

What does the Higgs boson tell us about New Physics?

Béranger Dumont

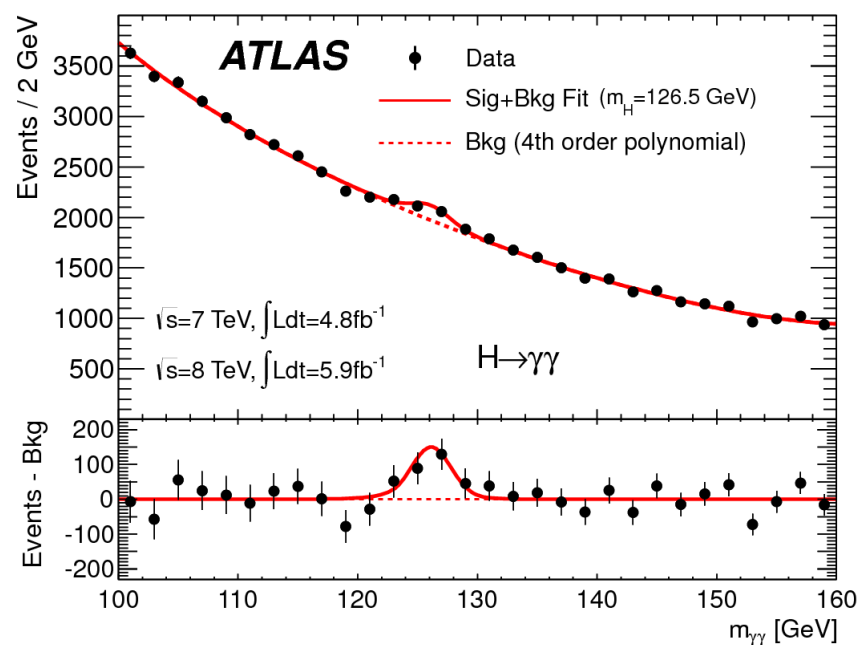
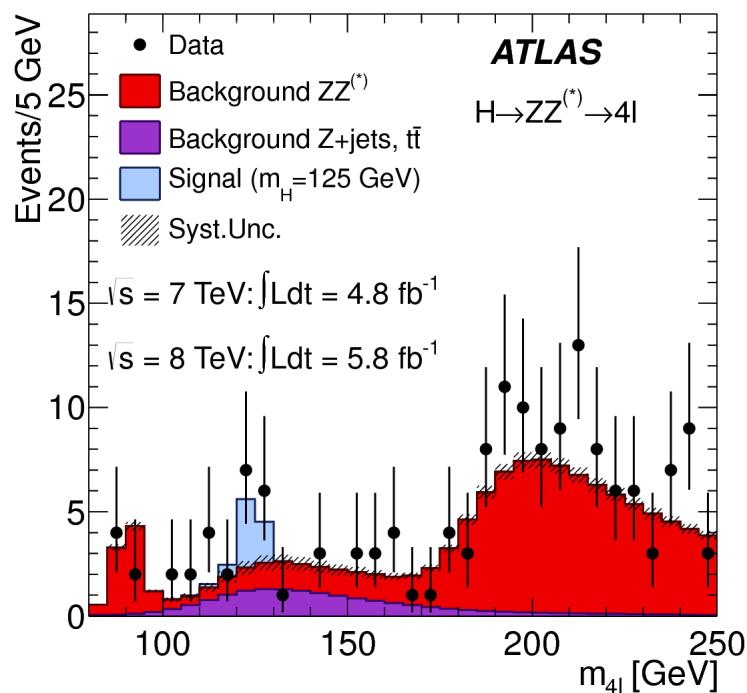


Séminaire des doctorants
May 28, 2013

Discovery of a new particle

on July 4th:

discovery of a new particle “consistent with the Higgs boson” at the LHC
mass ≈ 125 GeV



the Higgs boson has been theorized in 1964
long-awaited discovery after 30 years of search!

Standard Model Lagrangian

$$\mathcal{L}_{\text{SM}} = \mathcal{L}_{\text{gauge}} + \mathcal{L}_{\text{Higgs}}$$


- Follows from the local $\underbrace{\text{SU}(3)_C}_{\text{QCD}} \times \underbrace{\text{SU}(2)_L \times \text{U}(1)_Y}_{\text{electroweak}}$ symmetry
- interaction between fermions and gauge bosons (g, W, Z, γ)
- nice & simple, experimentally tested with high accuracy

- we need to break the electroweak symmetry!
 - gives mass to the W and Z bosons
 - gives masses to the fermions (for free)
- but *ad hoc* and not so well-known...

Standard Model Higgs

$$\mathcal{L}_{\text{Higgs}} = |D_\mu H|^2 - \mu^2 |H|^2 + \lambda |H|^4 + Y^{ij} \bar{f}_L^i f_R^j H$$

quadratic divergences?

the hierarchy problem:

we want $\mu \sim \mathcal{O}(100 \text{ GeV})$

we obtain $\mu \sim \mathcal{O}(10^{18} \text{ GeV})$

vacuum instability?

what if $\lambda < 0$?

flavor structure?

lack of understanding

most of the problems of the Standard Model come from the Higgs sector!

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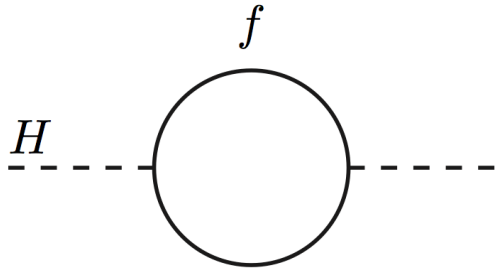
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The hierarchy problem



quantum corrections to the Higgs mass parameter from New Physics arising at the scale Λ_{UV}

$$\hookrightarrow \Delta\mu^2 \propto \Lambda_{UV}^2 \sim \mathcal{O}(M_{Pl}^2)$$

remain true if the new particles are only indirectly coupled to the Higgs



The Higgs mass parameter is naturally pushed to M_{Pl} ...but we know that $m_H \approx 125 \text{ GeV}$!

→ we **need** New Physics to explain the smallness of m_H compared to M_{Pl}

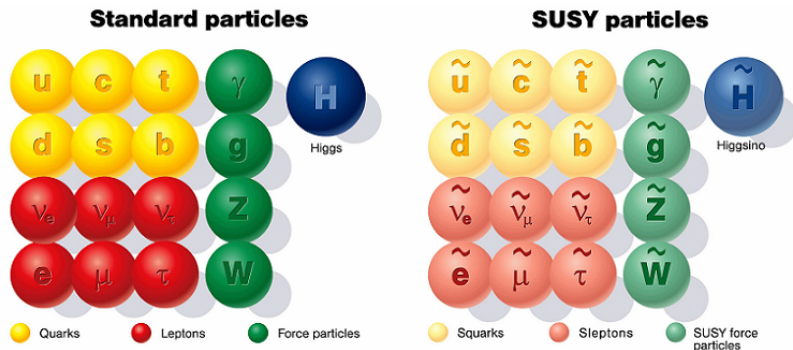
Possible solutions to the hierarchy problem

two popular solutions to the hierarchy problem

supersymmetry

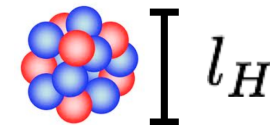
additional symmetry relating fermions and bosons

→ predicts “superpartners” of the existing particles (spin differing by $\frac{1}{2}$ unit)



composite Higgs

Higgs boson: hadron of new strong force
corrections to m_H screened above $1/l_H$

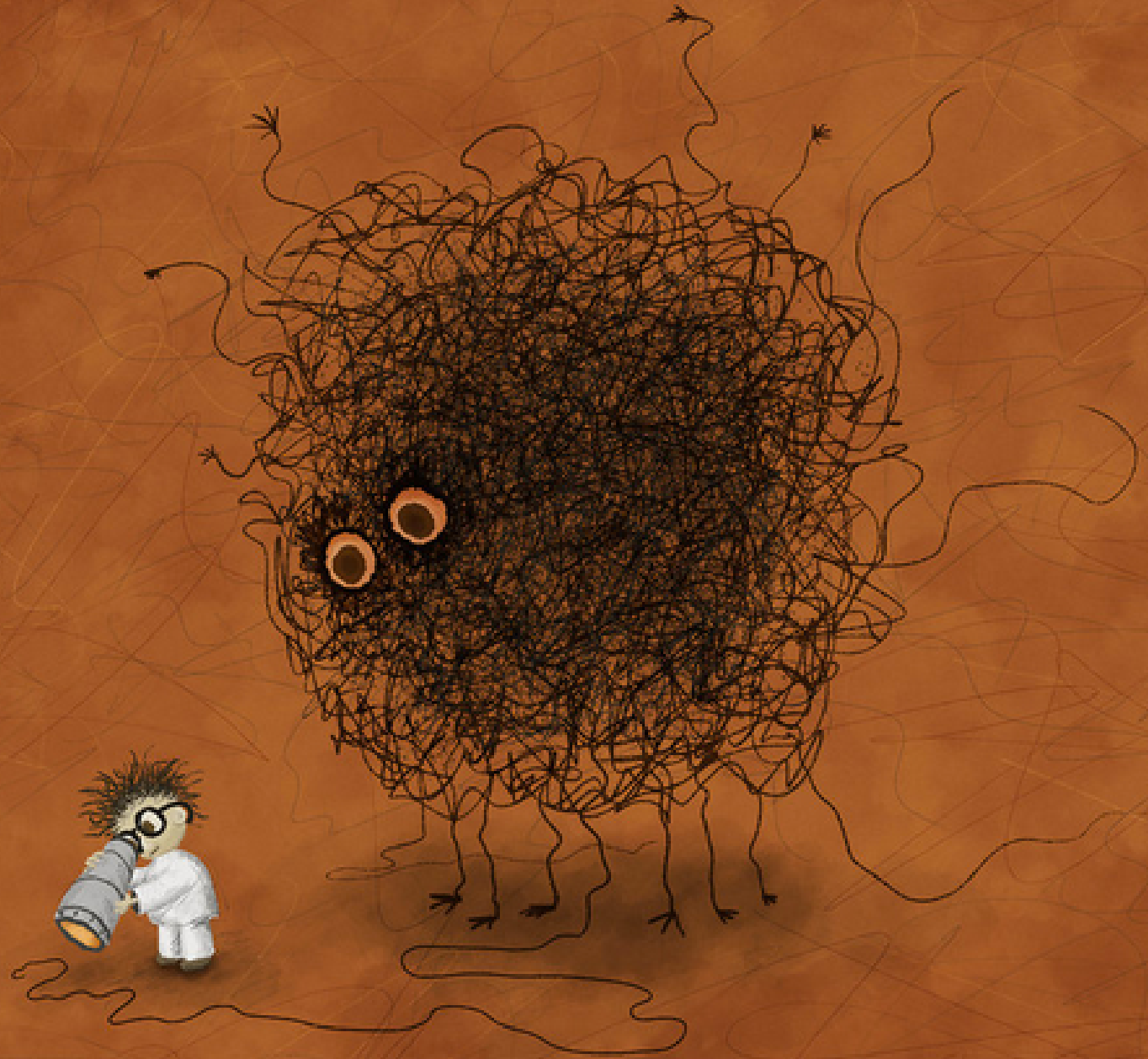


→ generically predicts (light) top partners

in both cases:

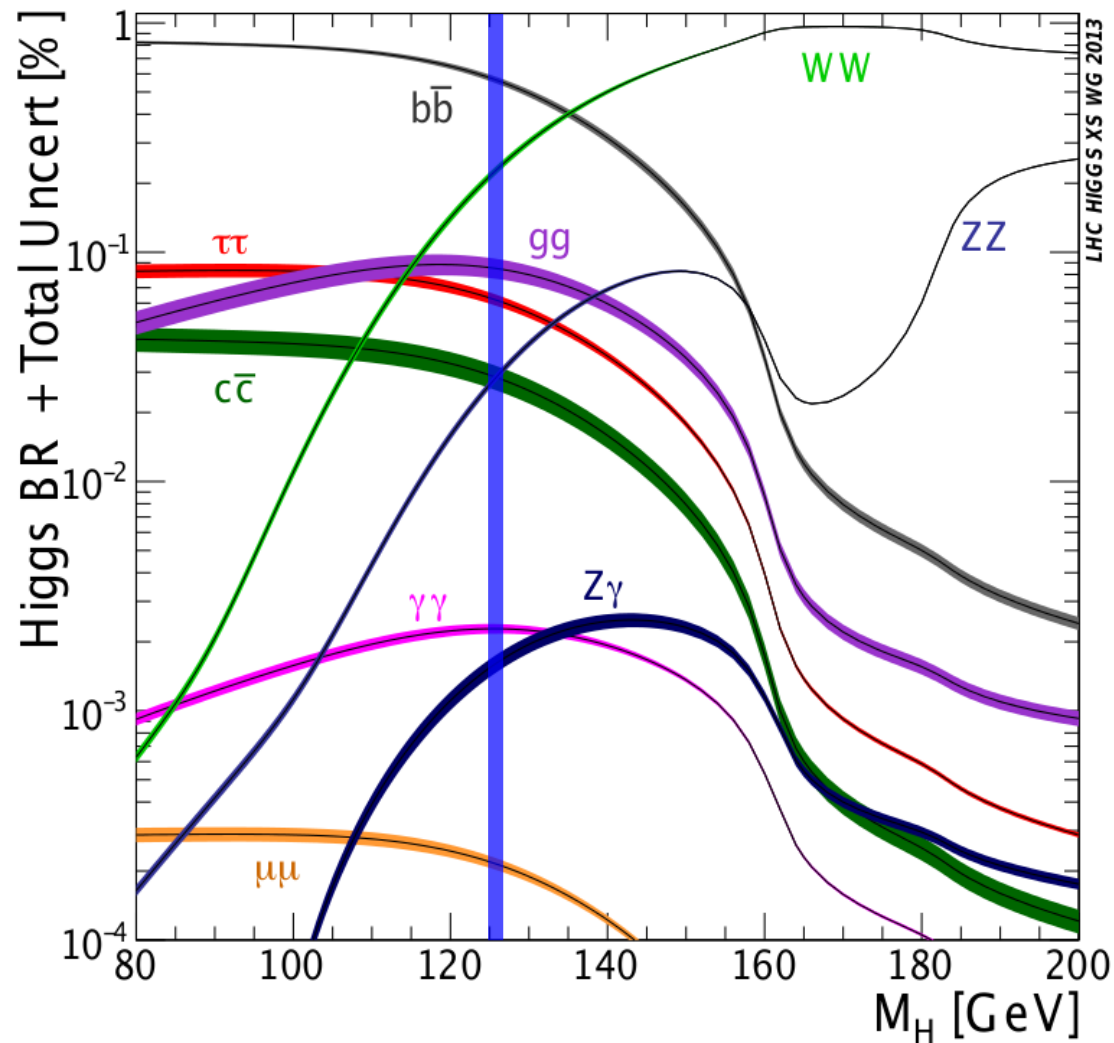
- natural to expect new particles at the TeV scale, i.e. **accessible at the LHC!**
- sizeable modifications of the properties of the Higgs are possible: **probed at the LHC!**

Standard Model Higgs... or New Physics?



taken from Alexey Drozdetskiy's talk at HCP2012

Accessible decays of the Higgs boson

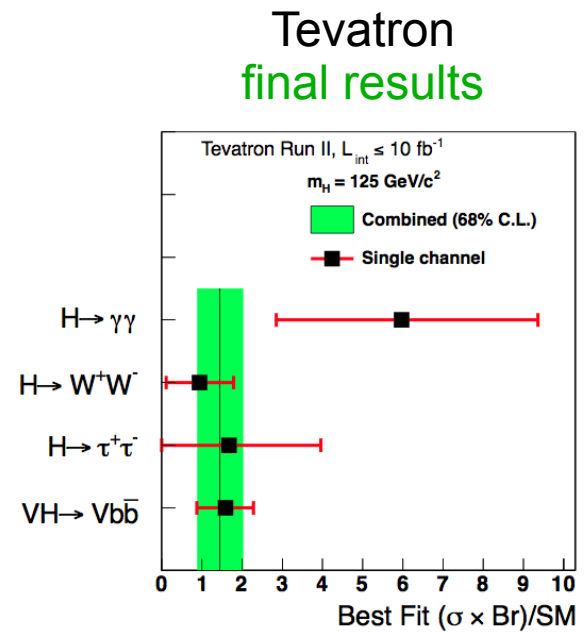
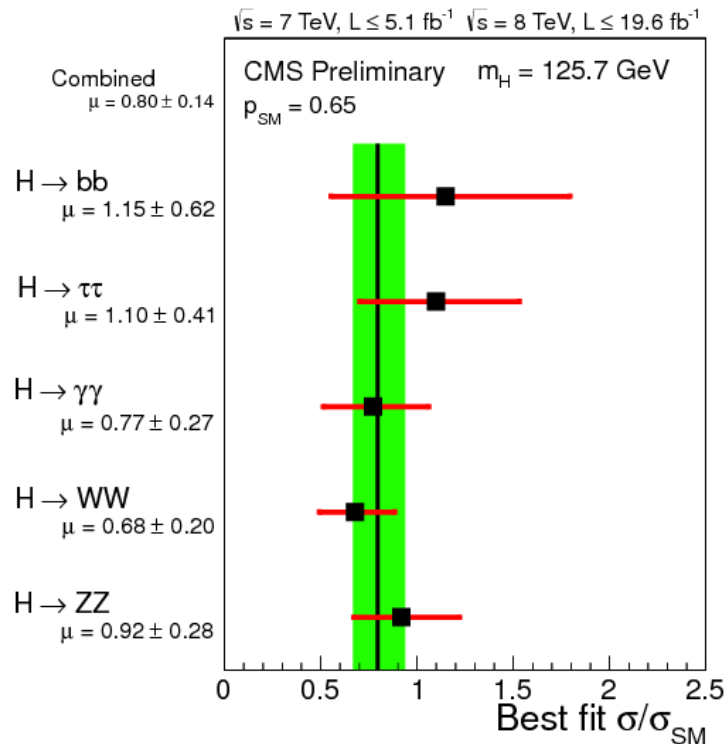
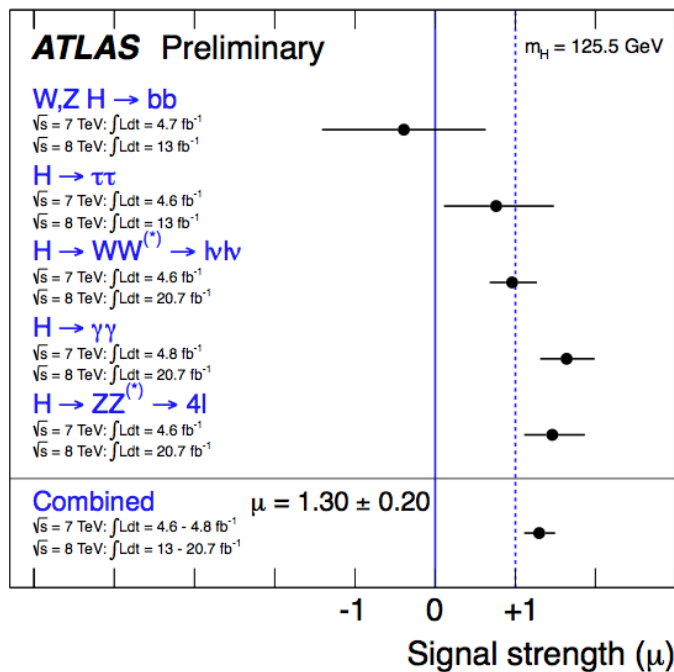


- $m_H \sim 125$ GeV is interesting: many channels are accessible!
- the mass of the new particle is inferred from the $\gamma\gamma$ and ZZ channels
- $H \rightarrow gg$ and $H \rightarrow cc$ channels: not accessible at the LHC (QCD background)

Searches for the Higgs boson

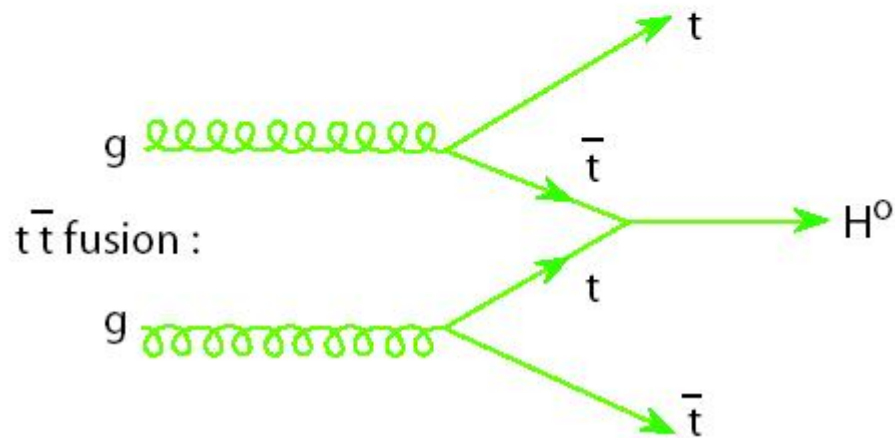
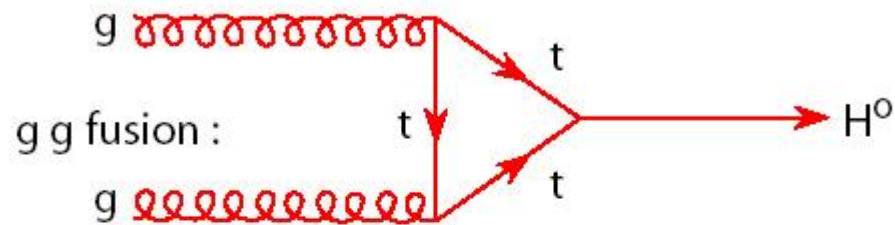
- new particle at observed 125.5 GeV with $> 7\sigma$ significance by ATLAS and CMS!
almost all bosonic searches have been updated with full luminosity

- results summarized in terms of signal strengths: $\mu_i = \frac{\left[\sum_j \sigma_{j \rightarrow h} \times \mathcal{B}(h \rightarrow i) \right]_{\text{observed}}}{\left[\sum_j \sigma_{j \rightarrow h} \times \mathcal{B}(h \rightarrow i) \right]_{\text{SM}}}$



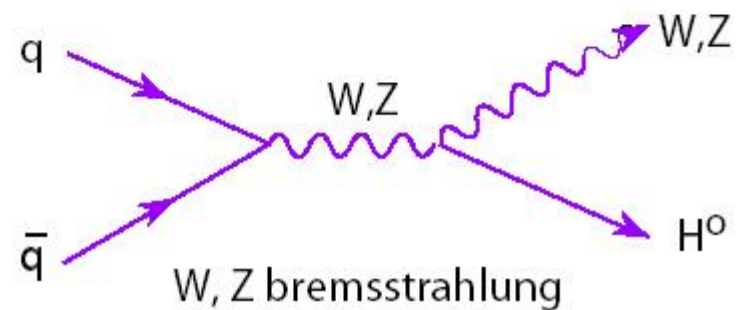
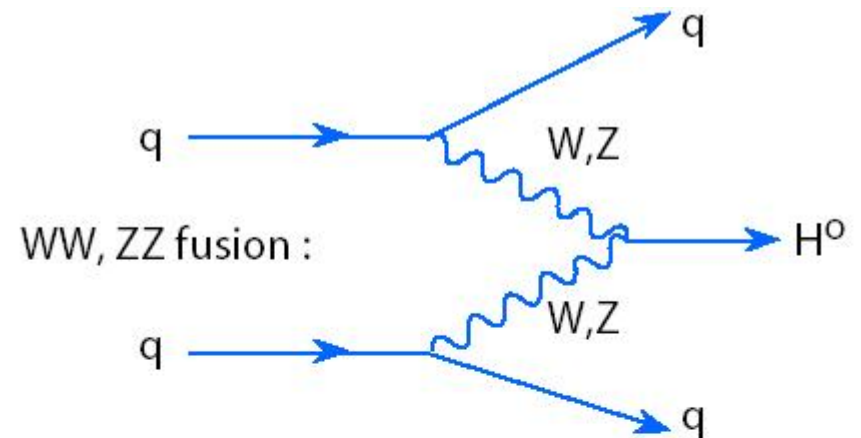
Production of a Higgs boson

...but New Physics modifies not only the Higgs decays but also its production!



(not observed yet)

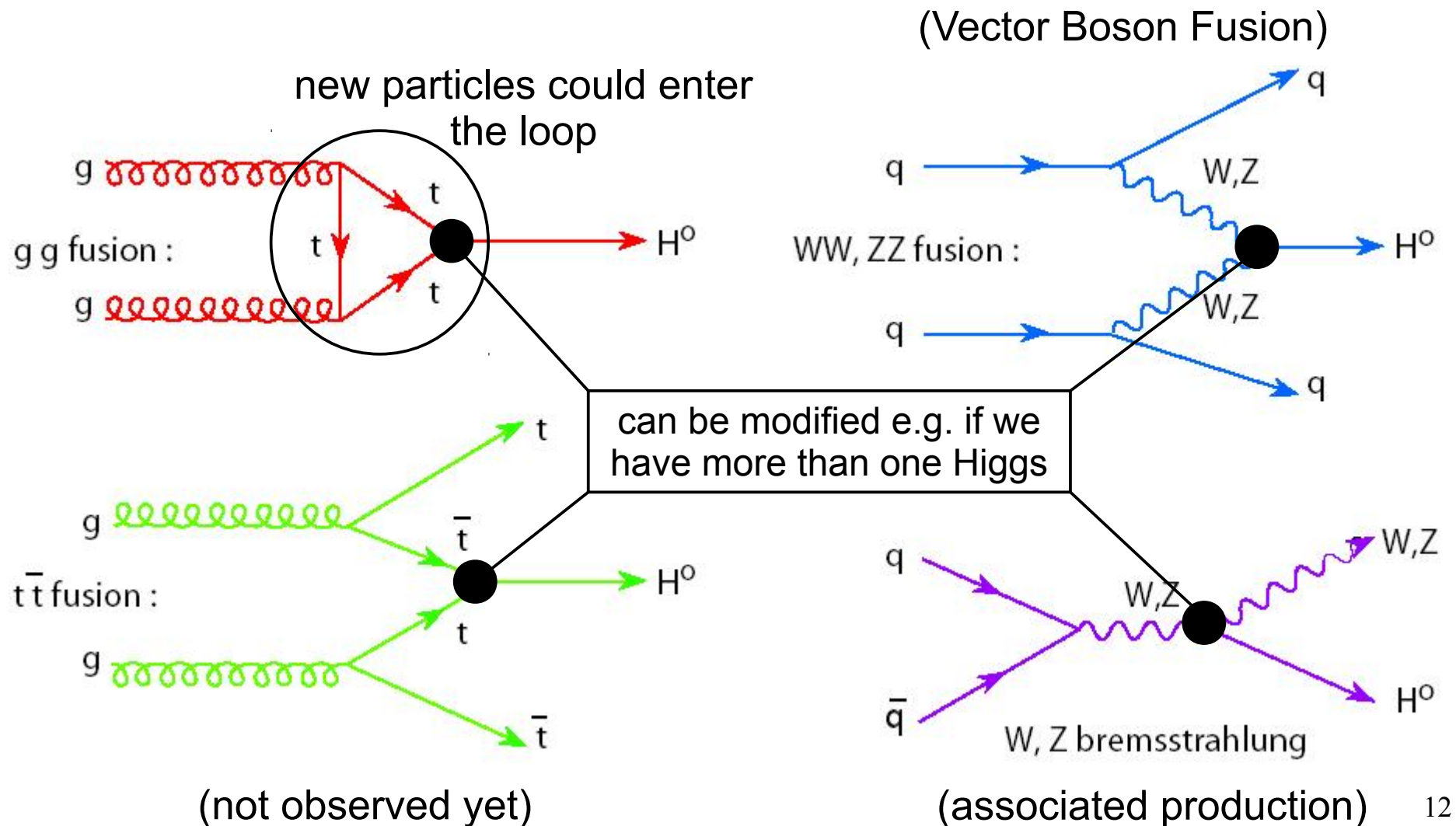
(Vector Boson Fusion)



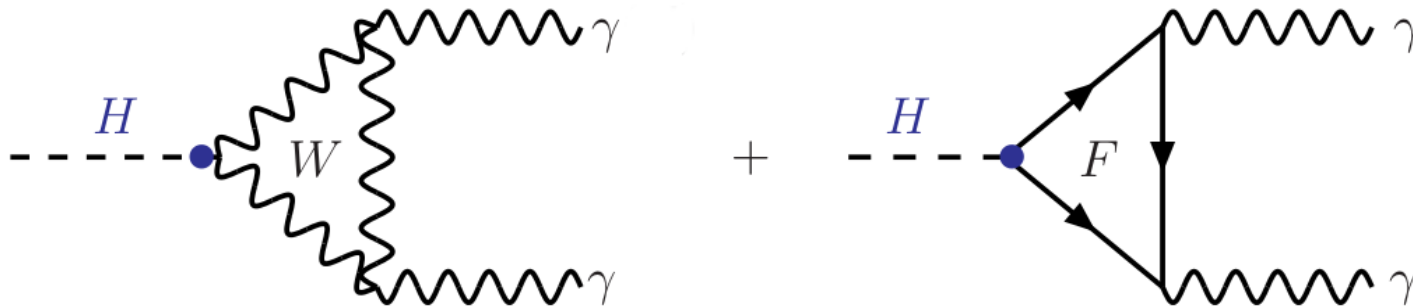
(associated production)

Production of a Higgs boson

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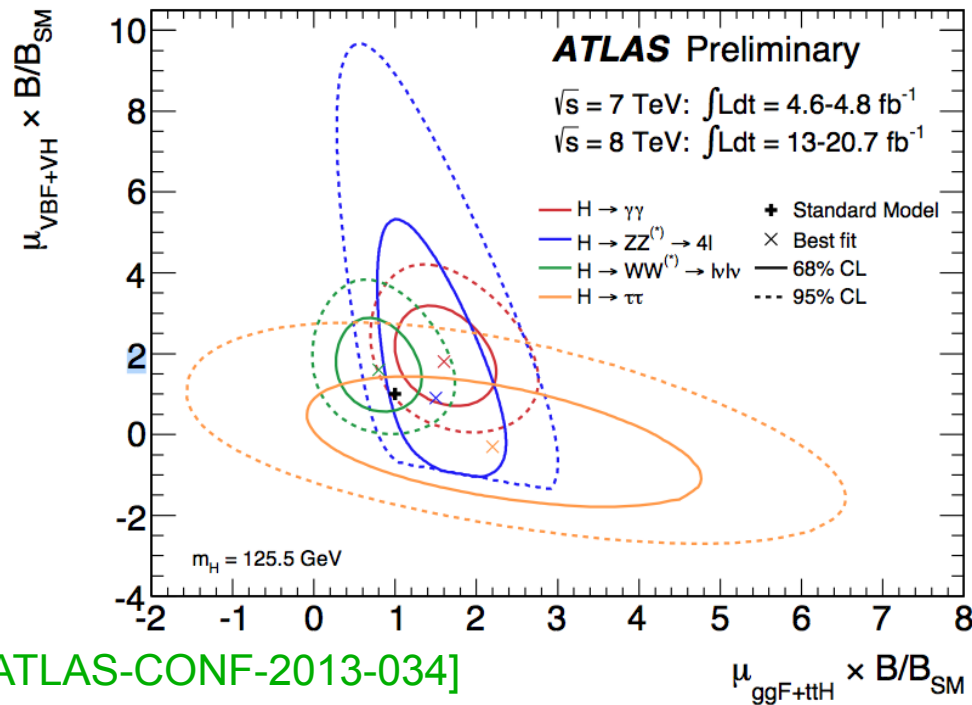
A word on $H \rightarrow \gamma\gamma$



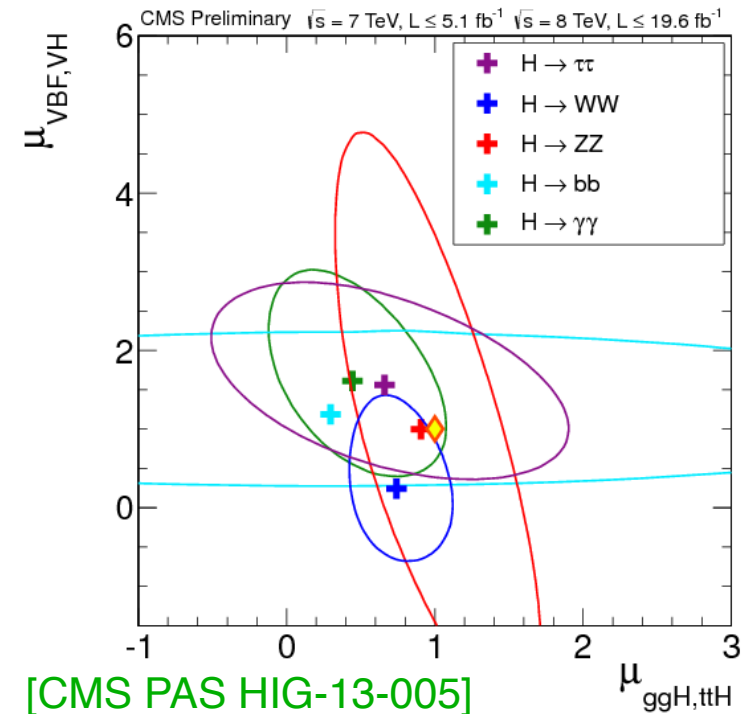
- contribution from the W is 5 times larger than from the top quark
- small contributions from bottom and lighter quarks
- new particles in the loop could change the $H\gamma\gamma$ rate!
(e.g. charged Higgses, charginos, staus, ...)

2D μ plots from ATLAS and CMS

ATLAS



CMS



we implement these results using the Gaussian approximation
 (validity checked ✓)

Understanding what we observe

The Higgs is Standard Model-like so far

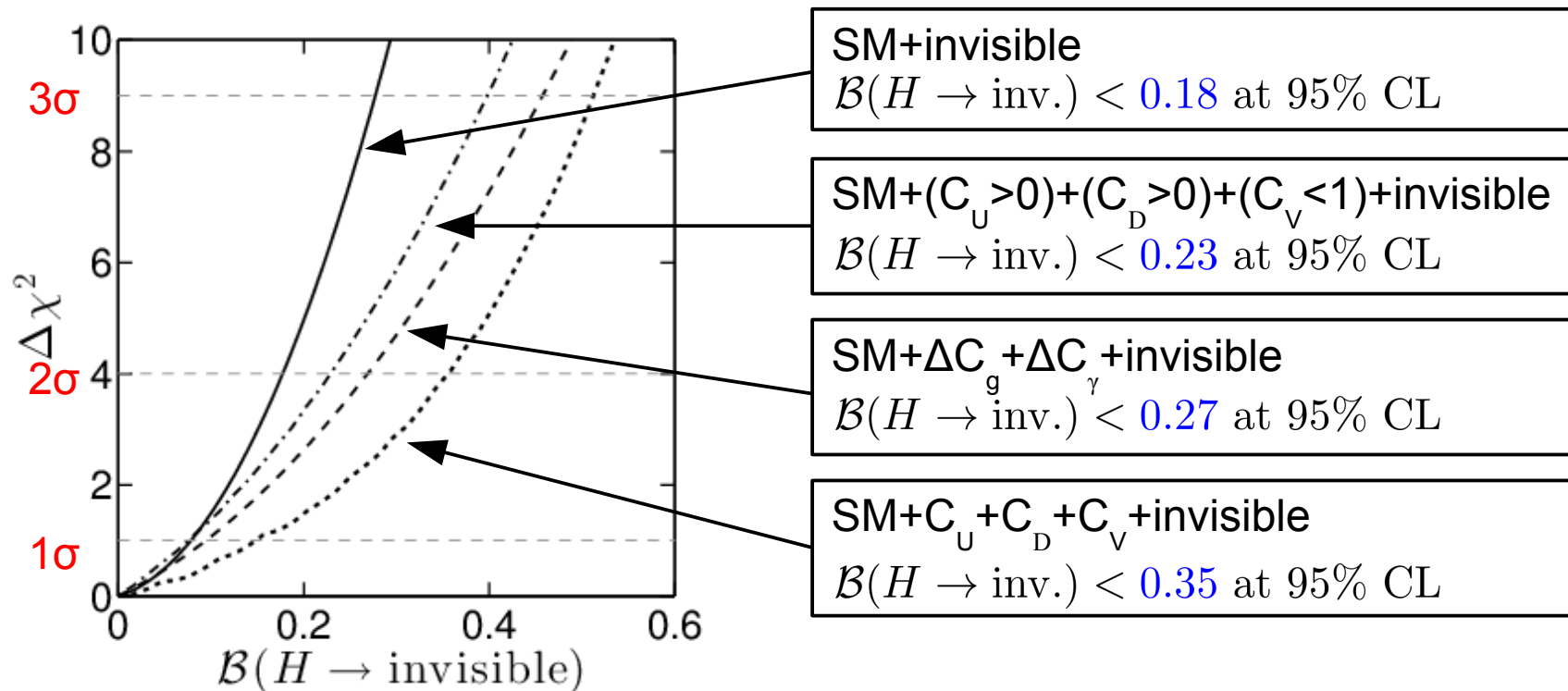
→ provides constraints on theories Beyond the Standard Model

- we have performed various couplings fits of the observed new particle

example of result:

G. Belanger, BD, U. Ellwanger, J. F. Gunion, and S. Kraml,

[arXiv:1212.5244] and [arXiv:1302.569]



Philosophy of our EFT approach

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \sum_i \frac{\alpha_i}{\Lambda^{n_i}} \mathcal{O}_i$$

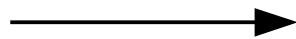
we consider dimension-6 operators only

BD, S. Fichet,
G. von Gersdorff
[\[arXiv:1304.3369\]](#)

underlying
assumptions

- the observed state at ~ 125 GeV is
 - CP-even
 - spin 0
 - and belongs to a $\text{SU}(2)_L$ doublet
- there is a mass gap between the SM and New Physics (arising at the scale Λ)

complementarity with the
“anomalous couplings”
approach



less general but...

- clear ordering between operators
- fully consistent theoretical framework

Our basis of relevant operators

$$\mathcal{O}_D = J_{H\mu}^a J_\mu^a, \quad \mathcal{O}_{D^2} = |H|^2 |D_\mu H|^2$$

$$\mathcal{O}_{WW} = H^\dagger H (W_{\mu\nu}^a)^2, \quad \mathcal{O}_{BB} = H^\dagger H (B_{\mu\nu})^2$$

$$\mathcal{O}_{WB} = H^\dagger W_{\mu\nu} H B_{\mu\nu}, \quad \mathcal{O}_{GG} = H^\dagger H (G_{\mu\nu}^a)^2$$

$$\mathcal{O}_t = 2y_t |H|^2 H \bar{t}_L t_R, \quad \mathcal{O}_b = 2y_b |H|^2 H \bar{b}_L b_R, \quad \mathcal{O}_\tau = 2y_\tau |H|^2 H \bar{\tau}_L \tau_R$$

no custodial-symmetry
violating operators

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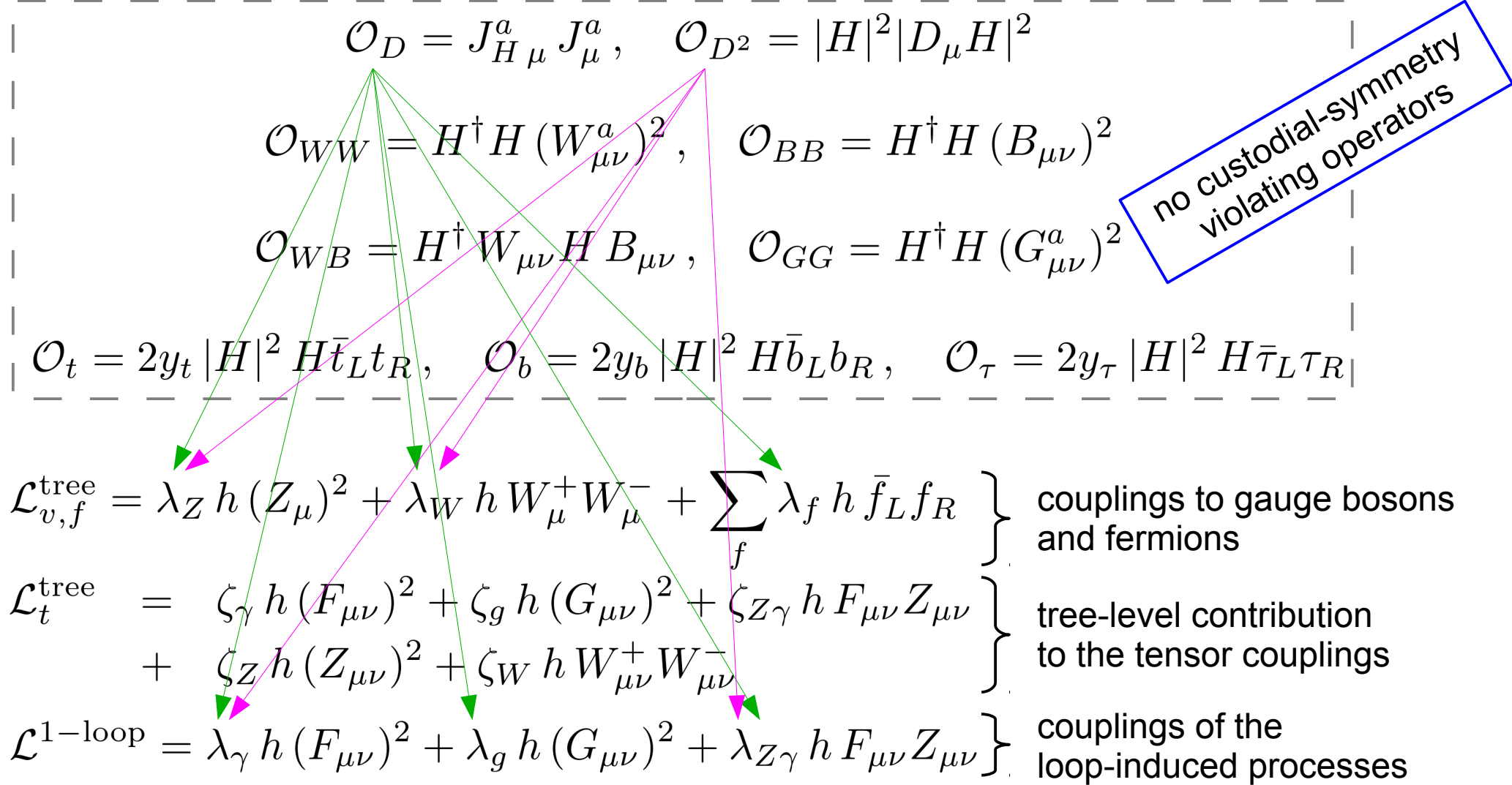
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$$\mathcal{L}_{v,f}^{\text{tree}} = \lambda_Z h (Z_\mu)^2 + \lambda_W h W_\mu^+ W_\mu^- + \sum_f \lambda_f h \bar{f}_L f_R \quad \left. \vphantom{\sum_f} \right\} \text{couplings to gauge bosons and fermions}$$

$$\begin{aligned} \mathcal{L}_t^{\text{tree}} &= \zeta_\gamma h (F_{\mu\nu})^2 + \zeta_g h (G_{\mu\nu})^2 + \zeta_{Z\gamma} h F_{\mu\nu} Z_{\mu\nu} \\ &+ \zeta_Z h (Z_{\mu\nu})^2 + \zeta_W h W_{\mu\nu}^+ W_{\mu\nu}^- \end{aligned} \quad \left. \vphantom{\zeta_W} \right\} \text{tree-level contribution to the tensor couplings}$$

$$\mathcal{L}^{1\text{-loop}} = \lambda_\gamma h (F_{\mu\nu})^2 + \lambda_g h (G_{\mu\nu})^2 + \lambda_{Z\gamma} h F_{\mu\nu} Z_{\mu\nu} \quad \left. \vphantom{\lambda_{Z\gamma}} \right\} \text{couplings of the loop-induced processes}$$

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these operators cannot be generated at tree-level
within a perturbative UV theory

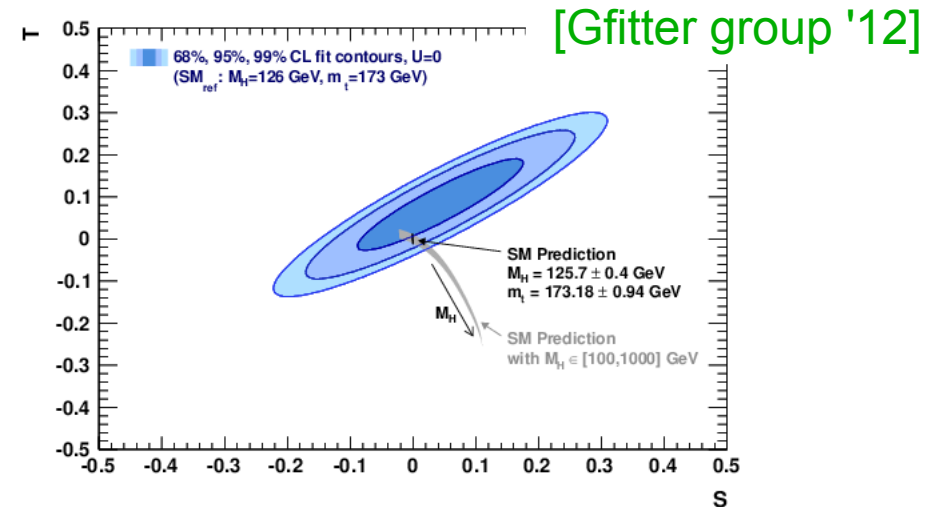
Scenario I) democratic HDOs
→ all operators on equal footing

Scenario II) loop suppressed $\mathcal{O}_{FF'}$'s
compared to the other operators
 $FF = WW, WB, BB, GG$

Experimental constraints

- Higgs properties:
all measurements up-to-date (incl. limits on $h \rightarrow Z\gamma$)
→ implemented as in the Higgs couplings fits

- electroweak precision observables:
Peskin–Takeuchi S & T parameters



- measurements of the Triple Gauge Vertices (TGV) $WW\gamma$, WWZ :
 $\kappa_\gamma = 0.973^{+0.044}_{-0.045}$, $g_1^Z = 0.984^{+0.022}_{-0.019}$ [LEPEWWG/TGC/2005-01]

Bayesian inference & MCMC

in a model M , having:

- parameters of interest ϕ ,
- other parameters ψ ,

the posterior probability on ϕ given the experimental data is:

$$\underbrace{p(\phi|d, M)}_{\text{marginal posterior on the parameters of interest}} \propto \int \underbrace{L(\phi, \psi)}_{\text{likelihood}} \underbrace{\pi(\phi, \psi|M)}_{\text{prior}} d\psi$$

with $L = L_{\text{Higgs}} \times L_{S,T} \times L_{\text{TGV}}$

$\pi(\alpha_i) = \pi(\beta_i) = 1$ (uniform prior)

we sample the posterior probability distribution using Markov Chain Monte Carlo (MCMC)

Setup of the analysis

scan ranges for the 2 scenarios we consider:

	I) Democratic HDOs	II) Loop-suppressed \mathcal{O}_{FF} 's
Λ	$4\pi v$	$4\pi v$
β_{FF}	$[-1, 1]$	$[-1/16\pi^2, 1/16\pi^2]$
Other β	$[-1, 1]$	$[-1, 1]$

where $\beta_i = \alpha_i v^2 / \Lambda^2$

$$FF = WW, WB, BB, GG$$

we fix $\Lambda = 4\pi v \approx 3$ TeV but the dependence in Λ is mild
→ results remain valid for TeV-scale New Physics

Results

1D probability distributions

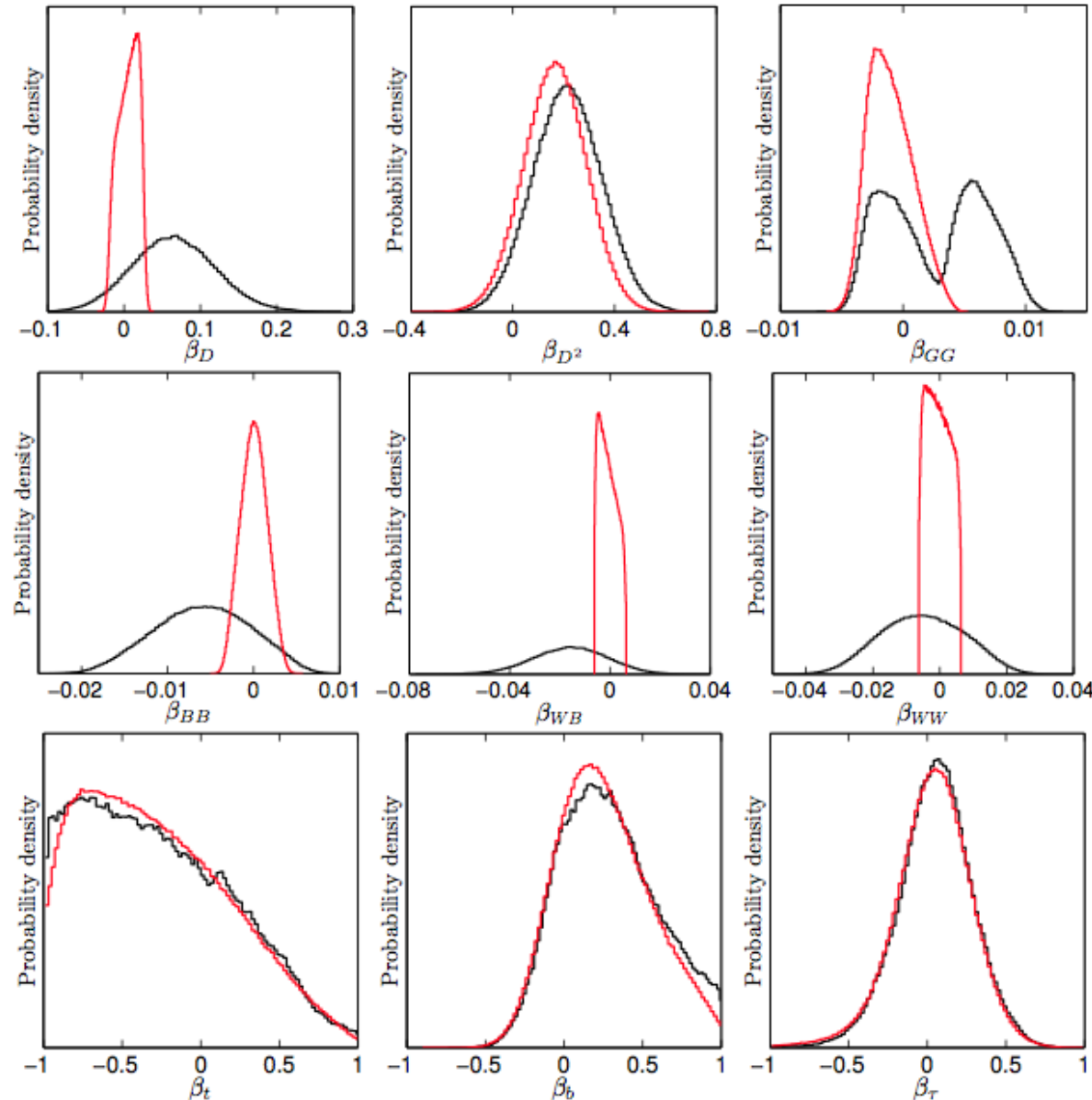
- black line
→ democratic HDOs
- red line
→ loop suppressed \mathcal{O}_{FF}' s

- β_{FF} are $\mathcal{O}(0.01)$
 $FF = WW, WB, BB, GG$

- β_f can go up to $\mathcal{O}(1)$

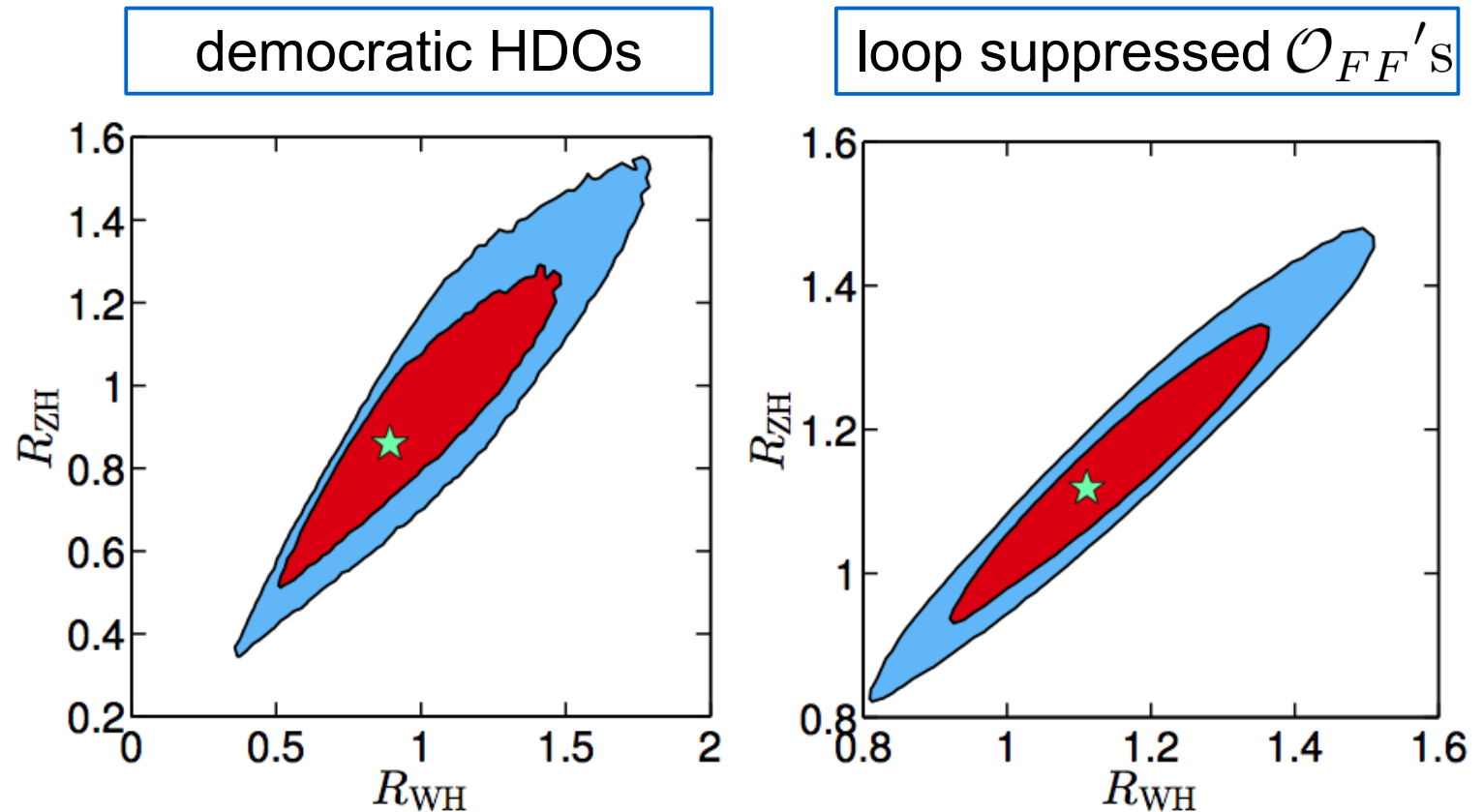
where $\beta_i = \alpha_i v^2 / \Lambda^2$

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \sum_i \frac{\alpha_i}{\Lambda^{n_i}} \mathcal{O}_i$$



Results associated production

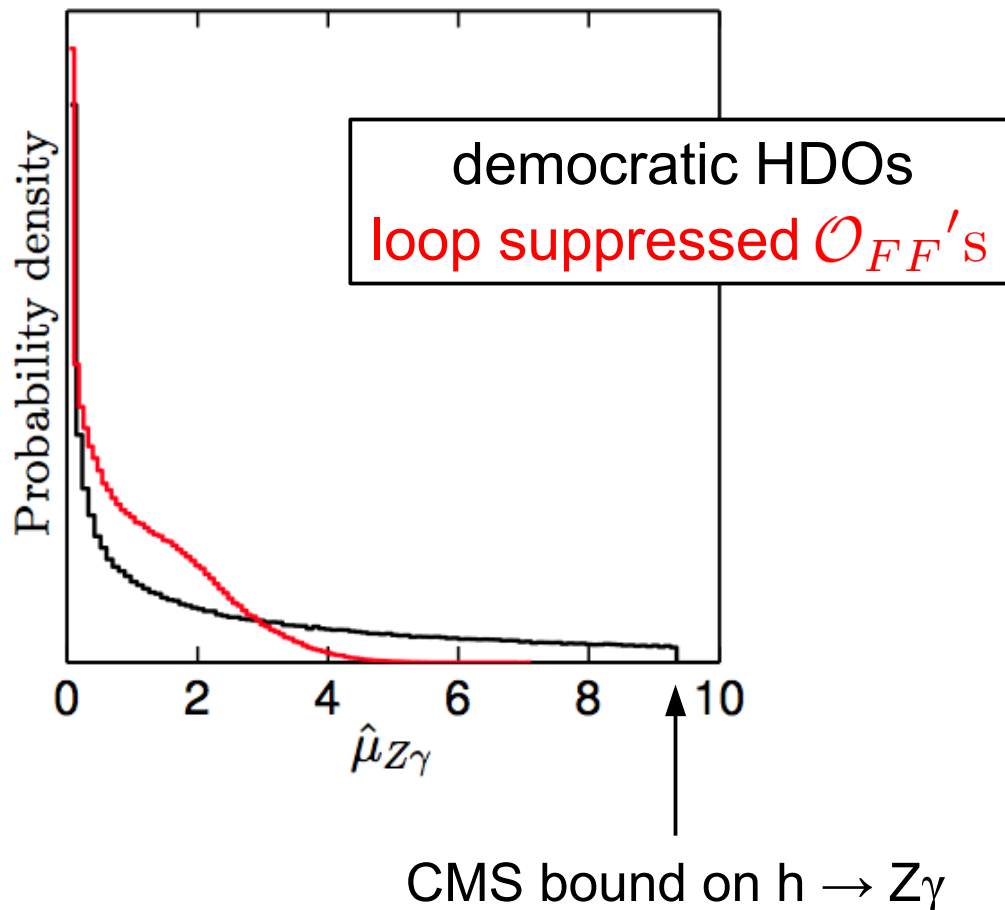
$$R_{\text{VH}} = \frac{\sigma_{\text{VH}}}{\sigma_{\text{VH}}^{\text{SM}}}$$



- tensorial couplings \rightarrow sizeable change in the rescaling of WH and ZH
- we plead for a clear separation of WH and ZH in the LHC Higgs results

Results

$h \rightarrow Z\gamma$



- possible large deviations from the SM value:
comes from tensorial coupling $\zeta_{Z\gamma}$
- future measurements of $h \rightarrow Z\gamma$
 \Rightarrow important constraints on our HDO parameters

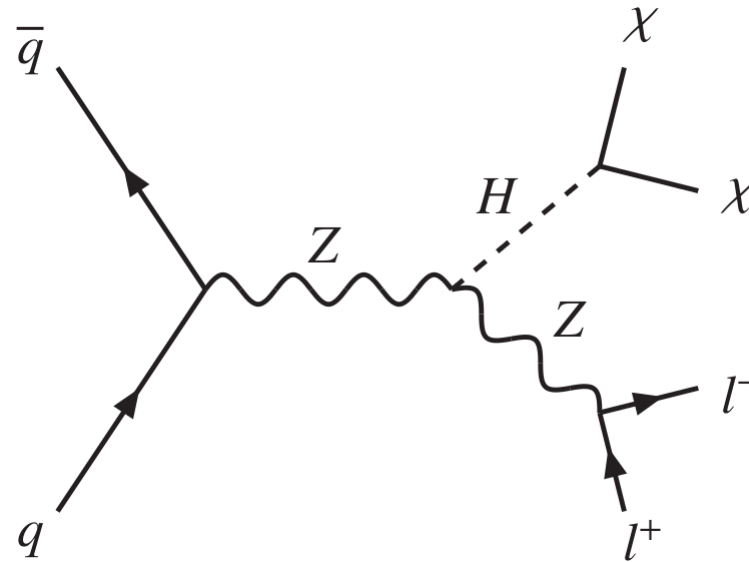
Conclusion

- overall, the observed Higgs boson seems very SM-like
(but still waiting for updates, especially in fermionic channels)
- precision era in Higgs physics has only just begun
however Higgs results are already a unique probe of New Physics
- model-independent studies as a first step in the study of the
implications of the new boson
→ time has come to fully explore the consequences for BSM models

Searches for invisible decays of the Higgs boson

ATLAS

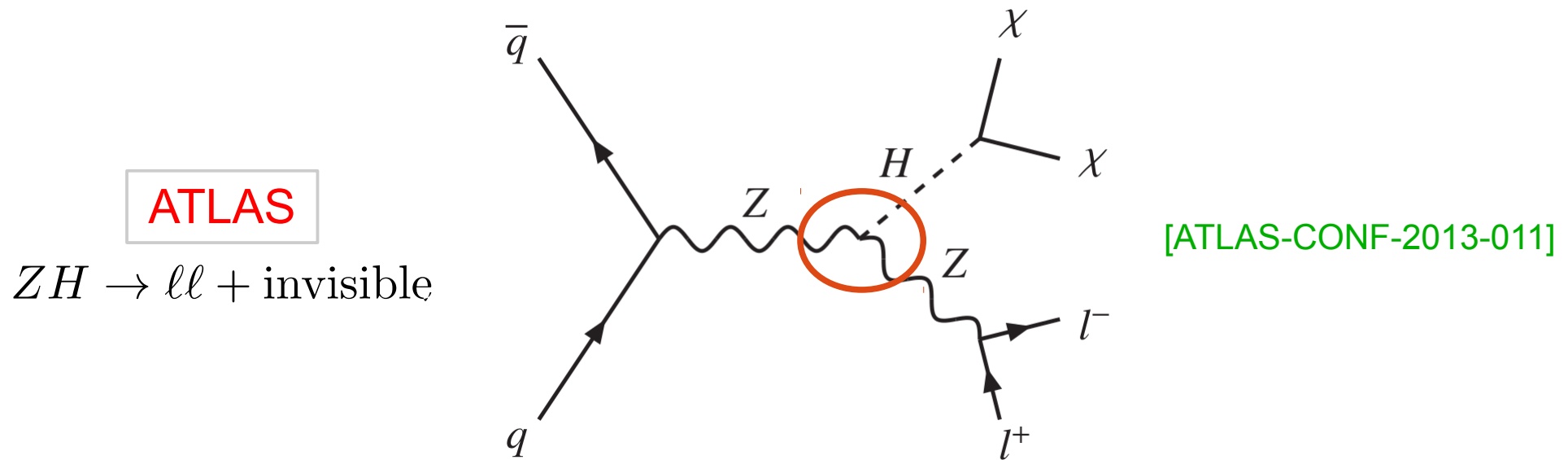
$ZH \rightarrow \ell\ell + \text{invisible}$



[ATLAS-CONF-2013-011]

$$\mathcal{B}(H \rightarrow \text{inv.}) < 0.65 \text{ at } 95\% \text{ CL}$$

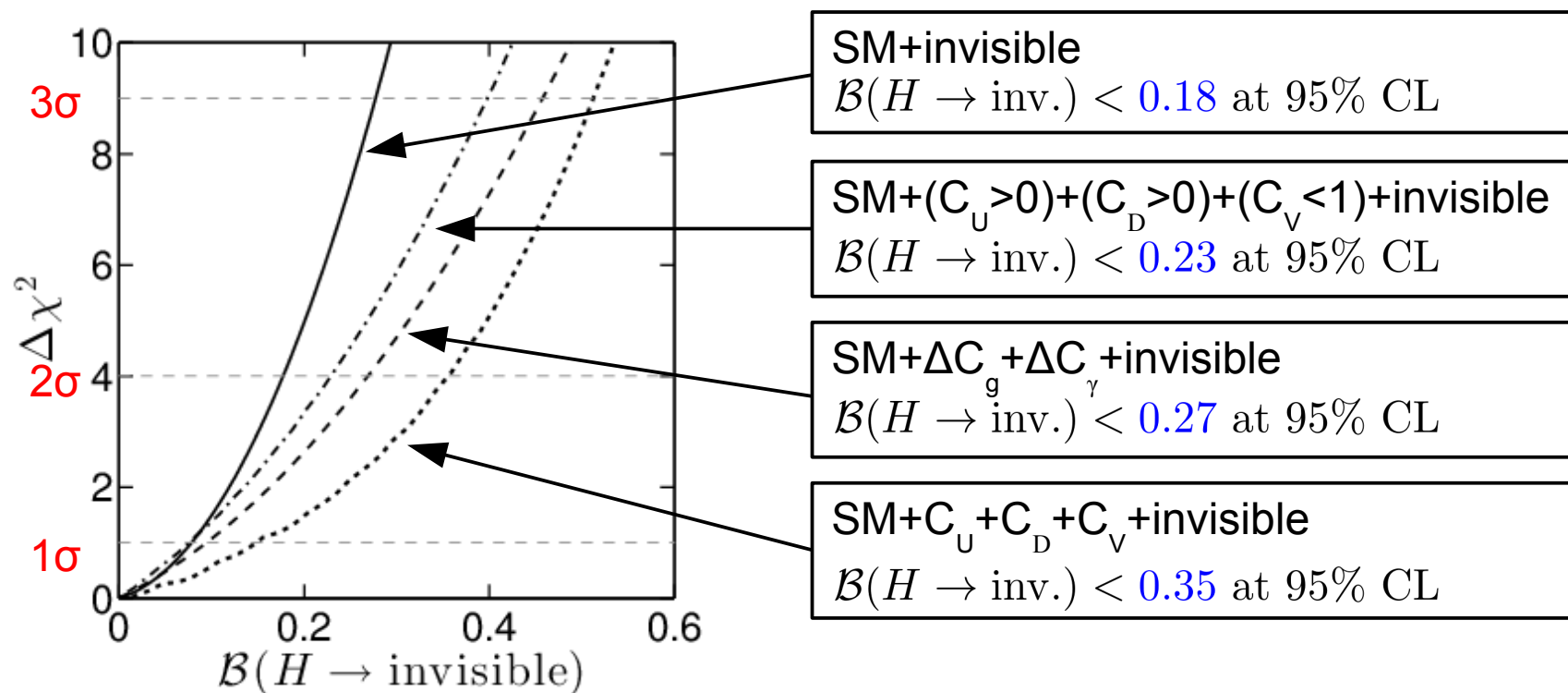
Searches for invisible decays of the Higgs boson



$$C_V^2 \mathcal{B}(H \rightarrow \text{inv.}) < 0.65 \text{ at } 95\% \text{ CL}$$

see also earlier studies based on e.g. monojet searches [Djouadi *et al.* '12]
...but so far global fits are more constraining

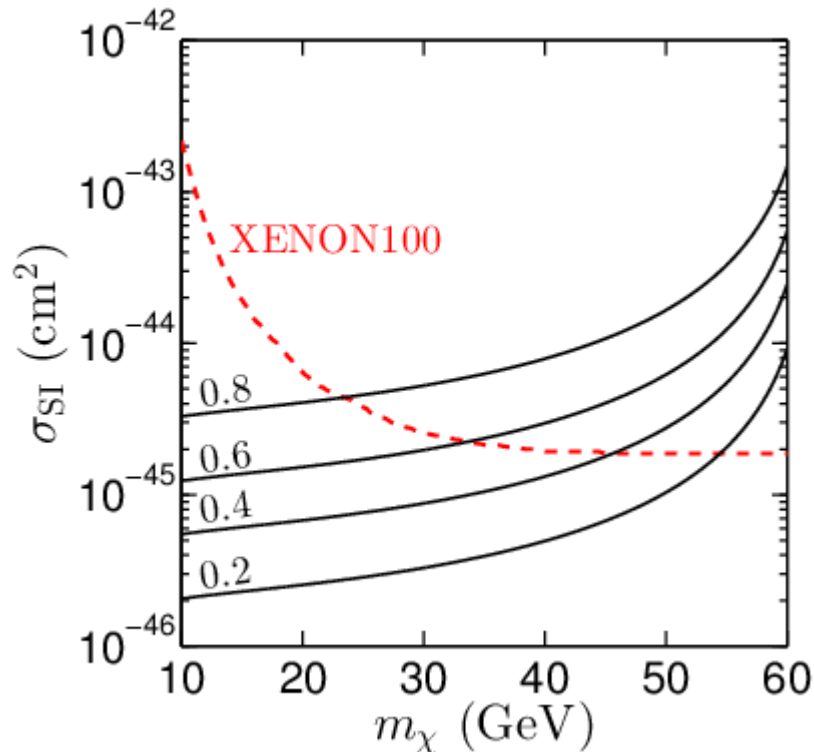
Invisible decays of the Higgs boson



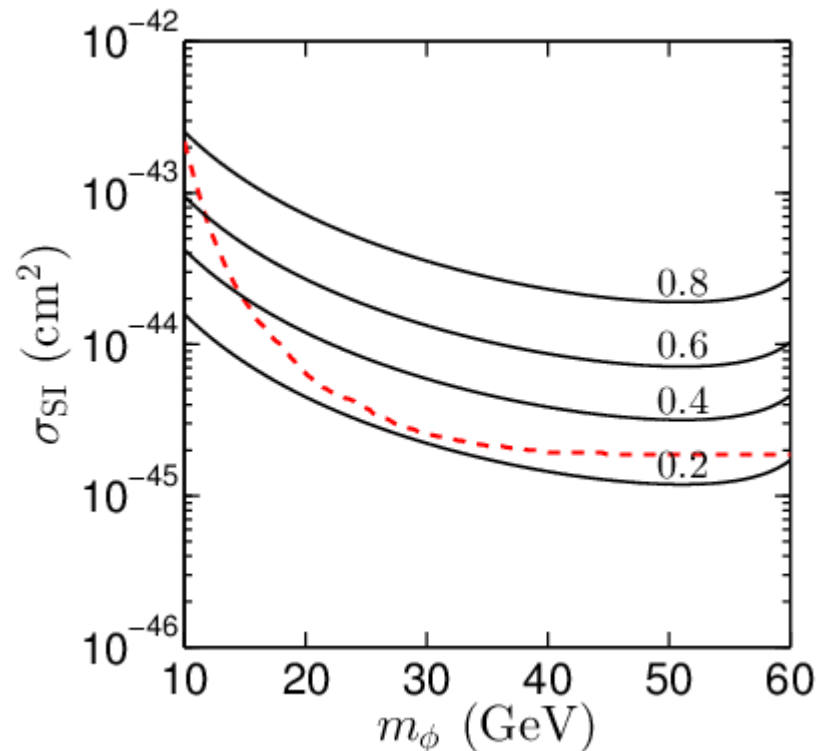
Invisible decays of the Higgs boson and dark matter

if invisible = dark matter:
interplay between direct searches and $H \rightarrow \text{invisible}$

Majorana dark matter



scalar dark matter



Tensorial couplings and kinematic distributions

The amplitude associated to a hVV vertex (with the V 's possibly off-shell) is in general

$$\mathcal{M}(hVV)^{\lambda_1, \lambda_2} = e_{\lambda_1}^{\mu(*)} e_{\lambda_2}^{\nu(*)} \left(i a_V \lambda_V^{\text{SM}} g^{\mu\nu} - i 2 \zeta_V q_1 \cdot q_2 \left[g^{\mu\nu} - \frac{q_1^\mu q_2^\nu}{q_1 \cdot q_2} \right] \right), \quad (4.1)$$

Experimental data we use

ATLAS

Channel	Signal strength μ	m_H (GeV)	Production mode			
			ggF	VBF	VH	ttH
$H \rightarrow \gamma\gamma$ (4.8 fb ⁻¹ at 7 TeV + 20.7 fb ⁻¹ at 8 TeV) [1, 2]						
$\mu(\text{ggF} + \text{ttH}, \gamma\gamma)$	1.60 ± 0.41	125.5	100%	–	–	–
$\mu(\text{VBF} + \text{VH}, \gamma\gamma)$	1.94 ± 0.82	125.5	–	60%	40%	–
$H \rightarrow ZZ$ (4.6 fb ⁻¹ at 7 TeV + 20.7 fb ⁻¹ at 8 TeV) [3, 2]						
$\mu(\text{ggF} + \text{ttH}, ZZ)$	1.50 ± 0.50	125.5	100%	–	–	–
$\mu(\text{VBF} + \text{VH}, ZZ)$	1.50 ± 2.52	125.5	–	60%	40%	–
$H \rightarrow WW$ (4.6 fb ⁻¹ at 7 TeV + 20.7 fb ⁻¹ at 8 TeV) [4, 5]						
$\mu(\text{ggF} + \text{ttH}, WW)$	0.79 ± 0.35	125.5	100%	–	–	–
$\mu(\text{VBF} + \text{VH}, WW)$	1.71 ± 0.76	125.5	–	60%	40%	–
$H \rightarrow b\bar{b}$ (4.7 fb ⁻¹ at 7 TeV + 13.0 fb ⁻¹ at 8 TeV) [6, 2]						
VH tag	-0.39 ± 1.02	125.5	–	–	100%	–
$H \rightarrow \tau\tau$ (4.6 fb ⁻¹ at 7 TeV + 13.0 fb ⁻¹ at 8 TeV) [2]						
$\mu(\text{ggF} + \text{ttH}, \tau\tau)$	2.31 ± 1.61	125.5	100%	–	–	–
$\mu(\text{VBF} + \text{VH}, \tau\tau)$	-0.20 ± 1.06	125.5	–	60%	40%	–

Table 1: ATLAS results, as employed in this analysis. The following correlations are included in the fit: $\rho_{\gamma\gamma} = -0.27$, $\rho_{ZZ} = -0.46$, $\rho_{WW} = -0.18$, $\rho_{\tau\tau} = -0.49$.

Experimental data we use

CMS

Channel	Signal strength μ	m_H (GeV)	Production mode			
			ggF	VBF	VH	ttH
$H \rightarrow \gamma\gamma$ (5.1 fb $^{-1}$ at 7 TeV + 19.6 fb $^{-1}$ at 8 TeV) [7, 8]						
$\mu(\text{ggF} + \text{ttH}, \gamma\gamma)$	0.46 ± 0.40	125.7	100%	–	–	–
$\mu(\text{VBF} + \text{VH}, \gamma\gamma)$	1.68 ± 0.87	125.7	–	60%	40%	–
$H \rightarrow ZZ$ (5.1 fb $^{-1}$ at 7 TeV + 19.6 fb $^{-1}$ at 8 TeV) [9]						
$\mu(\text{ggF} + \text{ttH}, ZZ)$	0.98 ± 0.46	125.8	100%	–	–	–
$\mu(\text{VBF} + \text{VH}, ZZ)$	1.07 ± 2.37	125.8	–	60%	40%	–
$H \rightarrow WW$ (up to 4.9 fb $^{-1}$ at 7 TeV + 19.5 fb $^{-1}$ at 8 TeV) [10, 11, 12, 8]						
$\mu(\text{ggF} + \text{ttH}, WW)$	0.78 ± 0.23	125.7	100%	–	–	–
$\mu(\text{VBF} + \text{VH}, WW)$	0.33 ± 0.70	125.7	–	60%	40%	–
$H \rightarrow b\bar{b}$ (up to 5.0 fb $^{-1}$ at 7 TeV + 12.1 fb $^{-1}$ at 8 TeV) [13, 14, 8]						
VH tag	$1.31^{+0.68}_{-0.61}$	125.7	–	–	100%	–
ttH tag	$-0.15^{+2.82}_{-2.90}$	125.7	–	–	–	100%
$H \rightarrow \tau\tau$ (4.9 fb $^{-1}$ at 7 TeV + 19.4 fb $^{-1}$ at 8 TeV) [15, 8]						
$\mu(\text{ggF} + \text{ttH}, \tau\tau)$	0.67 ± 0.79	125.7	100%	–	–	–
$\mu(\text{VBF} + \text{VH}, \tau\tau)$	1.59 ± 0.83	125.7	–	60%	40%	–

Table 2: CMS results, as employed in this analysis. The following correlations are included in the fit: $\rho_{\gamma\gamma} = -0.48$, $\rho_{ZZ} = -0.73$, $\rho_{WW} = -0.21$, $\rho_{\tau\tau} = -0.47$.

Experimental data we use

Tevatron

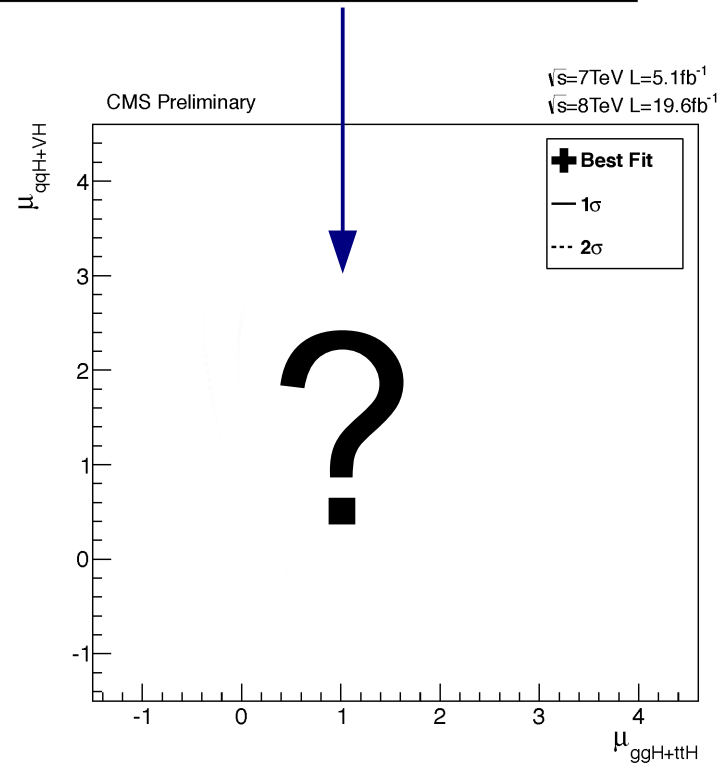
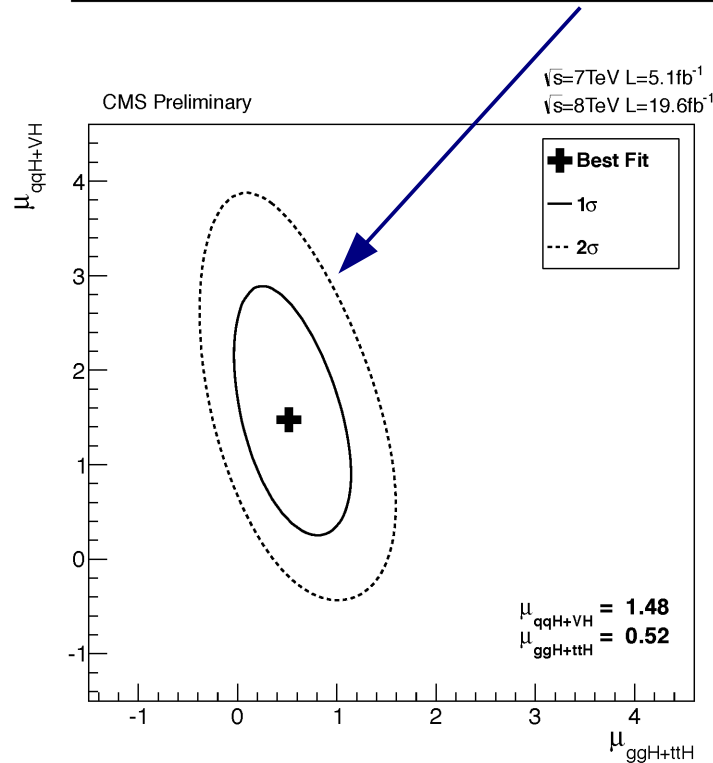
Channel	Signal strength μ	m_H (GeV)	Production mode			
			ggF	VBF	VH	ttH
$H \rightarrow \gamma\gamma$ [17]						
Combined	$5.97^{+3.39}_{-3.12}$	125	78%	5%	17%	–
$H \rightarrow WW$ [17]						
Combined	$0.94^{+0.85}_{-0.83}$	125	78%	5%	17%	–
$H \rightarrow b\bar{b}$ [17]						
VH tag	$1.59^{+0.69}_{-0.72}$	125	–	–	100%	–

Table 3: Tevatron results for up to 10 fb^{-1} at $\sqrt{s} = 1.96 \text{ TeV}$, as employed in this analysis.

- Tevatron $H \rightarrow \tau\tau$ is omitted (large uncertainties)
- $H \rightarrow \gamma\gamma$ and $H \rightarrow WW$ are approximated as inclusive searches (ratio of inclusive cross sections for $p\bar{p}$ collisions at 2 TeV)

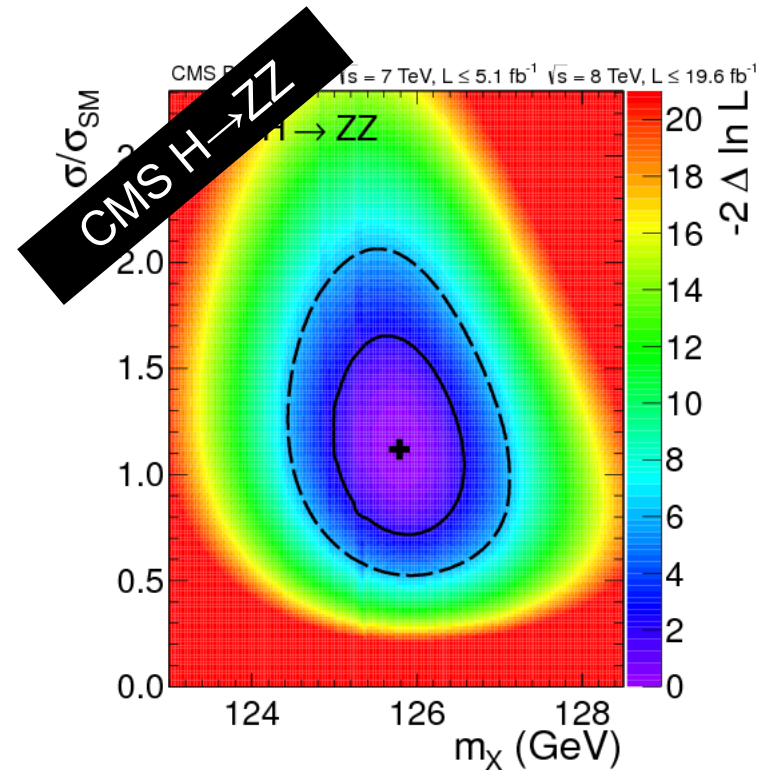
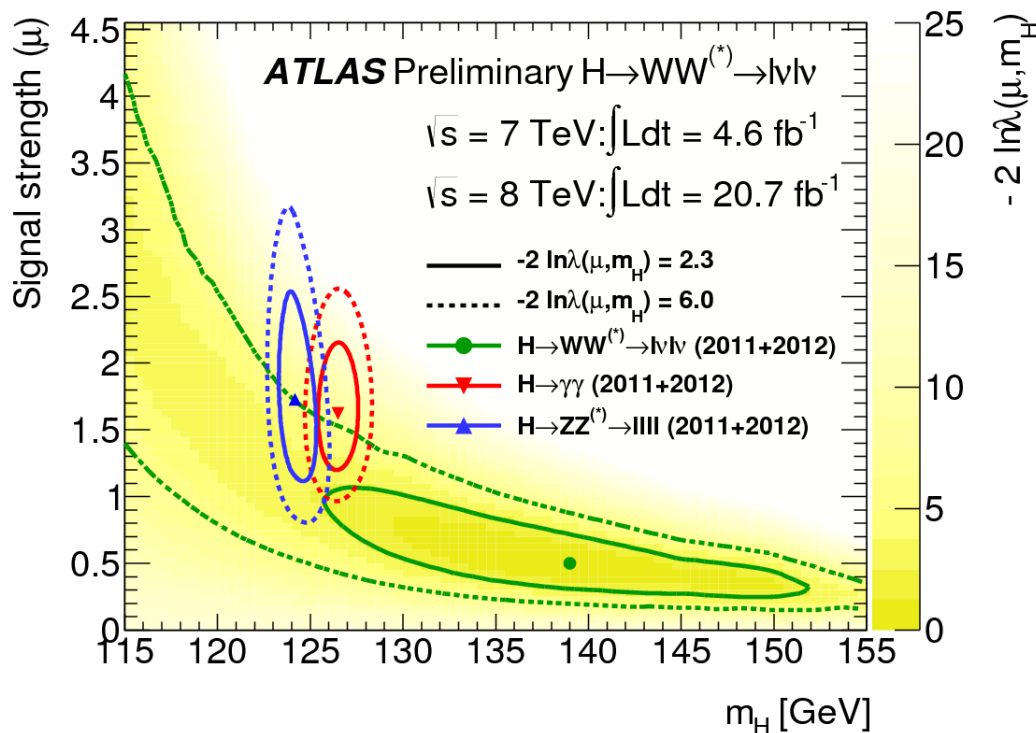
A word on CMS $H \rightarrow \gamma\gamma$

	MVA analysis (at $m_H=125$ GeV)	cut-based analysis (at $m_H=124.5$ GeV)
7 TeV	$1.69^{+0.65}_{-0.59}$	$2.27^{+0.80}_{-0.74}$
8 TeV	$0.55^{+0.29}_{-0.27}$	$0.93^{+0.34}_{-0.32}$
7 + 8 TeV	$0.78^{+0.28}_{-0.26}$	$1.11^{+0.32}_{-0.30}$



Dependence on m_H

- we would like to treat the Higgs mass as a nuisance parameter
- a priori important for the two high resolution channels ($H \rightarrow ZZ$ and $H \rightarrow \gamma\gamma$)



- unfortunately impossible to use together with the 2D μ information