

# "Who wants neutrons?"



MONTCLAIR  
STATE UNIVERSITY

# ”Who wants neutrons?”



Ecole de Physique des Houches does!

# Beyond Ramsey: nEDM techniques with Superfluid $^4\text{He}$ for reaching $10^{-28}$ e.cm

$\pi/2$  pulse

NEUTRONS

Ecole de Physique des Houches does!

Kent Leung, EDMs2016: WE-Heraeus Workshop

*Les Houches, France, March 03, 2026*



“In the sciences of observation, the work of the observer consists in making experiments, and in observing the results; but his/her **skill is shown in the selection of the experiments, and in the estimation of the accuracy of the results.**”

— James Clerk Maxwell, “On the Elementary Relations between Electrical Measurements,” *Philosophical Magazine* (1861).

“The history of science teaches only too plainly the lesson that **no single method is absolutely to be relied upon, that sources of error lurk where they are least expected, and that they may escape the notice of the most experienced and conscientious worker.**”

— Sir John William Strutt, 3rd Baron Rayleigh, *Report of the British Association for the Advancement of Science* (1883).



# Where the EDM field started?

Vol. 3—No. 13

OAK RIDGE, TENNESSEE



## Harvard University Conducts Important Research at ORNL

The growing importance of Oak Ridge National Laboratory as a research center is manifested particularly in its assistance to universities and technical schools on various projects in which nuclear research is involved. An example of such relationship is its present collaboration with Harvard University in an investigation to determine if neutrons have permanent electric dipole moments.

The work of the project is under the direction of Professors E. M. Purcell and Norman F. Ramsey of the Harvard University Physics Department and is being conducted on the Laboratory area by James H. Smith, a graduate student at Harvard. Dur-

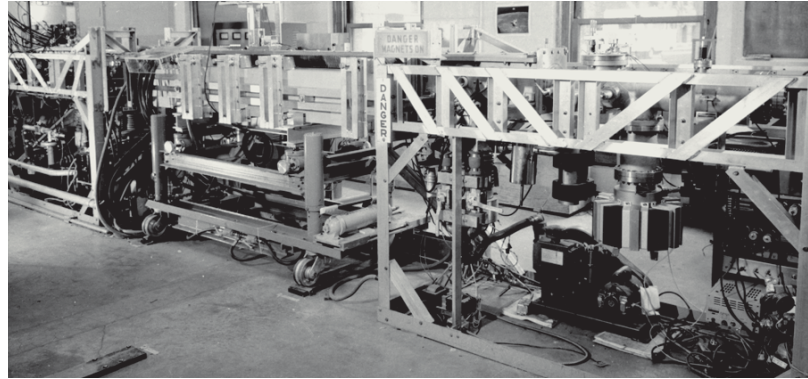
**HARVARD UNIVERSITY SPONSORS PROGRAM HERE** — James H. Smith, Harvard University graduate student in physics, is shown as he adjusts a neutron beam apparatus at the south face of the Oak Ridge Pile. Using the Pile as a source of neutrons, Mr. Smith is engaged in a project jointly sponsored by Harvard University and Oak Ridge National Laboratory for the purpose of determining if neutrons have permanent electric dipole moments.

Friday, September 29, 1950



- Oak Ridge National Laboratory's X-10 pile went critical in 1943, and civilian research began in 1945.
- Wu & Ambler's P-violation experiment in 1957
- nEDM Result published in 1957 by *Smith, Purcell & Ramsey*:  $D = (-0.1 \pm 2.4) \times 10^{-20} \text{ cm}$ ,

# Ramsey's method of separated oscillatory fields



## The Nobel Prize in Physics 1989



Norman F. Ramsey  
Prize share: 1/2



Hans G. Dehmelt  
Prize share: 1/4



Wolfgang Paul  
Prize share: 1/4

The Nobel Prize in Physics 1989 was divided, one half awarded to Norman F. Ramsey "for the invention of the separated oscillatory fields method and its use in the hydrogen maser and other atomic clocks",

## Norman F. Ramsey

In 1949 I was looking for a way to measure nuclear magnetic moments by the molecular-beam resonance method, but to do it more accurately than was possible with the arrangement developed by I. I. Rabi and his colleagues at Columbia University. The method I found<sup>1,2</sup> was that of separated oscillatory fields, in which the single oscillating magnetic field in the center of a Rabi device is replaced by two oscillating fields at the entrance and exit, respectively, of the space in which the nuclear magnetic

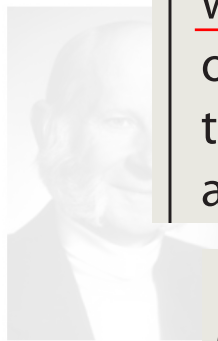
## Advice from Rabi

The first advice I received from Rabi in 1937 when I applied to him to begin my research was that I should not go into the field of molecular beams since the interesting problems amenable to that technique had already been solved and there was little future to the field. I have often wondered how I ... had the temerity to disregard this bit of advice from the master. However, I

(Quoted from Norman Ramsey's contribution to *A Tribute to Professor I. I. Rabi*, Columbia University, New York, 1970.)



Norman F. Ramsey  
Prize share: 1/2



Hans G. Dehmelt  
Prize share: 1/4

The Nobel Prize in Physics 1989 was divided equally between Norman F. Ramsey "for the invention of the separated oscillatory fields method and its use in the hydrogen maser and other atomic clocks",

## Ramsey

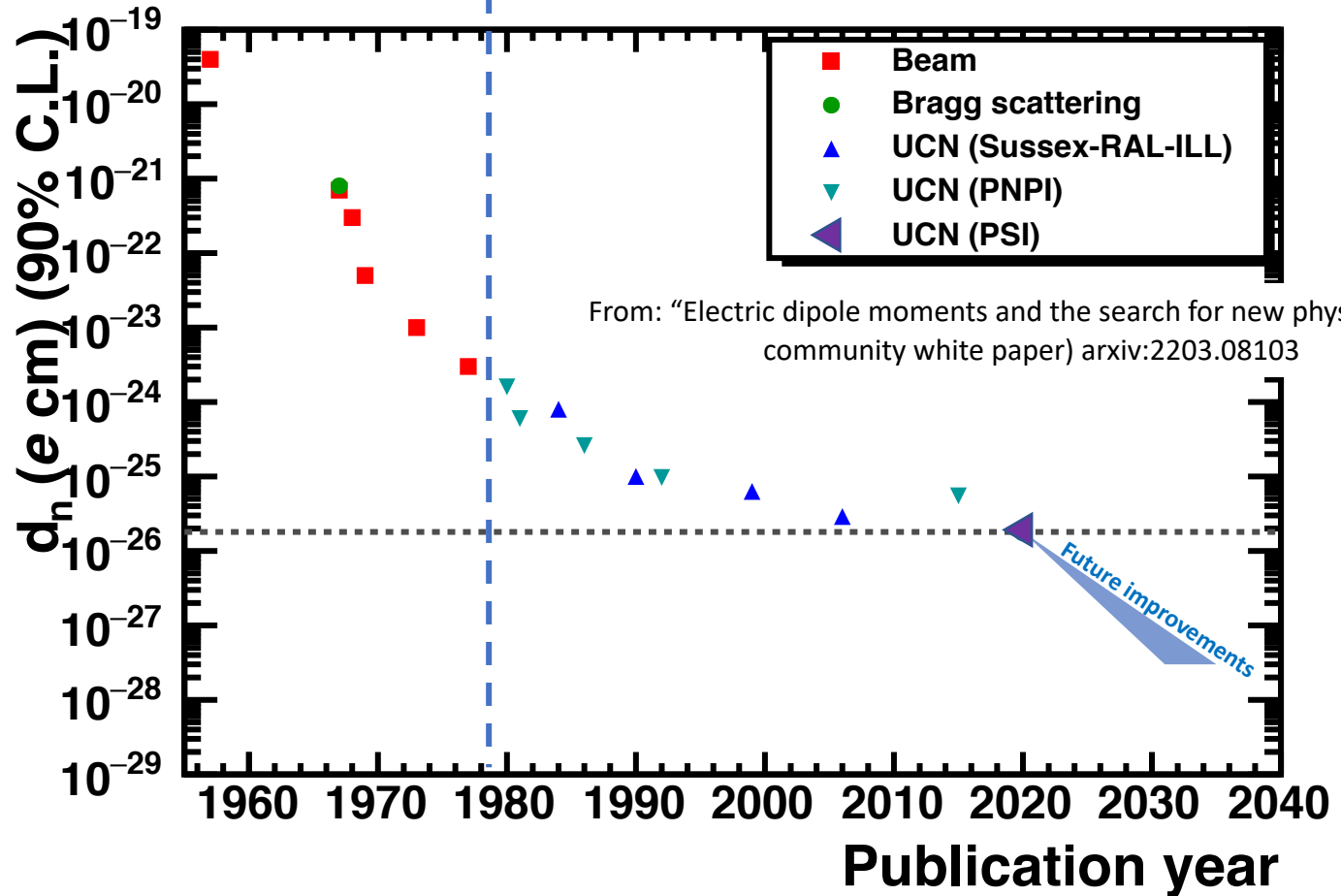
...king for a way to measure magnetic moments by the resonance method, but more accurately than was possible with the method developed by I. I. Rabi and his colleagues at Columbia University. The method I found<sup>1,2</sup> was based on the use of two separated oscillatory fields, in which the center of a Rabi device is flanked by two oscillating fields at the entrance and exit, respectively, of the space in which the nuclear magnetic

# Neutron EDM historical sensitivity

(High numbers/rates; short interaction/useful time)

(Low numbers/rates; long interaction/useful time)

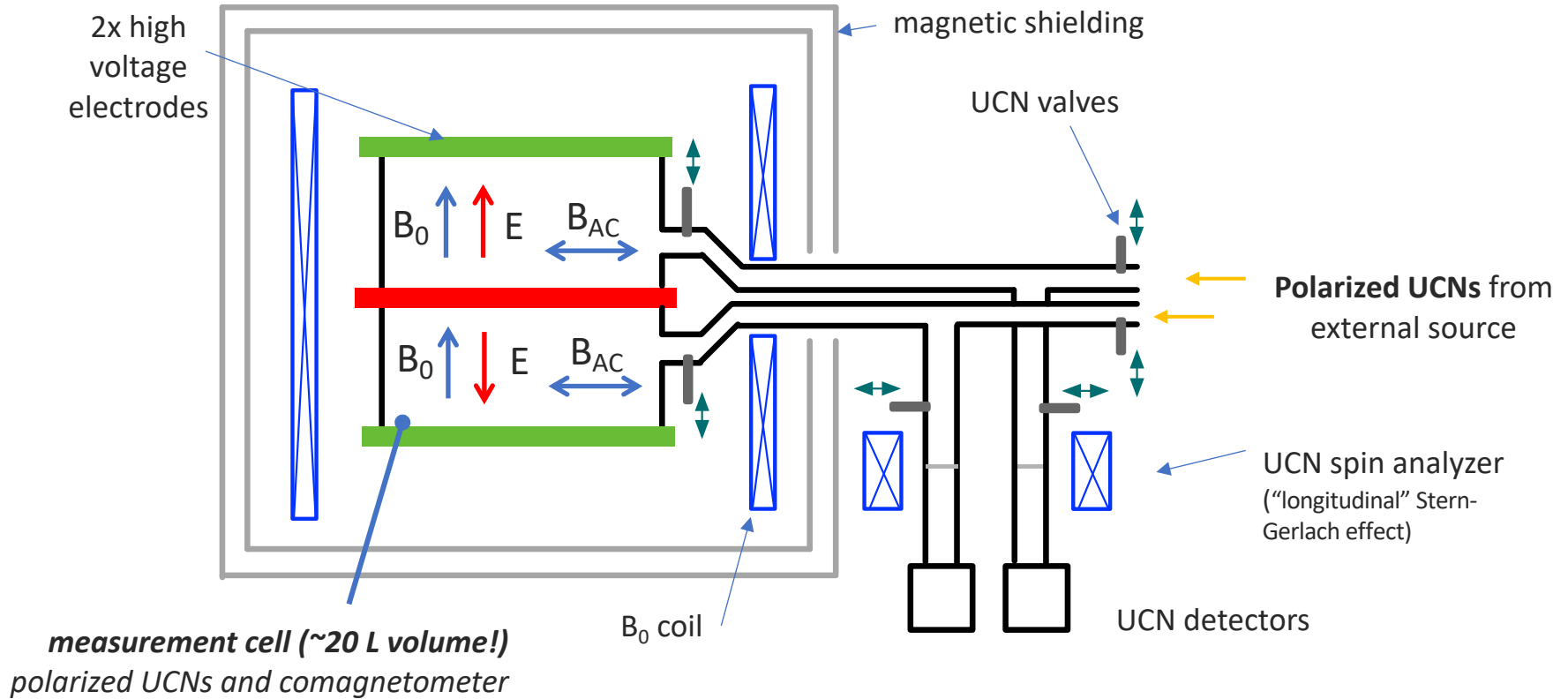
**Cold neutron beam era → ultracold neutrons era (but kept Ramsey's method)**



Experimental techniques with neutrons restricted by: neutral & no electronic states.  
(Excited  $\Delta$  state  $\sim 200$  MeV away...)



# Standard Ramsey technique with UCNs



- At the **end of free precession**, neutron population's **final phase** measured by counting number of spin down  $N_{\text{down}}$  (or  $N_{\text{up}}$ ). For UCNs, spin analysis typically done *ex-situ* for UCNs measurements). A full cycle takes 200-300s.
- Repeat cycle  $\sim 4x$  to extract a frequency in a cell
- Reverse E-field or, better, the "super ratio"

$$S = \frac{f_{\text{cell1}}^{\text{HV}\uparrow} f_{\text{cell2}}^{\text{HV}\downarrow}}{f_{\text{cell1}}^{\text{HV}\downarrow} f_{\text{cell2}}^{\text{HV}\uparrow}}$$

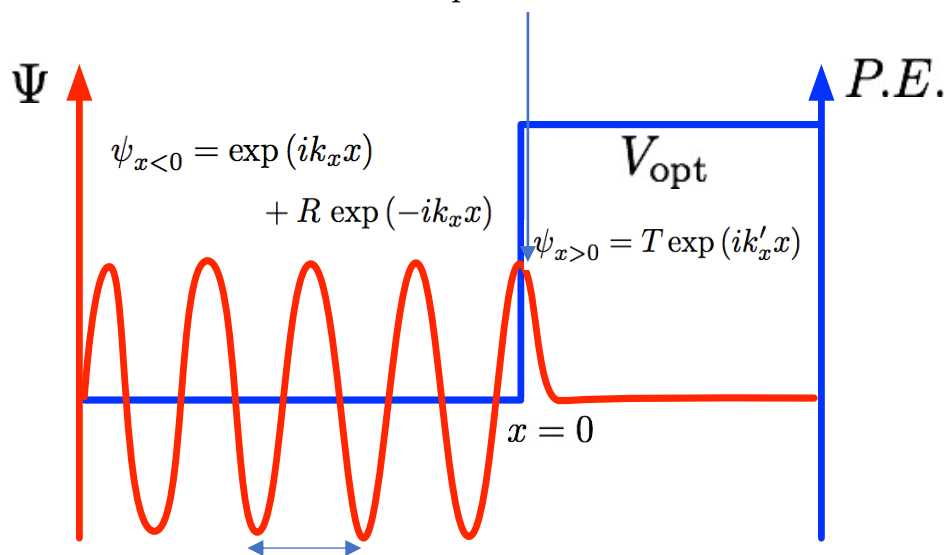
# “Ultracold Neutrons are Kitchen”

- L. N. Bondarenko, Kurchatov Institute

Transporting, especially polarized UCNs, can be tricky! Loss mechanisms not fully understood.

$$W = \frac{\hbar}{2} \sum_i n_i \sigma_{\text{loss}}(v) v$$

$$U = V_{\text{opt}} + iW$$



*Neutron incoherent scattering cross-section*

Isotope	Inc xs
H	80.26
1H	80.27
2H	2.05
3H	0.14
He	0
3He	1.6
4He	0
Li	0.92
6Li	0.46
7Li	0.78
Be	0.0018
B	1.7
10B	3
11B	0.21
C	0.001

Theoretical loss  $\sim 10^{-5}$  per reflection.  
But typically observe several  $10^{-4}$ .

- Hydrogen is everywhere! Especially on surfaces (and diffuses into materials).
- Depolarization losses could be magnetic impurities on surface.

# “Ultracold Neutrons are Kitchen”

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# Sensitivity

Comes from the uncertainty principle:  $\Delta E \Delta t \geq \hbar/2$

Get the “shot noise” limit:  $\sigma(d_n) \geq \frac{\hbar}{2\alpha ET\sqrt{N}}$  (often greater!)

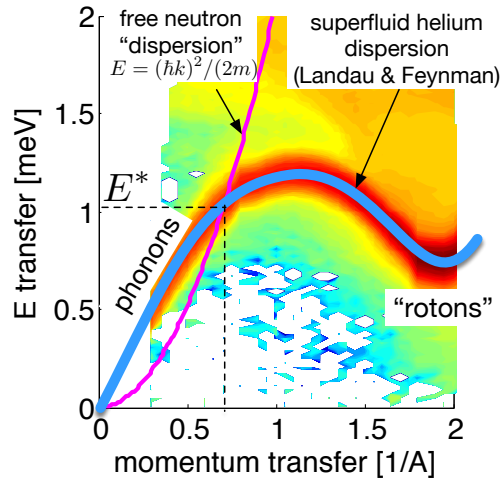
<b><math>\alpha</math> = polarization contrast</b>	~ 80 % → 95% if UCN spin analysis done in-situ
<b><math>E</math> = electric field strength</b>	~ 10 kV/cm in vacuum nEDM cells (due to side walls). Can degrade due to adding (gas) comagnetometer → 70 kV/cm if in insulating fluid
<b><math>N</math> = no. detected neutrons</b>	<ul style="list-style-type: none"> <li>• Current <math>10^{-27}</math> e.cm experiments detect 200,000-300,000 in large 20 - 30 L sized cell. (UCN density in cell ~ 5-10 cm<sup>-3</sup>).</li> <li>• Can't increase volumes much more; false EDM systematic ~ (cell length)<sup>2</sup></li> <li>• Also need to control B-field to <math>\lesssim 1</math> ppb over cell volume</li> </ul>
<b><math>T</math> = interaction/useful time</b>	Restricted by UCN loss time on walls. But even if get longer, need to control B-field field drifts. A lot can happen in 100-1000s.



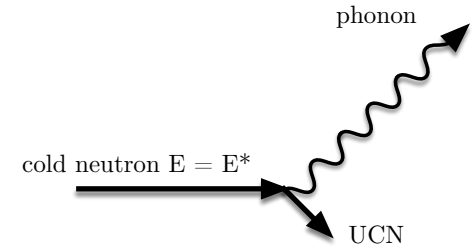
# Ultracold neutrons & Superfluid $^4\text{He}$

- **1 meV cold neutrons (11 K) scatter off phonons in superfluid  $^4\text{He}$  to become UCNs (< 160 neV, “2 mK”)**

Log contour plot of dynamic structure factor of He-II @ 1.2K  
from [Andersen et al. J. Phys. Condens. Matter (1994)]



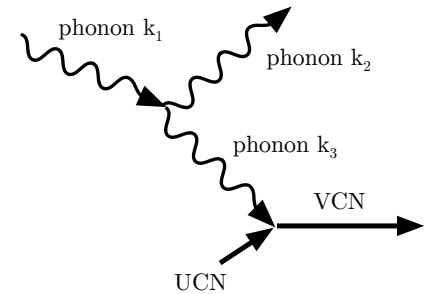
**Down-scattering** from single phonon  
( $\sim T$  independent):



**Dominant UCN loss** is two-phonon scattering:

$$\tau_{\text{up},2\text{-phonon}} = (100 \text{ s K}^7) T^{-7}$$

[Golub & Pendlebury, Physics Letters A (1977)]



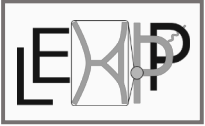
*Experimental studies & demonstrations:*

Golub et al, Z. Physik B (1983); Huffman et al. Nature (2000)

Zimmer et al. PRL (2011); Masuda et al. PRL (2012);

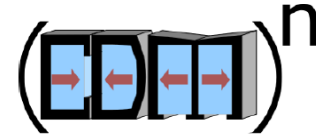
Piegsa, et al. PRC (2014); Schmidt-Wellenburg et al., PRC (2015); Leung et al. PRC (2016);

- “Super-thermal” because UCNs ( $\sim 2$  mK) not in thermal equilibrium with superfluid.
- At  $T = 0.4\text{K}$ , up-scattering (or “thermalization loss time”)  $\tau_{\text{up}} \approx 20$  hours.
- Neutron absorption by  $^4\text{He}$  is zero (when isotopically pure).
- **UCNs can be kept inside superfluid  $^4\text{He}$  and UCNs studied *in-situ***
- Superfluid also **scintillates** at  $\sim 80$  nm (EUV)  $\rightarrow$  used to detect n- $^3\text{He}$  capture events (later)
- $^3\text{He}$  atoms scatter off phonons  $\rightarrow$  mean-free-path  $\sim T^{7.5} \rightarrow$  important for key “false EDM” systematic control

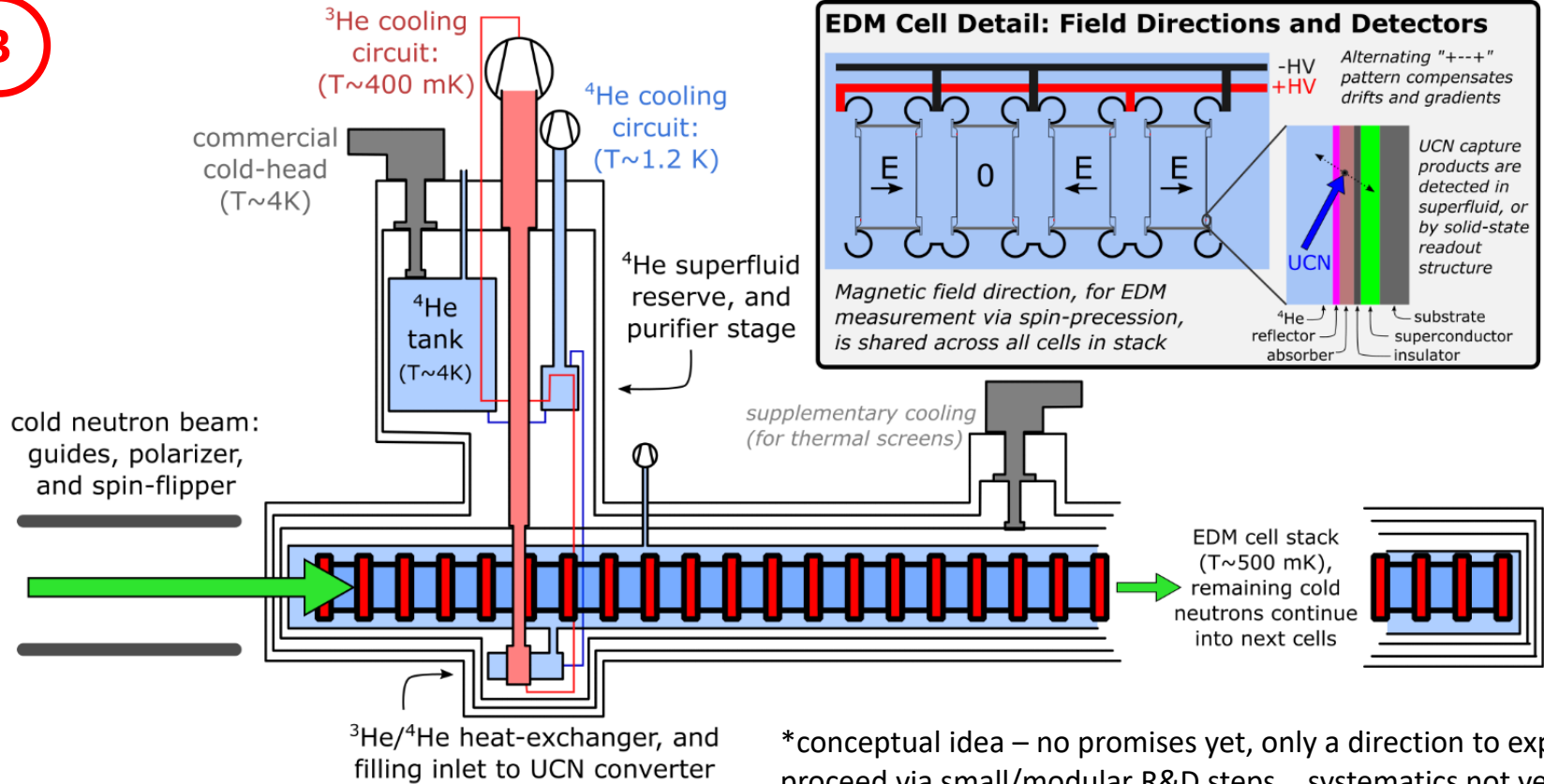


Skyler Degenkolb's idea:

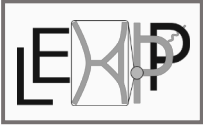
# The next generation\* ... scaling up!



3

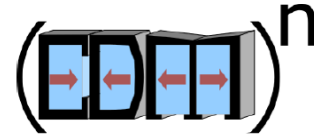


\*conceptual idea – no promises yet, only a direction to explore!  
proceed via small/modular R&D steps ...systematics not yet clear

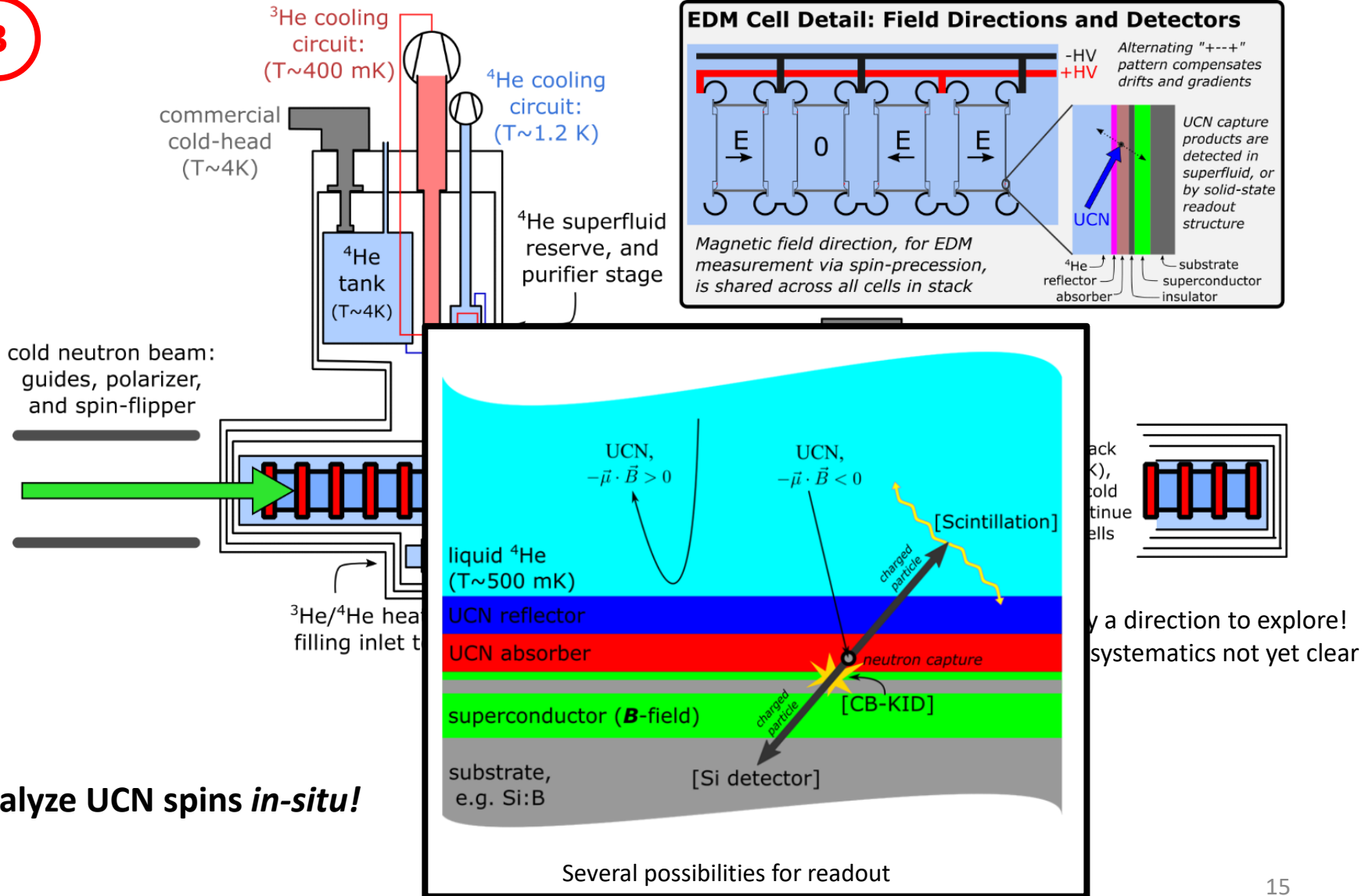


Skyler Degenkolb's idea:

# The next generation\* ... scaling up!



3



# Magnetic trapping of neutrons

**P. R. Huffman<sup>\*†</sup>, C. R. Brome<sup>\*</sup>, J. S. Butterworth<sup>\*‡</sup>, K. J. Coakley<sup>§</sup>,  
M. S. Dewey<sup>†</sup>, S. N. Dzhosyuk<sup>\*</sup>, R. Golub<sup>||</sup>, G. L. Greene<sup>¶</sup>, K. Habicht<sup>||</sup>,  
S. K. Lamoreaux<sup>¶</sup>, C. E. H. Mattoni<sup>\*</sup>, D. N. McKinsey<sup>\*</sup>, F. E. Wietfeldt<sup>\*†</sup>  
& J. M. Doyle<sup>\*</sup>**

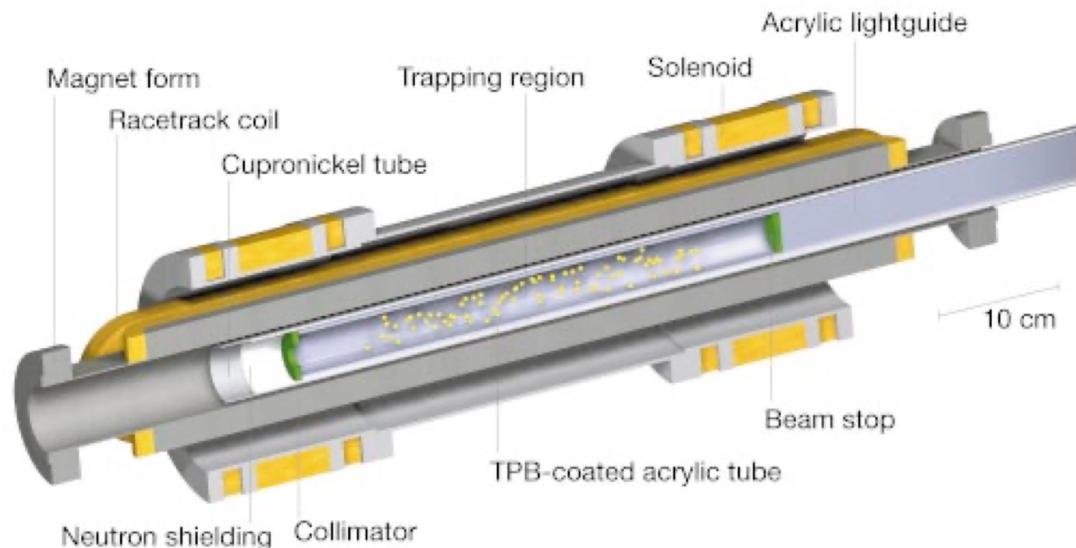
+ TRIUMF TUCAN UCN source

+ Institut Laue-Langevin's SUN-1,  
SUN-2, SuperSUN UCN sources

UCNs produced by cold neutron beam in Superfluid  $^4\text{He}$  and then watch decay via scintillation light ( $n \rightarrow p + e + \nu + 800 \text{ keV}$ ).

With magnetic trap: no wall losses, but still  $\beta$ -decay  $\tau_\beta = 880 \text{ s}$  and a bit more.

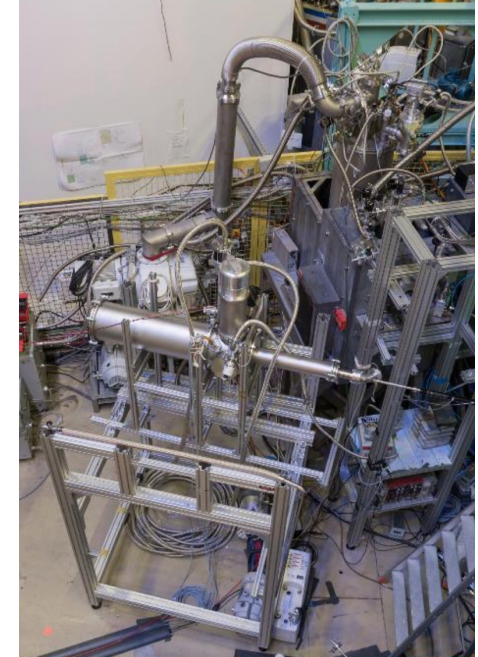
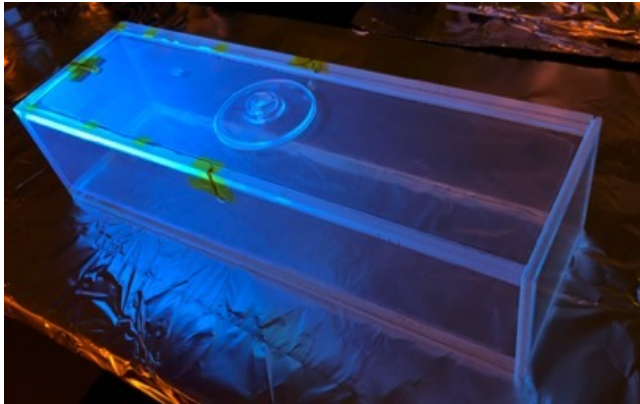
They observed  $\tau_{\text{loss total}} \sim 800 \text{ s}$ .





# Full-sized cryogenic cells demonstrated

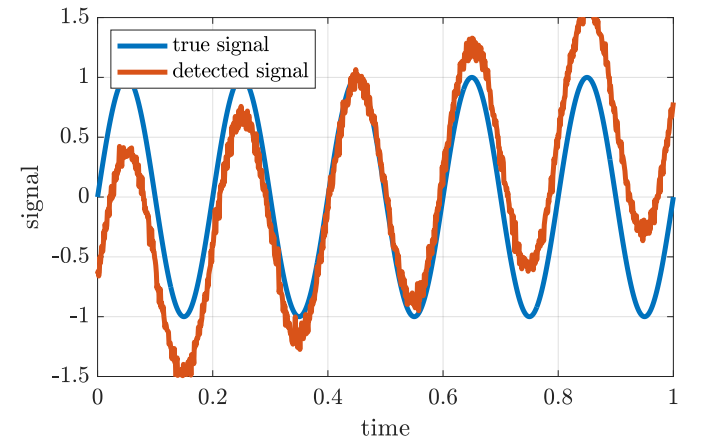
- Full-sized 3-L measurement cell (illuminated by 300 nm UV lamp)
- Deuterated polystyrene + deuterated tetraphenyl butadiene blend (170 neV) coating on acrylic. (Material known to be high-voltage friendly)
- **Measured**  $\tau_{\beta+\text{cell}}$  of  $570 \pm 20$  s at 30 K (Single exponential decay only observed)



- Neulinger et al. (panEDM): Fluoropolymer CYTOP (115 neV) coated 3-L volume cell observed  $\tau_{\beta+\text{cell}} = 560$  s (at 10 K)
- Long  $\tau_{\beta+\text{cell}}$  allows large number of UCNs accumulated per fill for in-situ production. **Polarized UCN density  $\sim 180$  UCN/cm<sup>3</sup>  $\Rightarrow$   $\sim 500,000$  UCNs per cell.**
- **UCN “usefulness” time restricted by  $\tau_{\beta+\text{cell}}$**  (transverse spin coherence times are  $\sim 20,000$  s)
- “Useful” UCN time of **1,000 s** use in a measurement cycle

# “Always measure frequency...” (Rabi? Ramsey? Wieman?)

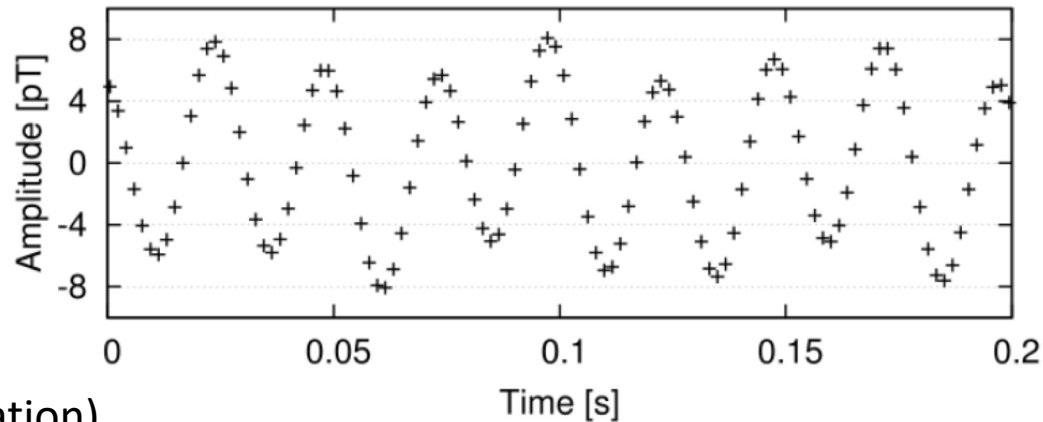
- Control and measurements of time/clocks can be done to a high precision.
  - ✓ used in Ramsey technique
- The power of a frequency measurement: need only **relative change** of a signal **over short time scales**. (Drifts on time scales greater than a few oscillations are suppressed)
- **BUT Ramsey’s technique only measures final neutron phase via. counting neutrons. Disadvantages:**
  - **Blind** to frequency fluctuations
  - Need two (possibly four) repeated Ramsey cycles (with different clock frequencies within the same cell) to determine a single EDM value.
  - Each cycle takes ~100-200s. **A lot can happen in this time.**
  - **Detector efficiency, magnetic field,  $\pi/2$  efficiency and other polarization, UCN source intensity drifts, etc..**
  - Can be corrected but each introduces additional statistics & systematics difficult for  $10^{-28}$  e.cm level



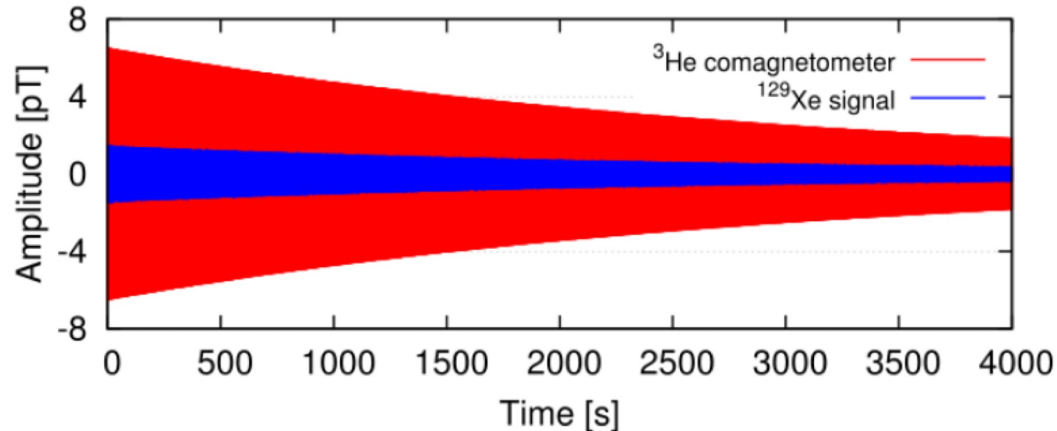
(dummy generated plots with arbitrary units to illustrate point only)

# Free precession decay observation

- Regularly done in NMR, in our field done for atomic systems, for example,  $^3\text{He}$ ,  $^{129}\text{Xe}$ ,  $^{199}\text{Hg}$  (but have to worry about unwanted Stark Shifts)



(From HeXe collaboration)



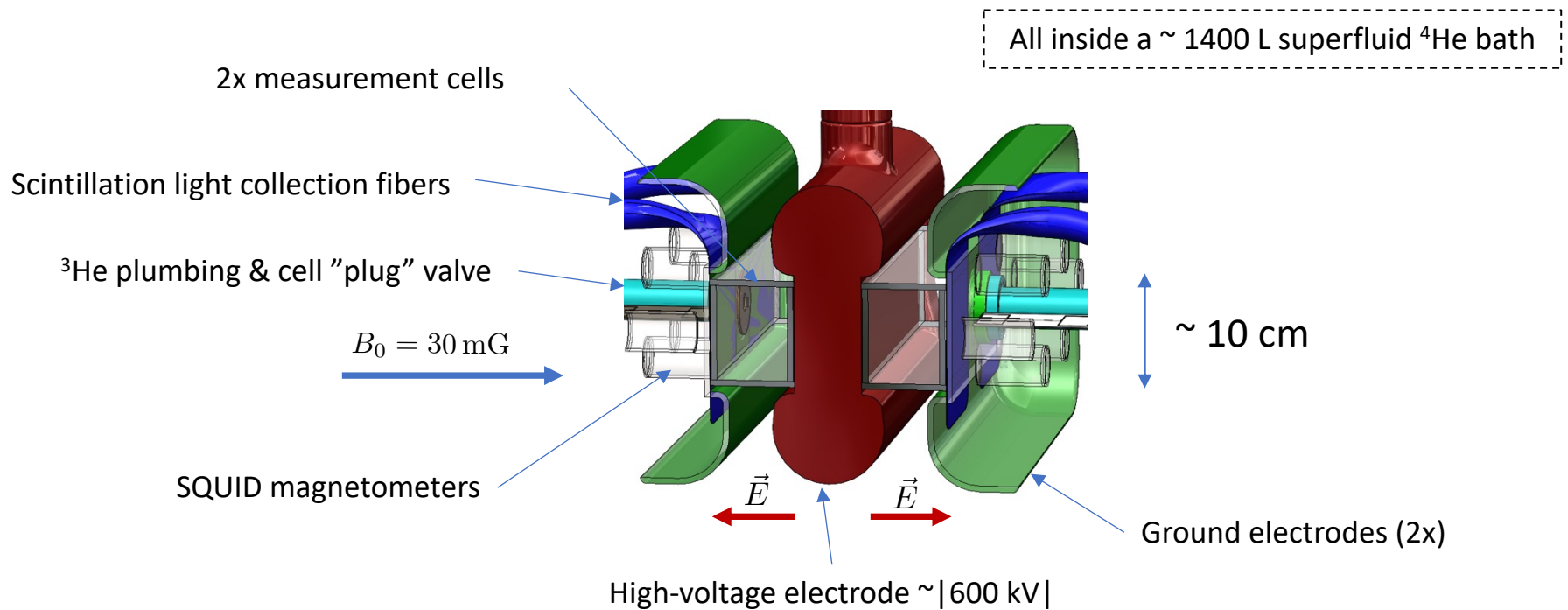
- Can it be done for neutrons?

# Experimental Overview



# nEDMSF: UCNs + superfluid $^4\text{He}$ + polarized $^3\text{He}$

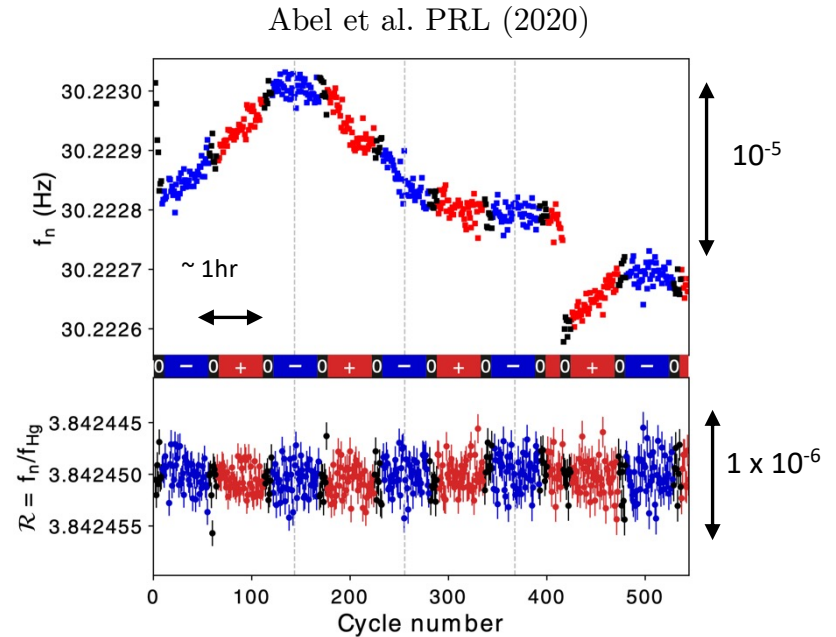
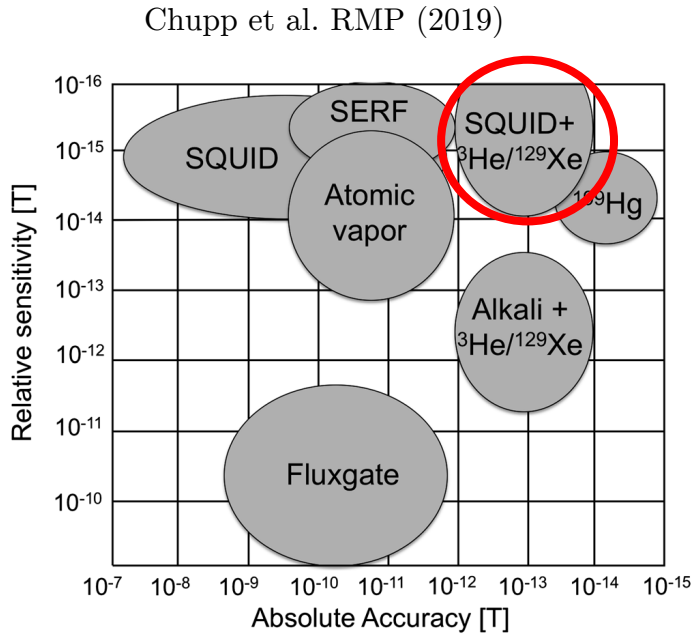
Based on Golub & Lamoreaux, Phys. Rep. (1994)



- **Double cell** setup with E-field relative to B-field opposite in each cell
- **>95% polarized  $^3\text{He}$**  loaded into cell (dissolved in isotopically-pure superfluid  $^4\text{He}$  at 0.4 K).
- $^3\text{He}$  serves as a **comagnetometer** and **UCN spin analyzer** (see later)
- **In-situ super-thermal UCN ( $\sim 100$  neV) production** and accumulation in **superfluid  $^4\text{He}$  with polarized 1 meV cold neutron beam** (direction into the page).
- Cold neutron beam  $\sim 100\%$  polarized  $\rightarrow$  produce  $\sim 100\%$  UCNs polarized.

# Comagnetometry

- A magnetometer species (typically a gas) that “cohabits” the same volume as the UCNs experience.
- Need sufficient comagnetometer atoms for precision; too high can cause electric breakdowns & UCN loss ( $\sim 10^{-6}$  mbar)



(for scale:  $10^{-28}$  e.cm sensitivity  $\rightarrow \sim 10^{-10}$  fractional frequency shift)

- nEDM experiments ultimately measure relative to comagnetometer’s EDM (usually “Schiff suppressed”)
- However, the two species still experiences **different (time & space) averaged magnetic fields** because:
  - UCNs have such low speeds, they “sag” in gravity a few mm.
  - Relativistic  $\vec{E} \times \vec{v}$  motional-field related effects are different. One produces a *false* EDM (most serious systematic error)

# Polarized $^3\text{He}$ as live and in-situ UCN spin analyzer

- Polarized  $^3\text{He}$  gas cells widely used as neutron beam spin analyzers (count survivors)
- Strong **spin-dependent capture cross-section**:  $^3\vec{\text{He}} + \vec{n} \rightarrow \text{p} + ^3\text{H} + 764 \text{ keV}$

Anti-parallel spins:  $\sigma_{\downarrow\uparrow,\text{thermal}} \approx 11 \text{ kb}$

Parallel spins:  $\sigma_{\uparrow\uparrow,\text{thermal}} \lesssim 0.1 \text{ kb}$

- **Capture rate** for polarized UCNs and  $^3\text{He}$  in same volume :

$$\dot{N}_3 = N_{\text{UCN}} \bar{\tau}_3^{-1} (1 - P_n P_3 \cos \theta_{n3})$$

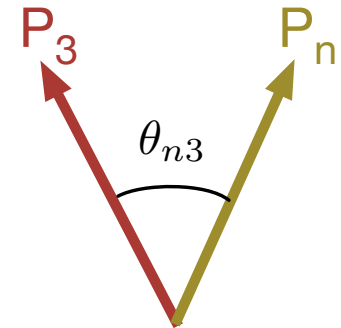
number of UCNs

polarizations

angle between spins

where  $\bar{\tau}_3^{-1} \approx [n_3 \sigma_{\downarrow\uparrow,\text{thermal}} \times (2200 \text{ m/s})]/2$

$^3\text{He}$  number density



or “continuous”

- Detect the 760 keV scintillation as generated  $\rightarrow$  in-situ and live UCN spin analyzer
  - 80 nm scintillation light  $\rightarrow$  blue photons (walls of cells)  $\rightarrow$  wavelength shifting fibers  $\rightarrow$  optical fibers  $\rightarrow$  Silicon Photomultipliers

# Live & in-situ UCN spin analysis

- Can't look **too intensely** or will kill UCNs too quickly



- Want  ${}^3\text{He}$ -n capture rate to be similar to UCN loss time (cell +  $\beta$ ), i.e.  $\bar{\tau}_3 \approx 500$  s
- So want number density  $n_3 \approx 2 \times 10^{12} \text{ cm}^{-3}$  (or  $10^{-10}$  concentration) with near  $P_3 \approx 100\%$   
→ **Achievable with atomic beam source.**
- This  $n_3$  &  $P_3$  of  ${}^3\text{He}$  produces  $\sim 6$  fT fields → **detectable by SQUIDs.**

**${}^3\text{He}$  can be used as comagnetometer AND live & in-situ UCN spin analyzer!**

# False EDM systematic effect

- Recall: nEDM is measured relative to the comagnetometer's EDM. Most effects cancel out with opposite E-field.
- Comagnetometer's EDM suppressed by Schiff screening but **can experience a false EDM**
- From *interaction* between the  $\vec{E} \times \vec{v}$  motional field and magnetic field gradients (“geometric-phase induced false EDM”)

Radius of cell

- “Discovered” in the nEDM field, false EDM for comagnetometer:  $d_{af} = \frac{J\hbar}{4} \left( \frac{\partial B_{0z}}{\partial z} \right) \frac{\gamma^2 R^2}{c^2} \left[ 1 - \frac{\omega_0^2}{\omega_r^2} \right]^{-1}$   
Pendlebury et al. Phys. Rev. C (2004)

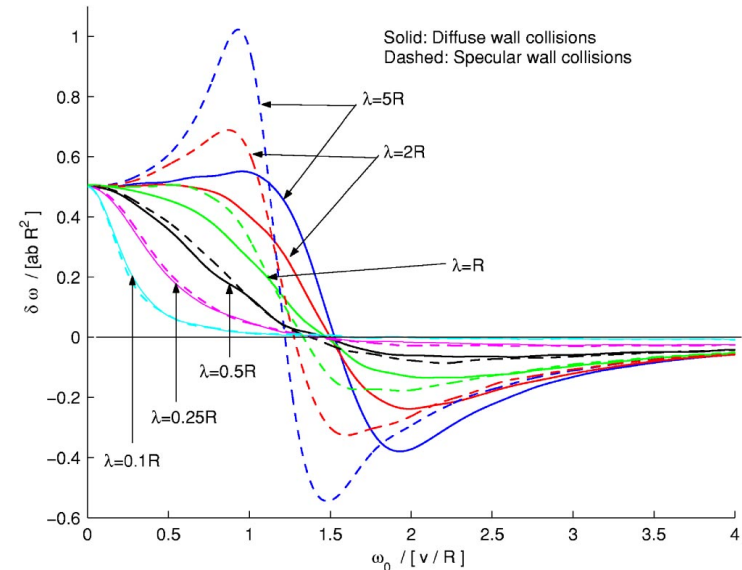
(Simplified cylindrical cell but general relationships hold. Rectangular cell work by Swank & Golub [Phys. Rev. A 93, 062703 \(2016\)](#))

- Note:** The transverse spin coherence time scales as  $R^4$ . Want small cells (with high UCN density)
- Can change  $^3\text{He}$ -phonon scattering mean-free-path by changes in superfluid temperature:

$$\lambda_{3\text{He}} \approx 0.077 \text{ cm} \times \left( \frac{0.45 \text{ K}}{T} \right)^{15/2}$$

**Can tune temperature to make false EDM zero by scanning T!**

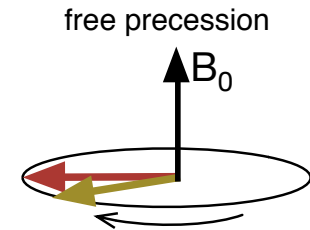
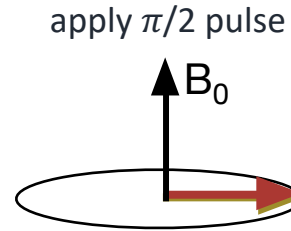
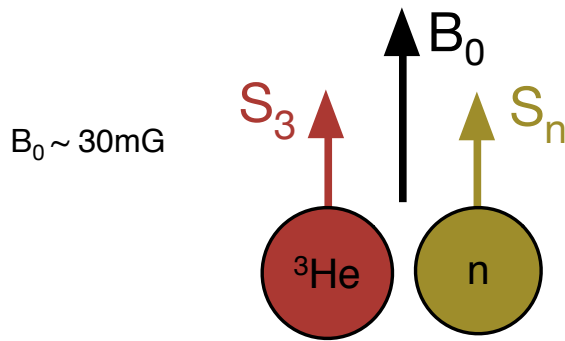
Numerical calculations from Golub:



**Two measurement modes of nEDMSF:  
double free precession & critical dressed spin**



# Double free precession mode



- Time evolution of phase:  $\theta_{3n}(t) = \theta_3(t) - \theta_n(t) = \left[ (\gamma_n - \gamma_3)B_0 \pm \frac{2d_n|E|}{\hbar} \right] t + \phi_0 \equiv \omega_{3n}^\pm t + \phi_0$ ,
- With  $B_0 = 30\text{ mG}$ :  $\gamma_3 B_0 / (2\pi) \approx 100\text{ Hz}$       $\gamma_3 \approx 1.1 \gamma_n$
- Scintillation light oscillation at “beat” frequency:  $|\gamma_n - \gamma_3| B_0 / (2\pi) \approx 10\text{ Hz}$
- The transverse spin coherence time (wall depolarization + gradient depolarization),  $T_2 \sim 10,000\text{ s}$
- Flipping high-voltage electrode often with known sequences to suppress 1<sup>st</sup> order drifts (e.g. + - - + - + + -) and analysis as a “super-asymmetry”

# Free precession scintillation signal

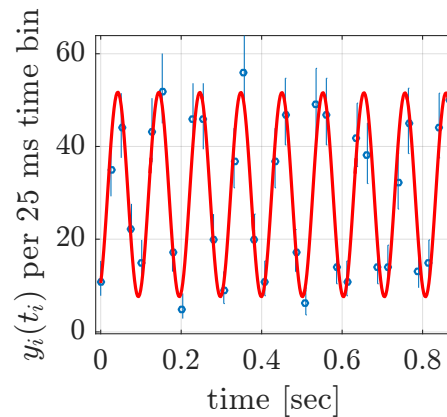
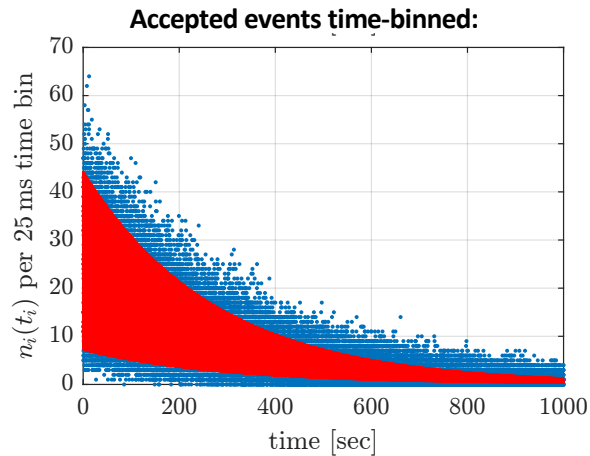
- “accepted” rate of scintillation light events:

$$\dot{N}_{\text{ac-s}}(t) = N_n(t) \left( \frac{\epsilon_\beta}{\tau_\beta} + \frac{\epsilon_3}{\bar{\tau}_3} \left\{ 1 - P_3(t)P_n(t) \cos[\theta_{3n}(t) + \phi_{3n0}] \right\} \right) + R_{\text{BG}}$$

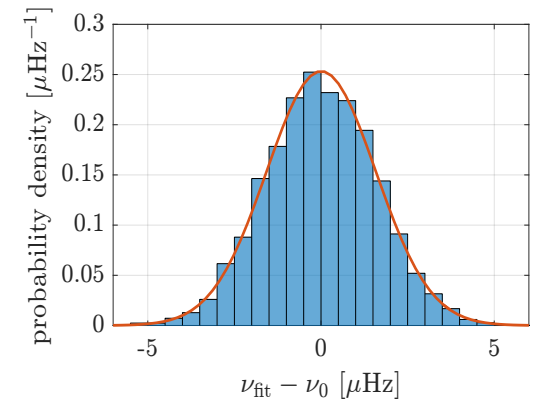
Acceptance probability for  $\beta$ -decay events ( $\sim 0.33$ )  
 Acceptance probability for  $n$ - $^3\text{He}$  capture ( $\sim 0.93$ )  
 Background rate (could be time-dependent)  
 no. UCNs in cell  
 recall: 500 s  
 polarizations  
 $P_3(t)P_n(t) = P_{30}P_{n0} \exp(-t/T_{2,\text{tot}}) \approx (0.98)(0.98) \exp(-t/[10,000 \text{ s}])$

$$N_n(t) = \int_0^{E_{\text{max}}} dE n_{n0}(E) \exp \left[ -\frac{t}{\tau_{\text{cell}}(E)} - \frac{t}{\bar{\tau}_\beta} + \frac{P_n(t)P_3(t)}{\bar{\tau}_3} \int_0^t \cos \phi_{3n}(t') dt' \right]$$

UCN spectrum  
 Oscillating term due to previous  $n$ - $^3\text{He}$  absorption



Repeat generation & fit:



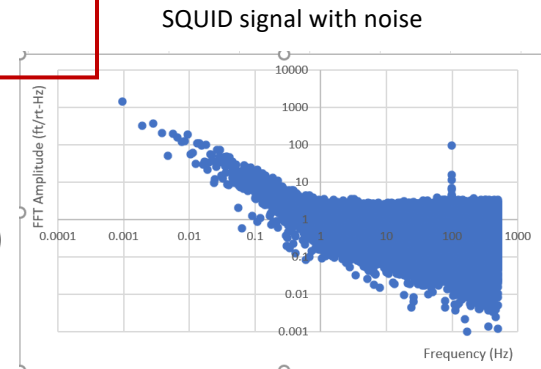
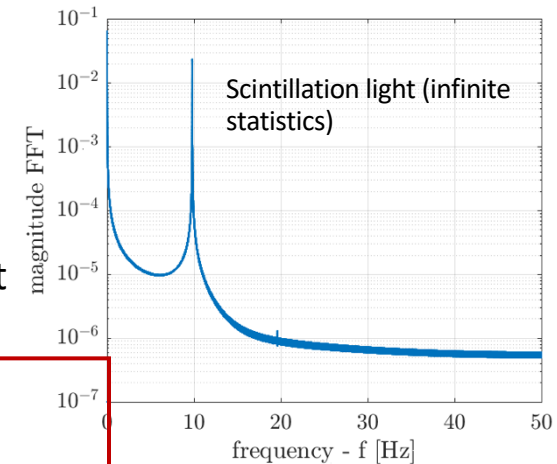
- Continuously measuring the UCN phase (relative to  $^3\text{He}$ )  $\rightarrow$  continuous frequency measurement (via derivative)!
- 300 live days of running (expected to take 3 years), get  $1\sigma$  nEDM error =  $3 \times 10^{-28} \text{ e.cm}$

# Data analysis simulations

- Neutron decay  $\beta$ -asymmetry
- **Spatial-variation scintillation light** detection efficiency
- **Oscillation in  $N_n(t)$**  due to history of  $n$ - $^3\text{He}$  absorption
- **Reduced parameter "contrast" fitting** to handle UCN energy-dependent wall loss

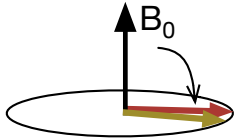
- Generation of scintillation light data with **magnetic field drifts**
- Generation of SQUID  $^3\text{He}$  signal with **noise and drifts** (UKy student: Mojtaba Behzadipour)
- **Simultaneously fitting SQUID+  $^3\text{He}$  magnetometer signal and scintillation light** signal with global likelihood parameter (UKy student: Mojtaba Behzadipour)
- Fit **temporal field drifts** with **orthogonal polynomials**

- **Particle-by-particle neutron** scintillation data generation code
- **Magnetic field noise** in spin-dressing mode
- **Novel** spin dressing field modulation modes (Caltech grad: Raymond Tat)
- UCN spin-tracking on **Graphics Processing Units** (NCSU/Caltech/ORNL: Morano, Tat, Matthews)
- UCN center-of-mass **gravitational offset time-evolution**

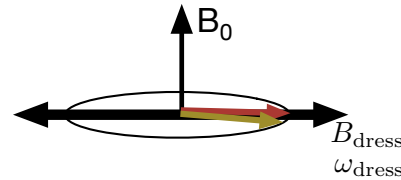


# Dressed-spin mode

apply  $\pi/2$  pulse



apply strong off-resonance dressing field perpendicular to  $B_0$  to alter precession of **both species**



Cohen-Tannoudji:

$$\hat{H} = -\gamma B_0 \hat{S}_z + \hbar \omega_d \hat{a}^\dagger \hat{a} + \lambda \hat{S}_x (\hat{a} + \hat{a}^\dagger)$$

$$\lambda = \gamma B_d / 2\sqrt{\hbar n}$$

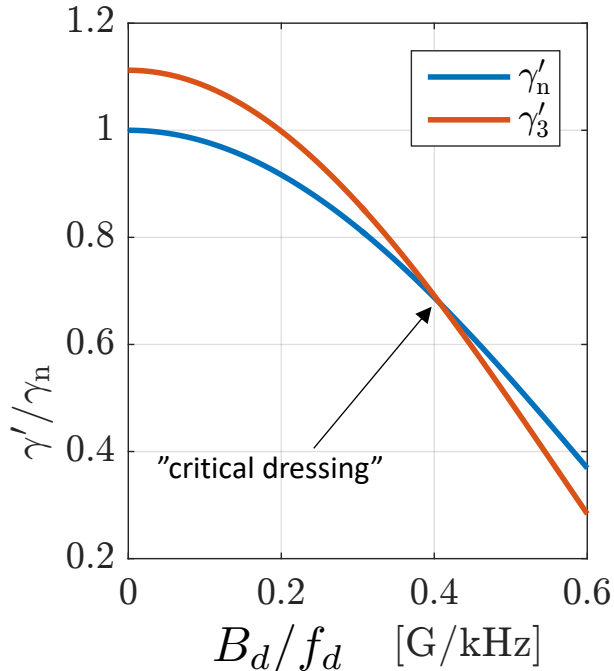
In the limit  $B_{\text{dress}} \gg B_0$

effective gyromagnetic ratio

0<sup>th</sup> order Bessel function

$$\gamma' = \gamma J_0 \left( \frac{\gamma B_{\text{dress}}}{\omega_{\text{dress}}} \right)$$

original



- Specific value of  $B_d/\nu_d$  can make  $\gamma'_3 = \gamma'_n$
- For instance, if  $B_d = 1 \text{ G}$  is chosen, then  $f_d \approx 2.5 \text{ kHz}$  is needed
- If above or below critical dressing condition, then **can make neutrons effectively precess faster or slower than  $^3\text{He}$**  as needed.

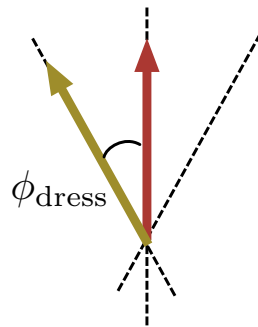
# Critical modulated dressed-spin mode

- The effect of **neutron EDM** with spin dressing:  $\gamma'_n B_0 \pm \frac{2ed_n E J_0 (\gamma_n B_{\text{dress}} / \omega_{\text{dress}})}{\hbar}$
- Example of modulation with “**square wave pulses**”. (Other modes possible.)

**Critical dressing field to sit at a fixed  $\phi_{\text{dress}}$**



(In rest frame of  $^3\text{He}$  spin,  
 $B_0$  coming out of screen)



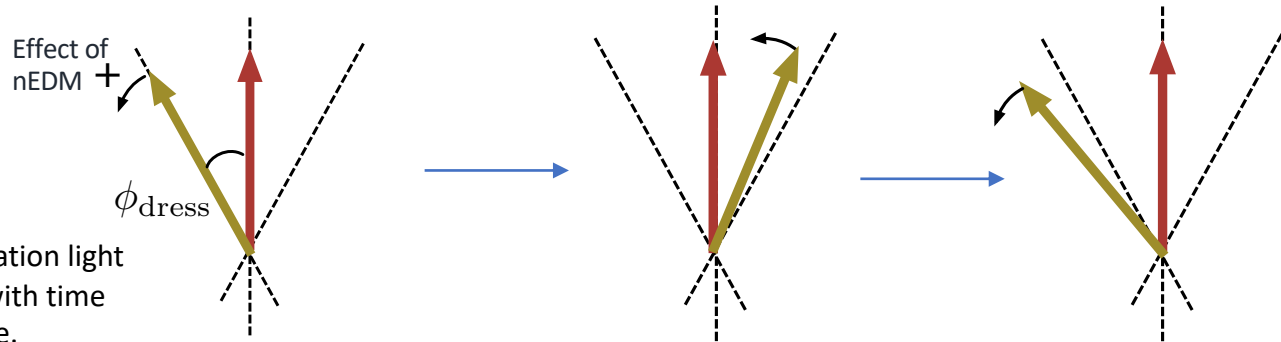
# Critical modulated dressed-spin mode

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**Critical dressing field to sit at a fixed  $\phi_{\text{dress}}$**



(In rest frame of  $^3\text{He}$  spin,  $B_0$  coming out of screen)



If there's a nEDM scintillation light increases or decreases with time depending on E and cycle.

- Can treat each pair as **asymmetry measurement**:  $A_d(t_i) = \frac{N_- - N_+}{N_- + N_+}$  (Ahmed et al. JINST 2019)

**Predicted sensitivity:** 300 live days of data (e.g. 3 years running)  $\sigma(d_n) = 1.7 \times 10^{-28} e \cdot \text{cm}$

- Less sensitive to static field in-homogeneities. Quality and control dressing field becomes main systematic.

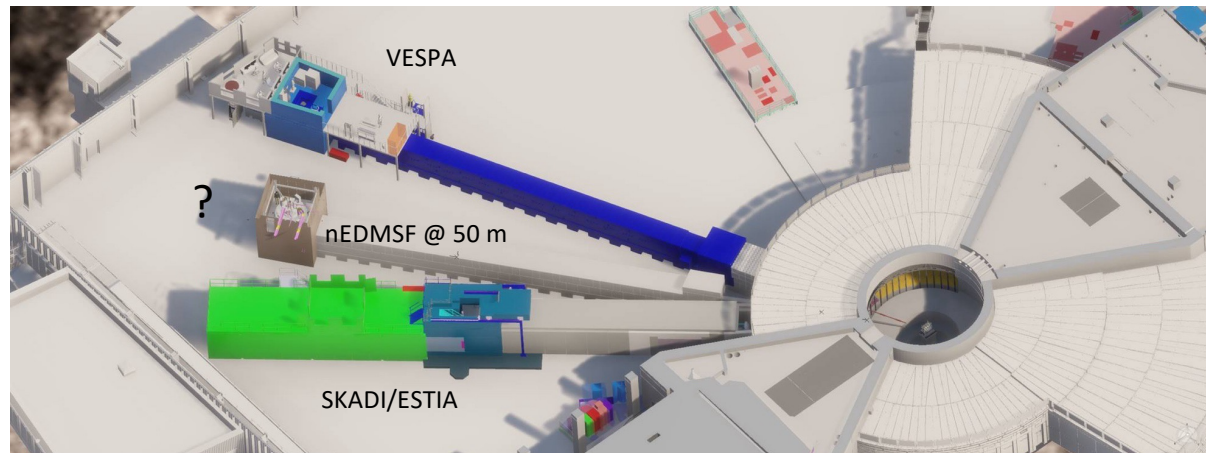


# nEDM-Superfluid Status

Early tests @ ILL?  
Small test apparatus



Strong support from ILL & ESS Directors.



# Hardware Status: Magnet System

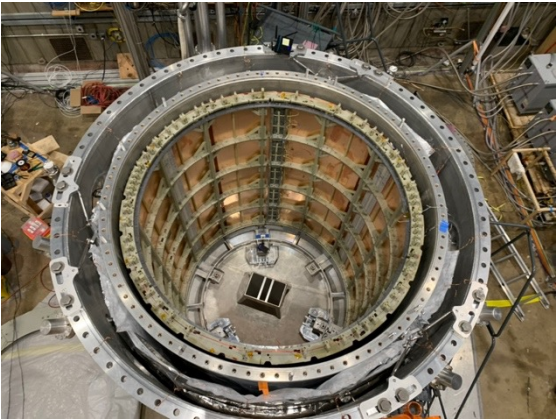
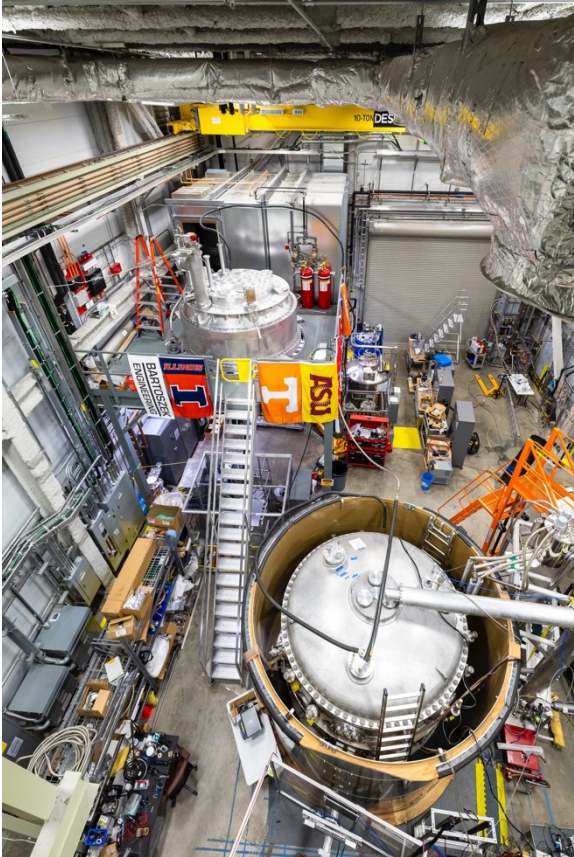


$B_0 \cos\theta$  coil



Magnet system closed for cold test at ORNL

## Oak Ridge National Lab



Completed  $B_0$  coil, sc shield, vacuum chamber at ORNL



Completed magnetic shielded room at Imedco: Hägendorf, CH



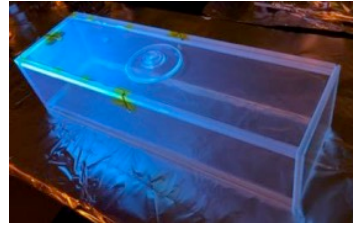
# Hardware Status: Central Detector System



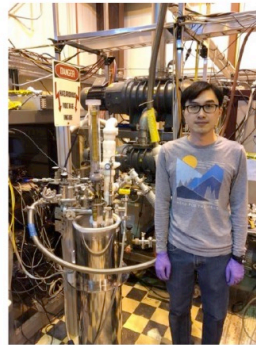
Completed cryostat top flange, tail for CDS testing



Half-scale HV test (T=0.4 K)



Measurement cell



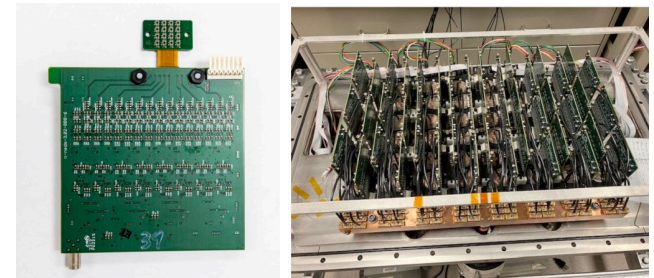
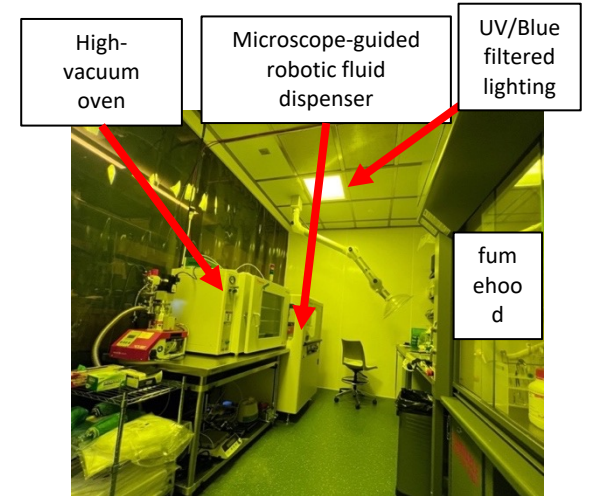
Small scale HV  
N. S. Phan, et al *J. Appl. Phys.*  
129, 083301 (2021)



OVC composite vessel



Full-size stainless Cavallo electrodes

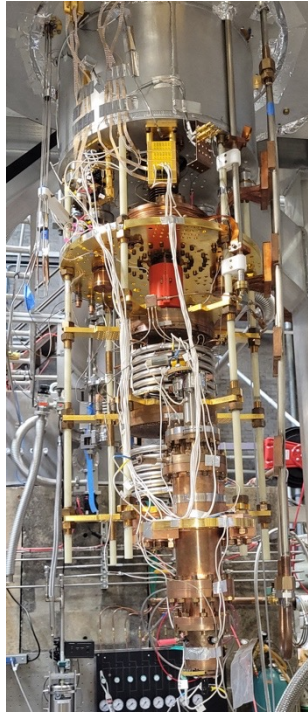


SiPM light collection

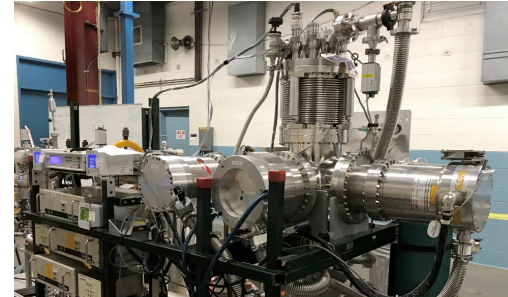
# Hardware Status: Polarized $^3\text{He}$ System



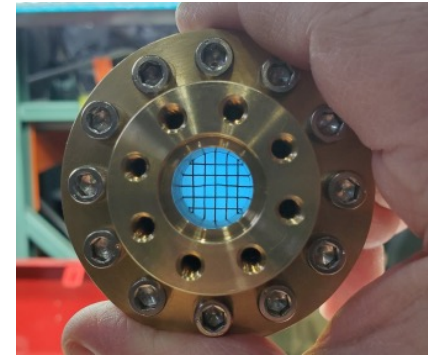
Dilution refrigerator/test cryostat



Dilution refrigerator + film burner  
(75 mW at 250 mK)



Atomic beam source



Nuclepore filter: magnetic particles

- To do: injection system commissioning

# Summary

- The live and in-situ UCN spin analysis allows **to control drifts** in short time scales in both free precession & critical dressed-spin mode.
- Sparked some different ideas for your favorite EDM system
- For the neutron: **cryogenic UCN +  $^3\text{He}$  + superfluid** scheme offers many advantages to **reach  $10^{-28}$  e.cm (especially systematics)**
  - In-situ produced UCN in **small cell with high-density and long storage times**
  - Supports high **electric field, SQUIDs** and **superconducting magnetic shielding**
  - Can vary motion of our  $^3\text{He}$  magnetometer with **small T changes to study key false EDM systematic effect**
  - **Two measurement modes** with different systematic effects for **self-checking our own results**

