Azimuthal anisotropy from quantum interference in ρ⁰ photoproduction in UPCs with ALICE



Andrea Giovanni Riffero¹ on behalf of the ALICE Collaboration

1. University and INFN Torino

DIS 2024, Grenoble, France, April 10th 2024



OUTLINE

PHYSICS MOTIVATION

DETECTOR AND DATA SAMPLE

DATA ANALYSIS

RESULTS

TAKE HOME AND OUTLOOK

PHYSICS MOTIVATIONS

ALICE webpage

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Purely hadronic interactions highly suppressed \rightarrow study of photon-induced reactions

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Coherent: the photon interacts with the nucleus as a whole Incoherent: the photon interacts with one nucleon

Interesting process: coherent photoproduction of a vector meson (e.g. $\rho^0)$















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The angular distribution of the pions is determined by the conservation of the total angular momentum











Each nucleus can act as the source of the photon or as the target in the interaction → two indistinguishable amplitudes contribute to the cross section

Interference between the amplitudes!

$$\sigma(p_{\rm T}, b, y = 0) = |A(p_{\rm T}, b) - A(p_{\rm T}, b) e^{i \vec{p} \cdot \vec{b}}|^2$$





cos(2φ) modulation



[4] PRL 84 (2000) 2330-2333

Why is it interesting?

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The pomeron exchange restricts the ho^0 production site within on<u>e of the nuclei</u>

[5] PRD 103 (2021) 3, 033007

Double-slit experiment at fm scale [5] $\rightarrow b$ = distance between the openings

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1) interaction dipole/target → Wilson lines
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→ account for different color charge configurations

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Depends on $b \rightarrow \text{ALICE}$

This talk! And paper in preparation!

DETECTOR AND DATA SAMPLE

Run 2 ALICE DETECTOR & DATA SAMPLE





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AD and VO used as veto \rightarrow suppression of purely hadronic interactions

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Trigger

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Topological trigger: events with at least two track segments in the ITS SPD with an opening angle θ >153°



DATA ANALYSIS





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Re-weighting procedure

Fit the MC generated p_T^2 distribution using the square of the nuclear form factor (1) to extract a_{Pb} and R_{Pb}

 $\frac{\mathrm{d}N}{\mathrm{d}p_{\mathrm{T}}^2} = c \mid F(t, a_{\mathrm{Pb}}, R_{\mathrm{Pb}}) \mid^2 (1)$

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Compute the weights using (2), where R_X is chosen to minimize discrepancies between data and reconstructed MC p_T distributions

$$w(p_{\rm T}) = \frac{|F(|t|, a_{\rm Pb}, R_{\rm X})|^2}{|F(|t|, a_{\rm Pb}, R_{\rm Pb})|^2} \quad (2)$$

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Build the MC mass distributions by weighting each event with $w(p_T)$ evaluated at the generated p_T

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 $w_{Y \rightarrow Z}$ = contribution of the physical class Y to the yield in the experimental class Z. Computed from measured cross sections ratios and migration probabilities [12].

[12] JHEP 06 (2020) 035

 a_2 = true amplitudes of the modulation



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> The effect of migrations is important especially in XnOn and XnXn

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RESULTS

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12/15

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The modulation strength strongly increases as *b* decreases



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Main systematic uncertainties: signal extraction, acceptance x efficiency, and ϕ angle definition
ASYMMETRY RESULTS

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TAKE HOME AND OUTLOOK -

Take home

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The same effect can be studied with other particles (e.g. J/ψ) where the model predictions are expected to be more precise

REFERENCES

[1] <u>ALICE webpage</u>

- [2] A. Baltz et al. The Physics of Ultraperipheral Collisions at the LHC, Phys.Rept. 458 (2008) 1171
- [3] <u>Talk by Daniel Brandenburg</u> *Linearly polarized photon-gluon collisions*
- [4] S. Klein, J. Nystrand, Interference in exclusive vector meson production in heavy ion collisions, Phys.Rev.Lett. 84 (2000) 2330-2333
- [5] W. Zha et al. *Exploring the double-slit interference with linearly polarized photons*, <u>Phys.Rev.D 103 (2021) 3, 033007</u>
- [6] H. Xing et al. *The cos 2φ azimuthal asymmetry in p0 meson production in ultraperipheral heavy ion collisions,* JHEP 10 (2020) 064
- [7] W. Zhao et al. *Effects of nuclear structure and quantum interference on diffractive vector meson production in ultra-peripheral nuclear collisions*, <u>arXiv:2310.15300 [nucl-th] (2023)</u>
- [8] STAR Collaboration, *Tomography of ultrarelativistic nuclei with polarized photon-gluon collisions*, Sci.Adv. 9 (2023) eabq3903
- [9] <u>Talk by Ashik Ikbal</u> (STAR), Exclusive J/ψ Photoproduction and Entanglement-Enabled Spin Interference in Ultra-Peripheral Collisions at STAR
- [10] M. Broz et al. A generator of forward neutrons for ultra-peripheral collisions: $n_0^0 n$, <u>Comput. Phys. Commun. (2020) 107181</u>
- [11] S. Klein, et al. STARlight: A Monte Carlo simulation program for ultra-peripheral collisions of relativistic ions, Comput. Phys. Commun. 212 (2017) 258–268
- [12] ALICE Collaboration, Coherent photoproduction of ρ0 vector mesons in ultra-peripheral Pb-Pb collisions at VsNN = 5.02 TeV, JHEP 06 (2020) 035

OTHER ALICE TALKS

Tue	09/04	09:50	A. Bylinkin	Exclusive four pion photoproduction in ultra-peripheral Pb–Pb collisions at VsNN = 5.02 TeV at ALICE
Tue	09/04	11:20	J. Park	Reacent heavy flavour measurements from ALICE
Tue	09/04	14:00	R. Guernane	Overview of ALICE Upgrades
Tue	09/04	14:20	L. Huhta	ALICE Forward Calorimeter upgrade (FoCal): Physics program and performance
Tue	09/04	16:00	A. Shatat	Coherent vector meson photoproduction and polarization in heavy- ion collisions with nuclear overlap in ALICE
Tue	09/04	16:20	G. Contreras	Energy dependence of coherent J/psi production off lead with ALICE
Wed	10/04	11:40	D. Grund	First study of the initial gluonic fluctuations using UPCs with ALICE
Wed	10/04	12:20	M. Kim	K+K- photoproduction in ultra-peripheral Pb-Pb collisions with ALICE

Check them out!

THANK YOU FOR YOUR ATTENTION!