

A unified description of DGLAP, CSS, and BFKL: TMD factorization bridging large and small x

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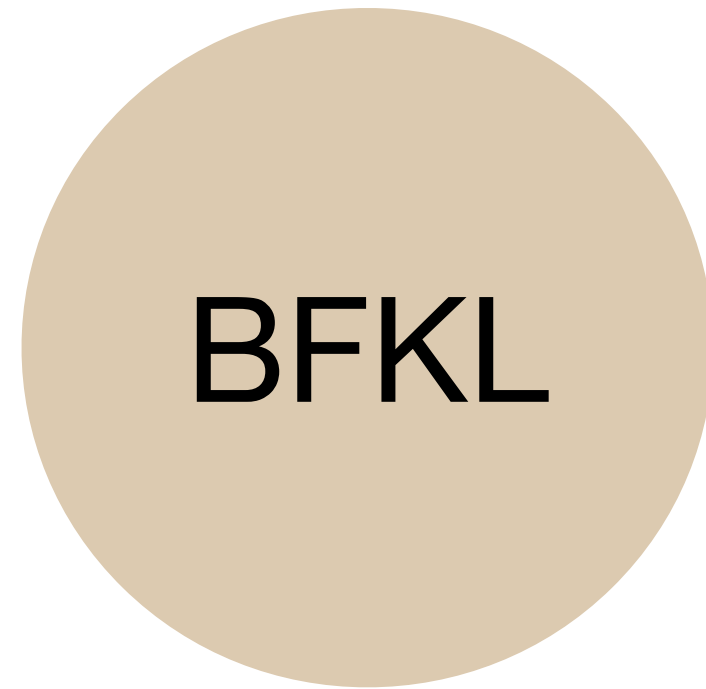
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initial condition for small-x evolution equation (BFKL) from lattice QCD / global fits ?

- ◆ small-x evolution — —> TMD physics
- ◆ evolutions of TMDPDF from lattice QCD / global fits: DGLAP+CSS, no BFKL
- ◆ different factorization, different IR structures, different evolutions, different nonperturbative TMDPDF — not universal
- ◆ need a universal TMDPDF / factorization that contains IR structures of both DGLAP and BFKL in the appropriate limits

high-energy limit



dipole($x_B = 0, b_\perp$)

all collinear twist

x_B -dep: resum all sub-eikonal

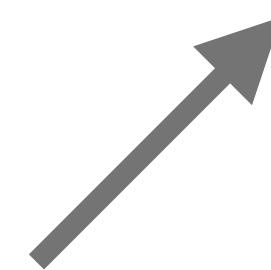
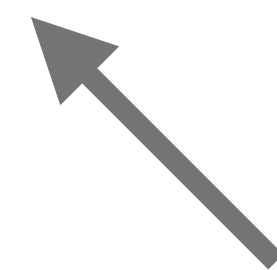
collinear limit



$\text{CSS}(b_\perp) \otimes \text{PDF}(x_B, b_\perp = 0)$

leading-twist

b_\perp -dep: resum all sub-lead twists



not known how

MSTT(-erious) factorization

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Unified description of DGLAP, CSS, and BFKL evolution: TMD factorization bridging large and small x

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a TMD factorization unifying IR structures of large and small x

summary

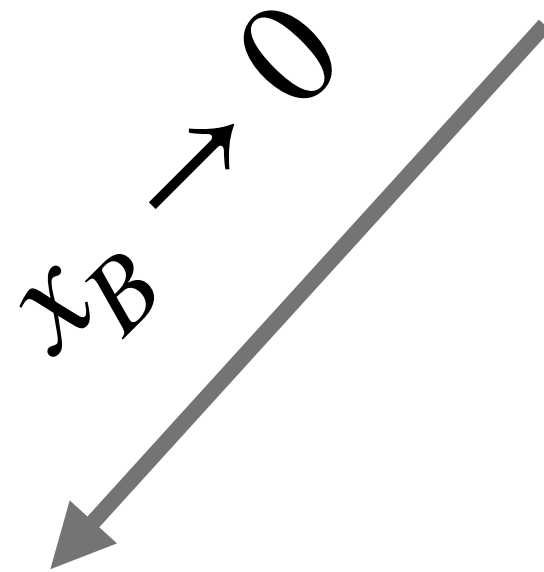
gluon TMDPDF(x_B, b_\perp) operator

NLO (2-gluon) corrections to
background fields

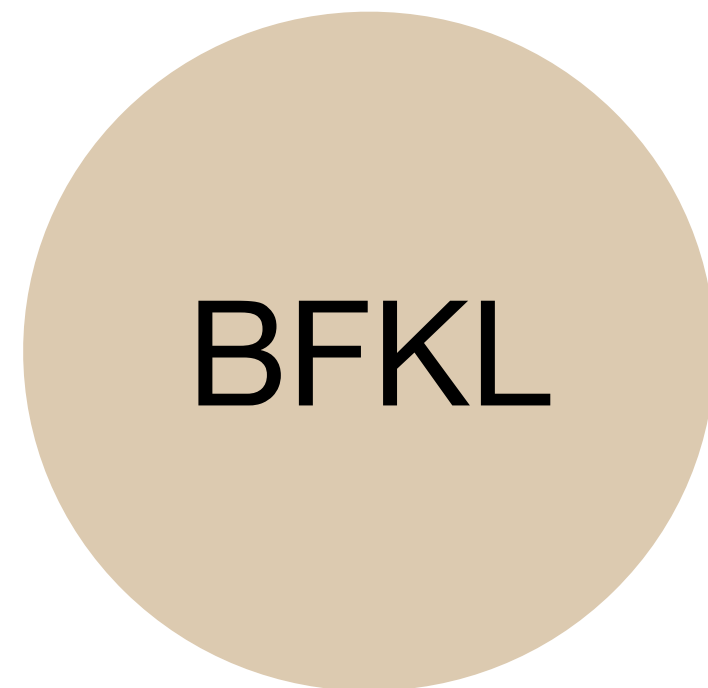
MSTT factorization



new general structure: IR & UV div. in trans. mom. + IR & UV in rapidity

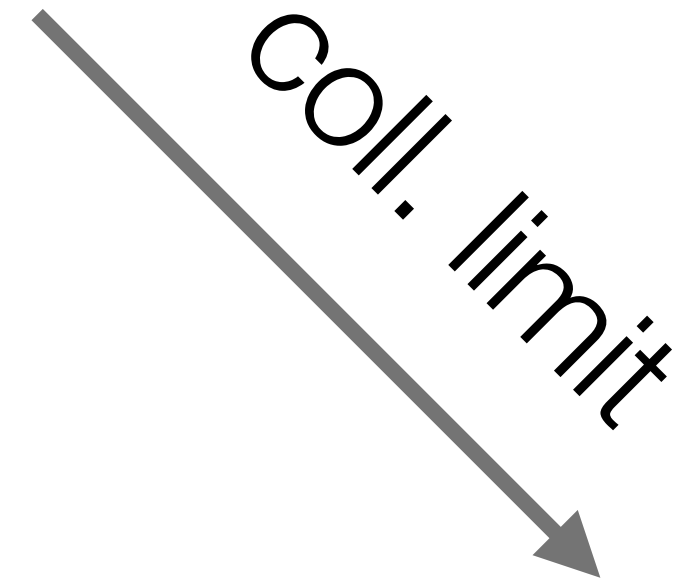


$x_B \rightarrow 0$



BFKL

$$WW(x_B = 0, b_\perp)$$



coll. limit



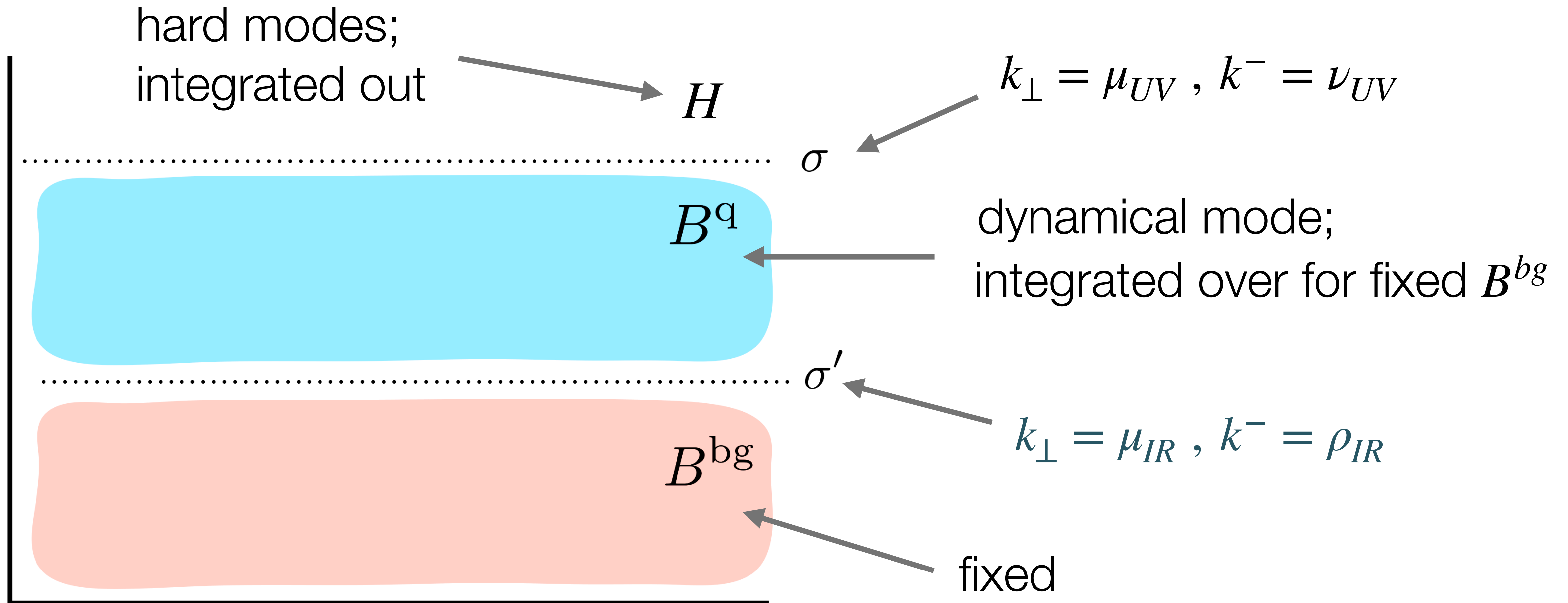
DGLAP,
CSS

$$\text{CSS}(b_\perp) \otimes \text{PDF}(x_B, b_\perp = 0)$$

few technicalities — background-field method

$$A \rightarrow H + B$$

$$B = B^q + B^{bg}$$



few technicalities — regularization schemes

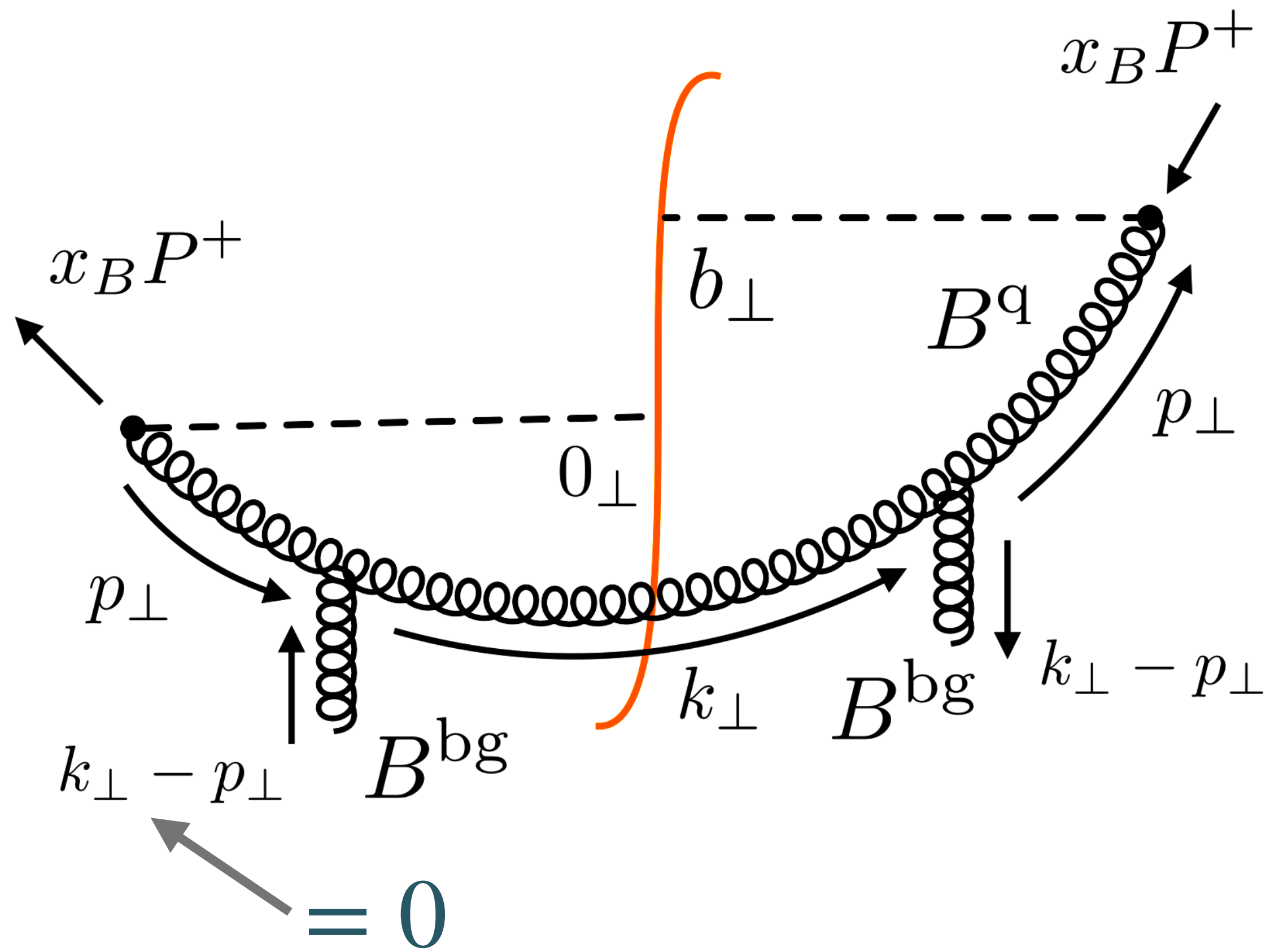
◆ divergences in $k_{\perp} \rightarrow \infty$ (μ_{UV}) & $k_{\perp} \rightarrow 0$ (μ_{IR}): dim-reg

◆ divergences in $k^{-} \rightarrow \infty$ (ν_{UV}) & $k^{-} \rightarrow 0$ (ρ_{IR}): η -scheme

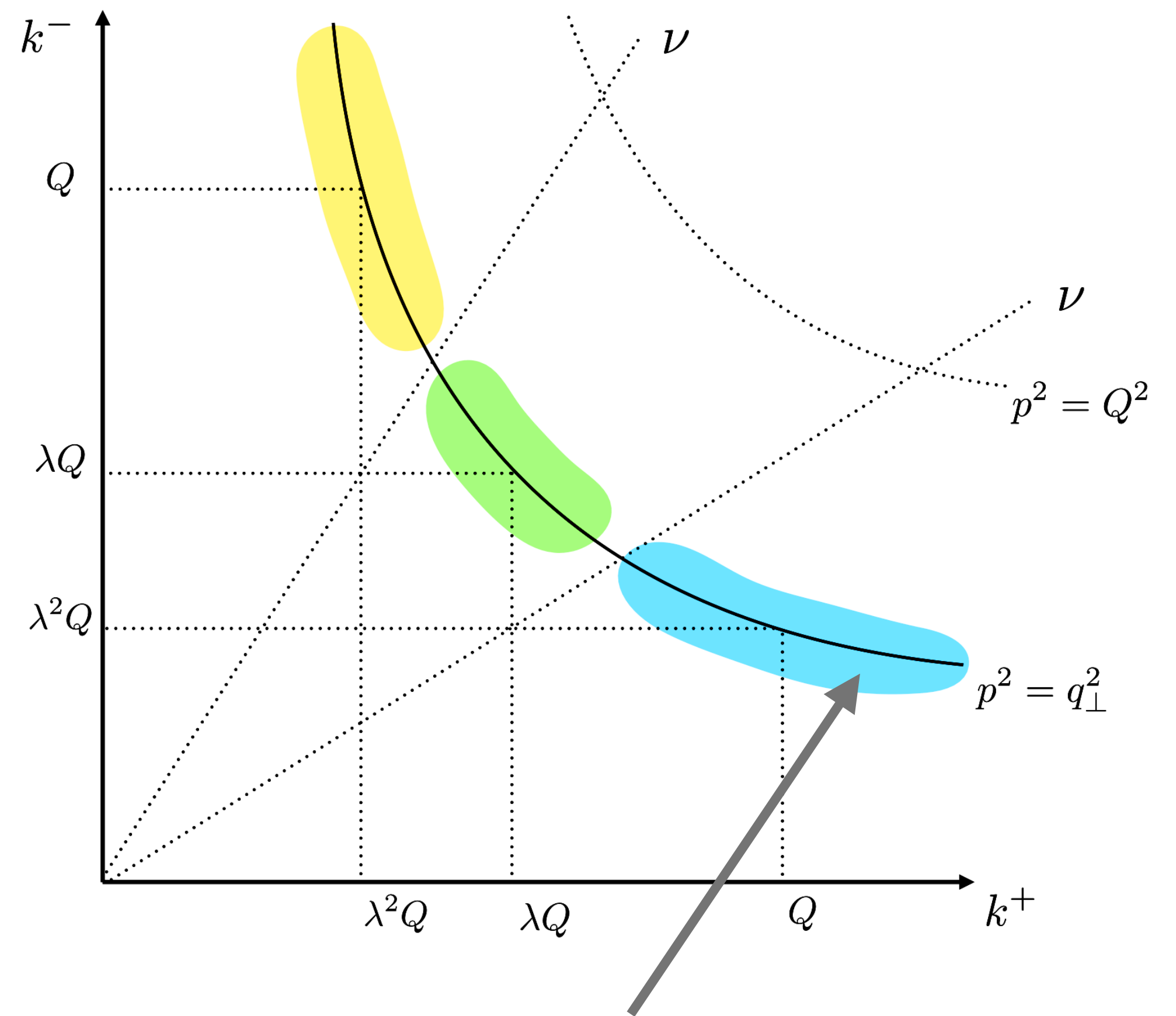
$$\int_0^{\infty} \frac{dk^{-}}{k^{-}} \rightarrow \nu^{\eta} \int_0^{\infty} \frac{dk^{-}}{k^{-}} |k^{+}|^{-\eta}$$

rapidity divergence in the calculation: $z \equiv \frac{x_B}{x_B + \frac{k_{\perp}^2}{2P^{+}k^{-}}}$ $k^{-} \rightarrow \infty : z \rightarrow 1$
 $k^{-} \rightarrow 0 : z \rightarrow 0$

CSS/SCET: missing ingredients



no trans. mom. supplied
by the background field

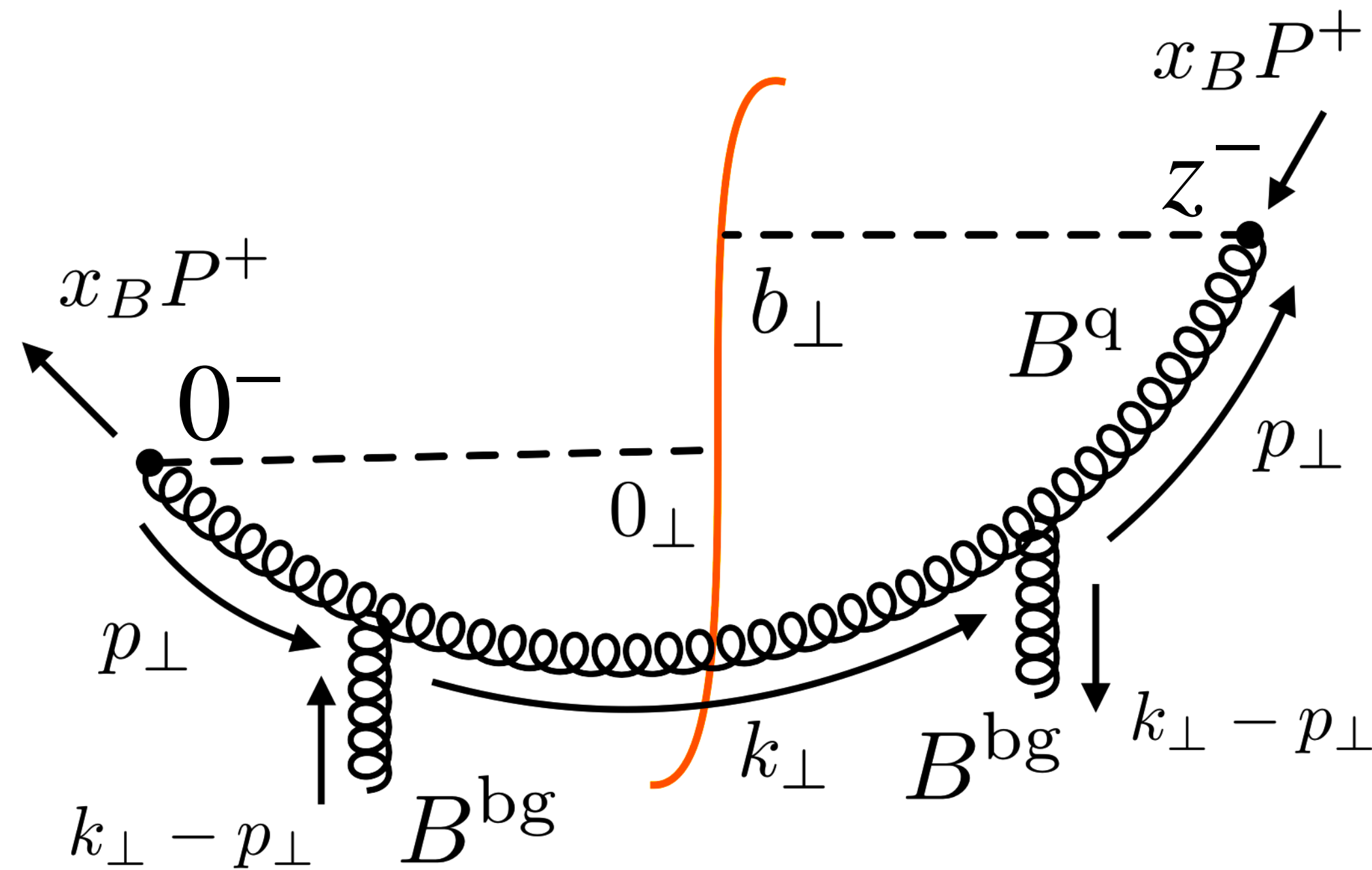


collinear modes are
approx. on mass-shell
no virtual corrections

MSTT

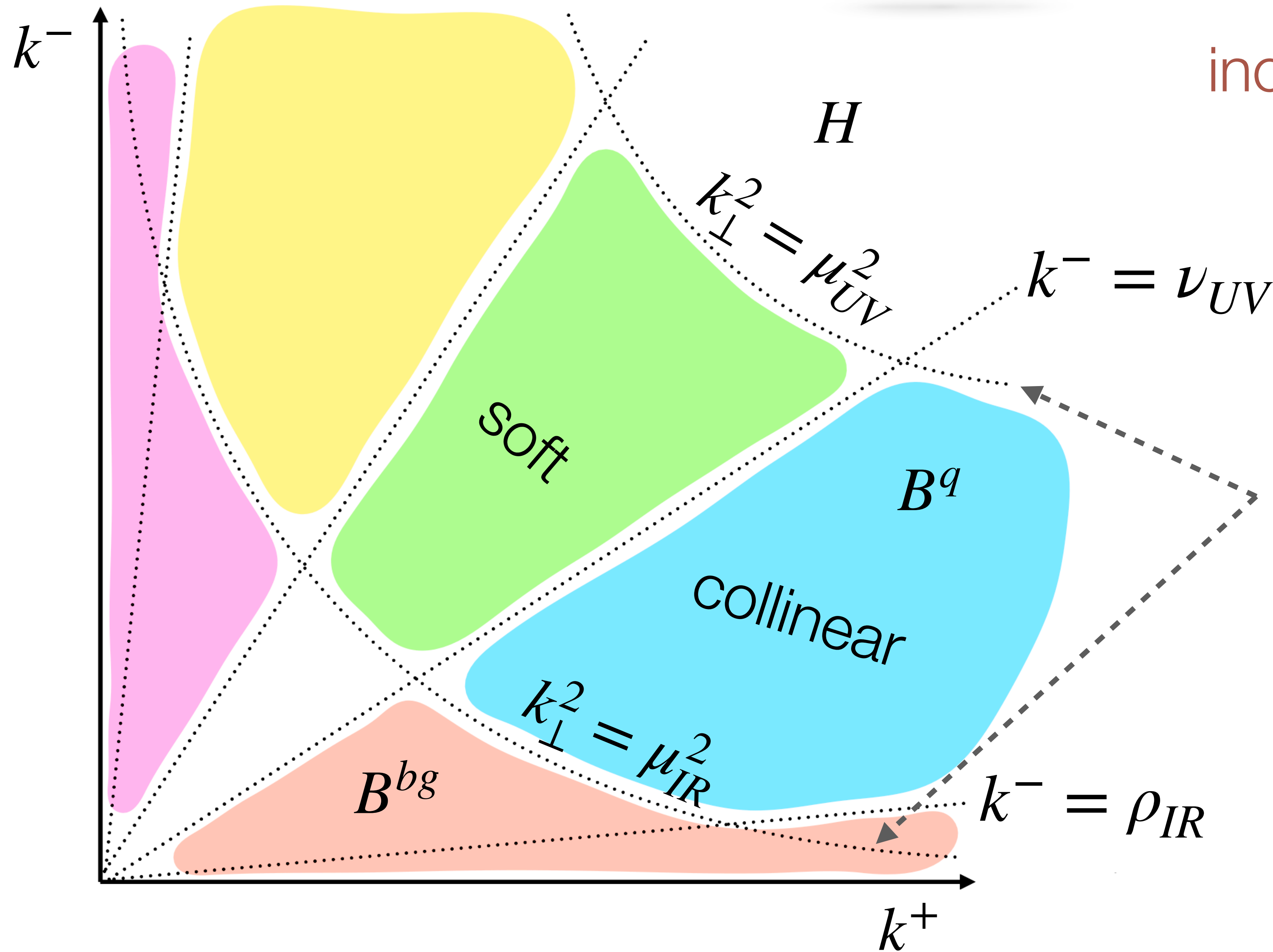
NLO (2-gluon) correction to
gluon TMDPDF operator

$$\mathcal{B}_{ij}(x_B, b_\perp) = \int_{-\infty}^{\infty} dz^- e^{-ix_B P^+ z^-} \langle P, S | \bar{T} \{ F_{-i}^m(z^-, b_\perp) \times [z^-, \infty]_b^{ma} \} T \{ [\infty, 0^-]_0^{an} F_{-j}^n(0^-, 0_\perp) \} | P, S \rangle$$



$$k_\perp - p_\perp > 0$$

MSTT



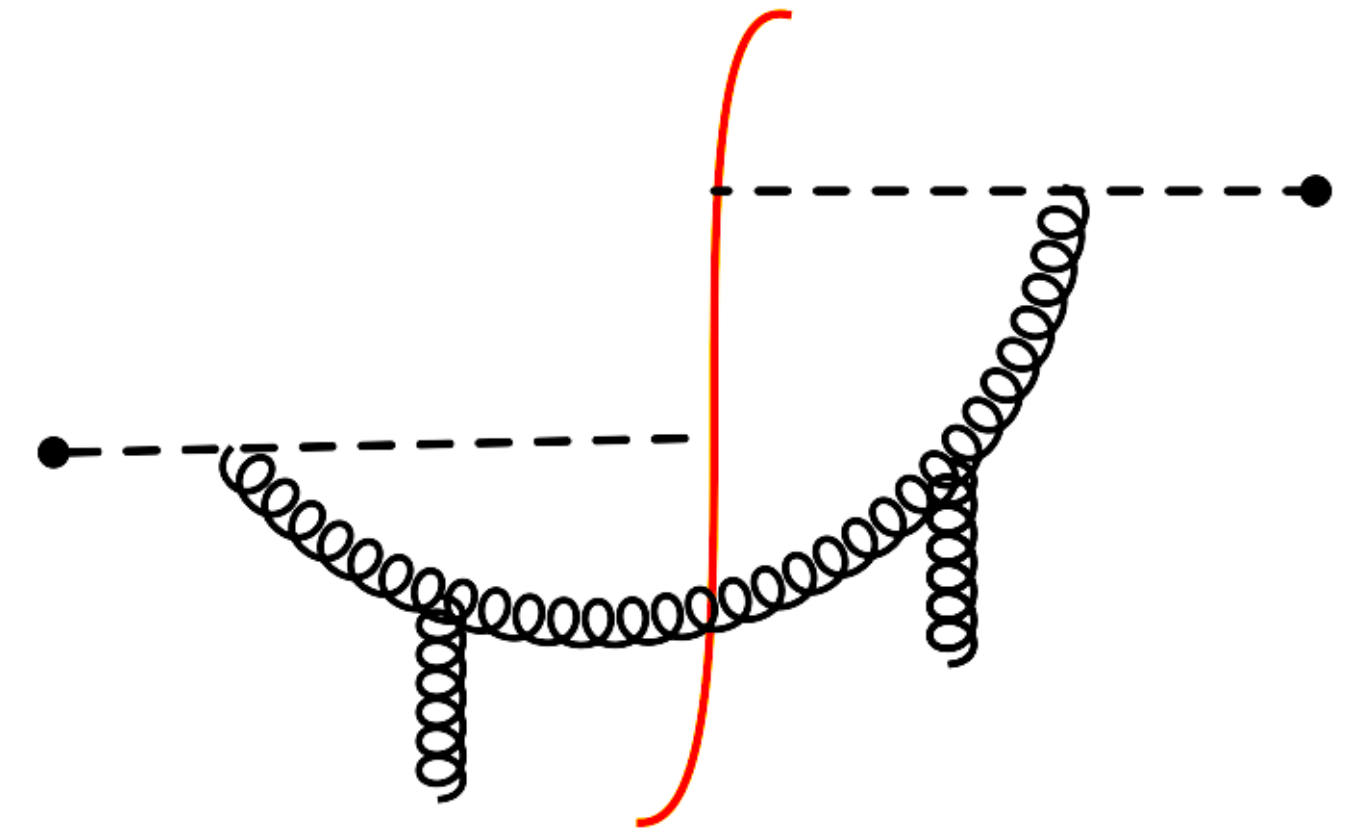
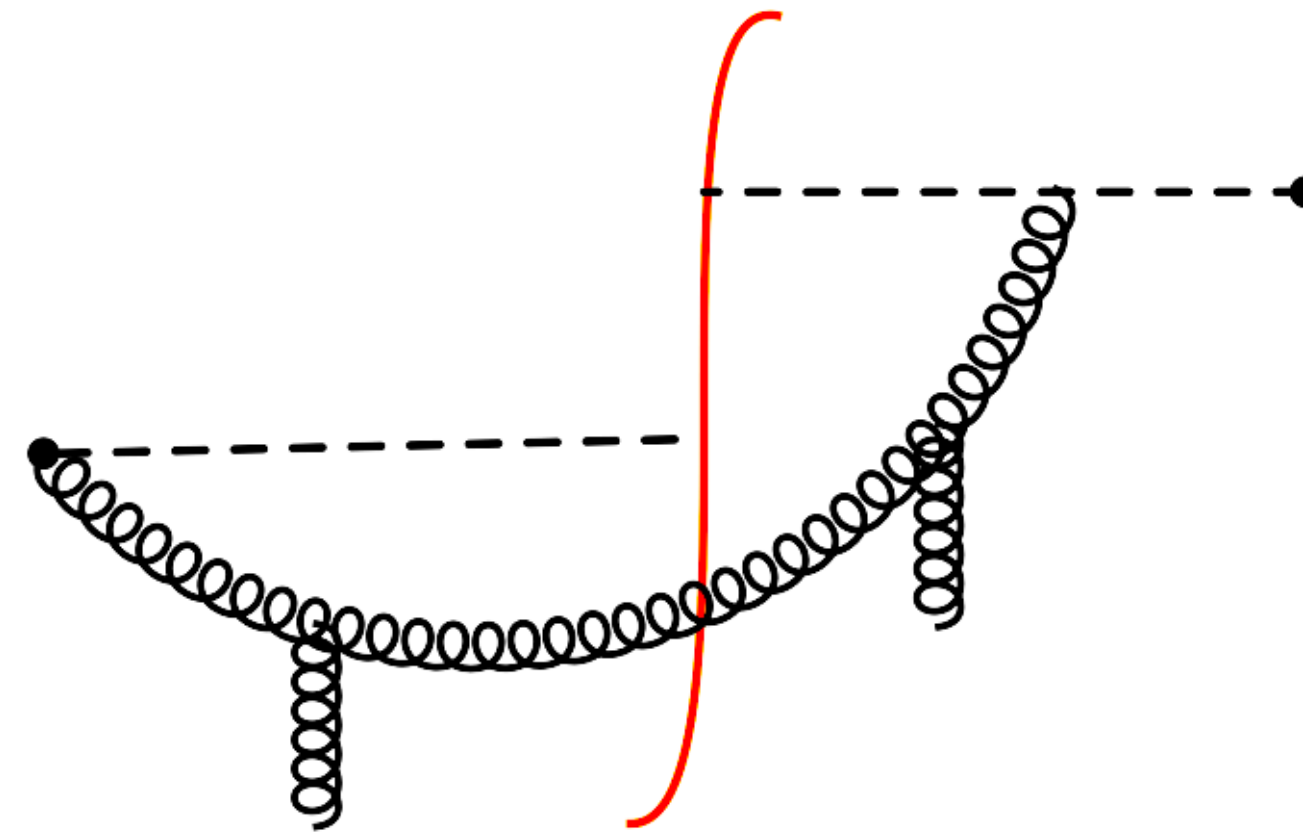
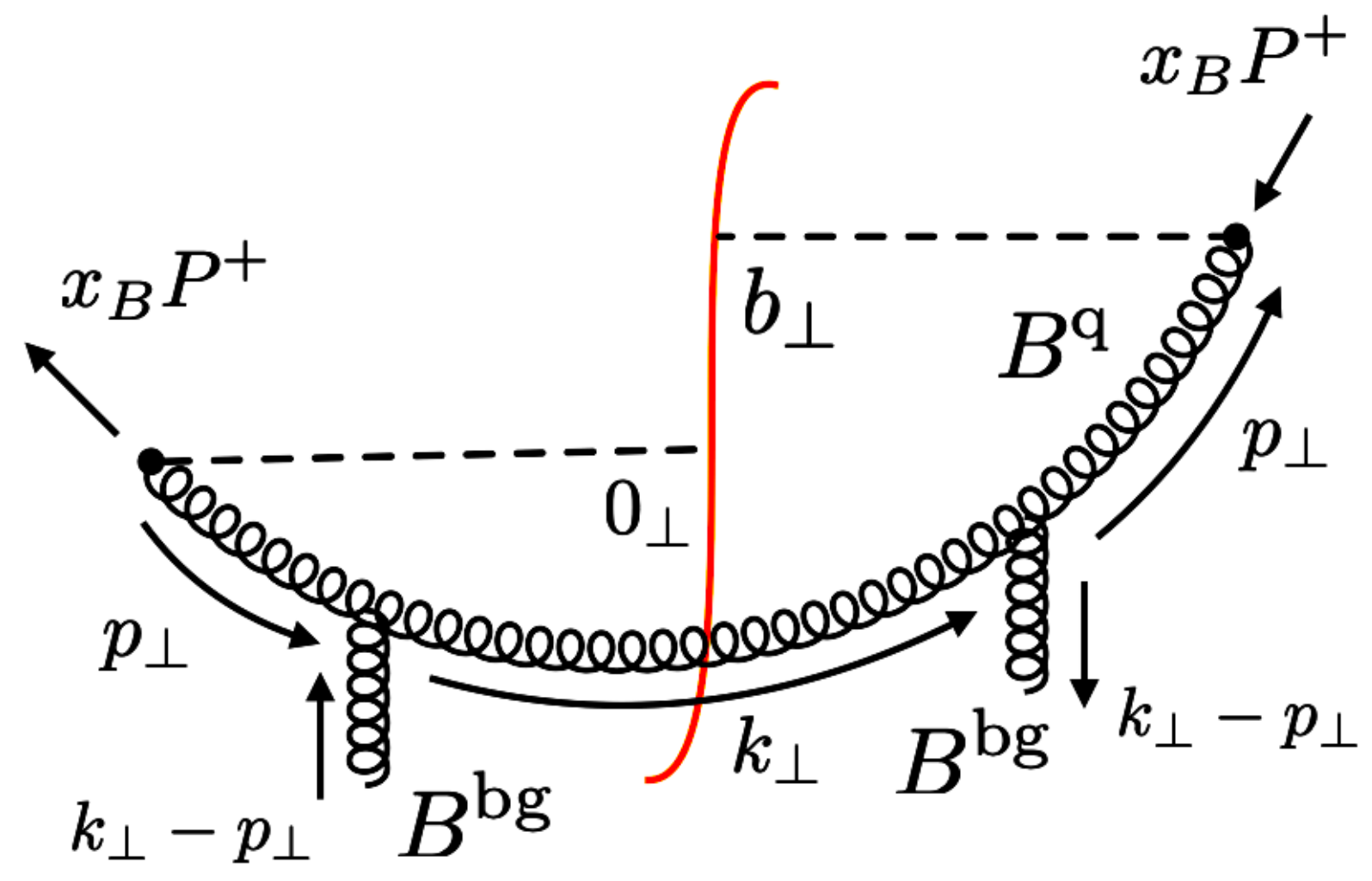
include virtual corrections

$$k_{\perp}^2 > 2k^+k^-$$

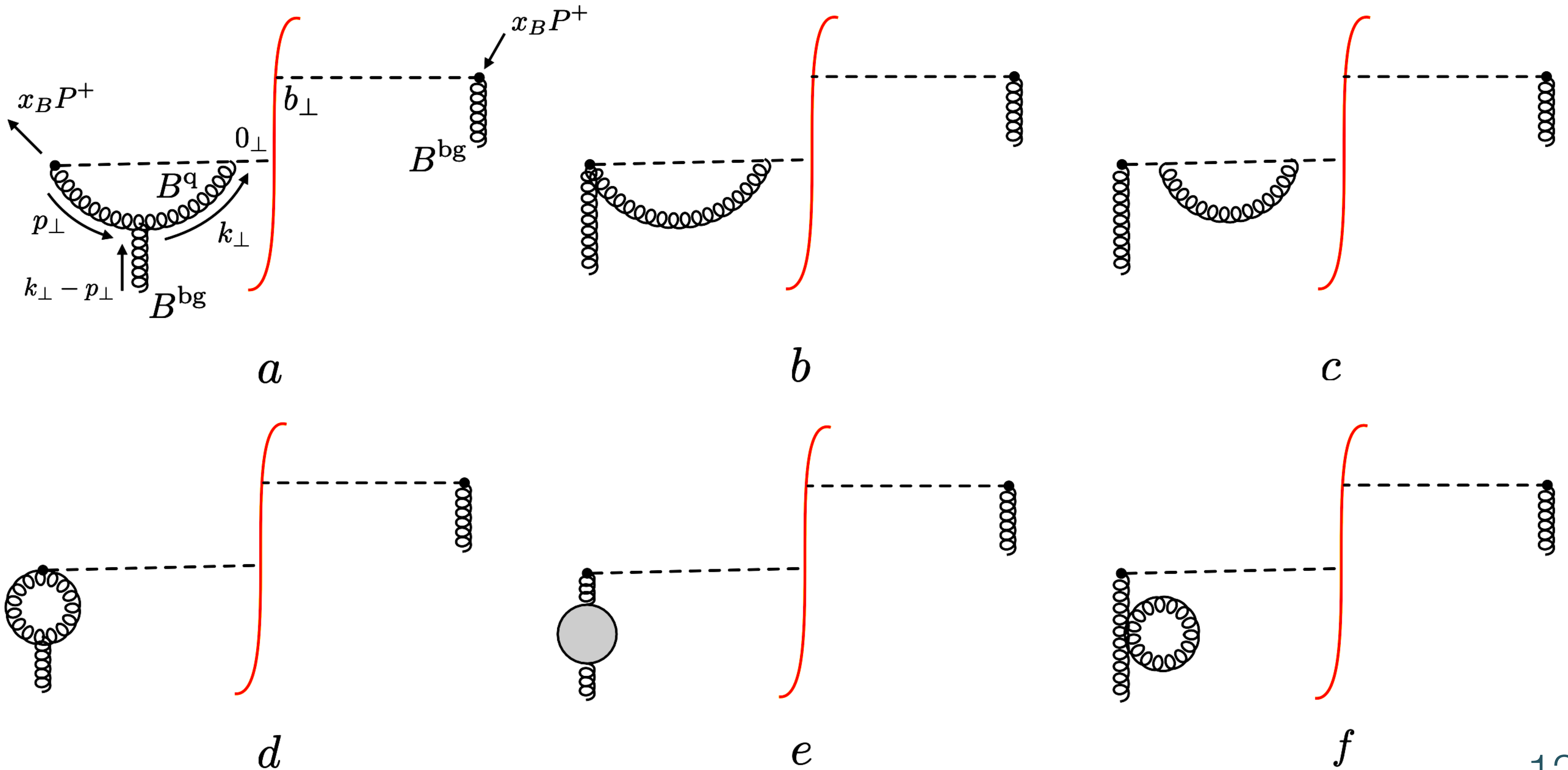
on mass-shell region

$$2k^+k^- = k_{\perp}^2 = \mu^2$$

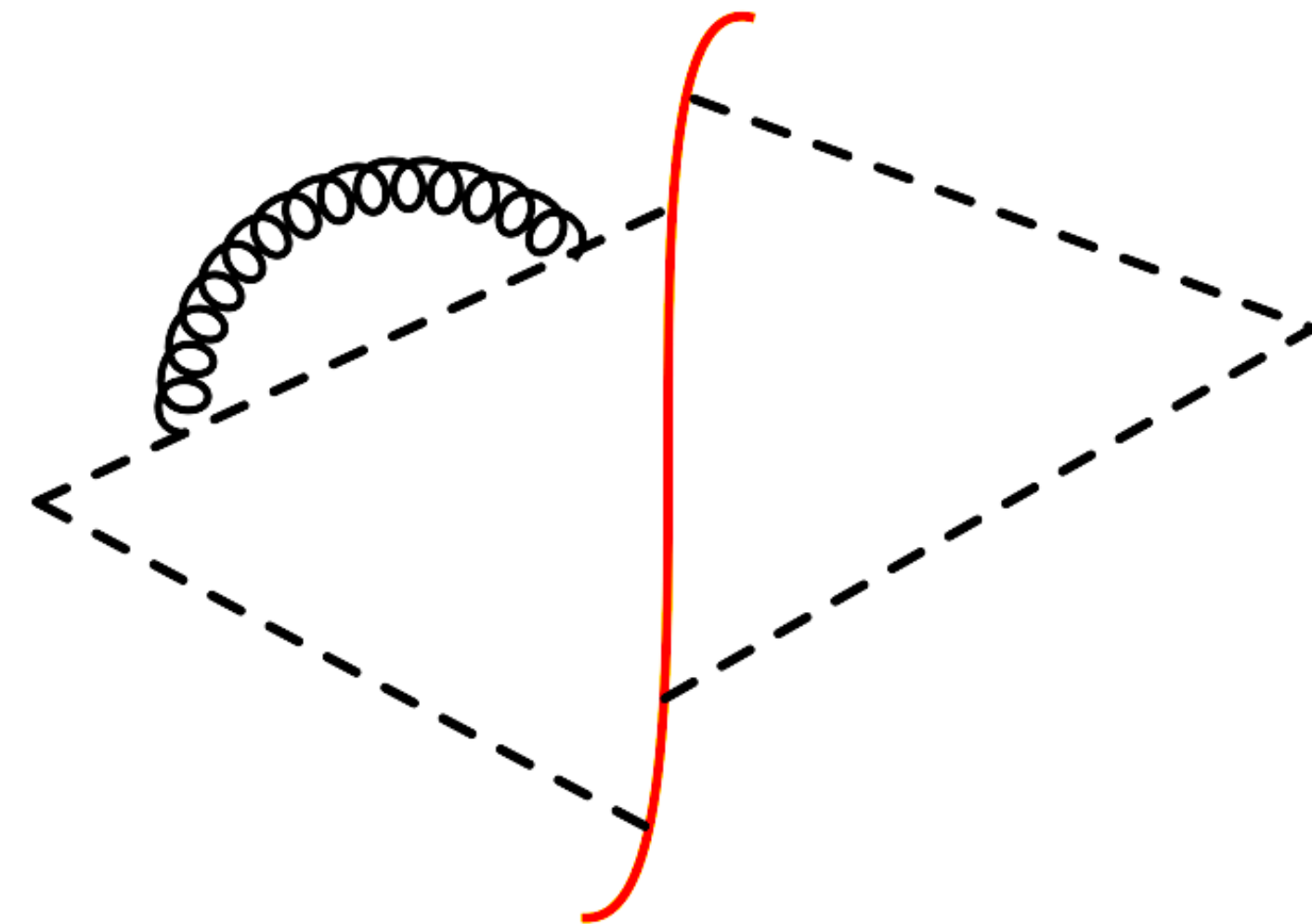
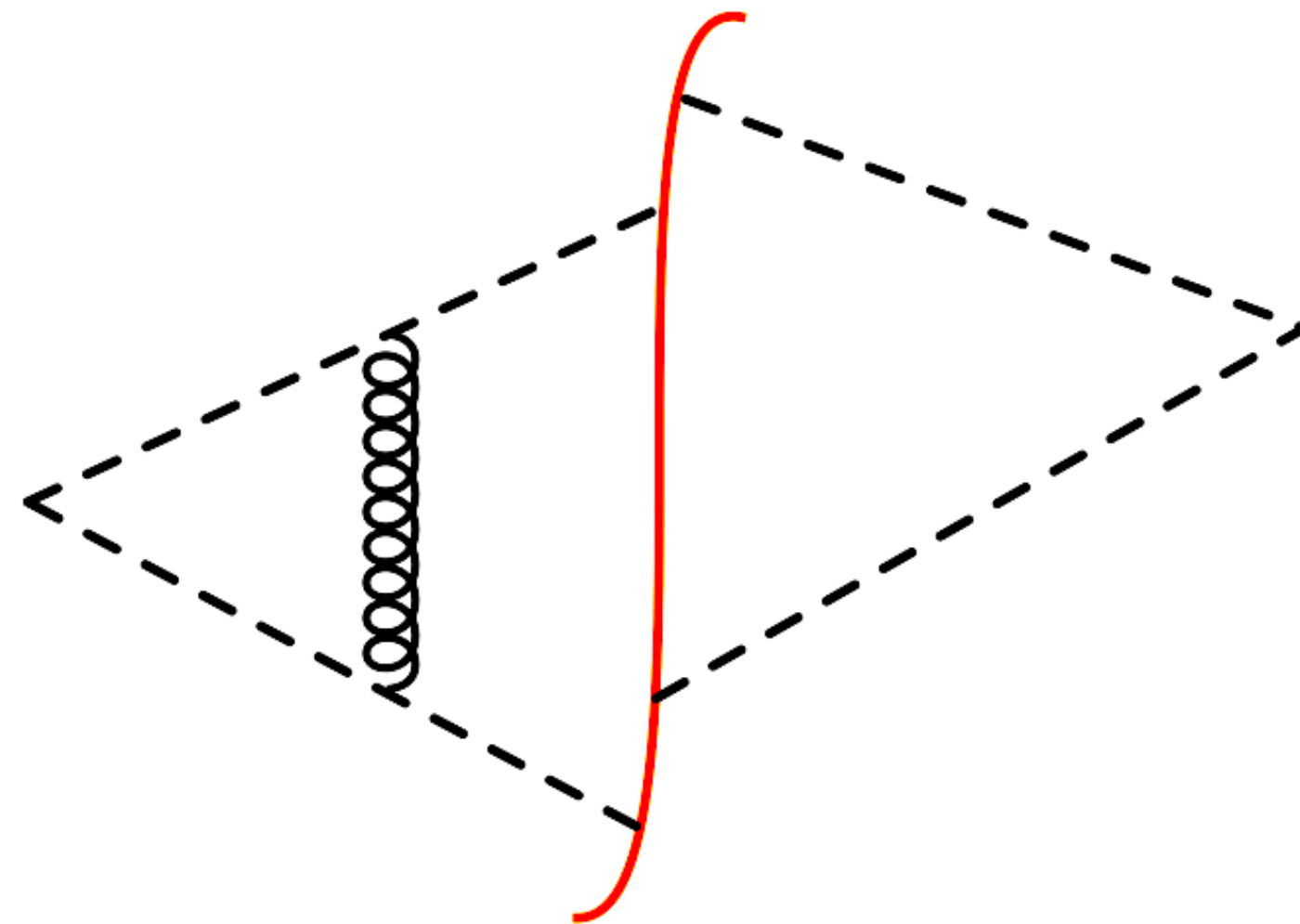
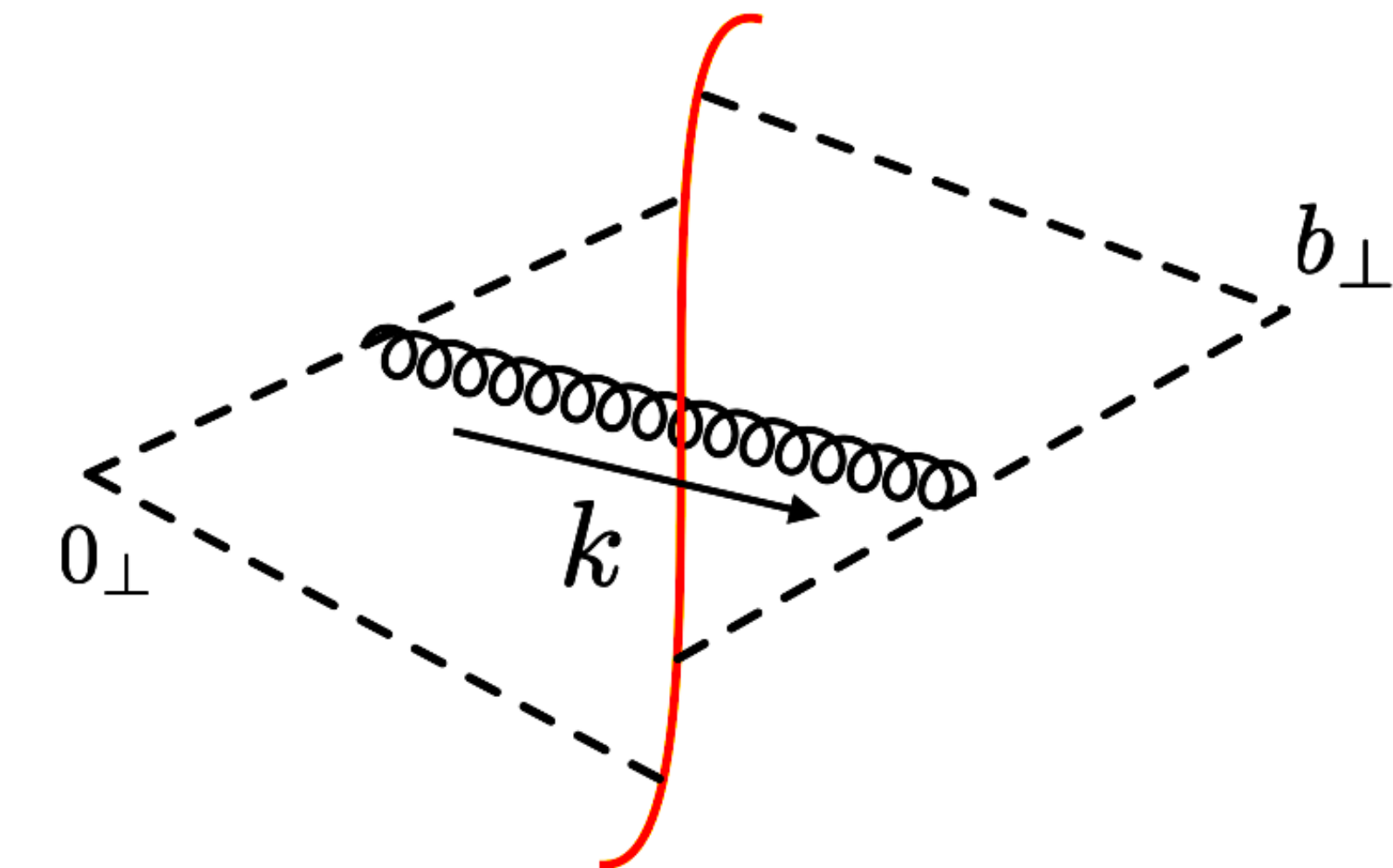
real emissions



virtual emissions



soft factor



MSTT factorization of gluon TMDPDF at NLO

$$\begin{aligned}
 f_{ij}(x_B, b_\perp, \mu_{\text{UV}}^2, \zeta) = & f_{ij}(x_B, b_\perp, \mu_{\text{IR}}^2, \rho) - 4\alpha_s N_c \int \vec{d}^2 p_\perp e^{ip_\perp b_\perp} \int_0^1 \frac{dz}{z(1-z)} \int \vec{d}^2 k_\perp \left[\mathcal{R}_{ij;lm}^a(z, p_\perp, k_\perp) \right. \\
 & \left. + \mathcal{R}_{ij;lm}^b(z, p_\perp, k_\perp) \right] \int d^2 z_\perp e^{-i(p_\perp - k_\perp)z_\perp} f_{lm}\left(\frac{x_B}{z}, z_\perp, \mu_{\text{IR}}^2, \rho\right) + \frac{\alpha_s N_c}{2\pi} \left(-\frac{1}{2} (L_b^{\mu_{\text{UV}}})^2 + L_b^{\mu_{\text{UV}}} \ln \frac{\mu_{\text{UV}}^2}{\zeta^2} - \frac{\pi^2}{12} \right) \\
 & \times f_{ij}(x_B, b_\perp, \mu_{\text{IR}}^2, \rho) - \frac{\alpha_s N_c}{\pi} L_b^{\mu_{\text{IR}}} \int_0^1 dz \left[\frac{1}{(1-z)_+} + \frac{1}{z} \right] f_{ij}\left(\frac{x_B}{z}, b_\perp, \mu_{\text{IR}}^2, \rho\right) - \frac{\alpha_s N_c}{2\pi} \int d^2 z_\perp \int \vec{d}^2 p_\perp e^{ip_\perp (b-z)_\perp} \\
 & \times \left(\frac{1}{2} \ln^2 \frac{\mu_{\text{IR}}^2}{p_\perp^2} + \ln \frac{\mu_{\text{IR}}^2}{p_\perp^2} \ln \frac{\rho}{\zeta} - \frac{\pi^2}{12} \right) \frac{g_{il} p_j p_m + p_i p_l g_{mj}}{p_\perp^2} f_{lm}(x_B, z_\perp, \mu_{\text{IR}}^2, \rho) \quad \text{part of DGLAP} \\
 & + \frac{\alpha_s N_c}{2\pi} \int d^2 z_\perp \int \vec{d}^2 p_\perp e^{ip_\perp (b-z)_\perp} \left(\frac{\beta_0}{2N_c} \ln \frac{\mu_{\text{UV}}^2}{p_\perp^2} + \frac{67}{18} - \frac{5N_f}{9N_c} \right) f_{ij}(x_B, z_\perp, \mu_{\text{IR}}^2, \rho) + O(\alpha_s^2). \quad \text{CSS}
 \end{aligned}$$

part of BFKL

$$L_b^\mu \equiv \ln \left(\frac{b_\perp^2 \mu^2}{4e^{-2\gamma_E}} \right)$$

$$\zeta = x_B P^+$$

gluon TMDPDF parametrization:

$$f_{ij}(x_B, b_\perp) = x_B P^+ \left[-\frac{g_{ij}}{2} f_1(x_B, b_\perp) + \left(\frac{g_{ij}}{2} + \frac{b_i b_j}{b_\perp^2} \right) h_1(x_B, b_\perp) \right]$$

collinear limit

$$f_{ij}(x_b, b_\perp \rightarrow 0) = f_1(x_b)$$

small-x limit

$$f_{ij}(x_b \rightarrow 0, b_\perp) = \frac{b_i b_j}{b_\perp^2} \mathcal{H}(b_\perp)$$

collinear limit of MSTT

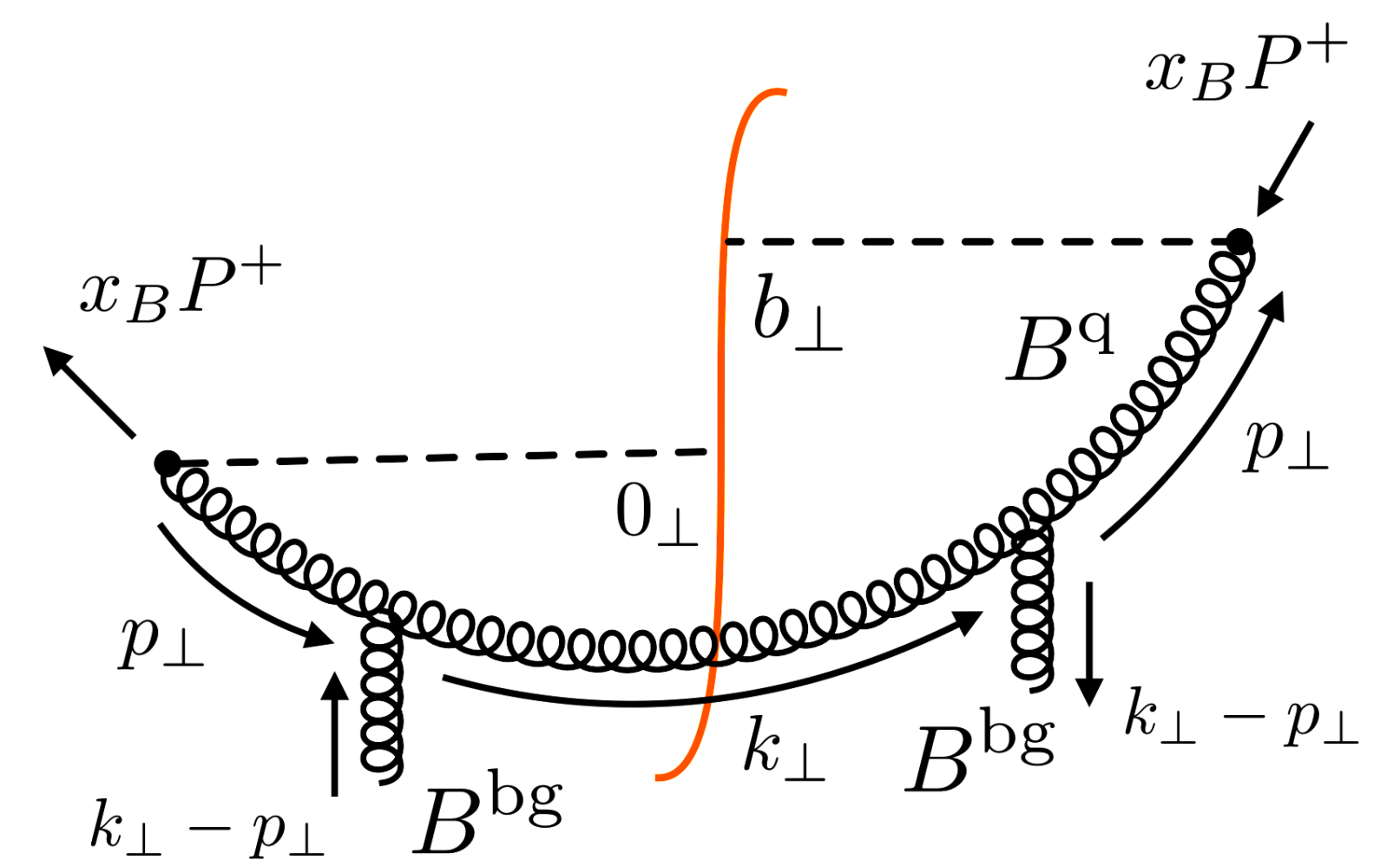
$$p_{\perp} \sim b_{\perp}^{-1} \gg k_{\perp} - p_{\perp} \sim \mu_{IR}$$

virtual corrections vanishes,
no rapidity IR div.

$$f_1(x_B, b_{\perp}, \mu_{UV}^2, \zeta) = f_1(x_B, 0_{\perp}, \mu_{IR}^2)$$

$$-\frac{\alpha_s N_c}{\pi} L_b^{\mu_{IR}} \int_0^1 \frac{dz}{z} P_{gg}(z) f_1\left(\frac{x_B}{z}, 0_{\perp}, \mu_{IR}^2\right) + \frac{\alpha_s N_c}{2\pi} \left(-\frac{1}{2} (L_b^{\mu_{UV}})^2 + L_b^{\mu_{UV}} \ln \frac{\mu_{UV}^2}{\zeta^2} - \frac{\pi^2}{12} \right) f_1(x_B, 0_{\perp}, \mu_{IR}^2)$$

DGLAP



$$k_{\perp} - p_{\perp} = 0$$

constant piece \rightarrow IR div. k_{\perp}

$$P_{gg}(z) = \frac{1}{(1-z)_+} + \frac{1}{z} - 2 + z - z^2$$

CSS

$x_B \rightarrow 0$ limit of MSTT

$1/z$ of $P_{gg}(z)$ generates additional rapidity IR div.

$$K_{\text{BFKL}}(p_{\perp}, k_{\perp}) = -\frac{\alpha_s N_c}{\pi} \int d^2 b_{\perp} e^{ik_{\perp} b_{\perp}} L_b^{\mu_{\text{IR}}} + \frac{\alpha_s N_c}{\pi} (2\pi)^2 \delta^2(k_{\perp}) \ln \frac{\mu_{\text{IR}}^2}{p_{\perp}^2}$$

BFKL

$$f_1(x_B, p_{\perp}, \mu_{\text{UV}}^2, \zeta) \simeq \mathcal{H}_1(p_{\perp}, \rho) + \ln \frac{\rho}{\zeta} \int d^2 k_{\perp} K_{\text{BFKL}}(p_{\perp}, k_{\perp}) \mathcal{H}_1(p_{\perp} - k_{\perp}, \rho)$$

CSS

$$+ \frac{\alpha_s N_c}{2\pi} \int d^2 b_{\perp} \left(-\frac{1}{2} (L_b^{\mu_{\text{UV}}})^2 + L_b^{\mu_{\text{UV}}} \ln \frac{\mu_{\text{UV}}^2}{\zeta^2} - \frac{\pi^2}{12} \right) \int d^2 k_{\perp} e^{ik_{\perp} b_{\perp}} \mathcal{H}_1(p_{\perp} - k_{\perp}, \rho)$$
$$+ \frac{\alpha_s N_c}{2\pi} \left(\frac{\beta_0}{2N_c} \ln \frac{\mu_{\text{UV}}^2}{p_{\perp}^2} + \frac{67}{18} - \frac{5N_f}{9N_c} \right) \mathcal{H}_1(p_{\perp}, \rho).$$

$x_B \rightarrow 0$ limit: MSTT vs Glauber SCET

$$k_{\perp}^2 \gg 2k^+k^{-1}$$

propagator $\sim 1/k_{\perp}^2$

Glauber SCET:	$k_{\perp}^2 \gg k^+ \sim k^{-1}$	$k^{\mu} \sim Q(\lambda^2, \lambda^2, \lambda)$	mid rapidity
MSTT:	$k^+ \gg k_{\perp}^2 \gg k^{-1}$	$k^{\mu} \sim Q(1, \lambda^4, \lambda)$	forward rapidity

$(+, -, \perp), \lambda \ll 1$

summary

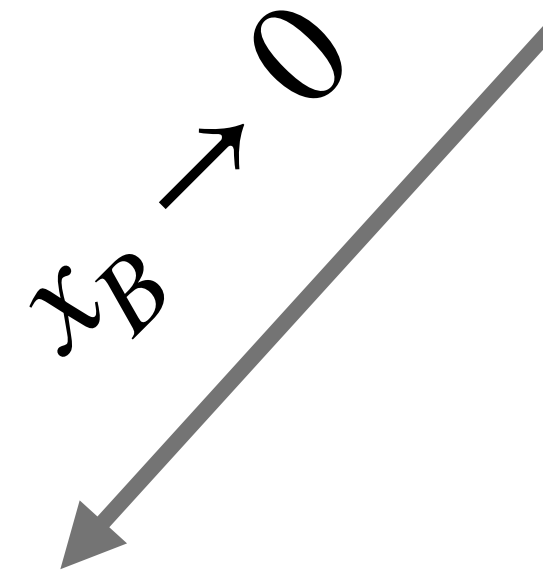
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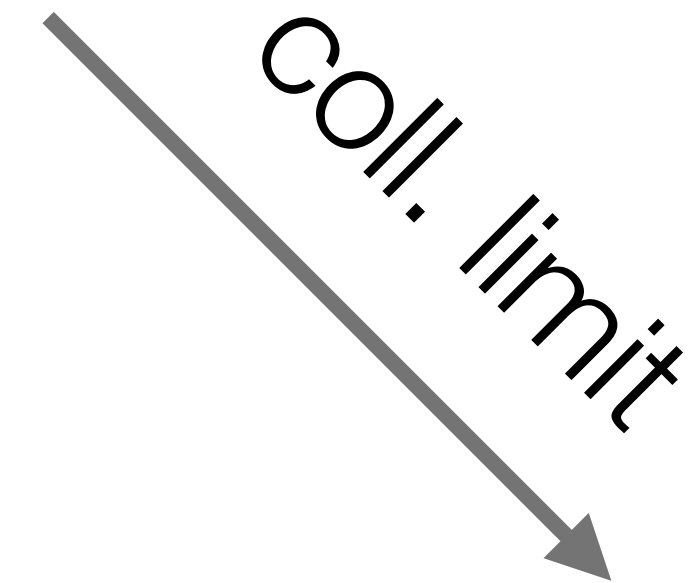


new general evolution



BFKL

$$WW(x_B = 0, b_\perp)$$



DGLAP,
CSS

$$CSS(b_\perp) \otimes PDF(x_B, b_\perp = 0)$$