

# Small- $x$ Quark and Gluon Helicity Contributions to the Proton Spin Puzzle

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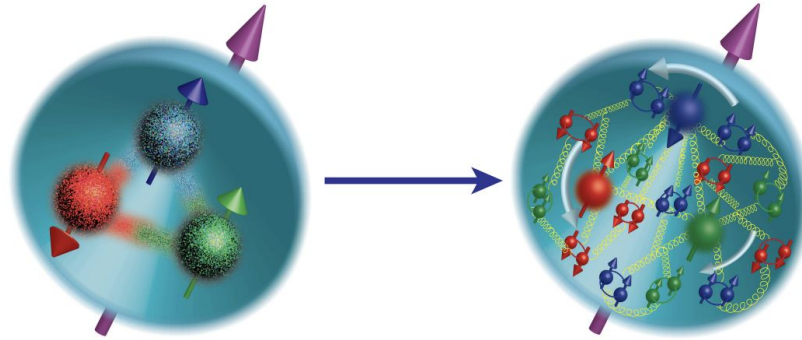
In collaboration with:

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D. Pitonyak, M. Sievert, N. Baldonado  
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Based on: 2204.11898, 2306.01651,  
2308.07461, and earlier publications

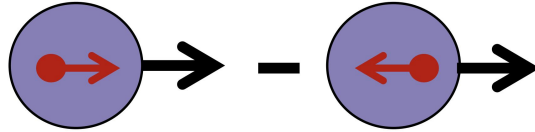


# Proton Spin



- In the past, proton spin was thought to be the sum of constituent quarks spins.
- Now, we believe it to be the sum of spins of valence quarks, sea quarks and gluons, together with their orbital angular momenta (OAM).

# Helicity PDF



- Helicity-dependent generalization of PDFs
- For each parton  $f$ ,

$$\Delta f(x, Q^2) = f(x, Q^2, +) - f(x, Q^2, -)$$

- For quarks, we often consider the “flavor singlet” quark hPDF:

$$\Delta\Sigma(x, Q^2) = \sum_{q=u,d,s} [\Delta q(x, Q^2) + \Delta\bar{q}(x, Q^2)]$$

and the “flavor non-singlet” quark hPDF:  $\Delta q^-(x, Q^2) = \Delta q(x, Q^2) - \Delta\bar{q}(x, Q^2)$

- Gluon hPDF:  $\Delta G(x, Q^2)$

# Proton Helicity Sum Rule

- Jaffe-Manohar sum rule:  $\frac{1}{2} = S_q + S_G + L_q + L_G$

where the helicity of quarks ( $S_q$ ) and gluons ( $S_G$ ) are

$$S_q(Q^2) = \frac{1}{2} \int_0^1 dx \Delta\Sigma(x, Q^2) \quad \text{and} \quad S_G(Q^2) = \int_0^1 dx \Delta G(x, Q^2)$$

- In the late 1980's, EMC measurement implied that  $S_q \approx 0.05$ , much lower than what would have been (1/2) had all the proton spin been carried by the constituent quarks.

# Current Knowledge of Proton Helicity

- More recently, the proton spin carried by quarks and gluon are estimated to be

$$S_q(Q^2 = 10 \text{ GeV}^2) \approx \frac{1}{2} \int_{0.001}^1 dx \Delta\Sigma(x, 10 \text{ GeV}^2) \in [0.15, 0.20]$$

$$S_G(Q^2 = 10 \text{ GeV}^2) \approx \int_{0.05}^1 dx \Delta G(x, 10 \text{ GeV}^2) \in [0.13, 0.26]$$

- They do not add to 1/2. The missing spin can come from:
  - Orbital angular momenta,  $L_q$  and  $L_G$ .
  - Small- $x$  region of  $\Delta\Sigma$  and  $\Delta G$ . Scattering experiments can only access finitely small  $x$ . The limit will improve with EIC.

$$\frac{1}{2} = S_q + S_G + L_q + L_G$$

$$S_q(Q^2) = \frac{1}{2} \int_0^1 dx \Delta\Sigma(x, Q^2)$$

$$S_G(Q^2) = \int_0^1 dx \Delta G(x, Q^2)$$

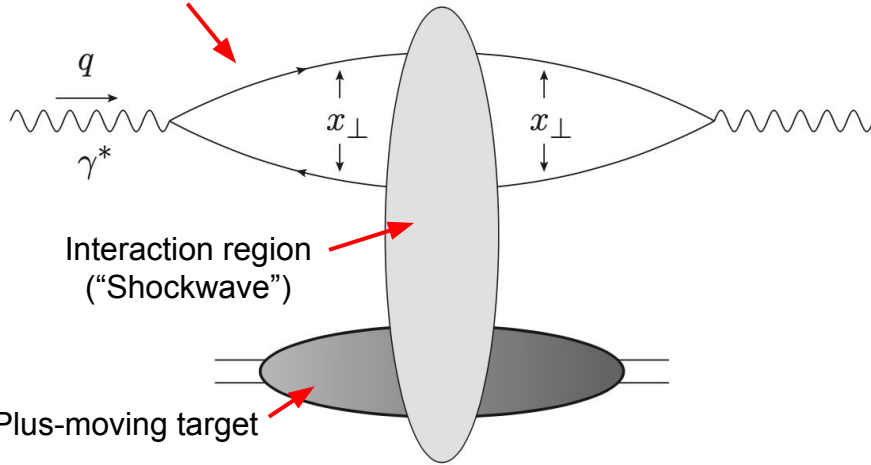
# DIS at Small x: The Dipole Picture

- Unpolarized PDF and structure functions,  $F_1(x, Q^2)$  and  $F_2(x, Q^2)$ , relate to the **s-matrix** of dipole-target scattering:

$$S(\underline{x}_1, \underline{x}_0, s) \equiv S_{10}(s) = \frac{1}{N_c} \left\langle \text{tr} \left[ V_{\underline{1}} V_{\underline{0}}^\dagger \right] \right\rangle (s)$$

Brackets: Averaging over target's state, including spin

Minus-moving dipole



where

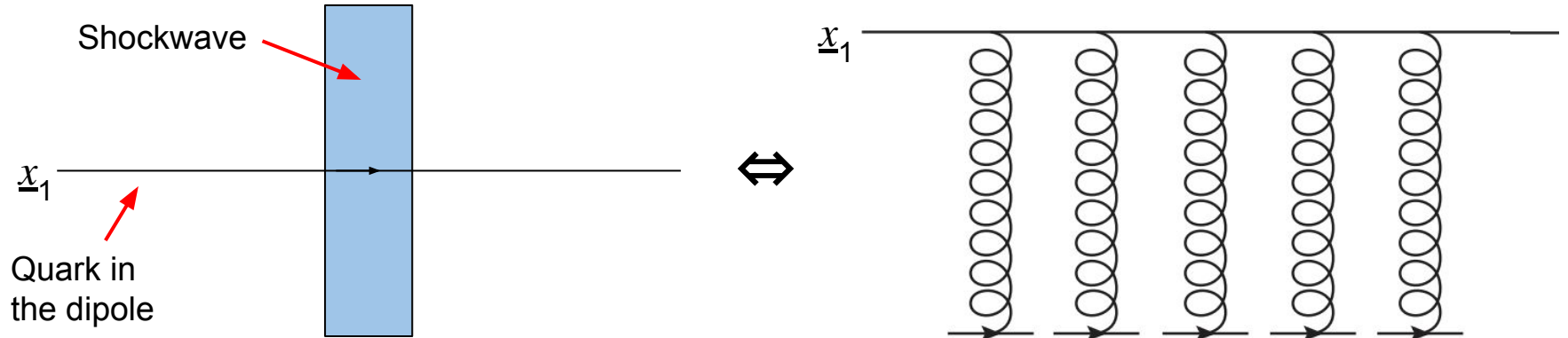
$$V_{\underline{1}}[x_f^-, x_i^-] \equiv V_{\underline{x}_1}[x_f^-, x_i^-] = \mathcal{P} \exp \left[ ig \int_{x_i^-}^{x_f^-} dx^- A^+(0^+, x^-, \underline{x}_1) \right]$$

$$V_{\underline{1}} \equiv V_{\underline{1}}[\infty, -\infty]$$

Lightcone (unpolarized) Wilson line

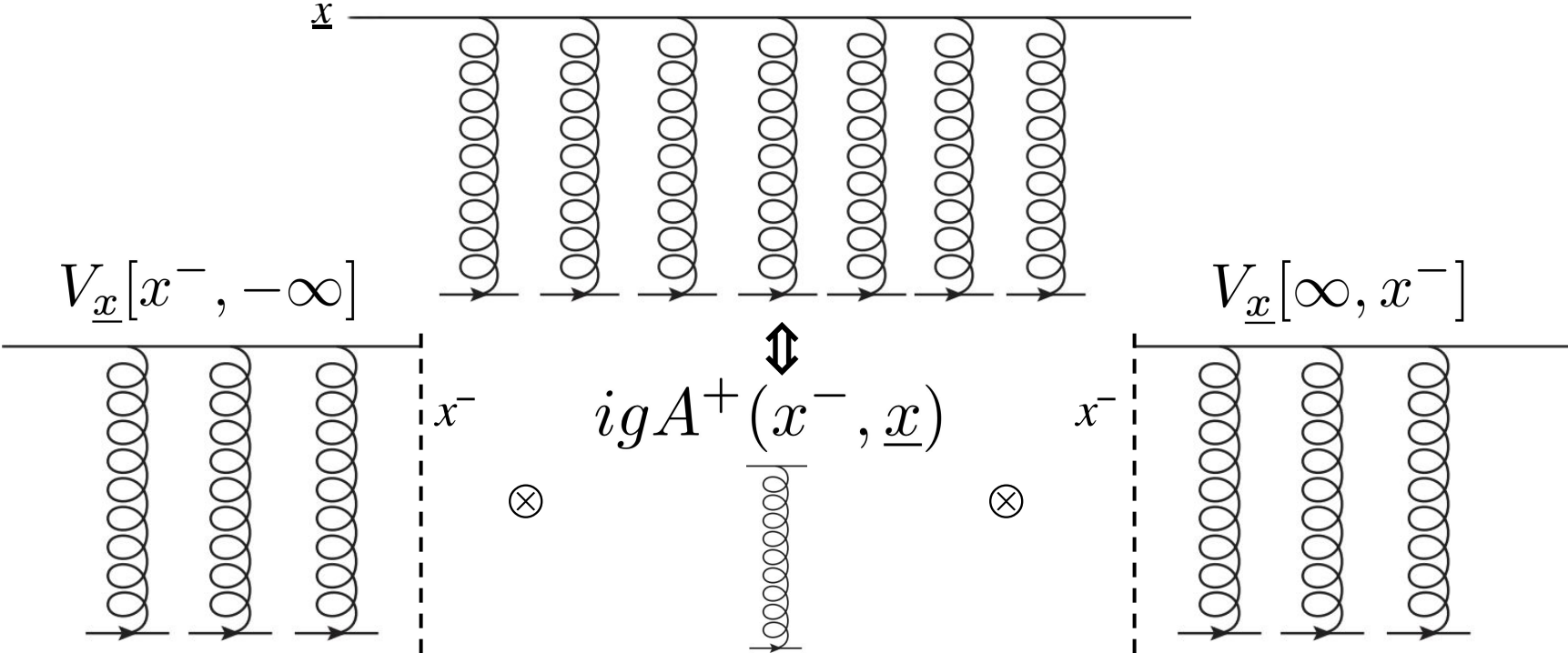
# Unpolarized Dipole Amplitude

- Parton **unpolarized PDF**,  $\Sigma(x, Q^2)$  and  $G(x, Q^2)$ , relate to **unpolarized dipole amplitude**,  $S_{10}(s) = \frac{1}{N_c} \left\langle \text{tr} \left[ V_{\underline{1}} V_{\underline{0}}^\dagger \right] \right\rangle (s)$ , which obeys BFKL/BK/JIMWLK evolution.
- Quark going through the shockwave at  $\underline{x}_1$ : unpolarized Wilson line, .
- Multiple parton exchanges at **eikonal** level (leading order in  $x$ ).



# Unpolarized Wilson Line

$$V_{\underline{x}_1}[x_f^-, x_i^-] = \mathcal{P} \exp \left[ ig \int_{x_i^-}^{x_f^-} dx^- A^+(0^+, x^-, \underline{x}_1) \right]$$



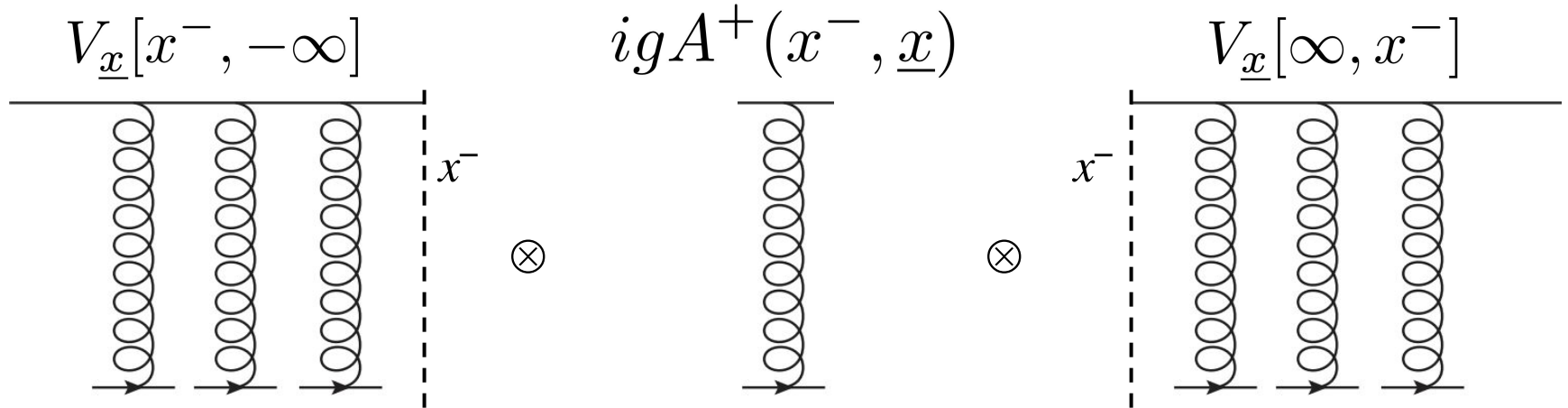


# Unpolarized Wilson Line

$$V_{\underline{x}_1}[x_f^-, x_i^-] = \mathcal{P} \exp \left[ ig \int_{x_i^-}^{x_f^-} dx^- A^+(0^+, x^-, \underline{x}_1) \right]$$

- Eikonal vertex insertion:

$$V_{\underline{x}} = ig \int_{-\infty}^{\infty} dx^- V_{\underline{x}}[\infty, x^-] A^+(x^-, \underline{x}) V_{\underline{x}}[x^-, -\infty]$$

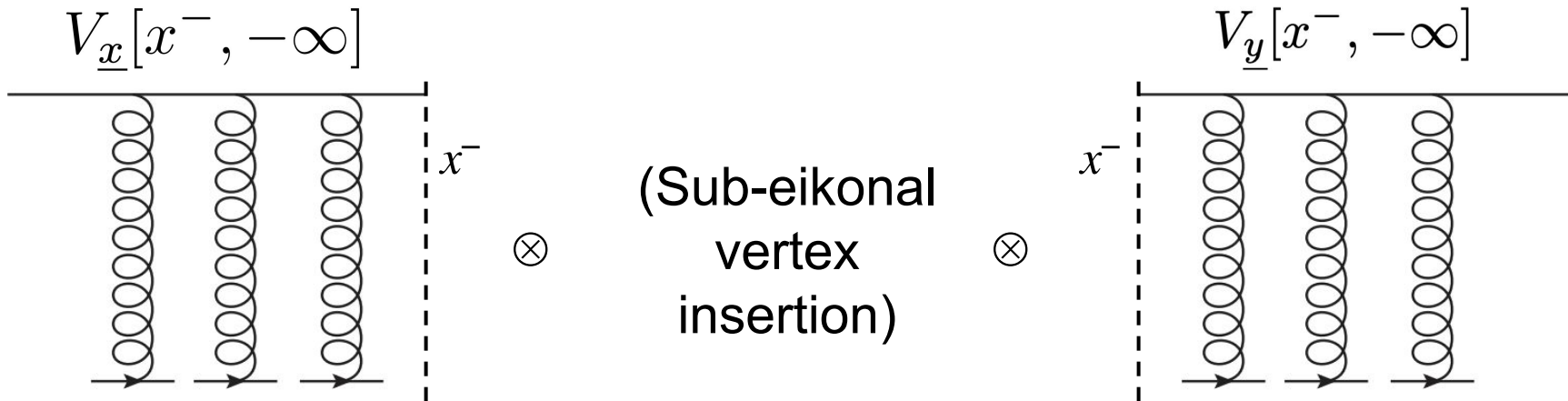


# Polarized Wilson Line

$$V_{\underline{x}_1}[x_f^-, x_i^-] = \mathcal{P} \exp \left[ ig \int_{x_i^-}^{x_f^-} dx^- A^+(0^+, x^-, \underline{x}_1) \right]$$

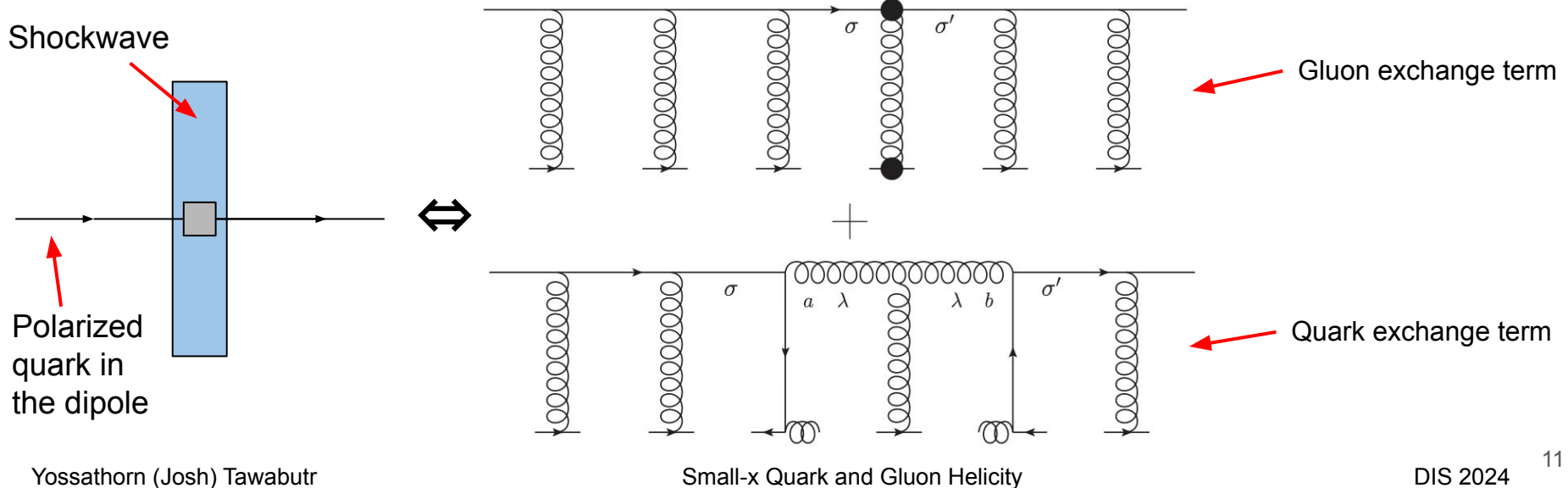
- Insertion of leading helicity-dependent vertex, which is
  - Sub-eikonal, i.e. (1/s)-suppressed
  - Not necessarily diagonal in transverse position
  - Denoted  $V_{\underline{x}, \underline{y}}^{\text{pol}}$

[Cougoulic, Kovchegov (YK),  
Tarasov, Tawabutr (JT),  
2204.11898 & predecessors]

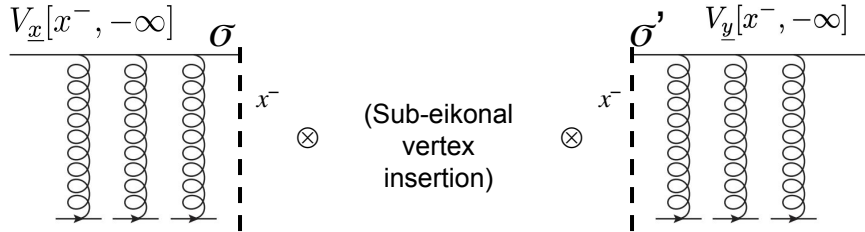


# Polarized Wilson Line

- Helicity-dependent quark line going through the shockwave corresponds to multiple eikonal parton exchanges, except for **one** helicity-dependent exchange, which is **sub-eikonal** (suppressed by an extra factor of  $x$ ).

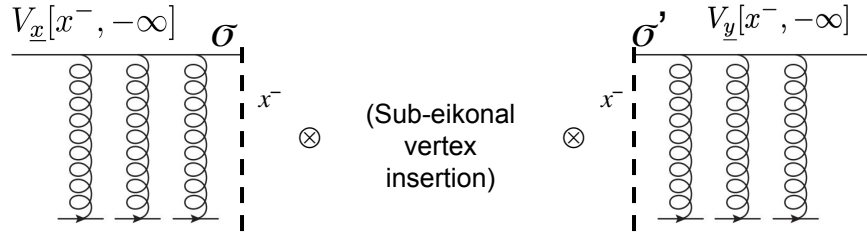


# Polarized Wilson Line



Polarized Wilson line	Type 1	Type 2
Helicity structure	$\sigma \delta_{\sigma, \sigma'}$	$\delta_{\sigma, \sigma'}$
Gluon exchange	$\sim F^{12} \delta^2(\underline{x} - \underline{y})$	$\sim \underline{\vec{D}} \cdot \underline{\vec{D}}$
Quark exchange	$\sim \psi (\gamma^+ \gamma_5) \bar{\psi} \delta^2(\underline{x} - \underline{y})$	N/A
Adjoint dipole	$\tilde{G}(x_{10}, zs)$	$G_2(x_{10}, zs)$
Fundamental dipole	$Q(x_{10}, zs)$	$G_2(x_{10}, zs)$

# Polarized Wilson Line



$$Q(x_{10}, z_s) \sim \left\langle \text{tr} \left[ V_{\underline{0}} V_{\underline{1}}^{\text{pol}[1]\dagger} \right] + \text{tr} \left[ V_{\underline{1}}^{\text{pol}[1]} V_{\underline{0}}^\dagger \right] \right\rangle$$

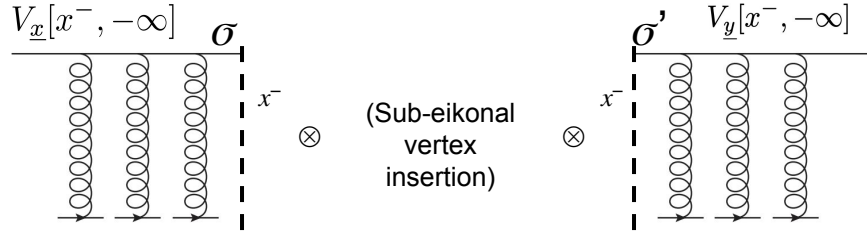
$$\tilde{G}(x_{10}, z_s) \sim \left\langle \text{Tr} \left[ U_{\underline{0}} U_{\underline{1}}^{\text{pol}[1]\dagger} \right] + \text{Tr} \left[ U_{\underline{1}}^{\text{pol}[1]} U_{\underline{0}}^\dagger \right] \right\rangle$$

$$G_2(x_{10}, z_s) \sim \left\langle \text{tr} \left[ V_{\underline{0}} V_{\underline{1}}^{\text{pol}[2]\dagger} \right] + \text{tr} \left[ V_{\underline{1}}^{\text{pol}[2]} V_{\underline{0}}^\dagger \right] \right\rangle$$

Brackets now include  $\frac{1}{2} \sum_S S$  of proton helicity

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# Polarized Wilson Line



$$\Delta\Sigma(x, Q^2) = -\frac{N_c N_f}{2\pi^3} \int_{\Lambda^2/s}^1 \frac{dz}{z} \int_{\frac{1}{zs}}^{\min\{\frac{1}{zQ^2}, \frac{1}{\Lambda^2}\}} \frac{dx_{10}^2}{x_{10}^2} [Q(x_{10}^2, zs) + 2G_2(x_{10}^2, zs)]$$

$$\Delta G(x, Q^2) = \frac{2N_c}{\alpha_s \pi^2} \left[ \left( 1 + x_{10}^2 \frac{\partial}{\partial x_{10}^2} \right) G_2 \left( x_{10}^2, zs = \frac{Q^2}{x} \right) \right]_{x_{10}^2 = \frac{1}{Q^2}}$$

$$g_1(x, Q^2) = -\sum_f \frac{N_c Z_f^2}{4\pi^3} \int_{\Lambda^2/s}^1 \frac{dz}{z} \int_{\frac{1}{zs}}^{\min\{\frac{1}{zQ^2}, \frac{1}{\Lambda^2}\}} \frac{dx_{10}^2}{x_{10}^2} [Q(x_{10}^2, zs) + 2G_2(x_{10}^2, zs)]$$

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# Small-x Asymptotics with Quark Exchanges (Large $N_c$ & $N_f$ )

- At small  $x$ , gluons dominate  $\rightarrow N_c \gg 1$
- Still important to include quark exchanges ( $\sim N_f/N_c$ ) for helicity evolution
- Flavor non-singlet hPDF:

$$\Delta q^-(x, Q^2) = \Delta q(x, Q^2) - \Delta \bar{q}(x, Q^2) \sim \left(\frac{1}{x}\right)^{\sqrt{\alpha_s N_c / \pi}} \quad [\text{YK, Pitonyak, Sievert, 1610.06197}]$$

- Flavor singlet hPDF:

$$\begin{aligned} \Delta \Sigma(x, Q^2) &= \sum_{q=u,d,s} [\Delta q(x, Q^2) + \Delta \bar{q}(x, Q^2)] \\ &\sim \Delta G(x, Q^2) \sim g_1(x, Q^2) \sim \left(\frac{1}{x}\right)^{3.43\sqrt{\alpha_s N_c / 2\pi}} \end{aligned} \quad [\text{Adamiak, YK, JT, 2306.01651}]$$

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Smaller than 1



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[YK, Pitonyak, Sievert, 1610.06197]

- Flavor singlet hPDF (with  $N_f = 3$ ):

$$\Delta \Sigma(x, Q^2) = \sum_{q=u,d,s} [\Delta q(x, Q^2) + \Delta \bar{q}(x, Q^2)]$$

$$\sim \Delta G(x, Q^2) \sim g_1(x, Q^2) \sim \left(\frac{1}{x}\right)^{3.43 \sqrt{\alpha_s N_c / 2\pi}}$$

Smaller than 1

Exceed 1 for  $\alpha_s \gtrsim 0.18$

Infinite spin from small  $x$ ???

[Adamiak, YK, JT, 2306.01651]

# Corrections to the DLA Evolution

- So far, helicity evolution resums  $\alpha_s \ln^2(1/x)$ .
- Potentially significant **single-log corrections**, resumming  $\alpha_s \ln(1/x)$ .
  - Convoluting with unpolarized dipoles, which obey BK evolution
  - Likely to include saturation mechanism
  - See [YK, Tarasov, JT, 2104.11765] and upcoming work

# Corrections to the DLA Evolution

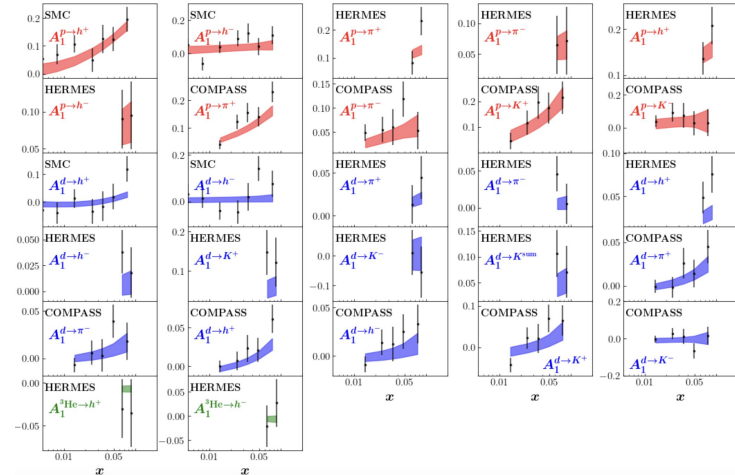
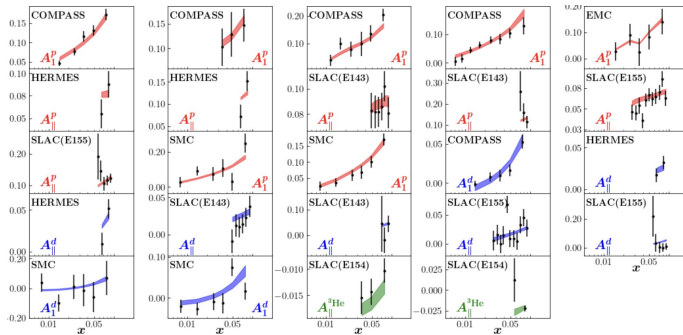
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  - Convoluting with unpolarized dipoles, which obey BK evolution
  - Likely to include saturation mechanism
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- Recently, a **running coupling correction** (daughter dipole prescription) is employed to the DLA evolution in a global fit with polarized DIS & SIDIS data.
- KPS-CTT evolution (with rc) starts at  $x_0 = 0.1$ .
- At larger  $x$ , employ generalized Born-level initial condition:  
Dipole  $\sim a \ln(\text{rapidity}) + b \ln(\text{dipole size}) + c$

[Adamiak et al,  
2308.07461]

# Global Fit

[Adamiak et al, 2308.07461]

- Polarized DIS and SIDIS data ( $A_1$ ,  $A_{\parallel}$ ,  $A_1^h$ ) from SLAC, EMC, SMC, COMPASS and HERMES at  $0.005 \leq x \leq 0.1$  and  $1.69 \text{ GeV}^2 \leq Q^2 \leq 10.4 \text{ GeV}^2$ .
  - Include proton, deuteron and helium-3 targets
  - For SIDIS, include  $\pi^{\pm}$ ,  $K^{\pm}$  and unidentified charged hadron productions
- In total,  $N_{\text{pts}} = 226$  data points
- Overall,  $\chi^2 / N_{\text{pts}} = 1.03$

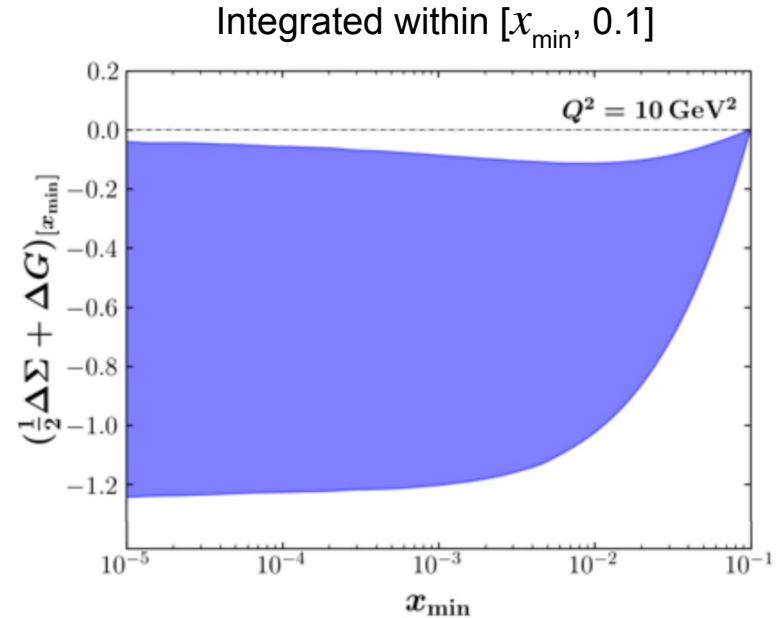
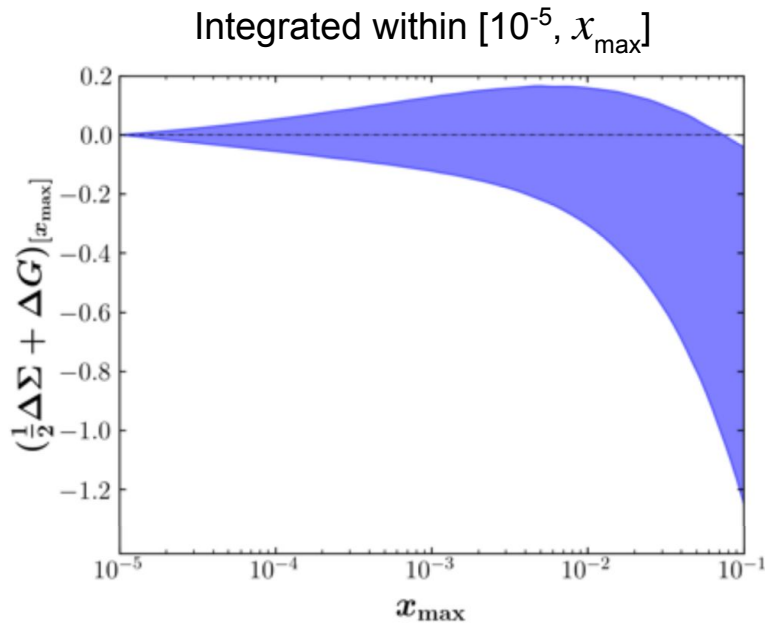


# Global Fit

[Adamiak et al, 2308.07461]

$$S_q + S_G \approx \int_{10^{-5}}^{0.1} dx \left( \frac{1}{2} \Delta\Sigma + \Delta G \right) (x) = -0.64 \pm 0.60$$

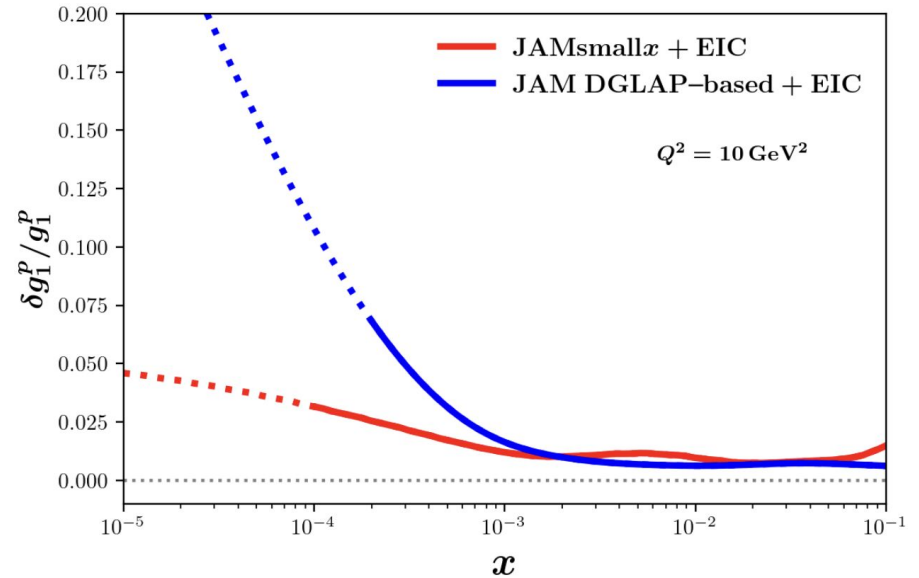
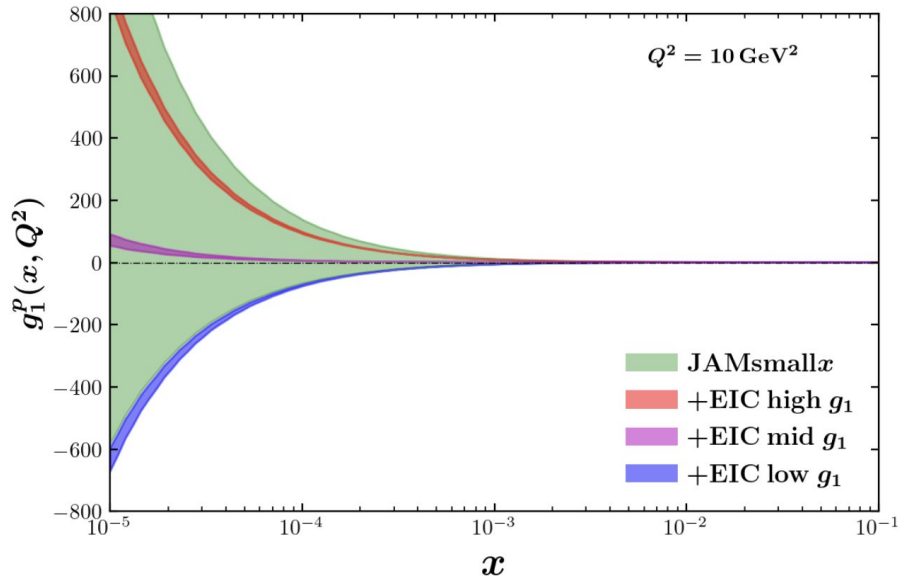
Significant spin from small  $x$



# Future EIC Impact

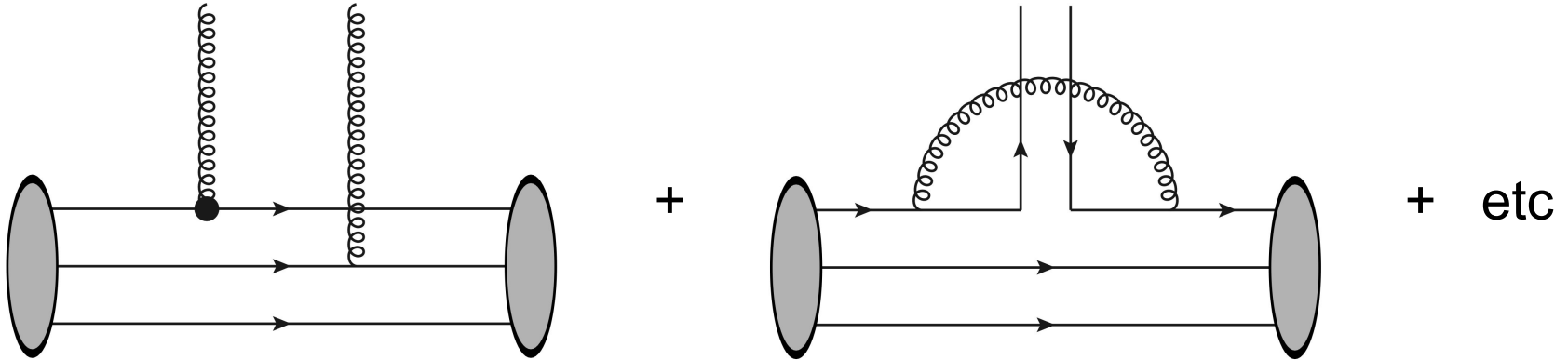
[Adamiak et al, 2308.07461]

- Significant reduction of uncertainty at small  $x$  with future EIC data.



# Global Fit: Next Step

- To allow the data to fix the total helicity, we need a more deterministic IC.
- Model proton target at moderate  $x$  with 3 valence quarks, c.f. [Dumitru et al 2010.11245, 2303.16339].
- Stay tuned

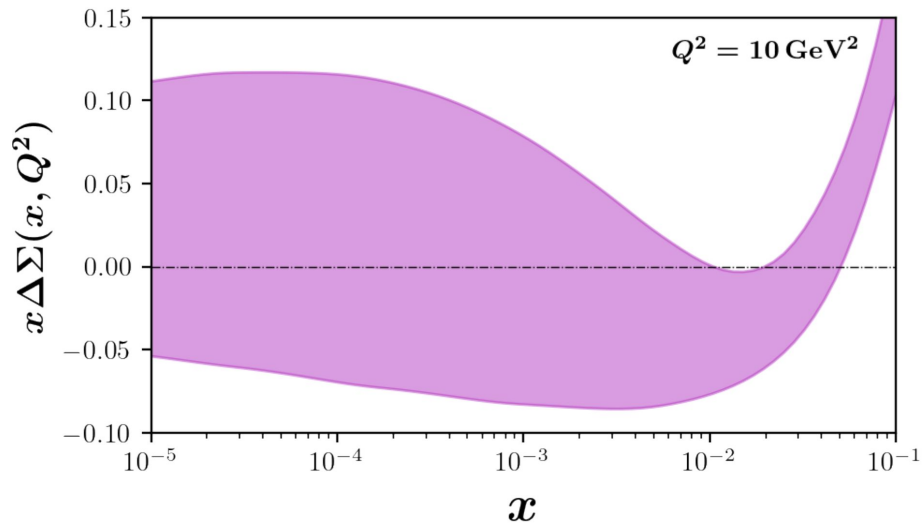
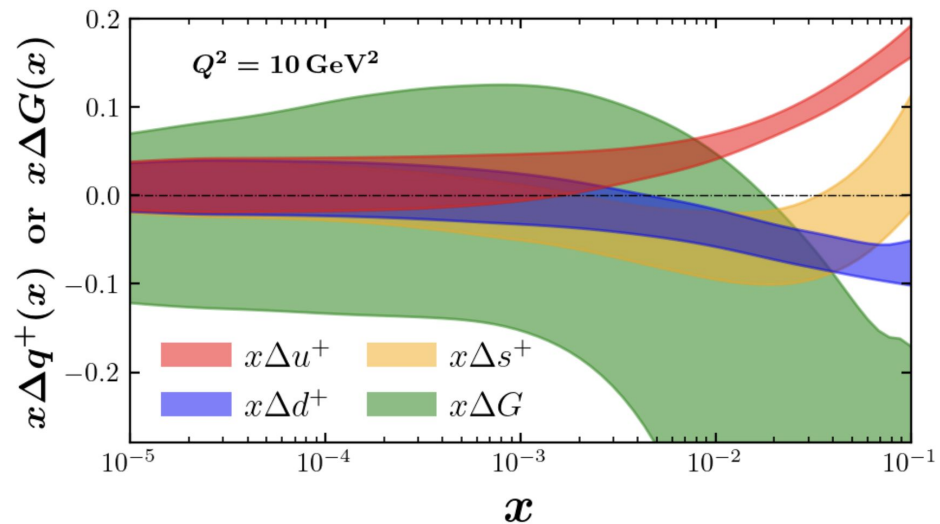


# Conclusion

- Already at DLA, KPS-CTT evolution provides a promising small- $x$  description of parton helicity, with potential improvement from future EIC results.
- Future work:
  - More deterministic initial condition using a valence-quark wave function
  - Improved global fit that includes  $pp$  particle production data
  - Complete single-logarithmic corrections, which will incorporate saturation
- The framework can be modified to calculate OAM's [YK, Manley, 1901.07453, 2310.18404] and other TMD's [YK, Santiago, 2108.03667, 2209.03538, 2310.02231].
- The framework has been generalized to helicity-JIMWLK evolution and helicity-dependent extension to MV model [Cougoulic, YK, 1910.04268, 2005.14688].



# Global Fit: hPDF Results



# Global Fit: Data Points

Data set ( $A_1$ )	Target	$N_{\text{pts}}$	$\chi^2/N_{\text{pts}}$
SLAC (E142) [141]	${}^3\text{He}$	1	0.60
EMC [146]	$p$	5	0.20
SMC [147, 149]	$p$	6	1.29
	$p$	6	0.53
	$d$	6	0.67
	$d$	6	2.26
COMPASS [150]	$p$	5	1.02
COMPASS [151]	$p$	17	0.74
COMPASS [152]	$d$	5	0.88
HERMES [153]	$n$	2	0.73
<b>Total</b>		59	0.91

Data set ( $A_{\parallel}$ )	Target	$N_{\text{pts}}$	$\chi^2/N_{\text{pts}}$
SLAC(E155) [144]	$p$	16	1.28
	$d$	16	1.62
SLAC (E143) [143]	$p$	9	0.56
	$d$	9	0.92
SLAC (E154) [142]	${}^3\text{He}$	5	1.09
HERMES [154]	$p$	4	1.54
	$d$	4	0.98
<b>Total</b>		63	1.19

Dataset ( $A_1^h$ )	Target	Tagged Hadron	$N_{\text{pts}}$	$\chi^2/N_{\text{pts}}$
SMC [148]	$p$	$h^+$	7	1.03
	$p$	$h^-$	7	1.45
	$d$	$h^+$	7	0.82
	$d$	$h^-$	7	1.49
HERMES [158]	$p$	$\pi^+$	2	2.39
	$p$	$\pi^-$	2	0.01
	$p$	$h^+$	2	0.79
	$p$	$h^-$	2	0.05
	$d$	$\pi^+$	2	0.47
	$d$	$\pi^-$	2	1.40
	$d$	$h^+$	2	2.84
	$d$	$h^-$	2	1.22
	$d$	$K^+$	2	1.81
	$d$	$K^-$	2	0.27
	$K^+ + K^-$	2	0.97	
HERMES [159]	${}^3\text{He}$	$h^+$	2	0.49
	${}^3\text{He}$	$h^-$	2	0.29
COMPASS [156]	$p$	$\pi^+$	5	1.88
	$p$	$\pi^-$	5	1.10
	$p$	$K^+$	5	0.42
	$p$	$K^-$	5	0.31
COMPASS [157]	$d$	$\pi^+$	5	0.50
	$d$	$\pi^-$	5	0.78
	$d$	$h^+$	5	0.90
	$d$	$h^-$	5	0.86
	$d$	$K^+$	5	1.50
	$d$	$K^-$	5	0.78
<b>Total</b>			104	1.01