

On the Interplay of Nuclear and Higher-Twist Corrections at Large x

S. Alekhin

University of Hamburg, Germany

S. Kulagin

INR Moscow

R. Petti

University of South Carolina, Columbia SC, USA

DIS 2024

April 10th, 2024, Grenoble, France

- ◆ *Precision studies of high-energy processes with nuclei require an understanding of nuclear effects at the parton level, which were observed to survive at $Q \gg 1 \text{ GeV}/c$.*
- ◆ *The study of nuclear corrections in ^2H , ^3H , ^3He provides insights into the mechanisms responsible for modifications of PDFs in the nuclear environment:*
 - *Dynamics of $A=2$ and $A=3$ nuclei better understood than the dynamics of many-particle nuclei;*
 - *Effects of the momentum distribution, nuclear binding and off-shell modification of bound nucleons driven by the wave/spectral function, which is directly related to the underlying N-N interaction.*

⇒ *Verify consistency with universal off-shell function from $A \geq 4$ nuclear targets*
- ◆ *Global QCD analysis with ^2H , ^3H , ^3He :*
 - *Determination of the off-shell function along with proton PDFs;*
 - *Interplay between nuclear corrections and higher-twist (HT) contributions in DIS at large x ;*
 - *Large n - p asymmetry of ^3H and ^3He allows to probe isospin effects in nuclear corrections.*

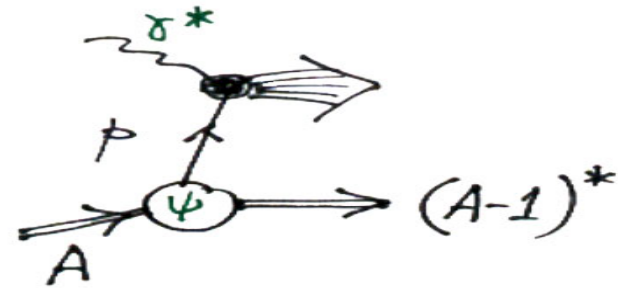
⇒ *Study flavor dependence of modifications of PDFs in the nuclear environment*

Microscopic Kulagin-Petti (KP) model [NPA 765 (2006) 126, PRC 90 (2014) 045204].
 At large x nuclear DIS dominated by incoherent scattering off bound nucleons:

- ◆ **FERMI MOTION AND BINDING** effects in nuclear PDFs from the *convolution* of nuclear spectral function with (bound) nucleon PDFs:

$$F_2^A = \sum_{i=p,n} \int d\varepsilon d^3\mathbf{p} \mathcal{P}_i(\varepsilon, \mathbf{p}) K_2 F_2^i(x', Q^2, p^2)$$

where $x' = Q^2 / (2p \cdot q)$ and $p = (M + \varepsilon, \mathbf{p})$ and K_2 kinematic factor ($K_2 \approx 1 + p_z/M$ for $Q \gg M$).



- ◆ Since bound nucleons are **OFF-MASS-SHELL** there appears dependence on the *nucleon virtuality* $p^2 = (M + \varepsilon)^2 - \mathbf{p}^2$ and expanding PDFs in the small $(p^2 - M^2)/M^2$:

$$F_2^i(x, Q^2, p^2) \approx F_2^i(x, Q^2, p^2 = M^2) (1 + \delta f(x)(p^2 - M^2)/M^2).$$

where we introduced a universal function for the NUCLEON: $\delta f(x)$

⇒ *Modification of bound nucleon partonic structure in the nuclear environment*

DESCRIPTION OF NUCLEON

Distribution of partons in a nucleon



STRUCTURE FUNCTIONS

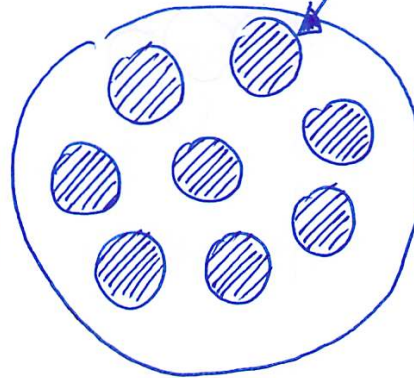
$F_i(x, Q^2)$: NNLO pQCD+TMC+HT from QCD fits

$\delta f(x)$: KP analysis of σ^A/σ^B DIS ratios

DESCRIPTION OF NUCLEUS

Distribution of bound nucleons

(A, Z)

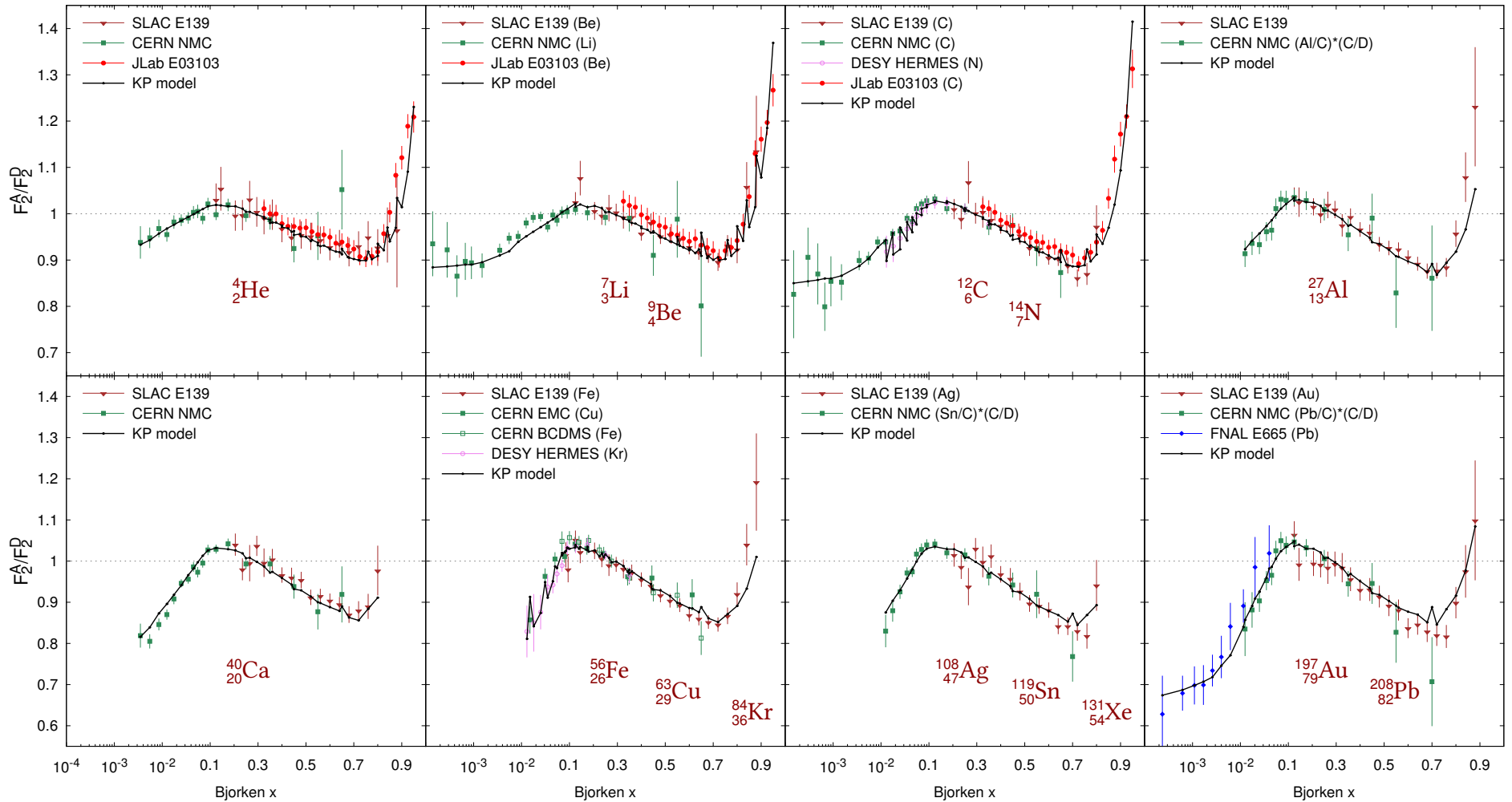


SPECTRAL FUNCTION

$\mathcal{P}(\varepsilon, \mathbf{p})$: mean field + N-N correlated part

Model includes meson-exchange current (MEC) correction balancing nuclear light-cone momentum and coherent multiple scattering effects responsible for nuclear shadowing

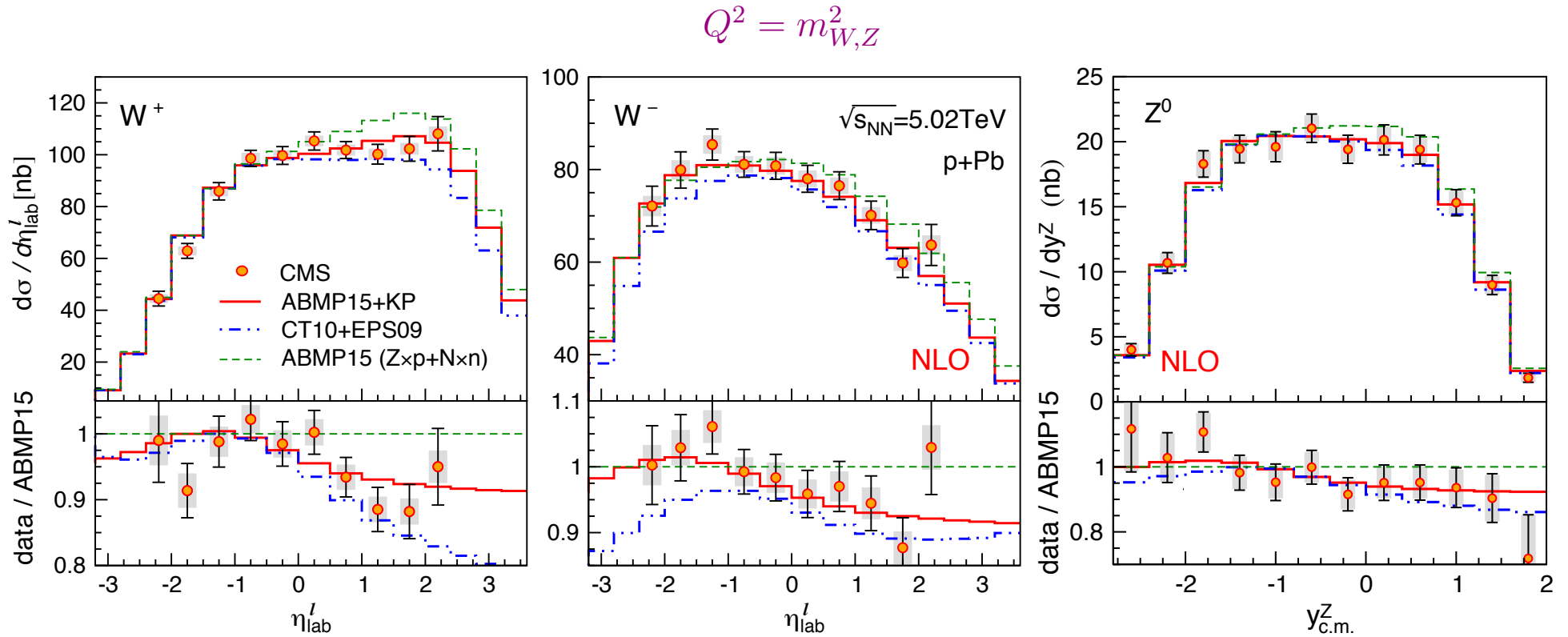
[NPA 765 (2006) 126, arXiv:hep-ph/0412425]



Microscopic KP model provides quantitative description of available data:

$$\chi^2/N_{\text{Data}} = 466.6/586 \text{ for DIS data with } Q^2 \geq 1 \text{ GeV}^2$$

\Rightarrow *Evidence for off-shell modification of bound nucleons from inclusive DIS*



*Predictions from KP model in excellent agreement with
Drell-Yan and W^\pm/Z boson production in pPb collisions up to $Q^2 = m_{W,Z}^2$*

(PRC 90 (2014) 045204; PRD 94 (2016) 113013)

- ◆ *Structure functions* are parameterized in the *NNLO QCD approximation*, supplemented by two (isoscalar) *High Twist (HT)* corrections to F_2 and F_T :

$$F_{2,T}(x, Q^2) = F_{2,T}^{\text{LT,TMC}}(x, Q^2) + \frac{H_{2,T}^N(x)}{Q^2}$$

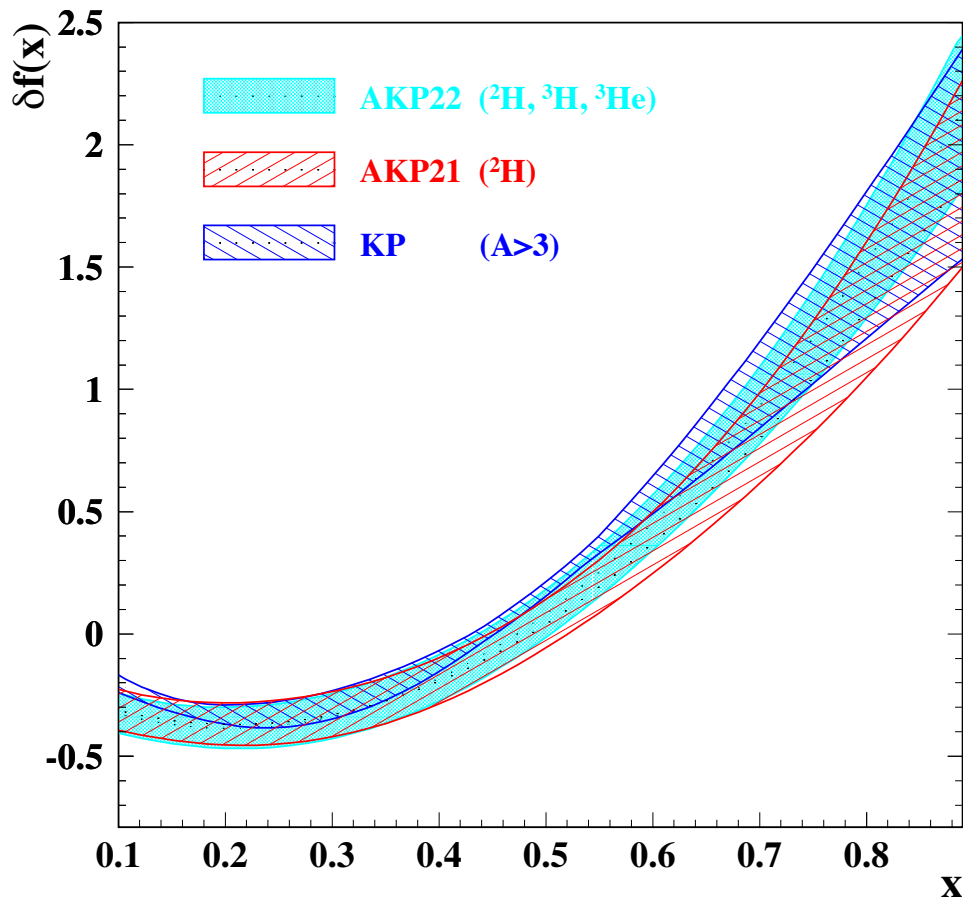
- *Target mass corrections (TMC)* in the *Leading Twist (LT)* term following *Georgi-Politzer*;
 - *Fixed flavor number scheme (FFNS)* with $n_f = 3$ and $\overline{\text{MS}}$ running masses for heavy quarks;
 - *PDFs* are parameterized following *ABMP16* at the initial scale $Q_0^2 = 9 \text{ GeV}^2$ [PRD 96 (2017) 014011];
 - Analysis performed in the region $Q^2 > 2.5 \text{ GeV}^2$ and $W^2 > 3 \text{ GeV}^2$.
- ◆ *Off-shell function* parameterized as *generic second order polynomial* to avoid model-dependent biases related to the functional form used:

$$\delta f(x) = a + b x + c x^2$$

- *Neglect nuclear effects* related to meson exchange currents and shadowing since focus at $x > 0.1$;
- ^2H *wave functions*: AV18 (default), Paris, CD-Bonn, WJC1, WJC2.
- ^3H and ^3He *spectral functions*: Rome with AV18 NN (default), Hannover with Paris NN.

⇒ *Simultaneous extraction of $\delta f(x)$, PDFs, and HT from global QCD analysis*

PRD 107 (2023) L051506; PRD 105 (2022) 114037; PRD 96 (2017) 054005



- ◆ *Different Q^2 dependence allows to disentangle off-shell correction from PDFs and HT*
 - ◆ *Results on $\delta f(x)$ agree with heavy target determination ($A \geq 4$) and our previous extraction from D data.*
 - ◆ *Determination of δf from QCD fits stable against all systematics studied.*
- ⇒ *Agreement with KP predictions based on δf universality*

- ◆ *Multiplicative vs. additive* implementation of High Twist (HT) terms:

$$F_{2,T}(x, Q^2) = F_{2,T}^{\text{LT,TMC}}(x, Q^2) + H_{2,T}^N(x)/Q^2$$
$$F_{2,T}(x, Q^2) = F_{2,T}^{\text{LT,TMC}}(x, Q^2) + F_{2,T}^{\text{LT}}(x, Q^2)h_{2,T}^N(x)/Q^2$$

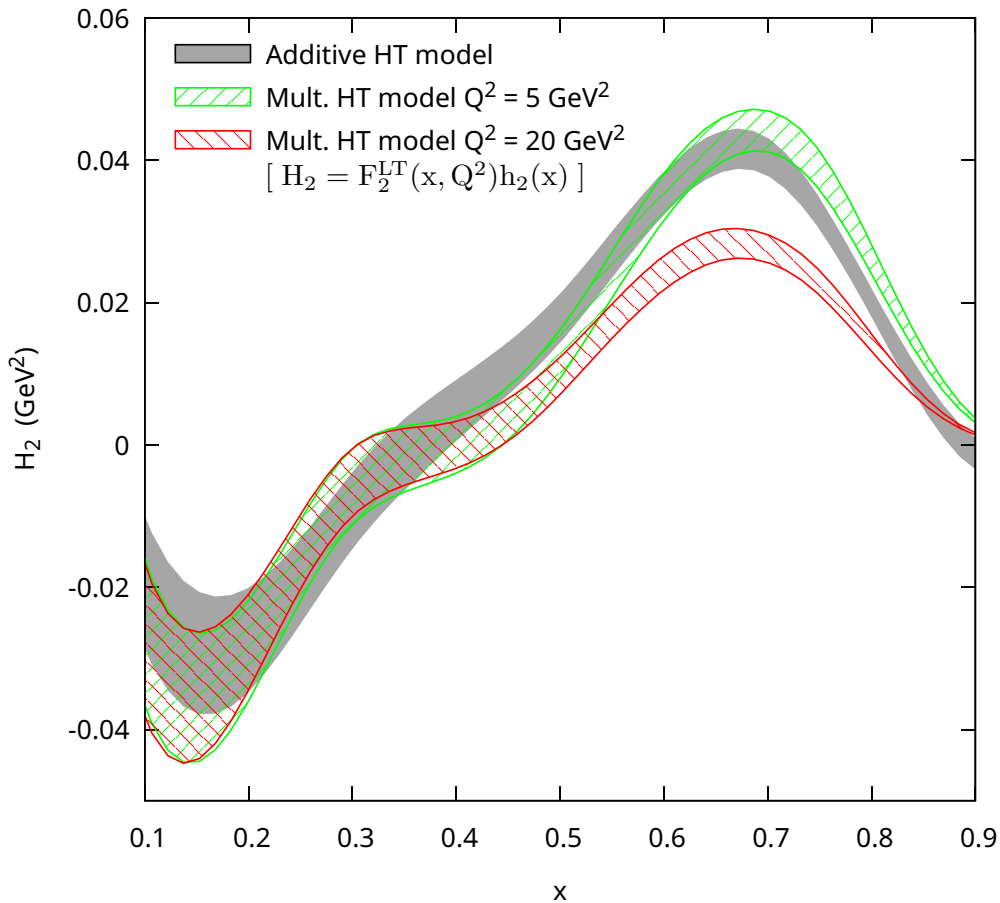
⇒ *Study impact of HT model on determination of LT and off-shell function δf*

- ◆ *Neutron-proton asymmetry in off-shell modification* of bound nucleons parameterized as generic first order polynomial:

$$\delta f^a(x) \equiv \delta f^n(x) - \delta f^p(x) = a_1 + b_1 x$$

- *Universality of isoscalar δf for all nuclei verified* with wide range of targets with $A \geq 4$ [NPA 765 (2006) 126], $A=3$ [PRC 82 (2010) 054614, PRL 128 (2022) 132002], and $A=2$ [PRD 105 (2022) 114037; PRD 96 (2017) 054005];
- *MARATHON $\sigma^{3\text{He}}/\sigma^{3\text{H}}$ data* [PRL 128 (2022) 132003] provides good sensitivity to n-p differences.

⇒ *Test/constrain a possible isospin dependence in the off-shell function δf*



◆ *Different Q^2 dependence for additive (aHT) and multiplicative (mHT) HT due to the interplay with LT.*

◆ *Consistent results from our fits with additive and multiplicative HT.*

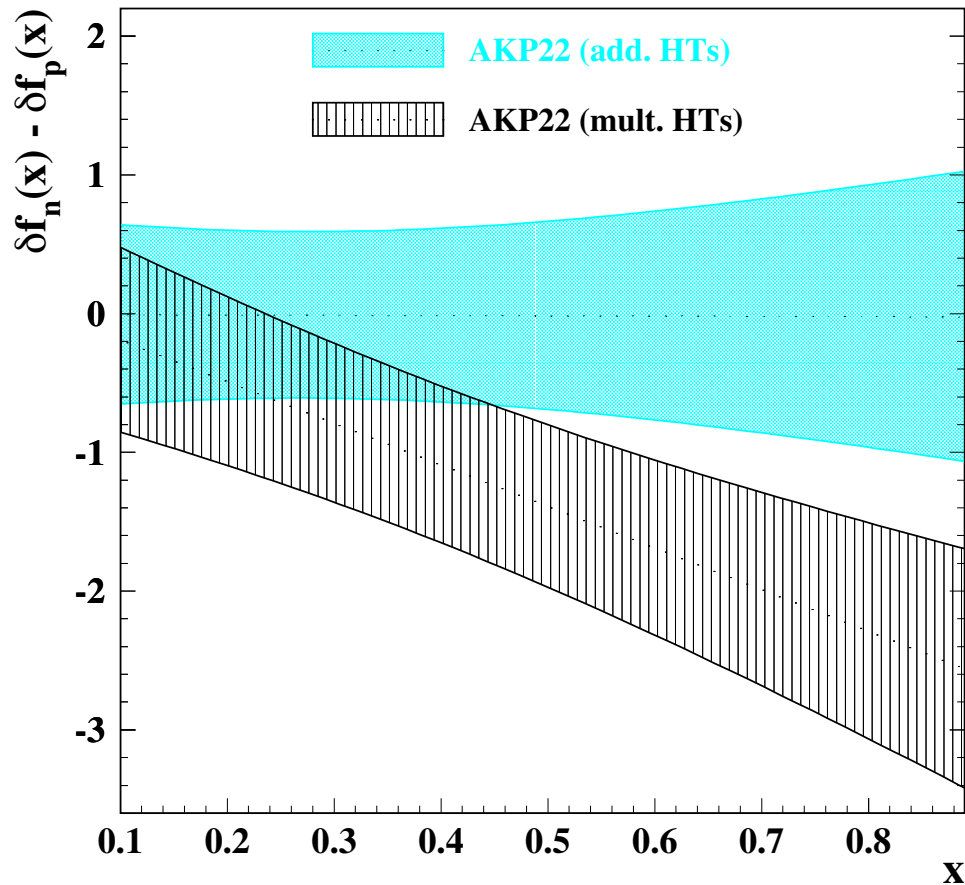
◆ *Intrinsic n-p difference in HT for multiplicative HT from the LT factor with isoscalar $h_{2,T}^N$ coefficients:*

$$F_{2,T}^n = F_{2,T}^{\text{LT},n} \left(1 + h_{2,T}^N / Q^2 \right)$$

$$F_{2,T}^p = F_{2,T}^{\text{LT},p} \left(1 + h_{2,T}^N / Q^2 \right)$$

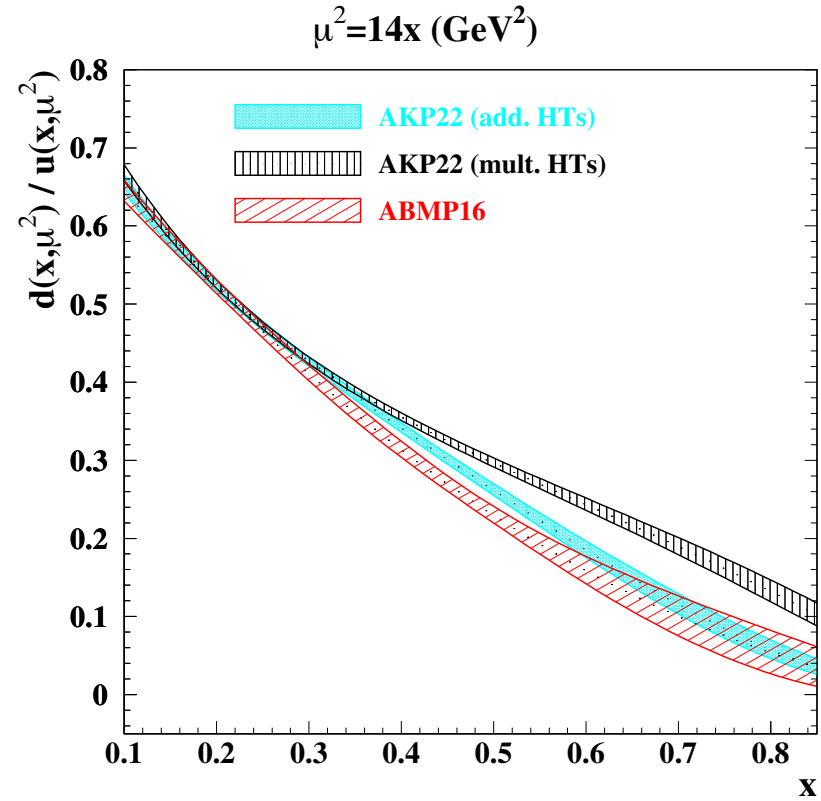
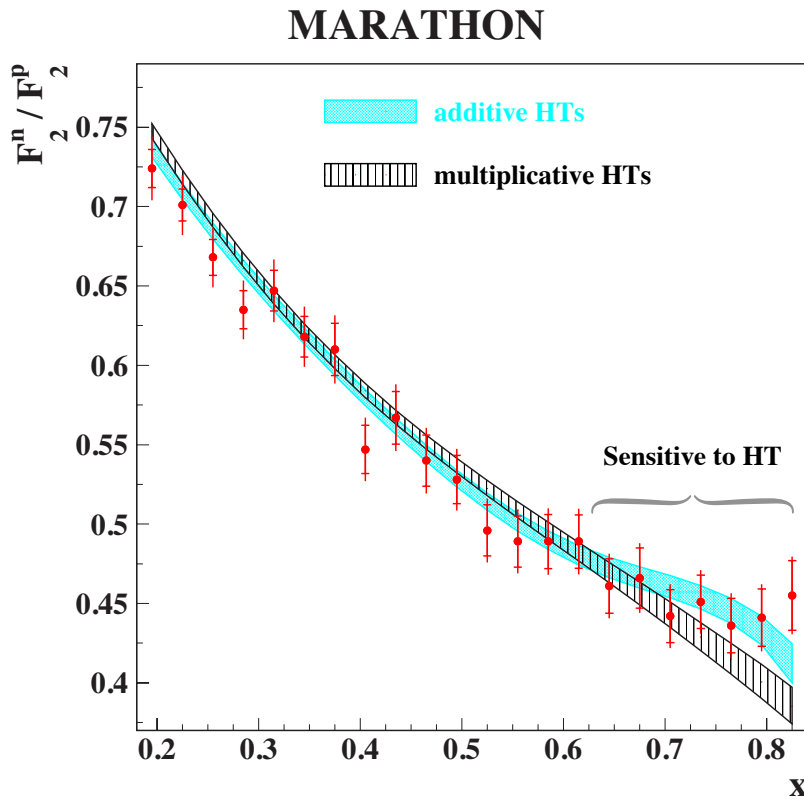
⇒ *Cancellation of HT terms in ratio*

$$F_{2,T}^n / F_{2,T}^p = F_{2,T}^{\text{LT},n} / F_{2,T}^{\text{LT},p}$$

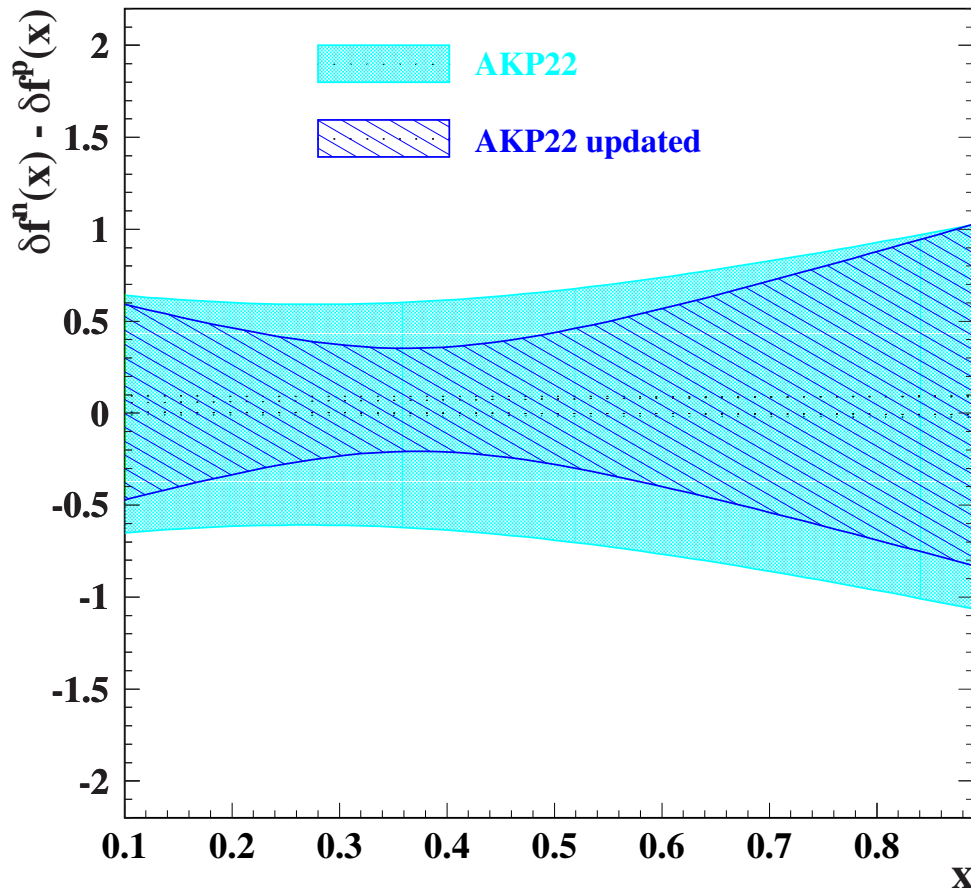


- ◆ *MARATHON $\sigma^{3\text{He}}/\sigma^{3\text{H}}$ allows an extraction of the n-p asymmetry $\delta f^a(x)$.*
 - ◆ *Same δf obtained for protons and neutrons with additive HT model.*
 - ◆ *Non-zero n-p asymmetry δf^a found with multiplicative HT model.*
- ⇒ *Bias introduced by LT-HT interplay in multiplicative HT model*

PREDICTIONS FOR F_2^n / F_2^p AND d/u RATIOS



- ◆ *MARATHON data prefers aHT over mHT model: $\chi^2/NDP = 20/22$ vs. $34/22$.*
- ◆ *Fitting MARATHON $\sigma^{3\text{He}}/\sigma^{3\text{H}}$ with mHT results in substantial d/u enhancement.*
- ⇒ *Enhancement of d/u correlated with the non-zero $n-p$ asymmetry δf^a with mHT*



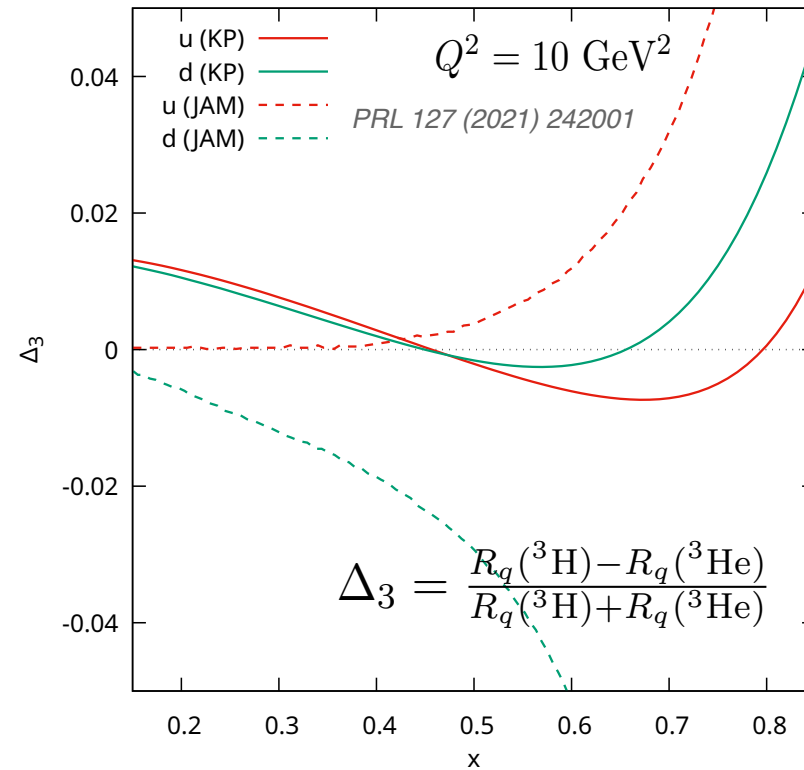
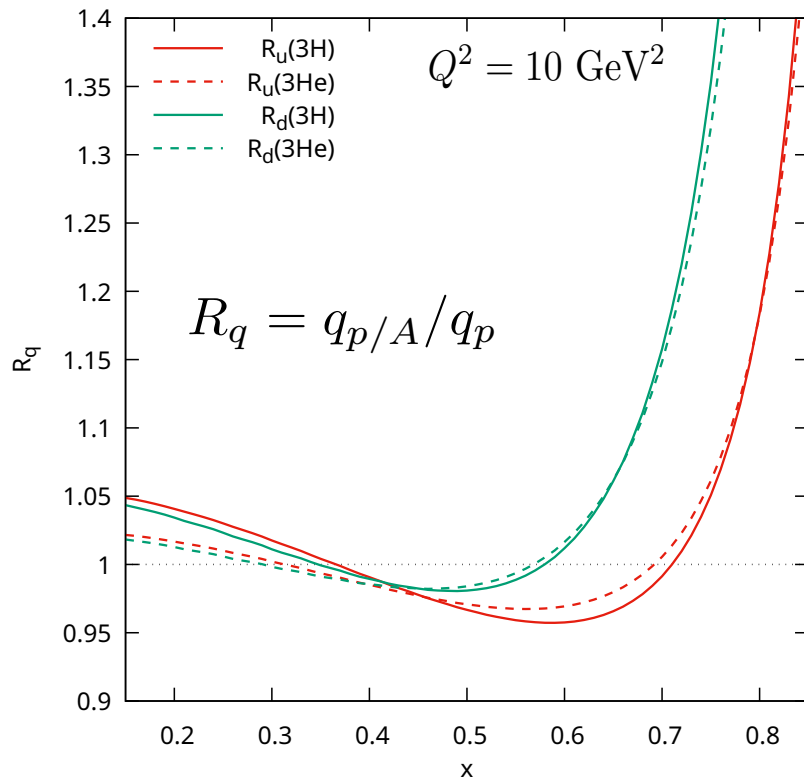
◆ Include $\sigma^{3\text{He}}/\sigma^{2\text{H}}$ data from JLab E03103 and HERMES in the analysis. [PRL 103 (2009) 202301; PLB 567 (2003) 339]

◆ Replace MARATHON $\sigma^{3\text{He}}/\sigma^{3\text{H}}$ with preliminary $\sigma^{3\text{He}}/\sigma^{2\text{H}}$ and $\sigma^{3\text{H}}/\sigma^{2\text{H}}$ data in analysis [F. Hauenstein, HiX2019].

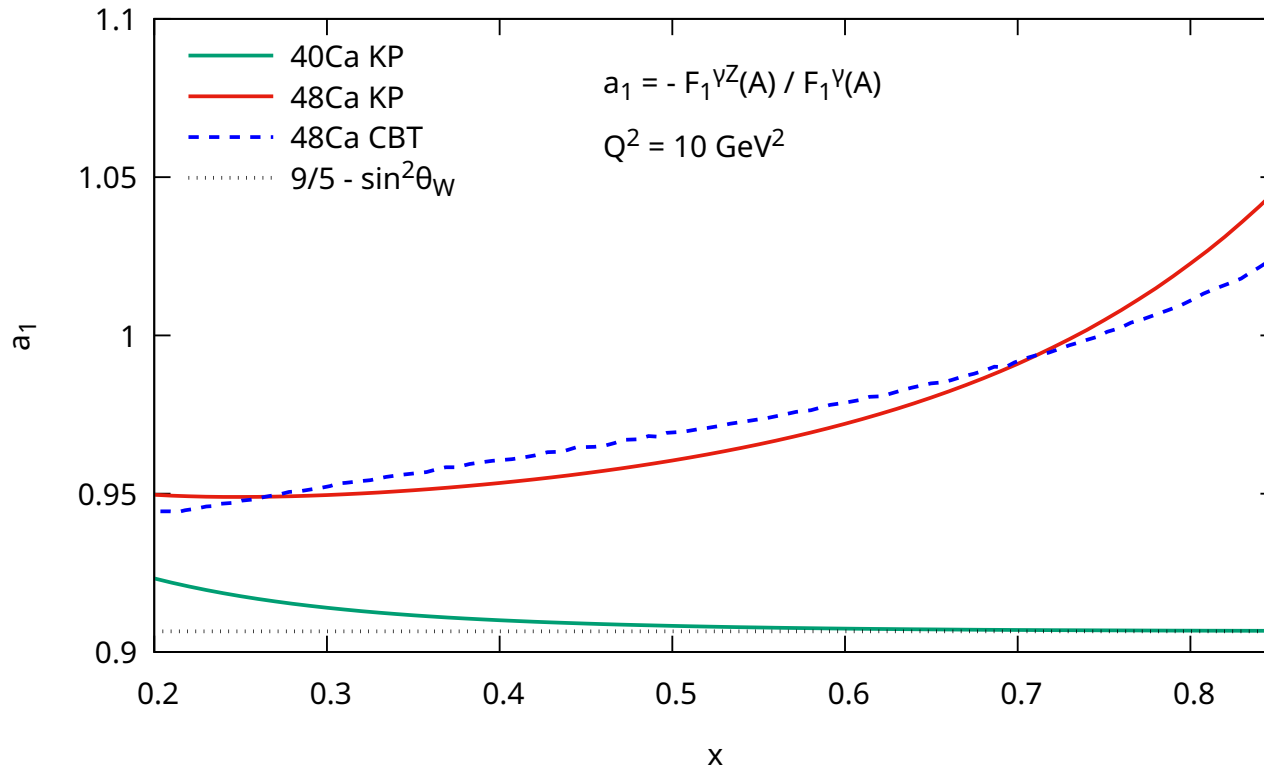
◆ Significant reduction of uncertainties on the n-p asymmetry $\delta f^a(x)$.

⇒ No evidence for isospin dependence of δf from all ^3He and ^3H data

NUCLEAR MODIFICATIONS OF d AND u IN ${}^3\text{H}$ AND ${}^3\text{He}$



- ◆ *Different nuclear corrections for d and u from convolution of their proton PDF shapes*
 - ◆ *Nuclear dependence of d, u modifications from p spectral functions in ${}^3\text{He}$ and ${}^3\text{H}$.*
- \implies *Significant isovector effects are present even with isoscalar off-shell $\delta f^n = \delta f^p$*

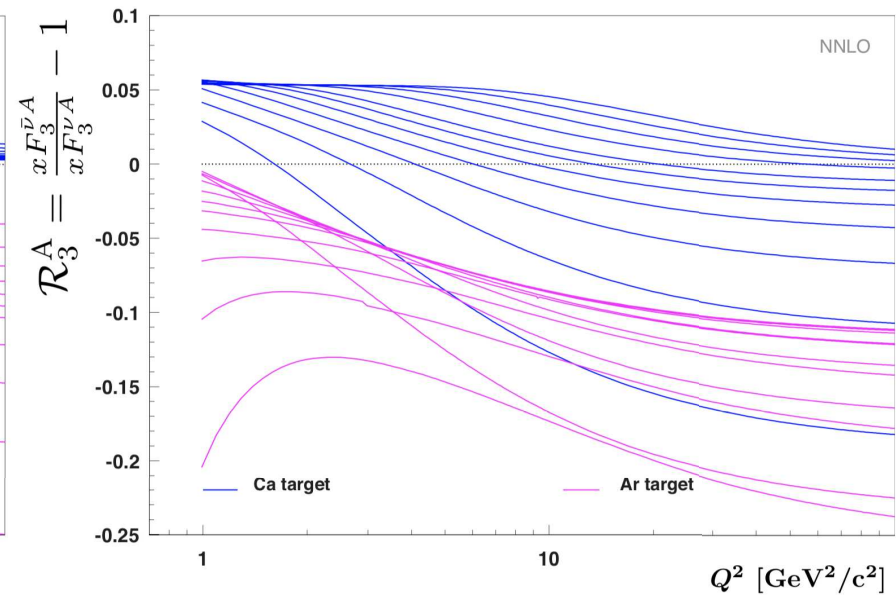
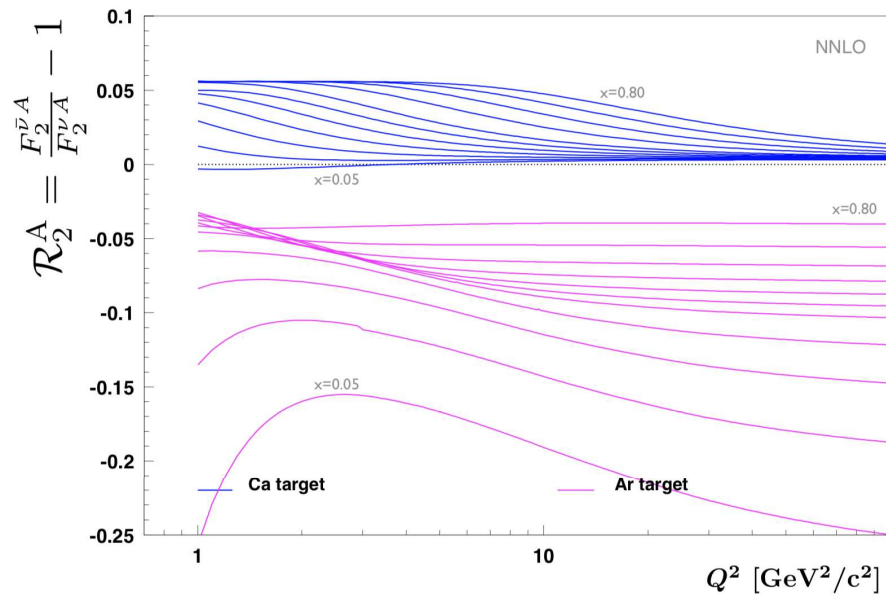


- ◆ *Isospin dependence of δf probed with parity-violating DIS in SoLID at JLab (PVEMC)*
- ◆ *Enhancement on asymmetry $a_1(x)$ in ^{48}Ca even with isoscalar off-shell $\delta f^n = \delta f^p$*
- ⇒ *Expected effects comparable to models with explicit isovector modifications*

EPJA 59 (2023) 194

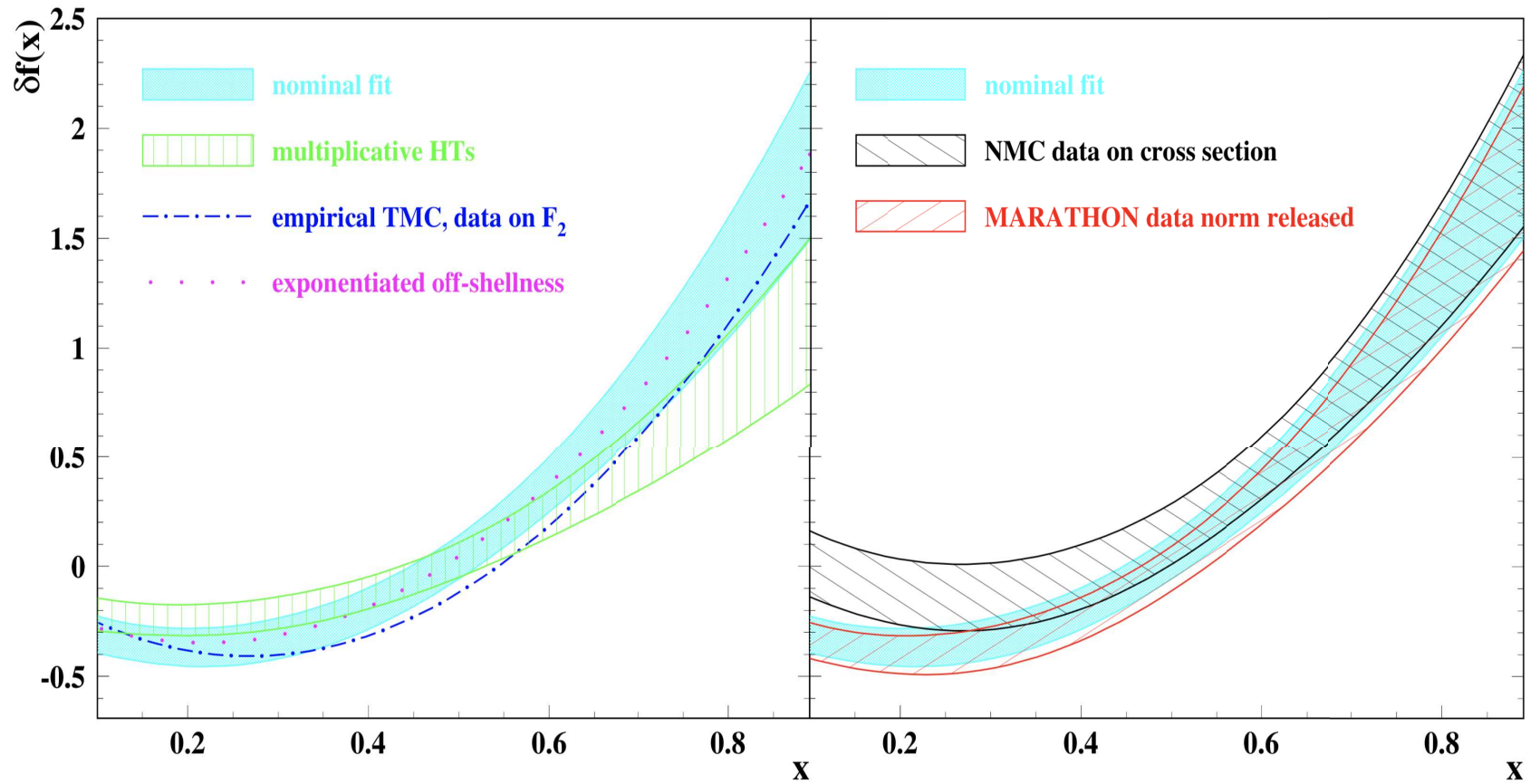
PREDICTIONS FOR $\bar{\nu}-\nu$ CC DIS AT LBNF

- ◆ Flavor selection of $\bar{\nu}-\nu$ CC DIS exploited in new program of *precision measurements on H,C,O,Ca,Ar targets with the intense LBNF beams at Fermilab.* PLB 834 (2022) 137469
- ◆ Comparison of $\mathcal{R}_{2,3}^A$ in ^{40}Ca and ^{40}Ar can probe isospin dependence of nuclear effects:
 - Same $A = 40$: neutron excess in Ar $\beta = (Z-N)/A \sim -0.1$, Ca mostly isoscalar $\beta \sim -2.6 \times 10^{-3}$;
 - Insights on physics mechanisms responsible for *isovector effects at both nucleon and nuclear level.*



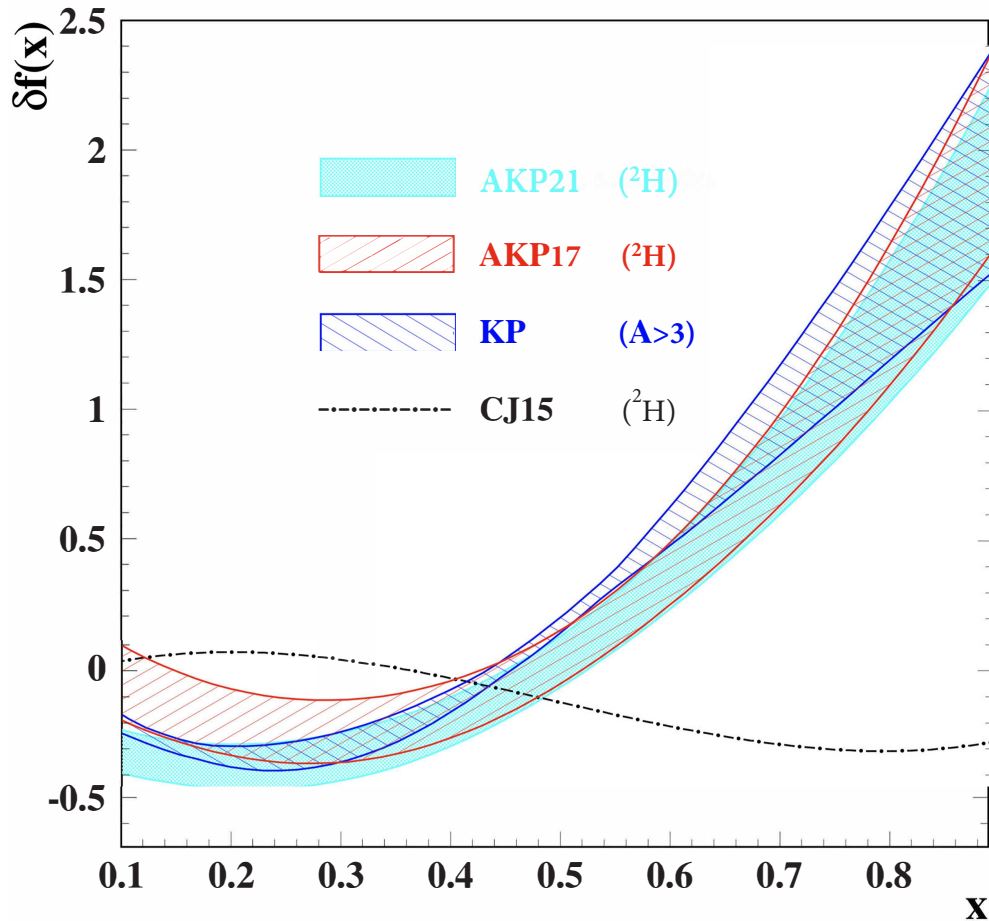
- ◆ *The off-shell modification of bound nucleons leads to an important nuclear correction which can be described by a universal function $\delta f(x)$ for all nuclei.*
- ◆ *The δf function determined from ${}^2\text{H}, {}^3\text{H}, {}^3\text{He}$ within our global QCD analysis is consistent with the one obtained from inclusive DIS data on nuclear targets with $A \geq 4$ (Kulagin and Petti) and from our earlier QCD analyses of ${}^2\text{H}$.*
- ◆ *We find no evidence for neutron-proton differences $\delta f^n - \delta f^p$ in the off-shell function from the QCD analysis of MARATHON $\sigma^{3\text{He}}/\sigma^{3\text{H}}$ data.
 \implies *Excellent agreement with MARATHON data with $\delta f^n = \delta f^p$ and isoscalar a_{HT}**
- ◆ *Our analysis indicates that the LT-HT interplay in the multiplicative HT model can bias both the n - p asymmetry δf^a and the d/u ratio extracted from MARATHON data.*
- ◆ *Substantial isovector effects on nuclear modifications of PDFs even with $\delta f^n = \delta f^p$ from convolution of different d/u shape with p and n spectral functions in nuclei.
 \implies *“Conventional” flavor dependence to be addressed before advocating $\delta f^n \neq \delta f^p$**

Backup slides



- ◆ *Determination of δf from QCD fits stable against all systematic variations studied*
- ◆ *Effect of model systematics comparable with the ones from use of different data sets*
- ⇒ *Consistency of results with nominal fit excludes model biases*

RESULTS ON $\delta f(x)$

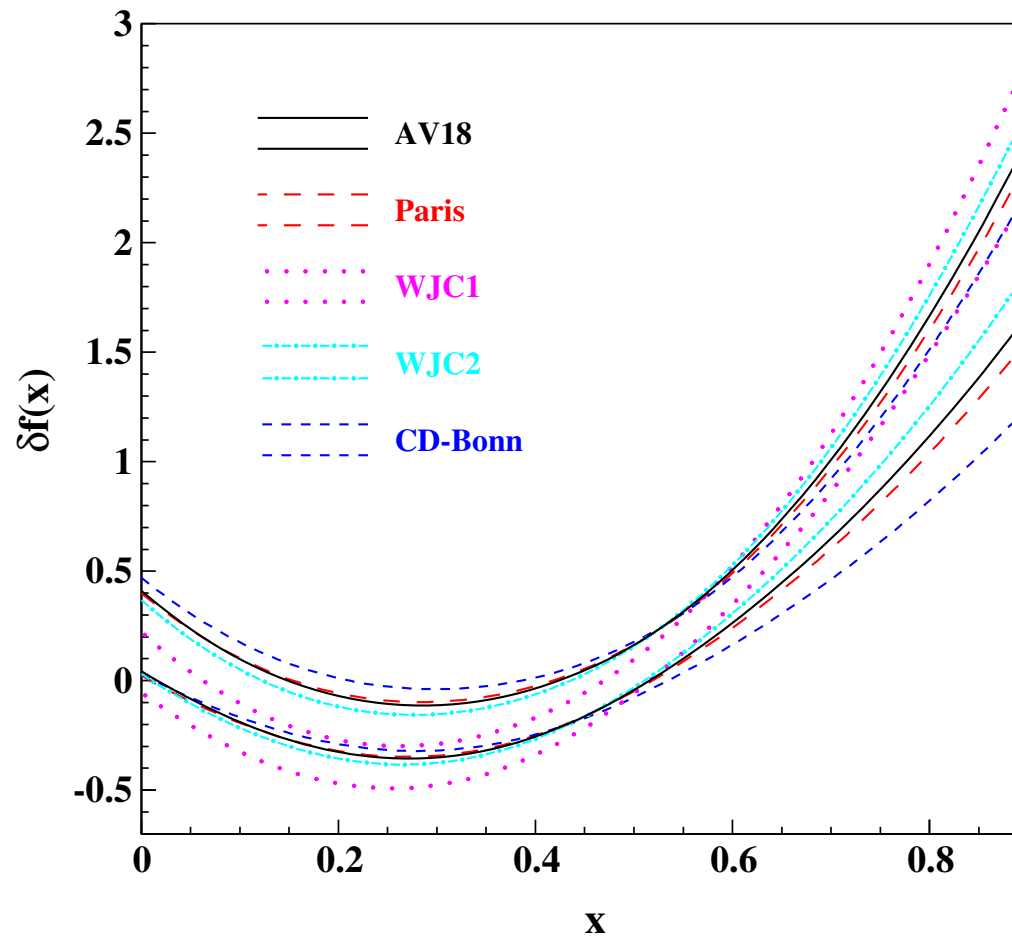


◆ *Our results on $\delta f(x)$ from ${}^2\text{H}$ consistent for all model variations agree with heavy target determination ($A \geq 4$).*

◆ *Clear disagreement with CJ15 results from ${}^2\text{H}$ in global QCD fits.*

◆ *Common meetings AKP-CJ held in 2020 and 2021 to try to understand differences*

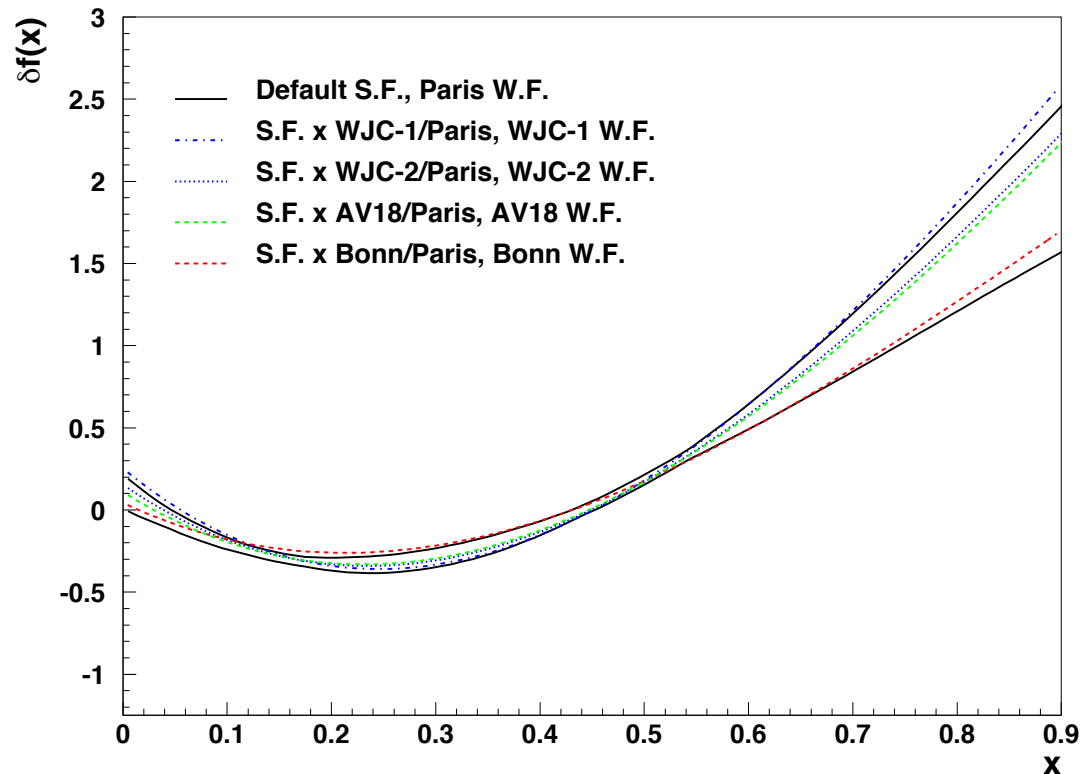
⇒ *Our QCD fits using CJ input data and model settings (with mHT and aHT) are consistent with AKP results*



Off-shell function determined from global QCD fits with different wave function models

PRD 96 (2017) 054005

OFF-SHELL FUNCTION FROM HEAVY TARGETS ($A \geq 4$)



- ◆ $\delta f(x)$ extracted phenomenologically from nuclear DIS ratios $\mathcal{R}_2(A, B) = F_2^A/F_2^B$:
 - Electron and muon scattering from BCDMS, EMC, E139, E140, E665 and NMC
 - Wide range of targets $^4\text{He}, ^7\text{Li}, ^9\text{Be}, ^{12}\text{C}, ^{27}\text{Al}, ^{40}\text{Ca}, ^{56}\text{Fe}, ^{64}\text{Cu}, ^{108}\text{Ag}, ^{119}\text{Sn}, ^{197}\text{Au}, ^{207}\text{Pb}$
 - Systematic uncertainties including modeling, functional form and spectral/wave function variations
- ⇒ Partial cancellation of systematics from spectral function in RATIOS $\mathcal{R}_2(A, B)$

NUCLEAR SPECTRAL FUNCTION

- ◆ Two-body ${}^2\text{H}$ spectral function determined by the wave function $\Psi_D(\mathbf{p})$:

$$\mathcal{P}(\varepsilon, \mathbf{p}) = 2\pi\delta\left(\varepsilon - \varepsilon_D + \frac{\mathbf{p}^2}{2M}\right) |\Psi_D(\mathbf{p})|^2$$

where $\varepsilon_D = M_D - 2M \approx -2.2$ MeV is the binding energy.

- ◆ Three-body ${}^3\text{He}$ and ${}^3\text{H}$ spectral functions from D bound state and continuum states:

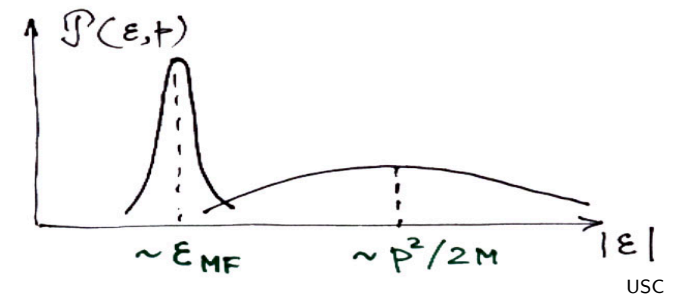
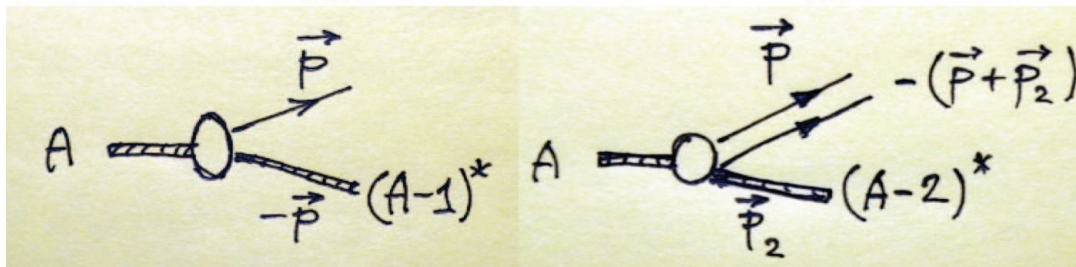
$$\mathcal{P}_{3\text{He}}^p = f_{3\text{He}}^D(\mathbf{p})\delta\left(\varepsilon + \varepsilon_{32} - \varepsilon_D + \frac{\mathbf{p}^2}{4M}\right) + f_{3\text{He}}^{pn}(\varepsilon, \mathbf{p}); \quad \mathcal{P}_{3\text{He}}^n = f_{3\text{He}}^{pp}(\varepsilon, \mathbf{p})$$

$$\mathcal{P}_{3\text{H}}^n = f_{3\text{H}}^D(\mathbf{p})\delta\left(\varepsilon + \varepsilon_{31} - \varepsilon_D + \frac{\mathbf{p}^2}{4M}\right) + f_{3\text{H}}^{pn}(\varepsilon, \mathbf{p}); \quad \mathcal{P}_{3\text{H}}^p = f_{3\text{H}}^{nn}(\varepsilon, \mathbf{p})$$

where $\varepsilon_{32} \approx -7.72$ MeV and $\varepsilon_{31} \approx -8.48$ MeV are the ${}^3\text{He}$ and ${}^3\text{H}$ binding energies.

- ◆ Spectral function for $A \geq 4$ nuclei with mean field \mathcal{P}_{MF} and NN correlated \mathcal{P}_{cor} parts:

$$\mathcal{P} = \mathcal{P}_{\text{MF}} + \mathcal{P}_{\text{cor}}$$



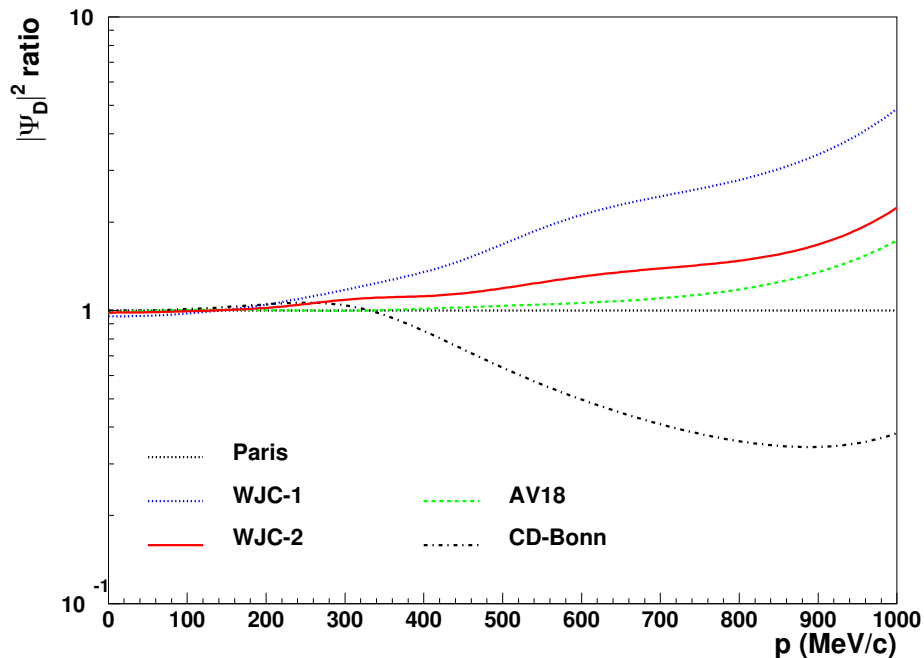
DEUTERON WAVE FUNCTION

- ◆ Two-body nucleus whose *spectral function* determined by the wave function $\Psi_D(\mathbf{p})$:

$$\mathcal{P}(\varepsilon, \mathbf{p}) = 2\pi\delta\left(\varepsilon - \varepsilon_D + \frac{\mathbf{p}^2}{2M}\right) |\Psi_D(\mathbf{p})|^2$$

where $\varepsilon_D = M_D - 2M \approx -2.2$ MeV is the binding energy.

- ◆ The deuteron is a superposition of *s*- and *d*-wave states. Different models of $\Psi_D(\mathbf{p})$ based on the corresponding underlying *N-N* interaction potentials, which are constrained at low momentum ($p < 300$ MeV/c) by *pp*, *pn* and *nn* scattering data.



$|\Psi_D(\mathbf{p})|^2$ gives deuteron momentum distribution

Different N-N potentials used

Paris: PRC 21 (1980) 861

CD-Bonn: PRC 63 (2001) 024001

AV18: PRC 84 (2011) 034003

WJC-1,2: PRC 82 (2010) 034004

[AKP, PRD 96 (2017) 054005]

Facility	Experiment	Beam	Beam energy (GeV)	Observable	Normalization factor	Normalization error(s) (%)	$\frac{\chi^2}{\text{NDP}}$
SLAC	E49a	e	11 ÷ 19.5	$\frac{d^2\sigma^d}{dE'd\Omega}$	0.988(10)	2.1 ^a	25/59
"	E49b	"	4.5 ÷ 18	"	0.996(10)	"	187/145
"	E87	"	8.7 ÷ 20	"	1.000(9)	"	114/109
"	E89b	"	10.4 ÷ 19.5	"	0.987(9)	"	52/72
"	E139	"	8 ÷ 24.5	"	1.002(9)	"	8/17
"	E140	"	3.7 ÷ 19.5	"	1	1.7	25/26
CERN	BCDMS	μ	100 ÷ 280	$\frac{d^2\sigma^d}{dx dQ^2}$	0.989(7)	3	273/254
"	NMC	"	90 ÷ 280	F_2^d/F_2^p	1	< 0.15	155/165
DESY	HERMES	e	27.6	σ^d/σ^p	1	1.4	21/30
JLab	E00-116	e	5.5	$\frac{d^2\sigma^d}{dE'd\Omega}$	0.981(10)	1.75	208/136
"	BONuS	"	4.2, 5.2	F_2^n/F_2^d	0.97(9)	7 ÷ 10	90/63
"	MARATHON	"	10.6	σ^d/σ^p	1	0.55	8/7
"	MARATHON	"	10.6	$\sigma^{^3\text{He}}/\sigma^{^3\text{H}}$	1	0.7	20/22
Total							1186/1105

List of ^2H , ^3H , and ^3He data used in the global QCD analysis