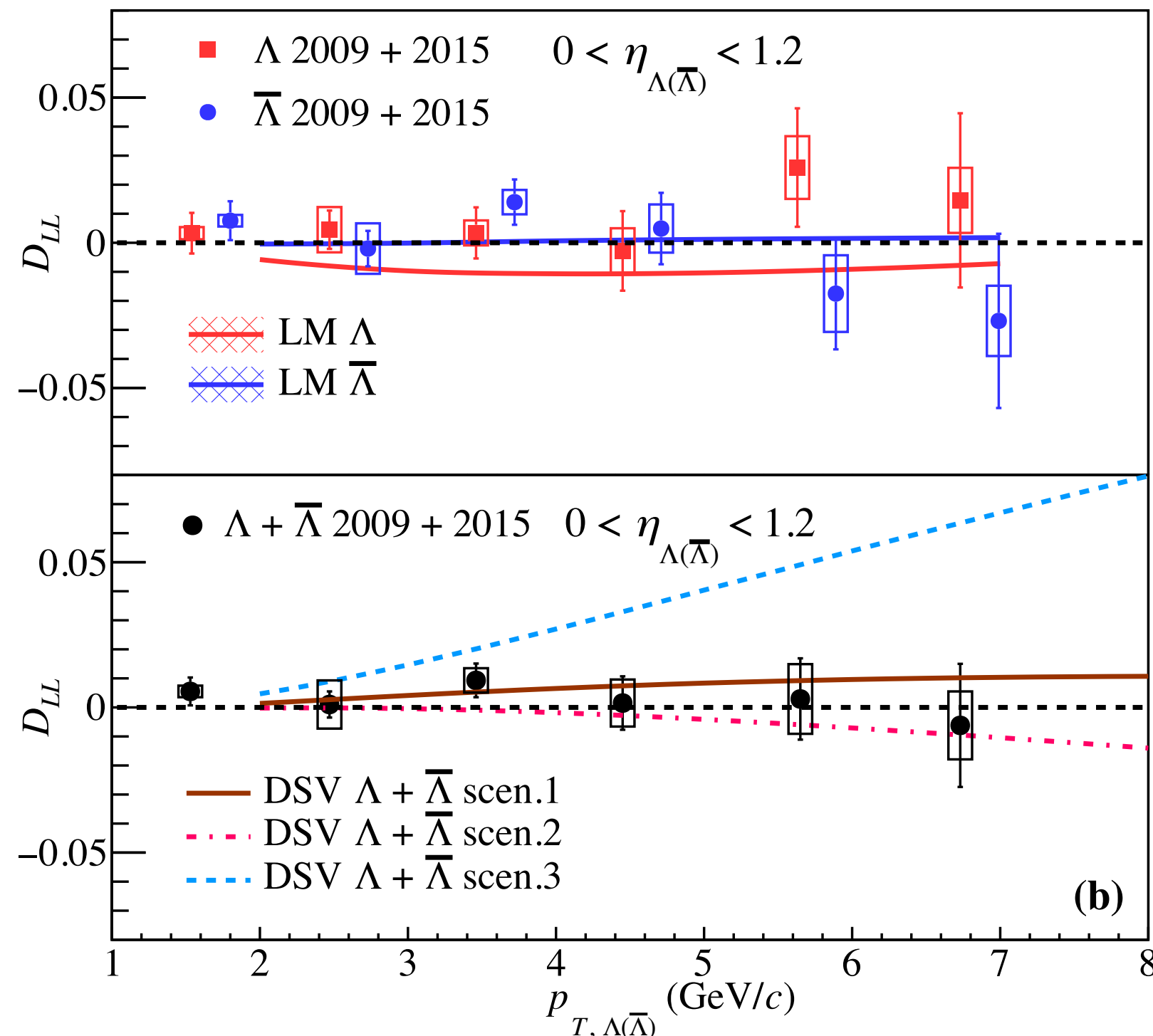
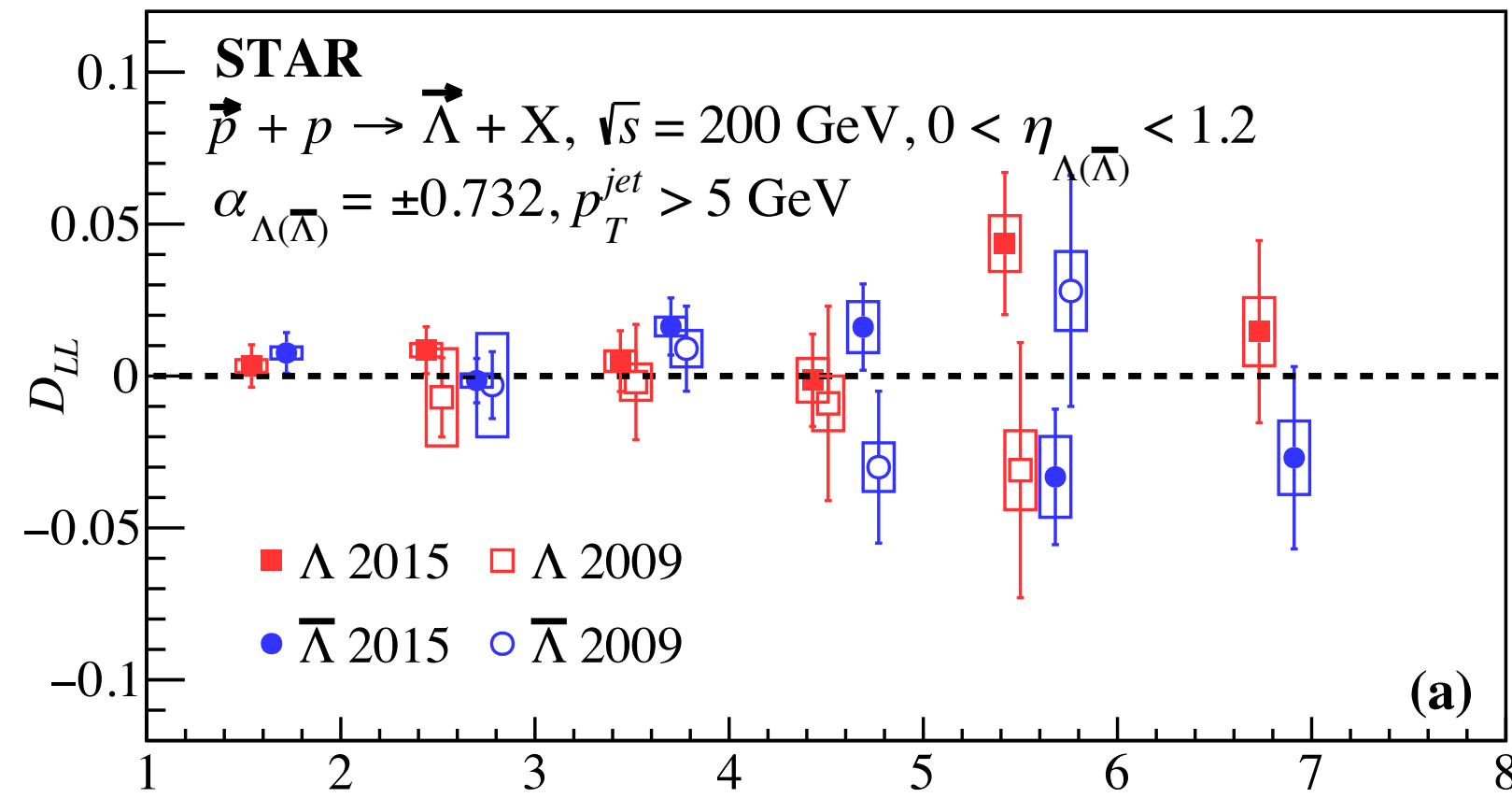
X.N. Liu, B.Q. Ma. *Eur. Phys. J. C* 10 (2019).





# New results of $D_{LL}$ vs $p_T$

[STAR], Phys. Rev. D **109**, 012004 (2024).



- Twice statistics larger as STAR 2009 data
- Most precise measurements up to date.
- Consistent results between  $\Lambda$  and  $\bar{\Lambda}$
- Two year's results are consistent
- Results are consistent with LM calculation
- Strong disfavor of the scenario 3 for the polarized FFs

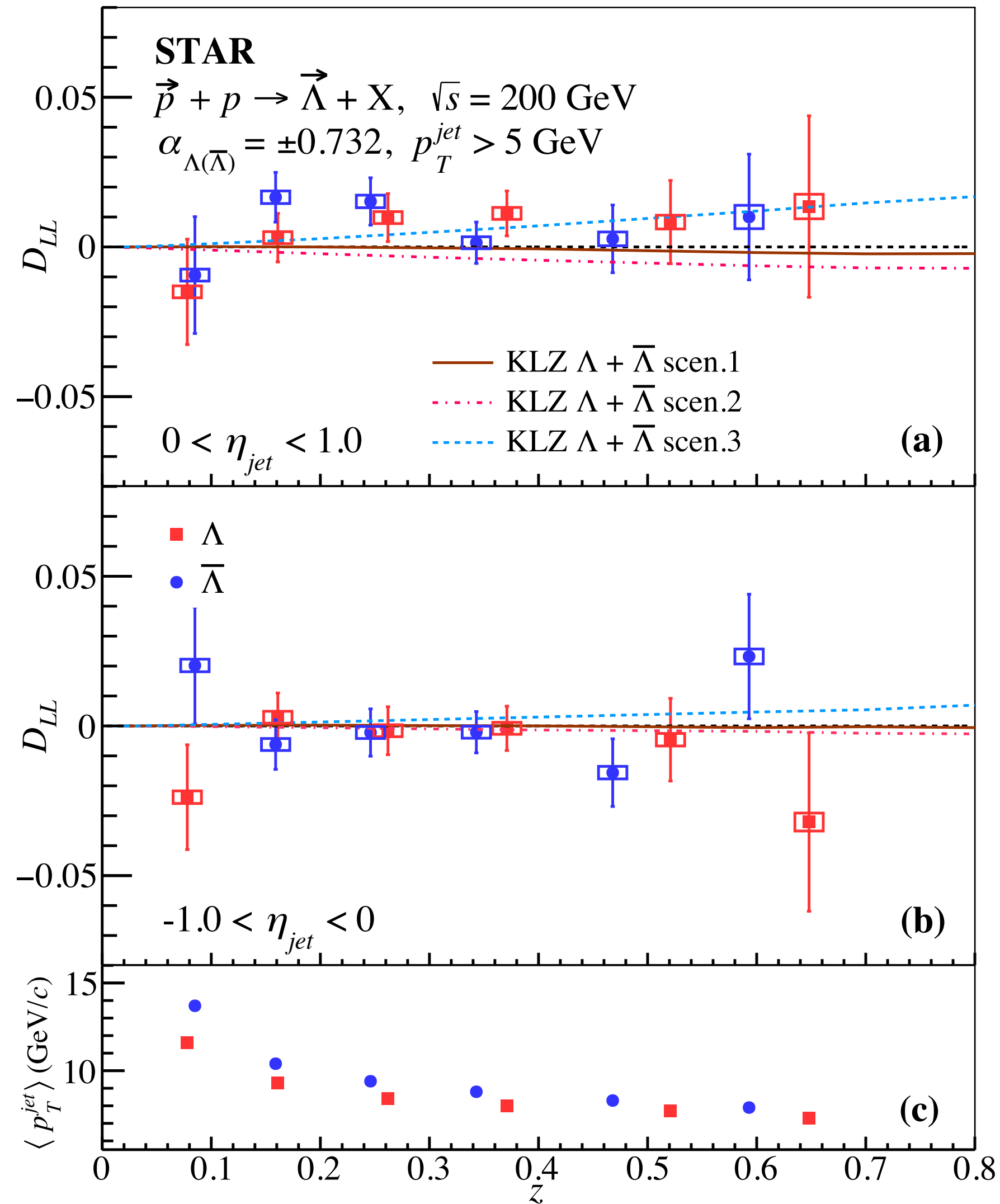
## Model predictions:

X.N. Liu, B.Q. Ma. *Eur. Phys. J. C* 10 (2019).

D. de Florian, M. Stratmann, and W. Vogelsang, *Phys. Rev. Lett.* 81, 530 (1998).

# First measurement of $D_{LL}$ vs $z$

[STAR], Phys. Rev. D **109**, 012004 (2024).



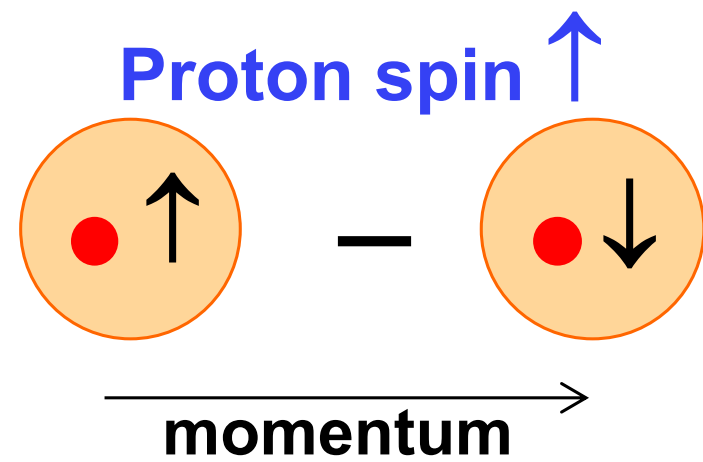
- The results directly probe the polarized fragmentation functions
- Results are comparable to model predictions within uncertainties
- Indication of small helicity distributions of (anti-) strange quark and/or small polarized fragmentation functions

**Model predictions:** Z.-B. Kang, K. Lee, and F. Zhao, Physics Letters B 809, 135756 (2020).



# Transverse spin transfer $D_{TT}$ in p+p collisions

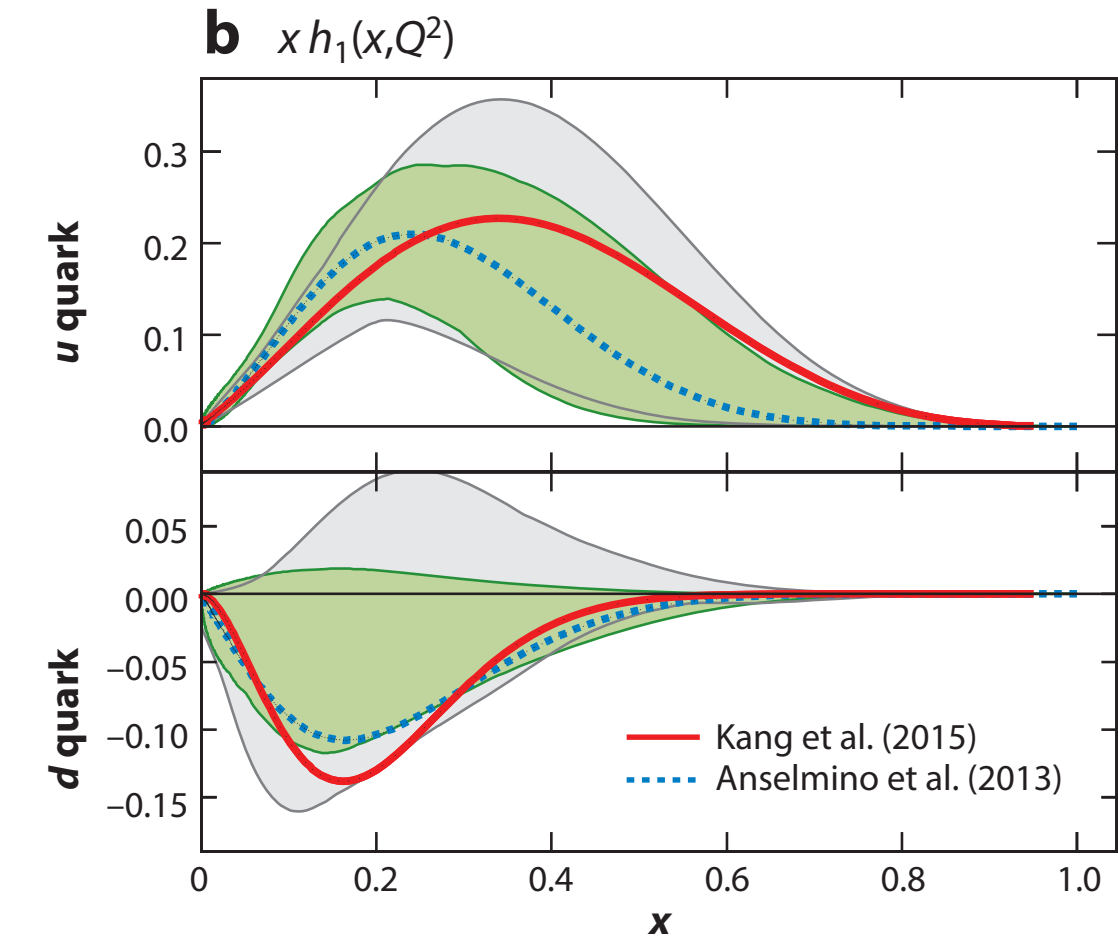
- Transversity**: least known leading twist parton distribution functions



$$\delta f(x) = f^\uparrow(x) - f^\downarrow(x)$$

Poor constraints on  $u$  and  $d$

Little info about  $s$  quark



- $D_{TT}$  can probe transversity

$$D_{TT}^\Lambda \equiv \frac{d\sigma^{p^\uparrow p \rightarrow \Lambda^\uparrow X} - d\sigma^{p^\uparrow p \rightarrow \Lambda^\downarrow X}}{d\sigma^{p^\uparrow p \rightarrow \Lambda^\uparrow X} + d\sigma^{p^\uparrow p \rightarrow \Lambda^\downarrow X}} = \frac{d\Delta_T \sigma}{d\sigma}$$

$$d\Delta_T \sigma \propto \delta f_a(x_a) f_b(x_b) \delta\sigma_T^{ab \rightarrow cd} \Delta_T D_c^\Lambda(z)$$

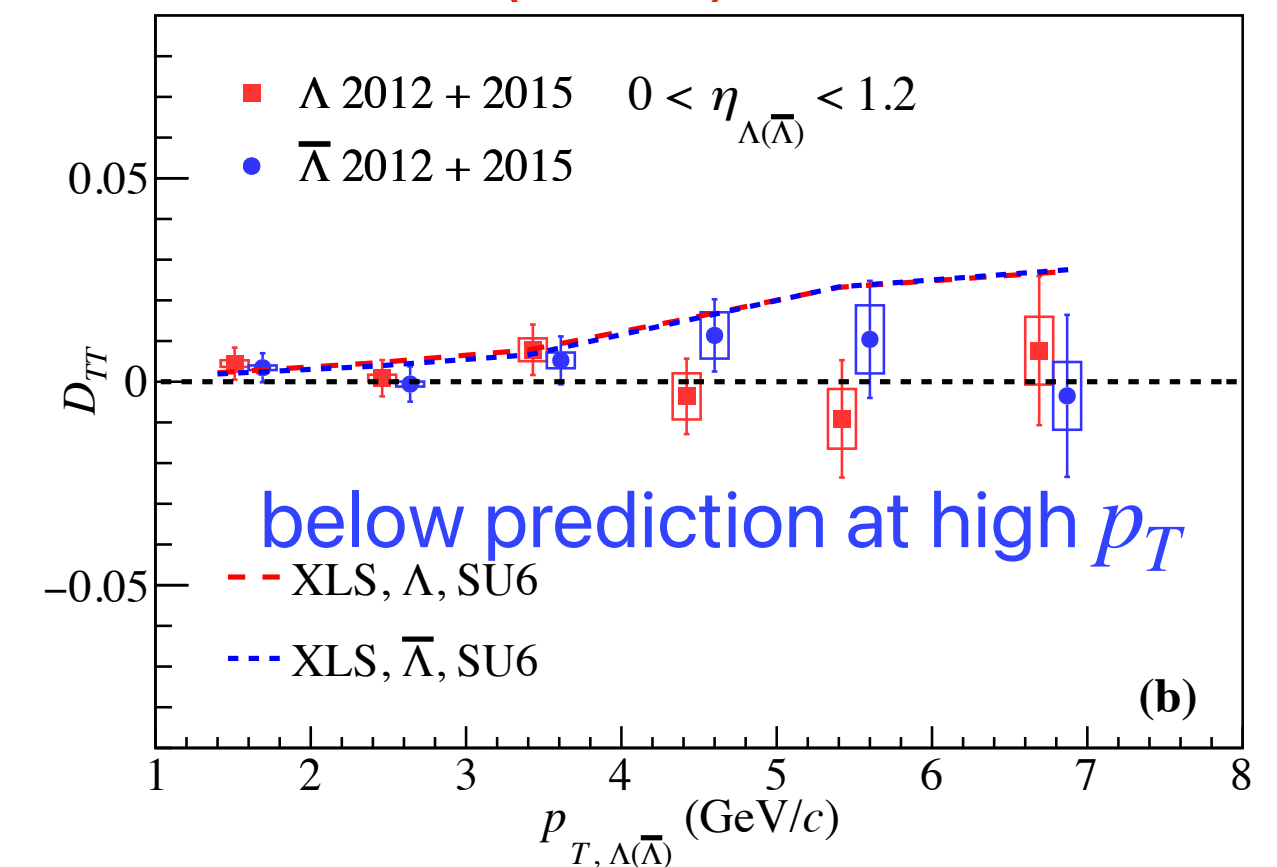
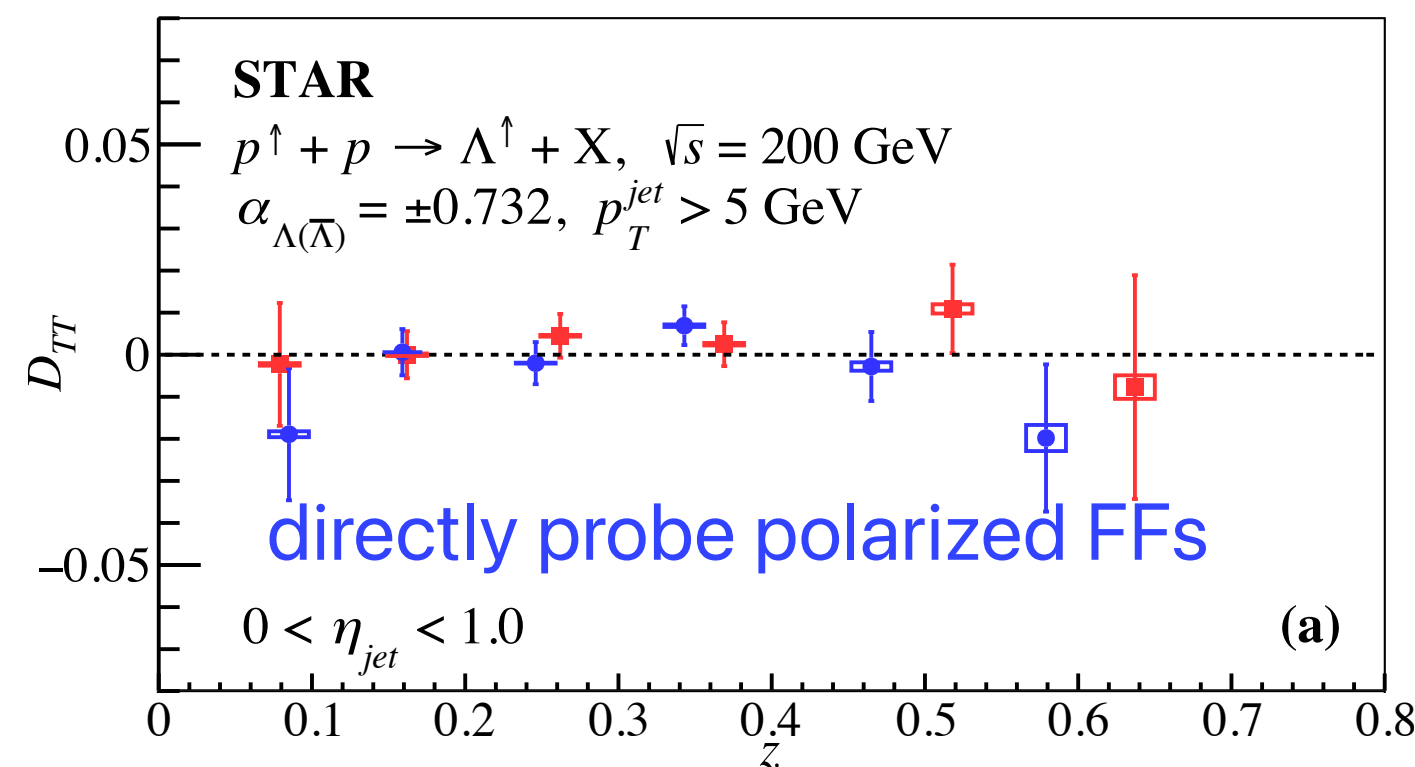
transversity

pQCD  
calculable

transversely  
polarized FFs

## new $D_{TT}$ measurements vs $p_T$ and $z$ at STAR

[STAR], Phys. Rev. D **109**, 012004 (2024).



model prediction

Q. H. Xu, Z. T. Liang, Phys. Rev. D 70, 034015 (2004).

Q. H. Xu, Z. T. Liang, and E. Sichterann, Phys. Rev. D 73, 077503 (2006).



# Summary

- $\pi^\pm$ -tagged jet  $A_{LL}$  in p+p collisions at  $\sqrt{s} = 200$  GeV at STAR
  - ▶ The results support positive  $\Delta g$
  - ▶  $A_{LL}$  is consistent with the prediction with NNPDFpol1.1 ( $\Delta g > 0$ )
  - ▶  $A_{LL}$  disfavors the prediction of JAM  $\Delta g < 0$
- $\Lambda$ ,  $\bar{\Lambda}$  and  $K_S^0$   $A_{LL}$  and  $D_{LL}$ 
  - ▶ First measurements of  $A_{LL}$  in polarized p+p collisions at  $\sqrt{s} = 200$  GeV
  - ▶ Indication of small strange quark and anti-quark helicity distribution
  - ▶  $D_{LL}$  disfavors the extreme scenario about the polarized FFs
  - ▶ First measurement of  $D_{LL}$  vs  $z$  provides direct access to the polarized FFs
- Larger data samples of p+p collisions at 510 GeV taken in 2012 and 2013 will improve the precision and extend to lower  $x$  region



# Backup



# Impact of the $\pi^\pm$ tagging

