# AMIDAS (Package and Website) for Direct Dark Matter Detection Experiments

### Chung-Lin Shan

Department of Physics, National Cheng Kung University

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in collaboration with M. Drees and M. Kakizaki based on JCAP 0706 011, JCAP 0806 012, 1103.0481, 1103.0482



### Model-independent data analyses

Motivation

Reconstruction of the WIMP velocity distribution

Determination of the WIMP mass

Estimation of the SI WIMP-nucleon coupling

Determinations of ratios of WIMP-nucleon cross sections

AMIDAS package and website

Summary



### Model-independent data analyses



## Motivation Motivation

O Differential event rate for elastic WIMP-nucleus scattering

$$\frac{dR}{dQ} = AF^{2}(Q) \int_{v_{\min}}^{v_{\max}} \left[ \frac{f_{1}(v)}{v} \right] dv$$

Here

$$v_{\min} = \alpha \sqrt{Q}$$

is the minimal incoming velocity of incident WIMPs that can deposit the recoil energy Q in the detector.

$$\mathcal{A} \equiv \frac{\rho_0 \sigma_0}{2m_\chi m_{\rm r,N}^2} \qquad \qquad \alpha \equiv \sqrt{\frac{m_{\rm N}}{2m_{\rm r,N}^2}} \qquad \qquad m_{\rm r,N} = \frac{m_\chi m_{\rm N}}{m_\chi + m_{\rm N}}$$

 $\rho_0$ : WIMP density near the Earth

 $\sigma_0$ : total cross section ignoring the form factor suppression

F(Q): elastic nuclear form factor

 $f_1(v)$ : one-dimensional velocity distribution of halo WIMPs



### Reconstruction of the WIMP velocity distribution

Normalized one-dimensional WIMP velocity distribution function

$$\begin{split} f_1(v) &= \mathcal{N} \left\{ -2Q \cdot \frac{d}{dQ} \left[ \frac{1}{F^2(Q)} \left( \frac{dR}{dQ} \right) \right] \right\}_{Q = v^2/\alpha^2} \\ \mathcal{N} &= \frac{2}{\alpha} \left\{ \int_0^\infty \frac{1}{\sqrt{Q}} \left[ \frac{1}{F^2(Q)} \left( \frac{dR}{dQ} \right) \right] dQ \right\}^{-1} \end{split}$$

Moments of the velocity distribution function

$$\begin{split} \langle v^n \rangle &= \mathcal{N}(Q_{\text{thre}}) \left( \frac{\alpha^{n+1}}{2} \right) \left[ \frac{2Q_{\text{thre}}^{(n+1)/2}}{F^2(Q_{\text{thre}})} \left( \frac{dR}{dQ} \right)_{Q=Q_{\text{thre}}} + (n+1)I_n(Q_{\text{thre}}) \right] \\ \mathcal{N}(Q_{\text{thre}}) &= \frac{2}{\alpha} \left[ \frac{2Q_{\text{thre}}^{1/2}}{F^2(Q_{\text{thre}})} \left( \frac{dR}{dQ} \right)_{Q=Q_{\text{thre}}} + I_0(Q_{\text{thre}}) \right]^{-1} \\ I_n(Q_{\text{thre}}) &= \int_0^\infty Q^{(n-1)/2} \left[ \frac{1}{F^2(Q)} \left( \frac{dR}{dQ} \right) \right] dQ \end{split}$$

[M. Drees and CLS, JCAP 0706, 011]

### Reconstruction of the WIMP velocity distribution

O Ansatz: the measured recoil spectrum in the nth Q-bin

$$\left(\frac{dR}{dQ}\right)_{\text{expt, }Q\simeq Q_n}\equiv r_n\,e^{k_n(Q-Q_{s,n})} \qquad \qquad r_n\equiv \frac{N_n}{b_n}$$

• Logarithmic slope and shifted point in the *n*th *Q*-bin

$$\overline{Q - Q_n}|_n \equiv \frac{1}{N_n} \sum_{i=1}^{N_n} (Q_{n,i} - Q_n) = \left(\frac{b_n}{2}\right) \coth\left(\frac{k_n b_n}{2}\right) - \frac{1}{k_n}$$

$$Q_{s,n} = Q_n + \frac{1}{k_n} \ln\left[\frac{\sinh(k_n b_n/2)}{k_n b_n/2}\right]$$

Reconstructing the one-dimensional WIMP velocity distribution

$$f_{1}(v_{s,n}) = \mathcal{N} \left[ \frac{2Q_{s,n}r_{n}}{F^{2}(Q_{s,n})} \right] \left[ \frac{d}{dQ} \ln F^{2}(Q) \right|_{Q=Q_{s,n}} - k_{n}$$

$$\mathcal{N} = \frac{2}{\alpha} \left[ \sum \frac{1}{\sqrt{Q_{a}} F^{2}(Q_{a})} \right]^{-1}$$

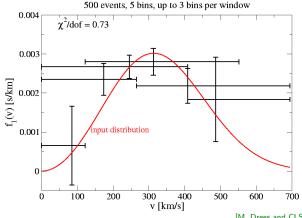
$$v_{s,n} = \alpha \sqrt{Q_{s,n}}$$

[M. Drees and CLS, JCAP 0706, 011]



### Reconstruction of the WIMP velocity distribution

O Reconstructed  $f_{1,rec}(v_{s,n})$  (76Ge, 500 events, 5 bins, up to 3 bins per window)





### Determination of the WIMP mass

Estimating the moments of the WIMP velocity distribution

$$\begin{split} \langle v^n \rangle &= \alpha^n \left[ \frac{2Q_{\min}^{1/2} r_{\min}}{F^2(Q_{\min})} + I_0 \right]^{-1} \left[ \frac{2Q_{\min}^{(n+1)/2} r_{\min}}{F^2(Q_{\min})} + (n+1)I_n \right] \\ I_n &= \sum_a \frac{Q_a^{(n-1)/2}}{F^2(Q_a)} \qquad \qquad r_{\min} = \left( \frac{dR}{dQ} \right)_{\text{expt}, \ Q = Q_{\min}} = r_1 \, e^{k_1(Q_{\min} - Q_{s,1})} \end{split}$$

[M. Drees and CLS, JCAP 0706, 011]

Determining the WIMP mass

$$\begin{split} m_{\chi}\big|_{\left\langle v^{n}\right\rangle} &= \frac{\sqrt{m_{\chi}m_{\gamma}} - m_{\chi}\mathcal{R}_{n}}{\mathcal{R}_{n} - \sqrt{m_{\chi}/m_{\gamma}}} \\ \mathcal{R}_{n} &= \left[\frac{2Q_{\min,\chi}^{(n+1)/2}r_{\min,\chi}/F_{\chi}^{2}(Q_{\min,\chi}) + (n+1)l_{n,\chi}}{2Q_{\min,\chi}^{1/2}r_{\min,\chi}/F_{\chi}^{2}(Q_{\min,\chi}) + l_{0,\chi}}\right]^{1/n} (X \longrightarrow Y)^{-1} \qquad (n \neq 0) \end{split}$$
 [CLS and M. Drees, arXiv:0710.429]

With the assumption of a dominant SI WIMP-nucleus interaction

$$m_{X}|_{\sigma} = \frac{(m_{X}/m_{Y})^{5/2} m_{Y} - m_{X} \mathcal{R}_{\sigma}}{\mathcal{R}_{\sigma} - (m_{X}/m_{Y})^{5/2}} \qquad \qquad \mathcal{R}_{\sigma} = \frac{\mathcal{E}_{Y}}{\mathcal{E}_{X}} \left[ \frac{2Q_{\min,X}^{1/2} r_{\min,X} / F_{X}^{2}(Q_{\min,X}) + l_{0,X}}{2Q_{\min,Y}^{1/2} r_{\min,X} / F_{Y}^{2}(Q_{\min,Y}) + l_{0,Y}} \right]$$

[M. Drees and CLS, JCAP 0806, 012]



### Determination of the WIMP mass

 $\circ$   $\chi^2$ -fitting

$$\chi^{2}(m_{\chi}) = \sum_{i,j} \left(f_{i,\chi} - f_{i,\gamma}\right) C_{ij}^{-1} \left(f_{j,\chi} - f_{j,\gamma}\right)$$

where

$$f_{i,X} = \alpha_X^i \left[ \frac{2Q_{\min,X}^{(i+1)/2} r_{\min,X} / F_X^2(Q_{\min,X}) + (i+1)I_{i,X}}{2Q_{\min,X}^{1/2} r_{\min,X} / F_X^2(Q_{\min,X}) + I_{0,X}} \right] \left( \frac{1}{300 \text{ km/s}} \right)^i$$

$$f_{n_{\text{max}}+1,X} = \mathcal{E}_X \left[ \frac{A_X^2}{2Q_{\text{min},X}^{1/2} r_{\text{min},X} / F_X^2(Q_{\text{min},X}) + I_{0,X}} \right] \left( \frac{\sqrt{m_X}}{m_X + m_X} \right)$$

$$\mathcal{C}_{ij} = \mathsf{cov}\left(f_{i,X}, f_{j,X}\right) + \mathsf{cov}\left(f_{i,Y}, f_{j,Y}\right)$$

 $\bigcirc$  Algorithmic  $Q_{\sf max}$  matching

$$Q_{\mathsf{max},Y} = \left(\frac{lpha_X}{lpha_Y}\right)^2 Q_{\mathsf{max},X}$$

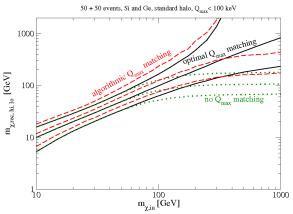
$$\left( v_{\text{cut}} = \alpha \sqrt{Q_{\text{max}}} \right)$$

[M. Drees and CLS, JCAP 0806, 012]



### Determination of the WIMP mass

O Reconstructed  $m_{\chi, \rm rec}$  (28Si + <sup>76</sup>Ge,  $Q_{\rm max}$  < 100 keV, 2 × 50 events)



[M. Drees and CLS, JCAP 0806, 012]



### Estimation of the SI WIMP-nucleon coupling

Spin-independent (SI) WIMP-nucleus cross section

$$\begin{split} \sigma_0^{\text{SI}} &= \left(\frac{4}{\pi}\right) m_{r,N}^2 \big[ Z f_p + (A-Z) f_n \big]^2 \simeq \left(\frac{4}{\pi}\right) m_{r,N}^2 A^2 |f_p|^2 = A^2 \left(\frac{m_{r,N}}{m_{r,p}}\right)^2 \sigma_{\chi p}^{\text{SI}} \\ \sigma_{\chi p}^{\text{SI}} &= \left(\frac{4}{\pi}\right) m_{r,p}^2 |f_p|^2 \end{split}$$

 $f_p$ ,  $f_n$ : effective SI WIMP-proton/neutron couplings

Estimating the SI WIMP-nucleon coupling

$$|f_{\rm p}|^2 = \frac{1}{\rho_0} \left[ \frac{\pi}{4\sqrt{2}} \left( \frac{1}{\mathcal{E}_Z A_Z^2 \sqrt{m_Z}} \right) \right] \left[ \frac{2Q_{\rm min,Z}^{1/2} f_{\rm min,Z}}{F_Z^2 (Q_{\rm min,Z})} + I_{0,Z} \right] (m_{\rm X} + m_Z)$$

[M. Drees and CLS, arXiv:0809.2441]



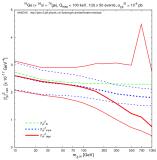
### Estimation of the SI WIMP-nucleon coupling

Estimating the SI WIMP-nucleon coupling

$$|f_{\rm p}|^2 = \frac{1}{\rho_0} \left[ \frac{\pi}{4\sqrt{2}} \left( \frac{1}{\mathcal{E}_Z A_Z^2 \sqrt{m_Z}} \right) \right] \left[ \frac{2Q_{\rm min,Z}^{1/2} r_{\rm min,Z}}{F_Z^2 (Q_{\rm min,Z})} + I_{0,Z} \right] (m_\chi + m_Z)$$

 $[\mathsf{M.\ Drees\ and\ CLS},\ \mathsf{arXiv}{:}0809.2441]$ 

 $O |f_p|^2 (^{76} ext{Ge} (+^{28} ext{Si} +^{76} ext{Ge}), Q_{ ext{max}} < 100 \text{ keV}, \sigma_{\chi p}^{ ext{SI}} = 10^{-8} \text{ pb}, 1(3) imes 50 \text{ events})$ 



[CLS, arXiv:1103.0481, submitted to JCAP]



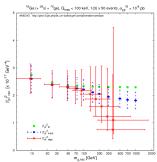
### Estimation of the SI WIMP-nucleon coupling

Estimating the SI WIMP-nucleon coupling

$$|f_{\rm p}|^2 = rac{1}{
ho_0} \left[ rac{\pi}{4\sqrt{2}} \left( rac{1}{\mathcal{E}_Z A_Z^2 \sqrt{m_Z}} 
ight) 
ight] \left[ rac{2Q_{
m min,Z}^{1/2} r_{
m min,Z}}{F_Z^2 (Q_{
m min,Z})} + I_{0,Z} 
ight] (m_\chi + m_Z)$$

[M. Drees and CLS, arXiv:0809.2441]

 $\bigcirc \ |f_{\rm p}|^2 \ {\rm vs.} \ m_\chi \ (^{76}{\rm Ge} \, (+^{28}{\rm Si} +^{76}{\rm Ge}), \ Q_{\rm max} < 100 \ {\rm keV}, \ \sigma_{\chi \rm p}^{\rm SI} = 10^{-8} \ {\rm pb}, \ 1(3) \times 50 \ {\rm events})$ 



[CLS, arXiv:1103.0481, submitted to JCAP]





Spin-dependent (SD) WIMP-nucleus cross section

$$\begin{split} \sigma_0^{\rm SD} &= \left(\frac{32}{\pi}\right) \, G_F^2 \, m_{\rm r,N}^2 \left(\frac{J+1}{J}\right) \left[\langle S_{\rm p} \rangle a_{\rm p} + \langle S_{\rm n} \rangle a_{\rm n}\right]^2 \\ \sigma_{\chi \rm p/n}^{\rm SD} &= \left(\frac{32}{\pi}\right) \, G_F^2 \, m_{\rm r,p/n}^2 \cdot \left(\frac{3}{4}\right) a_{\rm p/n}^2 \end{split}$$

J: total nuclear spin

 $\langle S_p \rangle$ ,  $\langle S_n \rangle$ : expectation values of the proton/neutron group spin

 $a_p$ ,  $a_n$ : effective SD WIMP-proton/neutron couplings

Determining the ratio of two SD WIMP-nucleon couplings

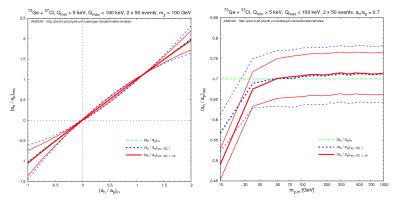
$$\begin{pmatrix} \frac{a_{n}}{a_{p}} \end{pmatrix}_{\pm,n}^{SD} = -\frac{\langle S_{p} \rangle_{X} \pm \langle S_{p} \rangle_{Y} \mathcal{R}_{J,n}}{\langle S_{n} \rangle_{X} \pm \langle S_{n} \rangle_{Y} \mathcal{R}_{J,n}} 
\mathcal{R}_{J,n} \equiv \left[ \left( \frac{J_{X}}{J_{Y}+1} \right) \left( \frac{J_{Y}+1}{J_{Y}} \right) \frac{\mathcal{R}_{\sigma}}{\mathcal{R}_{n}} \right]^{1/2} \qquad (n \neq 0)$$

[M. Drees and CLS, arXiv:0903.3300]



### Determination of the ratio of two SD WIMP-nucleon couplings

O Reconstructed  $(a_{\rm n}/a_{\rm p})_{\rm rec,1}^{\rm SD}$  $(^{73}{\rm Ge}+^{37}{\rm Cl},~Q_{\rm min}>5~{\rm keV},~Q_{\rm max}<100~{\rm keV},~2\times50~{\rm events},$   $m_{\chi}=100~{\rm GeV}~{\rm or}~a_{\rm n}/a_{\rm p}=0.7)$ 

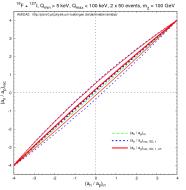


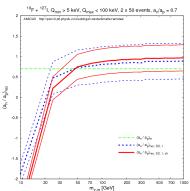
[M. Drees and CLS, arXiv:0903.3300; CLS, arXiv:1103.0482, submitted to JCAP]



### Determination of the ratio of two SD WIMP-nucleon couplings

O Reconstructed  $(a_{\rm n}/a_{\rm p})_{\rm rec,1}^{\rm SD}$  $(^{19}{\rm F}+^{127}{\rm I},~Q_{\rm min}>5~{\rm keV},~Q_{\rm max}<100~{\rm keV},~2\times50~{\rm events},$   $m_{\rm x}=100~{\rm GeV}~{\rm or}~a_{\rm n}/a_{\rm p}=0.7)$ 





[CLS, arXiv:1103.0482, submitted to JCAP]



O Differential rate for a combination of the SI and SD cross sections

$$\begin{split} \left(\frac{dR}{dQ}\right)_{\text{expt, }Q=Q_{\text{min}}} &= \mathcal{E}\left(\frac{\rho_0\sigma_0^{\text{SI}}}{2m_\chi m_{\text{r,N}}^2}\right) \left[F_{\text{SI}}^2(Q) + \left(\frac{\sigma_{\chi p}^{\text{SD}}}{\sigma_{\chi p}^{\text{SI}}}\right) \mathcal{C}_p F_{\text{SD}}^2(Q)\right] \int_{\nu_{\text{min}}}^{\nu_{\text{max}}} \left[\frac{f_1(\nu)}{\nu}\right] d\nu \\ \mathcal{C}_p &\equiv \frac{4}{2}\left(\frac{J+1}{\nu}\right) \left[\frac{\langle S_p \rangle + (a_n/a_p)\langle S_n \rangle}{a}\right]^2 \end{split}$$

Determining the ratio of two WIMP-proton cross sections

$$\begin{split} & \frac{\sigma_{XP}^{\text{SD}}}{\sigma_{XP}^{\text{SI}}} = \frac{F_{\text{SI},Y}^2(Q_{\text{min},Y})\mathcal{R}_{m,XY} - F_{\text{SI},X}^2(Q_{\text{min},X})}{\mathcal{C}_{\text{p},X}F_{\text{SD},X}^2(Q_{\text{min},X}) - \mathcal{C}_{\text{p},Y}F_{\text{SD},Y}^2(Q_{\text{min},Y})\mathcal{R}_{m,XY}} \\ & \mathcal{R}_{m,XY} \equiv \left(\frac{r_{\text{min},X}}{\mathcal{E}_X}\right) \left(\frac{\mathcal{E}_Y}{r_{\text{min},Y}}\right) \left(\frac{m_Y}{m_X}\right)^2 \end{split}$$

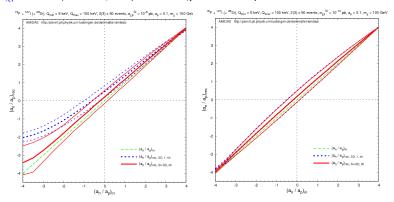
Determining the ratio of two SD WIMP-nucleon couplings

$$\begin{split} \left(\frac{a_{n}}{a_{p}}\right)_{\pm}^{\text{SI+SD}} &= \frac{-\left(c_{p,X}s_{n/p,X} - c_{p,Y}s_{n/p,Y}\right) \pm \sqrt{c_{p,X}c_{p,Y}} \left|s_{n/p,X} - s_{n/p,Y}\right|}{c_{p,X}s_{n/p,X}^{2} - c_{p,Y}s_{n/p,Y}^{2}} \\ c_{p,X} &\equiv \frac{4}{3}\left(\frac{J_{X}+1}{J_{Y}}\right) \left[\frac{\langle S_{p}\rangle_{X}}{J_{Y}}\right]^{2} \left[F_{\text{SI},Z}^{2}(Q_{\text{min},Z})\mathcal{R}_{m,YZ} - F_{\text{SI},Y}^{2}(Q_{\text{min},Y})\right]F_{\text{SD},X}^{2}(Q_{\text{min},X}) \end{split}$$

[M. Drees and CLS, arXiv:0903.3300]



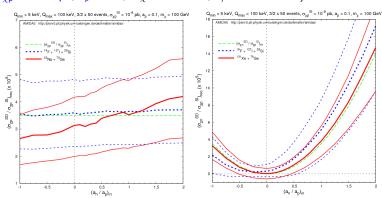
O Reconstructed  $(a_n/a_p)_{\text{rec}}^{\text{SI+SD}}$  vs  $(a_n/a_p)_{\text{rec},1}^{\text{SD}}$  ( $^{19}\text{F} + ^{127}\text{I} + ^{28}\text{Si}$ ,  $Q_{\text{min}} > 5$  keV,  $Q_{\text{max}} < 100$  keV,  $3 \times 50$  events,  $\sigma_{\text{Vp}}^{\text{SI}} = 10^{-8} / 10^{-10}$  pb,  $a_p = 0.1$ ,  $m_\chi = 100$  GeV)



[CLS, arXiv:1103.0482, submitted to JCAP]



O Reconstructed 
$$(\sigma_{\chi p}^{SD}/\sigma_{\chi p}^{SI})_{rec}$$
 and  $(\sigma_{\chi n}^{SD}/\sigma_{\chi p}^{SI})_{rec}$   $(^{19}\text{F} + ^{127}\text{I} + ^{28}\text{Si vs.} ^{23}\text{Na}/^{131}\text{Xe} + ^{76}\text{Ge}, Q_{min} > 5 \text{ keV}, Q_{max} < 100 \text{ keV},$   $\sigma_{\chi p}^{SI} = 10^{-8} \text{ pb}, a_p = 0.1, m_\chi = 100 \text{ GeV}, 3/2 \times 50 \text{ events})$ 



[CLS, arXiv:1103.0482, submitted to JCAP]





- A Model-Independent Data Analysis System for direct Dark Matter detection experiments
  - DAMNED Dark Matter Web Tool (ILIAS Project)
    http://pisrv0.pit.physik.uni-tuebingen.de/darkmatter/amidas/
    [CLS, arXiv:0909.1459, 0910.1971]
  - Online interactive simulation/data analysis system
  - > Full Monte Carlo simulations
  - > Theoretical estimations
  - Real/user-uploaded data analyses
- Further projects/ideas
  - > Analyze events with directional information
  - User account/security system (AMIDAS-II)
  - Combine with other simulation/data analysis packages for (in)direct detections



 $\circ$  Reconstructed results:  $m_{\chi}$  (with Woods-Saxon form factor)

#### Results

#### Reconstructed WIMP mass

$m_{\chi} (\text{GeV}/c^2)$	m <sub>χ</sub> , combined	m <sub>χ, 2</sub>	<i>m</i> χ, 1	<i>m</i> <sub>χ, −1</sub>	$m_{\chi, \sigma}$
m <sub>χ, rec</sub>	122.72	48.283	32.431	1.4779	36.369
<i>m</i> <sub>χ, lo</sub>	83.309	12.565	-1.0121	-57.873	-10.601
m <sub>χ, hi</sub>	196.31	92.53	75.3	107.25	125.66

The reconstructed WIMP mass and the upper and lower bounds of its 1e statistical uncertainty,  $m_{E,n}$  and  $m_{K,\sigma}$  have been estimated by Eqs. (34) and (40) of Ref. [2], respectively; while  $m_{\chi,c}$  ombined by the  $\chi^2$ -fitting defined in Eq. (51) of Ref. [2], which combines the estimators for  $m_{\chi,n}$  and  $m_{\chi,\sigma}$  with each other. The reconstructed WIMP mass  $m_{\chi,c}$  ombined as well as  $m_{\chi,n}$  and  $m_{\chi,\sigma}$  shown here have been corrected by the algorithmic  $Q_{\max}$  matching described in Ref. [2].

#### Data



 $\circ$  Reconstructed results:  $m_{\chi}$  (with exponential form factor)

#### Results

#### Reconstructed WIMP mass

$m_{\chi} (\text{GeV}/c^2)$	$m_{\chi}$ , combined	m <sub>χ, 2</sub>	<i>m</i> <sub>χ, 1</sub>	<i>m</i> <sub>χ, -1</sub>	<i>m</i> χ, σ
m <sub>χ, rec</sub>	113.26	50.279	36.874	0.048119	35.655
<i>m</i> <sub>χ, lo</sub>	75.343	12.413	2.3131	-57.236	-10.108
m <sub>χ, hi</sub>	180.44	89.574	72.862	106.59	127.92

The reconstructed WIMP mass and the upper and lower bounds of its 1e statistical uncertainty,  $m_{Z,R}$  and  $m_{Z,G}$  have been estimated by Eqs. (34) and (40) of Ref. [2], respectively; while  $m_{Z,Combined}$  by the  $\chi^2$ -fitting defined in Eq. (51) of Ref. [2], which combines the estimators for  $m_{Z,R}$  and  $m_{Z,G}$  with each other. The reconstructed WIMP mass  $m_{Z,Combined}$  as well as  $m_{Z,R}$  and  $m_{Z,G}$  shown here have been corrected by the algorithmic  $Q_{\max}$  matching described in Ref. [2].

#### Data



O Reconstructed results:  $|f_p|^2$  (with W-S form factor and input  $\rho_0$ )

#### Results

#### Reconstructed SI WIMP-nucleon coupling

$ f_{\rm p} ^2 (c^6/{\rm GeV}^4)$	fp  <sup>2</sup> input	fp 2recon
f <sub>p</sub>   <sup>2</sup> rec	8.1686E-18	7.8443E-18
fp  <sup>2</sup> lo	6.7381E-18	5.7005E-18
fp  <sup>2</sup> hi	9.5293E-18	1.1103E-17

The reconstructed squared spin-independent WIMP-nucleon coupling  $|f_p|^2$  and the lower and upper bounds of its lo statistical uncertainty estimated by Eqs. (17) and (18) of Ref. [3] with the input WIMP mass (with an overall uncertainty of 0.1%) and the reconstructed WIMP mass (estimated by the algorithmic process described in Ref. [2]).

#### Data



O Reconstructed results:  $|f_p|^2$  (with ex form factor and standard  $\rho_0$ )

#### Results

#### Reconstructed SI WIMP-nucleon coupling

$ f_{\rm p} ^2 (c^6/{\rm GeV}^4)$	fp  <sup>2</sup> input	fp 2recon
$ f_{\rm p} ^2_{\rm rec}$	9.3348E-18	1.0157E-17
fp  <sup>2</sup> lo	7.6996E-18	7.5775E-18
fp  <sup>2</sup> hi	1.0874E-17	1.4234E-17

The reconstructed squared spin-independent WIMP-nucleon coupling  $|f_p|^2$  and the lower and upper bounds of its lo statistical uncertainty estimated by Eqs. (17) and (18) of Ref. [3] with the input WIMP mass (with an overall uncertainty of 0.1%) and the reconstructed WIMP mass (estimated by the algorithmic process described in Ref. [2]).

#### Data



 $\circ$  Reconstructed results:  $a_n/a_p$  (with Woods-Saxon form factor)

#### Results

#### Reconstructed ratio between two SD WIMP-nucleon couplings

<i>a</i> <sub>n</sub> / <i>a</i> <sub>p</sub>	$(a_n / a_p)_{+, 1, sh}^{SD}$	$(a_n / a_p)_{-, 1, sh}^{SD}$	$(a_n/a_p)_{+, sh}^{SI+SD}$	$(a_n / a_p)_{-, sh}^{SI + SD}$
$(a_n / a_p)_{rec}$	16.441	0.97458	19.75	0.80541
$(a_n / a_p)_{lo}$	9.9344	0.59298	11.496	0.48606
$(a_n / a_p)_{hi}$	22.652	1.3532	27.389	1.1488

The reconstructed  $a_0 / a_p$  ratios and the lower and upper bounds of their  $1\sigma$  statistical uncertainties estimated by Eqs. (31) and (36) of Ref. [4] with n = -1, 1, and 2 as well as by Eqs. (57) and (60) of Ref. [4] at the shifted energy points (given by Eq. (38) of Ref. [41)).

#### Data



 $\circ$  Reconstructed results:  $a_n/a_p$  (with exponential form factor)

#### Results

#### Reconstructed ratio between two SD WIMP-nucleon couplings

<i>a</i> <sub>n</sub> / <i>a</i> <sub>p</sub>	$(a_n / a_p)_{+, 1, sh}^{SD}$	(a <sub>n</sub> / a <sub>p</sub> )-, 1, sh <sup>SD</sup>	$(a_n/a_p)_{+, sh}^{SI+SD}$	$(a_n / a_p)_{-, sh}^{SI + SD}$
$(a_n / a_p)_{rec}$	16.441	0.97458	19.745	0.80563
$(a_n / a_p)_{lo}$	9.9344	0.59298	11.496	0.48636
$(a_n / a_p)_{hi}$	22.652	1.3532	27.376	1.149

The reconstructed  $a_0 / a_p$  ratios and the lower and upper bounds of their  $1\sigma$  statistical uncertainties estimated by Eqs. (31) and (36) of Ref. [4] with n = -1, 1, and 2 as well as by Eqs. (57) and (60) of Ref. [4] at the shifted energy points (given by Eq. (38) of Ref. [41)).

#### Data



O Reconstructed results:  $\sigma_{(p,n),\chi}^{SD}/\sigma_{\chi p}^{SI}$  (with W-S form factor)

#### Results

#### Reconstructed ratio between the SD and SI WIMP-proton cross sections

$\sigma_{\chi p}^{SD} / \sigma_{\chi p}^{SI}$	$(\sigma_{\chi p}^{SD} / \sigma_{\chi p}^{SI})_{+, sh}^{XYZ}$	$(\sigma_{\chi p}^{SD} / \sigma_{\chi p}^{SI})_{-, sh}^{XYZ}$	$(\sigma_{\chi p}^{SD} / \sigma_{\chi p}^{SI})_{sh}^{XY}$
$(\sigma_{\chi p}^{SD} / \sigma_{\chi p}^{SI})_{rec}$	39390	913780	1008000
$(\sigma_{\chi p}^{SD} / \sigma_{\chi p}^{SI})_{lo}$	6036.9	615690	632860
$(\sigma_{\chi p}^{SD} / \sigma_{\chi p}^{SI})_{hi}$	74207	1216200	1373300

The reconstructed  $\sigma_{gp}^{\rm SI}/\sigma_{gp}^{\rm SI}$  ratios and the lower and upper bounds of their  $1\sigma$  statistical uncertainties estimated by Eqs. (50) and (61) (with  $a_n/a_p$  estimated by Eq. (57)) as well as by Eqs. (65) and (69) of Ref. [4] at the shifted energy points (given by Eq. (38) of Ref. [4]).

#### Reconstructed ratio between the SD and SI WIMP-neutron cross sections

$\sigma_{\chi n}^{SD} / \sigma_{\chi p}^{SI}$	$(\sigma_{\chi n}^{SD} / \sigma_{\chi p}^{SI})_{+, sh}^{XYZ}$	$(\sigma_{\chi n}^{SD} / \sigma_{\chi p}^{SI})_{-, sh}^{XYZ}$	$(\sigma_{\chi n}^{SD} / \sigma_{\chi p}^{SI})_{sh}^{XY}$
$(\sigma_{\chi n}^{SD} / \sigma_{\chi p}^{SI})_{rec}$	15164000	602750	487250
$(\sigma_{\chi n}^{SD} / \sigma_{\chi p}^{SI})_{lo}$	10220000	72093	269840
$(\sigma_{\chi n}^{SD} / \sigma_{\chi p}^{SI})_{hi}$	20193000	1152700	700660



O Reconstructed results:  $\sigma_{(p,n),\chi}^{SD}/\sigma_{\chi p}^{SI}$  (with exponential form factor)

#### **Results**

#### Reconstructed ratio between the SD and SI WIMP-proton cross sections

$\sigma_{\chi p}^{SD} / \sigma_{\chi p}^{SI}$	$(\sigma_{\chi p}^{SD} / \sigma_{\chi p}^{SI})_{+, sh}^{XYZ}$	$(\sigma_{\chi p}^{SD} / \sigma_{\chi p}^{SI})_{-, sh}^{XYZ}$	$(\sigma_{\chi p}^{SD} / \sigma_{\chi p}^{SI})_{sh}^{XY}$
$(\sigma_{\chi p}^{SD} / \sigma_{\chi p}^{SI})_{rec}$	39334	911960	1001900
$(\sigma_{\chi p}^{SD} / \sigma_{\chi p}^{SI})_{lo}$	6044.5	614460	629110
$(\sigma_{\chi p}^{SD} / \sigma_{\chi p}^{SI})_{hi}$	74084	1213800	1365300

The reconstructed  $\sigma_{gp}^{SD}/\sigma_{gp}^{SI}$  ratios and the lower and upper bounds of their  $1\sigma$  statistical uncertainties estimated by Eqs. (50) and (61) (with  $a_n/a_p$  estimated by Eq. (57)) as well as by Eqs. (65) and (69) of Ref. [4] at the shifted energy points (given by Eq. (38) of Ref. [4]).

#### Reconstructed ratio between the SD and SI WIMP-neutron cross sections

$\sigma_{\chi n}^{SD} / \sigma_{\chi p}^{SI}$	$(\sigma_{\chi n}^{SD} / \sigma_{\chi p}^{SI})_{+, sh}^{XYZ}$	$(\sigma_{\chi n}^{SD} / \sigma_{\chi p}^{SI})_{-, sh}^{XYZ}$	$(\sigma_{\chi n}^{SD} / \sigma_{\chi p}^{SI})_{sh}^{XY}$
$(\sigma_{\chi n}^{SD} / \sigma_{\chi p}^{SI})_{rec}$	15132000	601930	484270
$(\sigma_{\chi n}^{SD} / \sigma_{\chi p}^{SI})_{lo}$	10199000	72257	268430
$(\sigma_{\chi n}^{SD} / \sigma_{\chi p}^{SI})_{hi}$	20153000	1150800	696420



- Input setup for generating pseudodata
  - $m_{Y} = 130 \text{ GeV}$
  - $ightharpoonup \sigma_{\chi p}^{\rm SI} = 4 imes 10^{-9} \ \mathrm{pb}$
  - $> a_p = 0.1, a_n/a_p = 0.7$
  - > Woods-Saxon and thin-shell elastic nuclear form factors
  - $ho_0 = 0.4 \text{ GeV/cm}^3$
  - $> v_0 = 230 \text{ km/s}$
  - $> t_p = 140 \text{ d}, t_{expt} = 300 \text{ d}$
  - $|f_p|^2 = 9.305 \times 10^{-18} \text{ GeV}^{-4} \text{ (for 130 GeV } m_\chi \text{)}$
  - $\sigma_{\chi p}^{SD}/\sigma_{\chi p}^{SI} = 8.77 \times 10^5, \ \sigma_{\chi n}^{SD}/\sigma_{\chi p}^{SI} = 4.30 \times 10^5 \ (\text{for } 130 \text{ GeV } m_{\chi})$



### Summary



### Summary

- Once two or more experiments with different target nuclei observe positive WIMP signals, we could estimate
  - $ightharpoonup WIMP mass <math>m_{\chi}$
  - > SI WIMP-proton coupling  $|f_p|^2$
  - $\rightarrow$  ratio between the SD WIMP-nucleon couplings  $a_n/a_p$
  - $\succ$  ratios between the SD and SI WIMP-nucleon cross sections  $\sigma_{\chi({\bf p},{\bf n})}^{\rm SD}/\sigma_{\chi{\bf p}}^{\rm SI}$
- These analyses are independent of the velocity distribution, the local dentity, and the mass/couplings on nucleons of halo WIMPs (none of them is yet known).
- For a WIMP mass of 100 GeV, these quantities could be estimated with statistical uncertainties of 10% 40% with only  $\mathcal{O}(50)$  events from one experiment.



### Summary on background effects



### Summary on background effects

### References

- Y.-T. Chou and C.-L. Shan, "Effects of Residue Background Events in Direct Dark Matter Detection Experiments on the Determination of the WIMP Mass", JCAP 1008, 014 (2010).
- C.-L. Shan, "Effects of Residue Background Events in Direct Dark Matter Detection Experiments on the Reconstruction of the Velocity Distribution Function of Halo WIMPs", JCAP 1006, 029 (2010).
- C.-L. Shan, "Effects of Residue Background Events in Direct Dark Matter Detection Experiments on the Estimation of the Spin-Independent WIMP-Nucleon Coupling", arXiv:1103.4049 [hep-ph] (2011).
- C.-L. Shan, "Effects of Residue Background Events in Direct Dark Matter Detection Experiments on the Determinations of Ratios of WIMP-Nucleon Cross Sections", arXiv:1104.5305 [hep-ph] (2011).



### Summary on background effects

- For determining WIMP properties, the maximal acceptable background ratio is  $\sim 10\% 20\%$ .
- For determining ratios between different WIMP couplings/cross sections, between results reconstructed under different assumptions and/or with different moments of the WIMP velocity distribution function and/or by using different target nuclei there could be an (in)compatibility.
- For determining ratios between different WIMP couplings/cross sections, background events in low energy ranges would be (much) more problematic than those in high energy ranges.

[More detailed discussions will be given at TAUP 2011 in Munich]



Thank you very much for your attention [http://myweb.ncku.edu.tw/~clshan/Publications/Talks/]