A new concept to extract and collimate UCN without loosing phase-space density

A. Mietke^{1,2}, S. Baessler², V. Nesvizhevsky³

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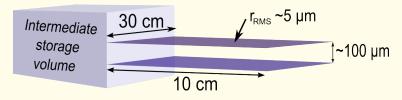


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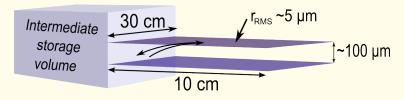
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- can be diffusely reflected, specularly reflected or absorbed
- ullet are accepted at the slit if: $E_{\perp} < m_n g h_{slit} pprox 10~peV$

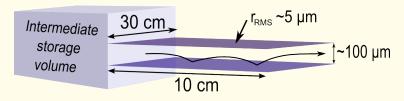




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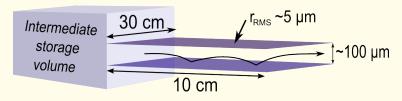




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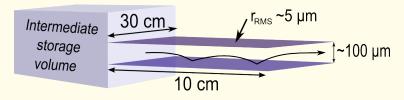




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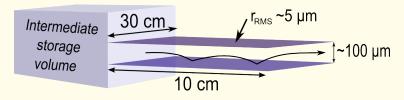




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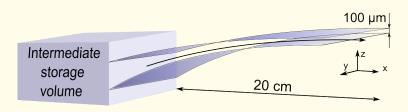




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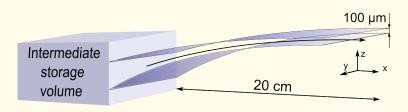




Collimator with parabola-shaped edges - the slope corresponds to $v_n \sim 7 \, {m \over s}$

Now we adapt the collimating system for the neutron trajectory and expect the following improvements:

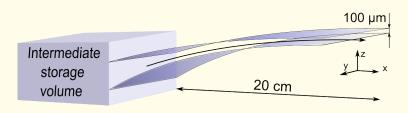
- Higher rate of reflections back (P_R) into the storage volume
 A larger aperture, that still provides a proper phase space density of the UCN, becomes possible
- $P_{diff} \propto v_{\perp}^2$ is essential for the effectiveness of the system \Rightarrow Once the neutron trajectory is adapted to the edge, specular reflections become more likely



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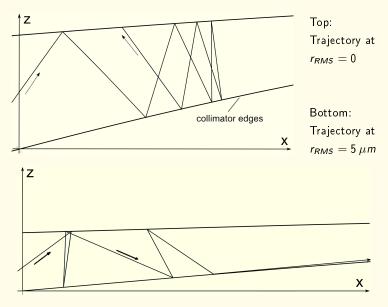
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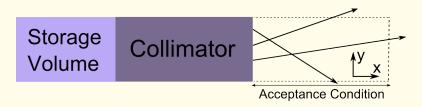


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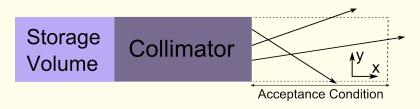


Collimationsystem top view - Visualization Acceptance Condition

 Neutrons are counted with respect to their alignment
 What intensity can be expected in a certain distance from the exit slit

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$$P_{diff} = \left(\frac{k_{\perp} r_{RMS}}{2\pi}\right)^2 \stackrel{k_{\perp}}{\Longrightarrow} \stackrel{= \frac{2\pi p_{\perp}}{h}}{\Longrightarrow} P_{diff} = \frac{r^2 m_n^2}{h^2} V_{\perp}^2$$

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- The neutrons are then scatterd with an angle distribution, that is based on the scattering law¹ and yields: $d\Omega_{out} = \cos\theta_{out} d\cos\theta_{out} d\phi_{out} = \frac{1}{2} d\cos^2\theta_{out} d\phi_{out}$

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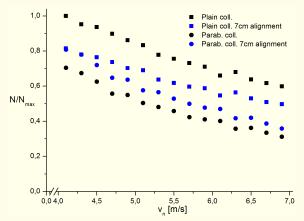
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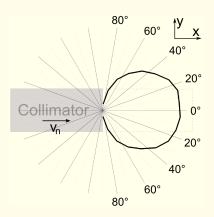
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4. Simulationdata - velocity distributions

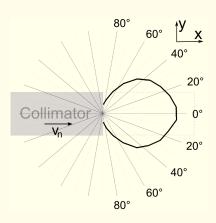


Velocity distributions for parabola-shaped and plain collimator

4. Simulationdata - angle distributions



Angle distribution plain collimator



Angle distribution parabola- shaped collimator

 A_0 : Aperture area, P_p : Passing probability, P_b : Back reflection probability

Flux relation:

$$\frac{f_{parabola}}{f_{plain}} = \frac{A_{0,par}P_{p,par}}{A_{0,plain}P_{p,plain}}$$

Effective aperture:

$$A_{eff} = A_0(1 - P_R)$$

Main goal:

$$A_{eff} < 1.5 \, mm^2$$



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Plain collimator:

Width	$A_{0,plain} [mm^2]$	$A_{eff} [mm^2]$	P_b	$P_{p,plain}$	
30 cm	30	1.05	96.5 %	0.7/0.45 %	

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Parabola-shaped collimator:

Width	$A_{0,par}$ $[mm^2]$	$A_{eff} [mm^2]$	P_b	$P_{p,par}$	f _{parabola} f _{plain}
26 cm	56.42	1.2	97.8 %	0.43/0.27 %	1.3
28 cm	60.76	1.3	97.8 %	0.46/0.28 %	1.4
30 cm	65.1	1.4	97.8 %	0.48/0.3 %	1.5

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- Inserting a mirror at the bottom decreases P_p rapidly
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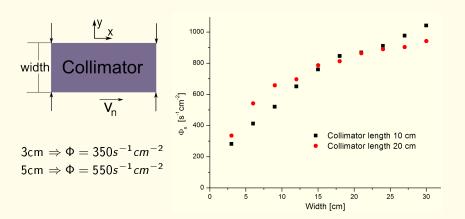
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6. Alternative application: Neutron reflectometry

ullet Only a width < 5~cm is useful for this measurement technique



- A Monte-Carlo simulation was successfully developed, to investigate and compare two collimation systems
- While taking a limit for the effective aperture into account, the parabola-shaped collimator was optimized

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Appendix - Scattering Law

The condition of detailed balance states, that the number of neutrons reflected from Ω_0 to Ω must be equal to the number of neutrons reflected from Ω to Ω_0 and can be expressed as:

$$\cos\theta_0 d\Omega_0 W_{ref}(\Omega_0, \Omega) d\Omega = \cos\theta d\Omega W_{ref}(\Omega, \Omega_0) d\Omega_0
\Rightarrow \cos\theta_0 W_{ref}(\Omega_0, \Omega) = \cos\theta W_{ref}(\Omega, \Omega_0) \tag{1}$$

Where θ und θ_0 are the angles between normal of the surface and direction of passing neutrons. Because of the required equilibrium, W_{ref} must be symmetric in Ω and Ω_0 and with respect to (1) can be satisfied by functions as: $W_{ref}(\Omega_i,\Omega_j)=f(\Omega_i,\Omega_j)\cos\theta_j$, where f is an arbitrary function, symmetric in its arguments. Assuming f=const. allows us to normalize W_{ref} to satisfy $\int W_{ref} d\Omega=1$ and leads to:

$$W_{ref}d\Omega = \frac{1}{\pi}\cos\theta d\Omega = \frac{1}{2\pi}d\cos^2\theta d\phi$$

