

# Azimuthal correlations between high $p_T$ heavy flavour (jets)

*9<sup>th</sup> International workshop on high  $p_T$  Physics at LHC*

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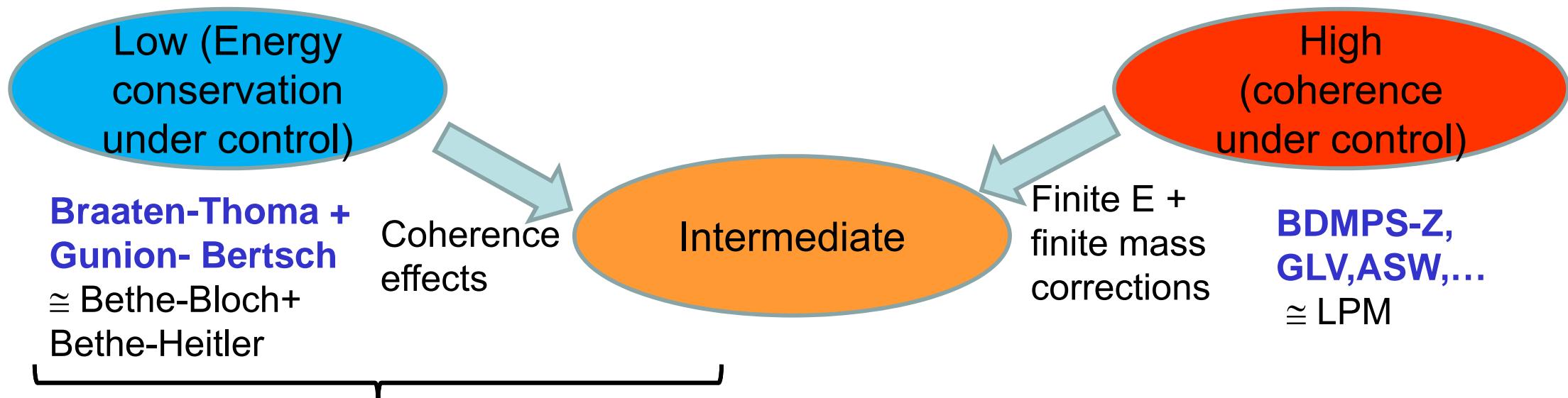
with

**M. Nahrgang, J. Aichelin, and K. Werner**

Inspired from arxiv 1305.6544

# Motivation and context

- Most of the *interesting* HF observables so far: located at *intermediate*  $p_T$
- Intermediate  $p_T$ : hope that pQCD (or pQCD inspired models) apply (as compared to low  $p_T$ )
- Intermediate  $p_T$ : mass effect still present and thus hope to learn something more as compared to large  $p_T$



Approach pursued in our **models**... Unfortunately too many of them

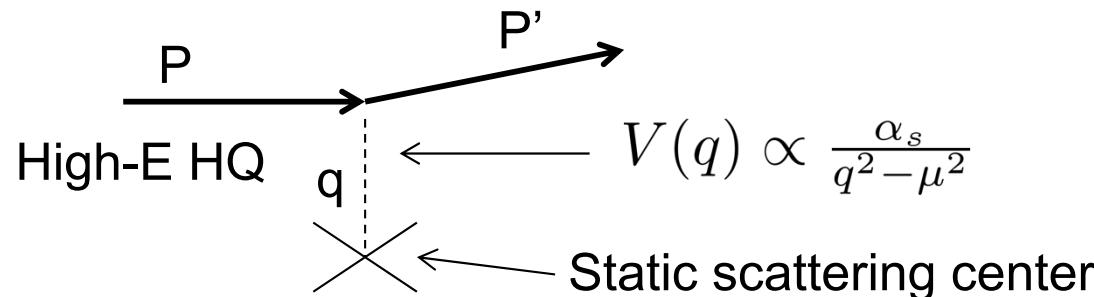
=>Need for falsification (more observables; IQCD): **Azimuthal correlations** ?

# Insufficient control on energy loss theory in QCD

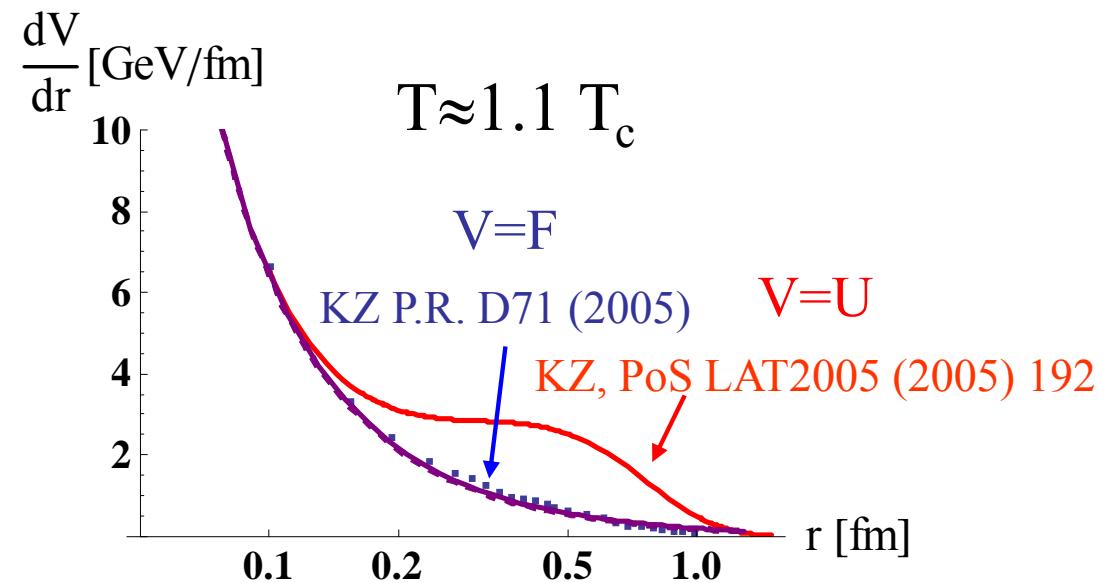
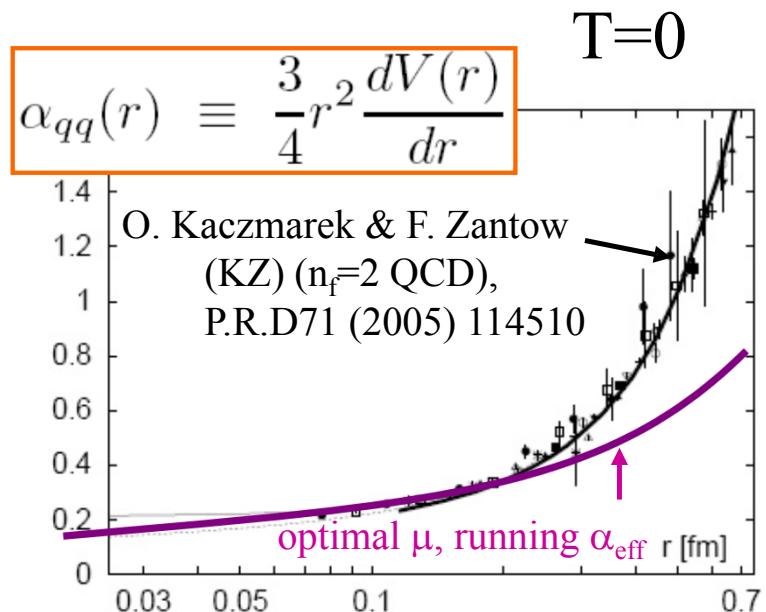
Basic ingredient in the derivation of QED collisional Eloss; transverse force

In QCD: non perturbative « corrections » even at large HQ energy

In most models:



Lattice QCD :



Significant r-tail in the transverse force acting on the high E HQ

# Our basic ingredients for HQ energy loss

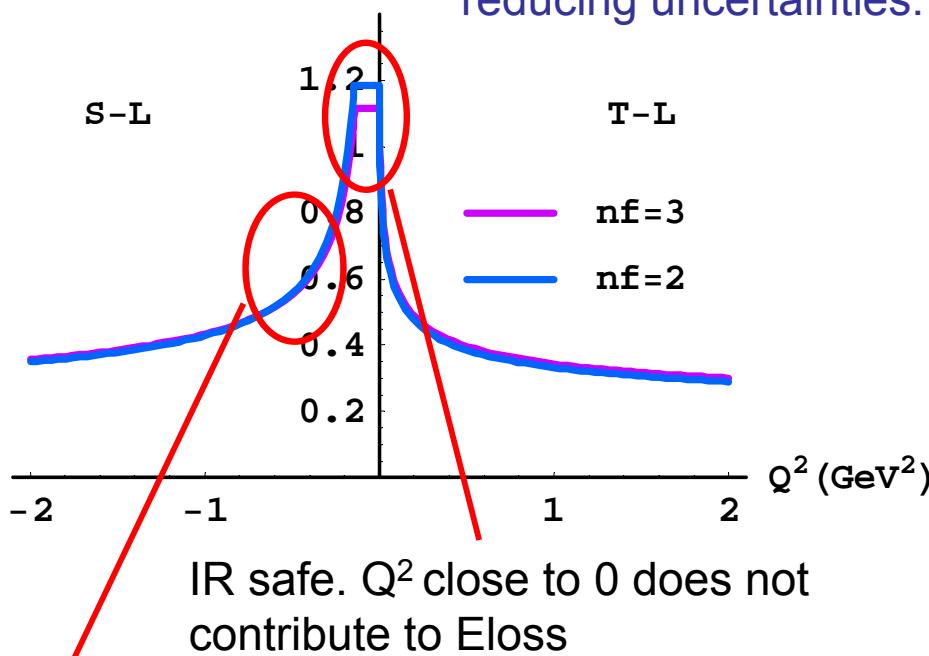
## Elastic

Motivation: Even a fast parton with the largest momentum  $P$  will undergo collisions with moderate  $q$  exchange and large  $\alpha_s(Q^2)$ . The running aspect of the coupling constant has been “forgotten/neglected” in most of approaches

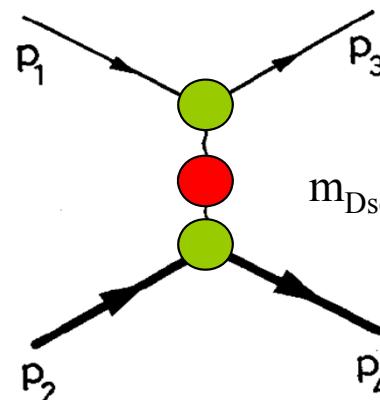
Effective  $\alpha_s(Q^2)$  (Dokshitzer 95, Brodsky 02)

$$\frac{1}{Q_u} \int_{|Q^2| \leq Q_u^2} dQ \alpha_s(Q^2) \approx 0.5$$

“Universality constrain” (Dokshitzer 02) helps reducing uncertainties:



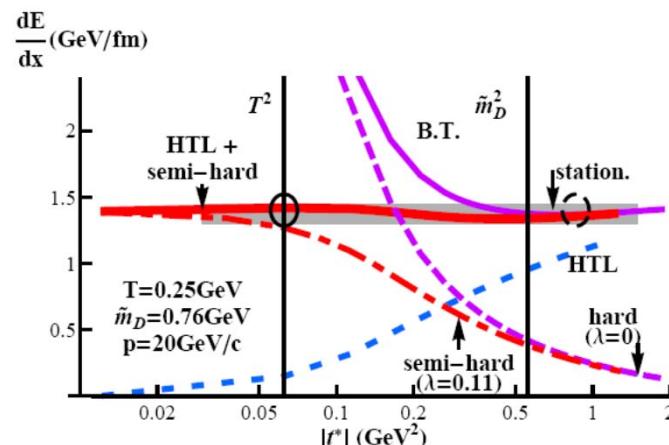
Large values for intermediate momentum-transfer => larger cross section



$$m_{\text{Dself}}^2(T) = (1+n_f/6) 4\pi \alpha_{\text{eff}}(m_{\text{Dself}}^2) T^2$$

+ u and s channels

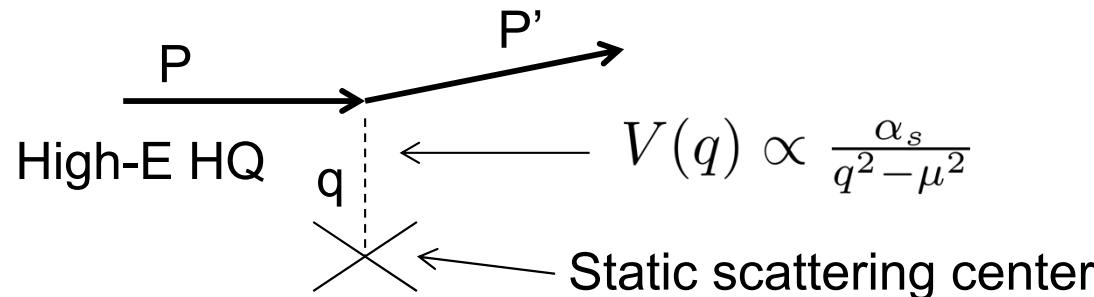
One gluon exchange effective propagator, designed in order to guarantee maximal insensitivity of  $dE/dx$  in Braaten-Thomas scheme



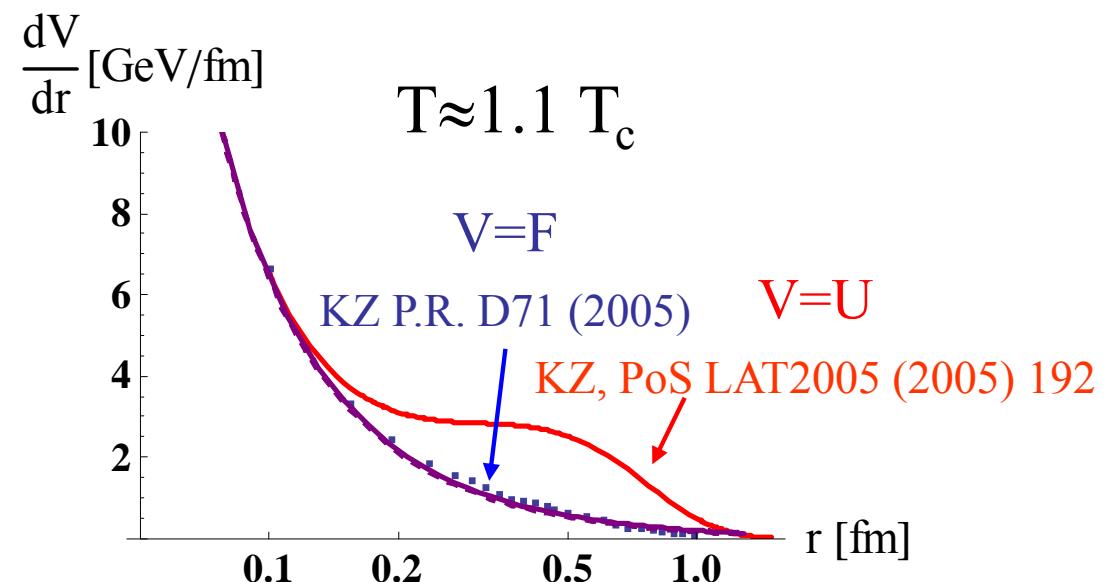
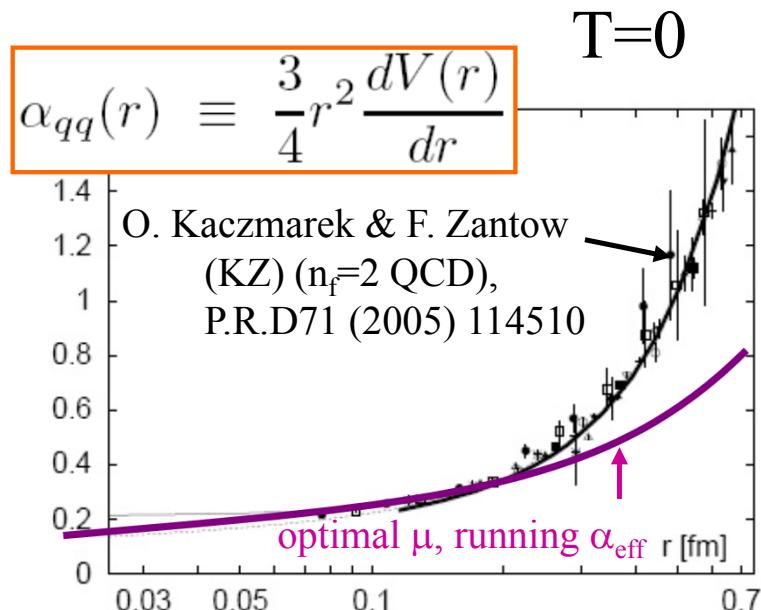
# Insufficient control on energy loss theory

Non perturbative « corrections » even at large HQ energy

In most models:



Lattice QCD :



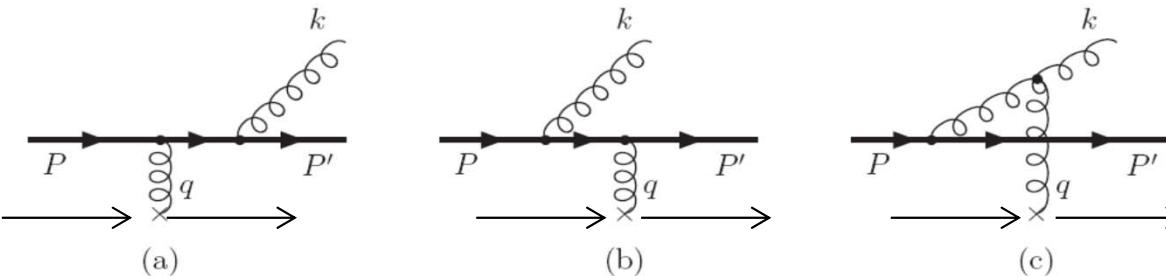
Our force is close to the one extracted from the free energy as a potential

=> Allow for some global rescaling of the rates: “K” fixed on experiment 5

# Our basic ingredients for HQ energy loss

## Incoherent Induced Radiative

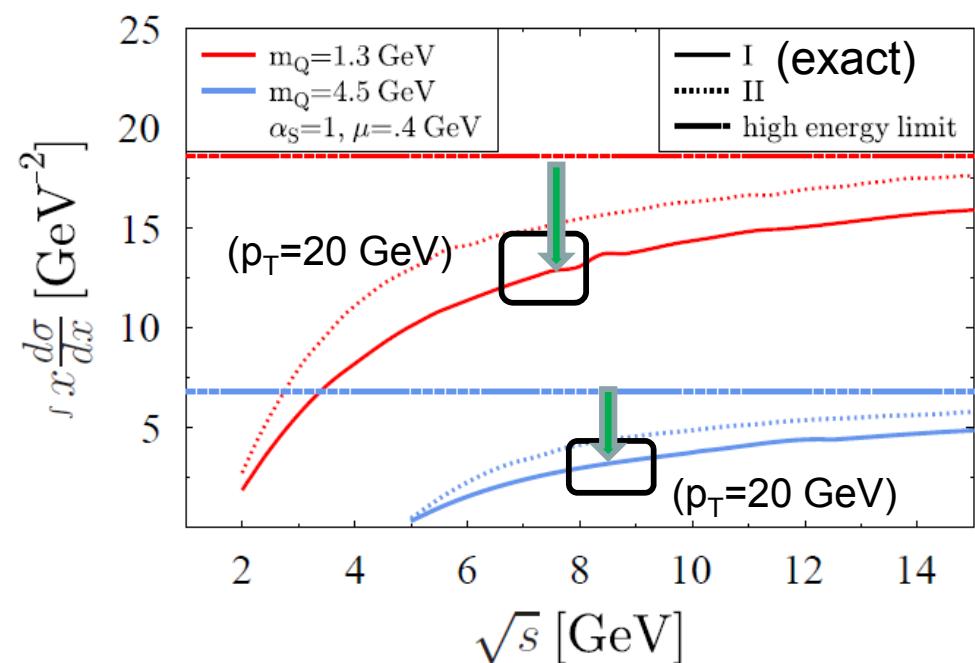
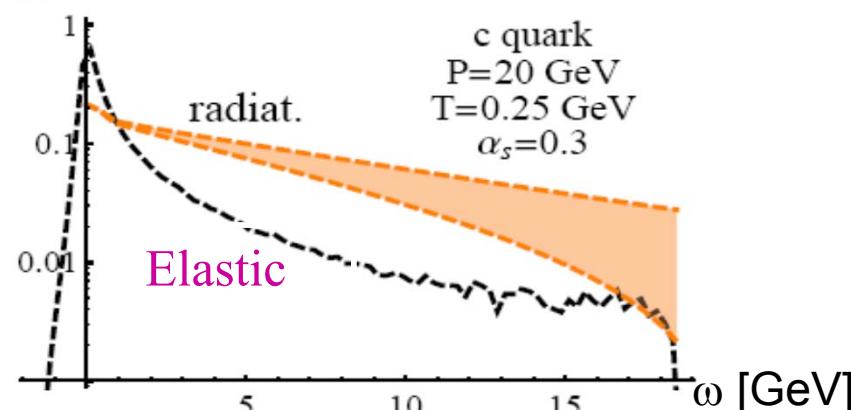
Generalized Gunion-Bertsch for finite mass, finite energy and dynamical light partons arxiv 1307.5270



$$\omega \frac{d^3 \sigma_{\text{rad}}^{x \ll 1}}{d\omega d^2 k_\perp dq_\perp^2} = N_c \alpha_s (1 - x) \times \frac{J_{\text{QCD}}^2}{\omega^2} \times \boxed{\frac{d\sigma_{\text{el}}^{Qq}}{dq_\perp^2}}$$

$$\frac{J_{\text{QCD}}^2}{\omega^2} = \left( \frac{\vec{k}_\perp}{k_\perp^2 + x^2 M^2 + (1-x)m_g^2} - \frac{\vec{k}_\perp - \vec{q}_\perp}{(\vec{k}_\perp - \vec{q}_\perp)^2 + x^2 M^2 + (1-x)m_g^2} \right)^2$$

$$\frac{dP(\omega)}{dz} [\text{fm}^{-1}]$$



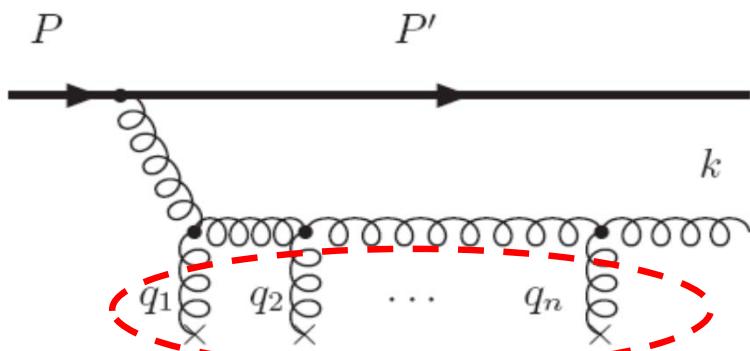
Finite energy lead to strong reduction of the radiative energy loss at intermediate  $p_T$

differential energy loss  $\omega$  per unit length ( $T, M, \dots$ ): big differences between the 2 contributions

# Our basic ingredients for HQ energy loss

## Coherent Induced Radiative

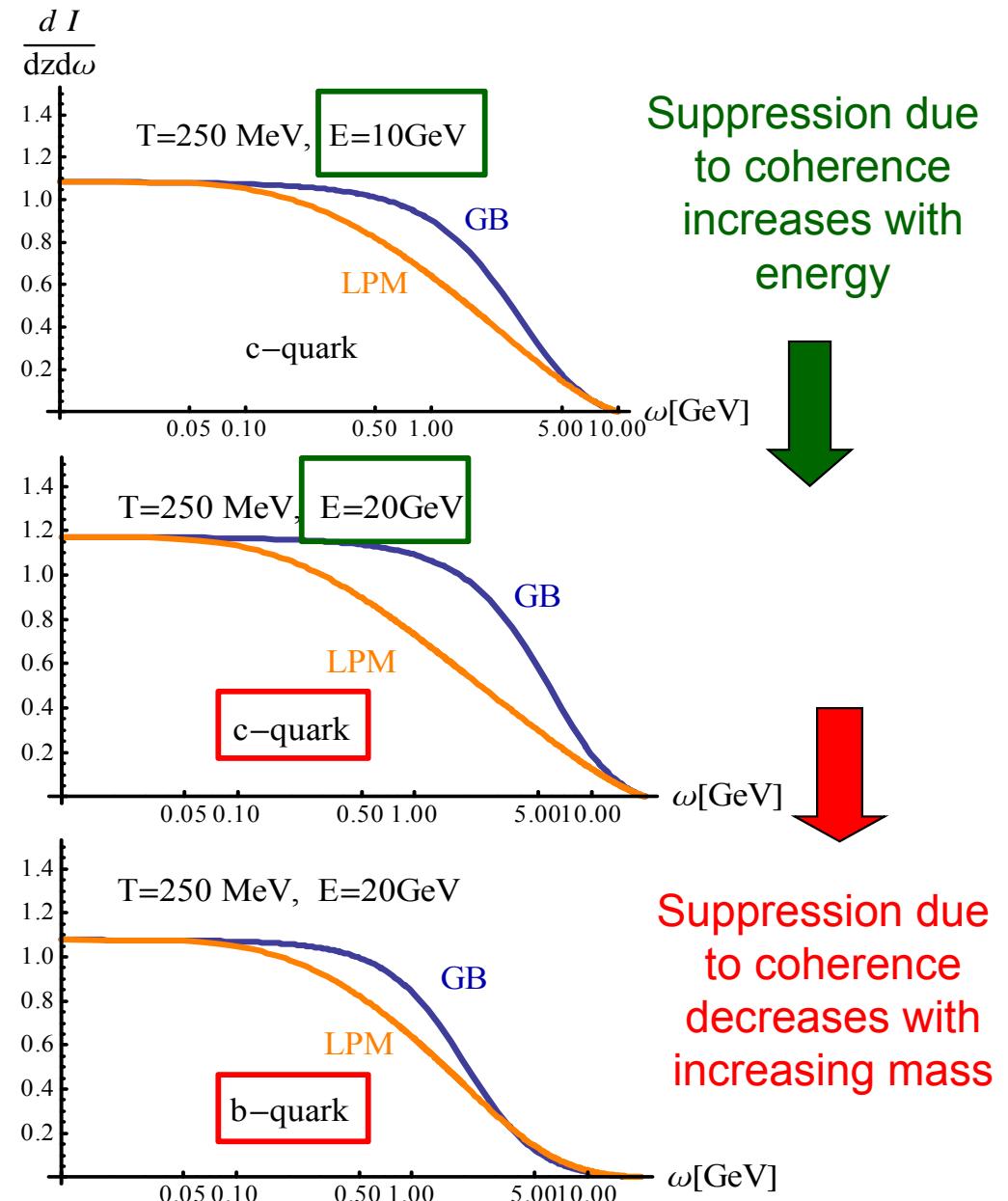
Formation time picture: for  $l_{f,mult} > \lambda$ , gluon is radiated coherently on a distance  $l_{f,mult}$



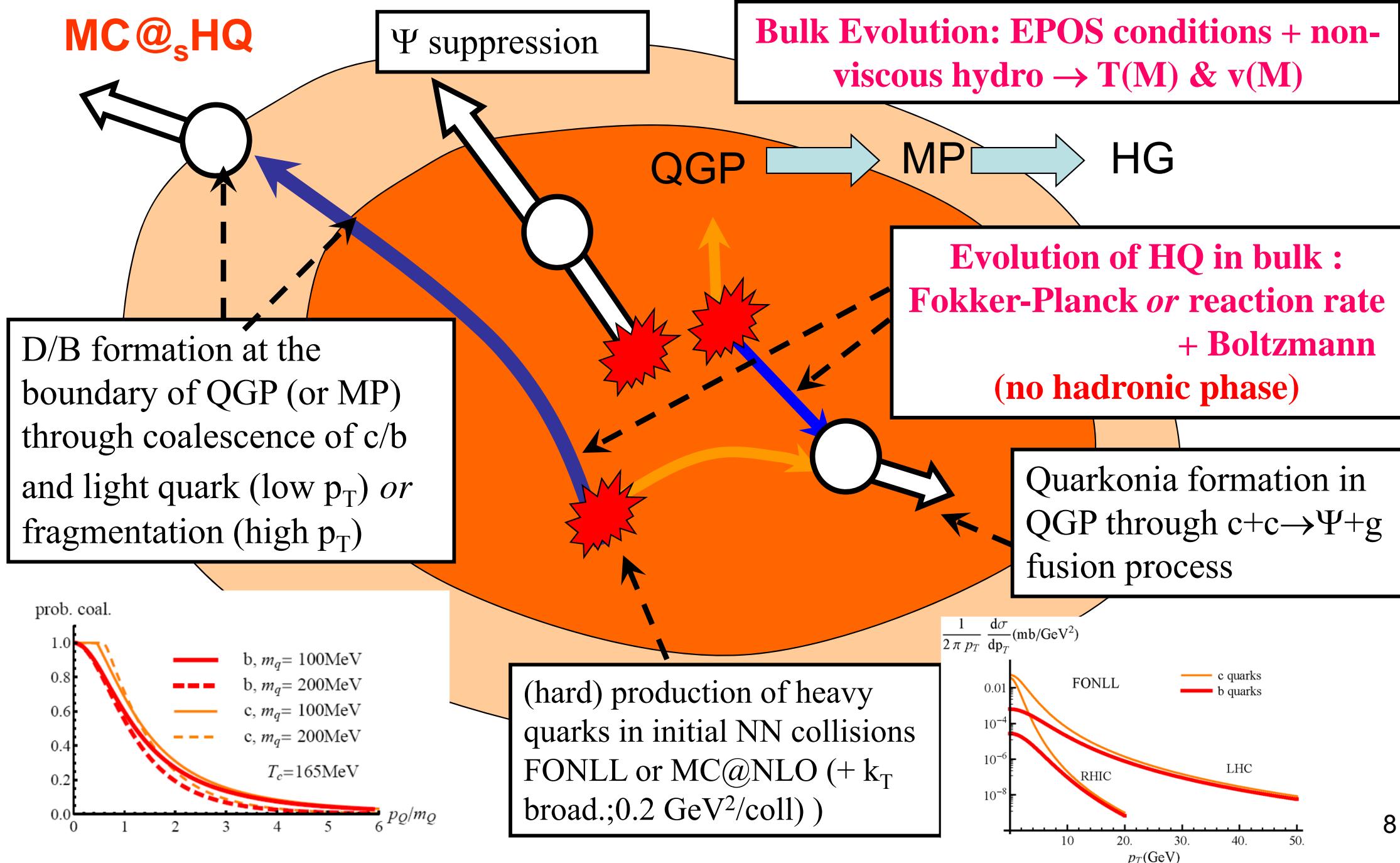
Model: all  $N_{coh}$  scatterers act as a single effective one with probability  $p_{Ncoh}(Q_\perp)$  obtained by convoluting individual probability of kicks

$$\frac{d^2 I_{\text{eff}}}{dz d\omega} \sim \frac{\alpha_s}{N_{coh} \tilde{\lambda}} \ln \left( 1 + \frac{N_{coh} \mu^2}{3 (m_g^2 + x^2 M^2 + \sqrt{\omega \hat{q}})} \right)$$

[arXiv:1209.0844] (Hard Probes 2012)

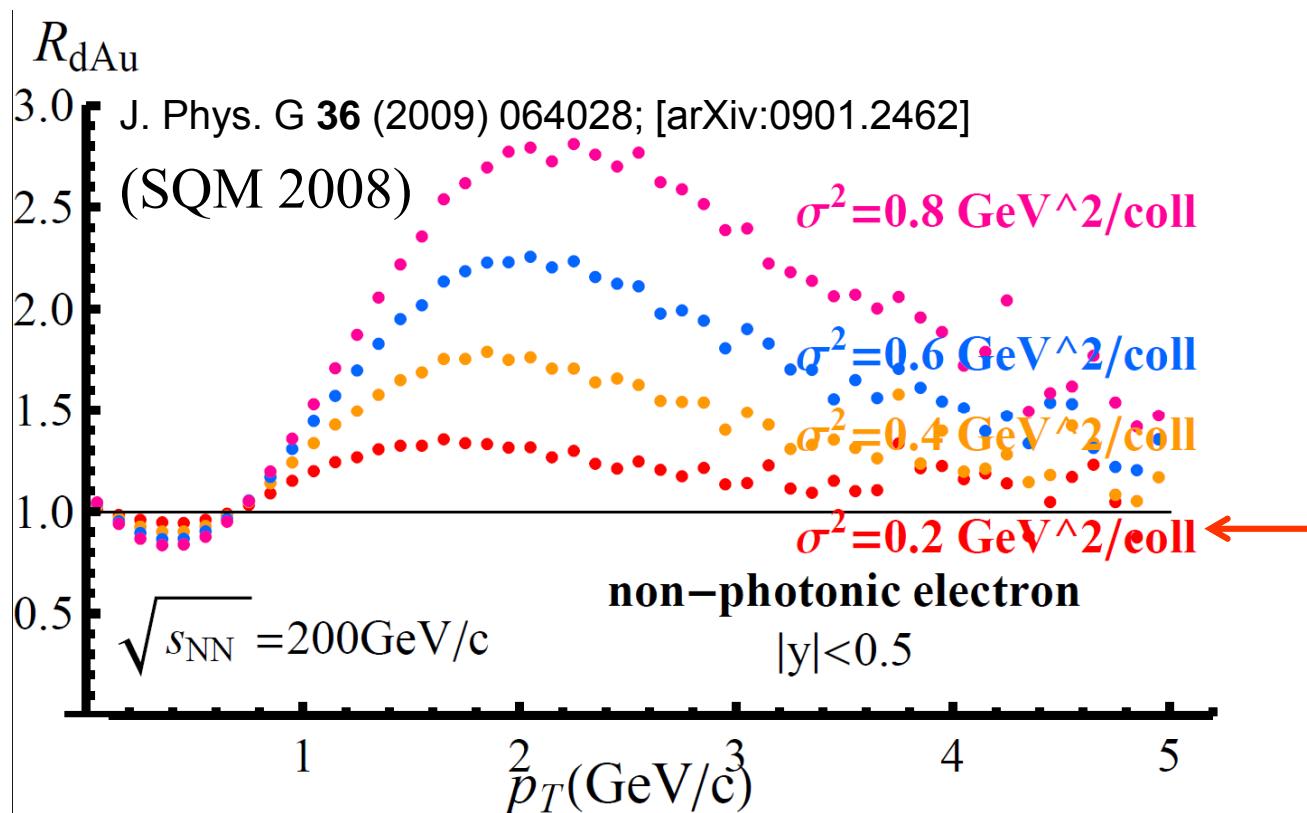


# Schematic view of « Monte Carlo @ Heavy Quark » generator



# A Lucky guess for the CNM effects ?

Inserting a phenomenological  $k_T$  broadening between HQ FONLL production and fragmentation:

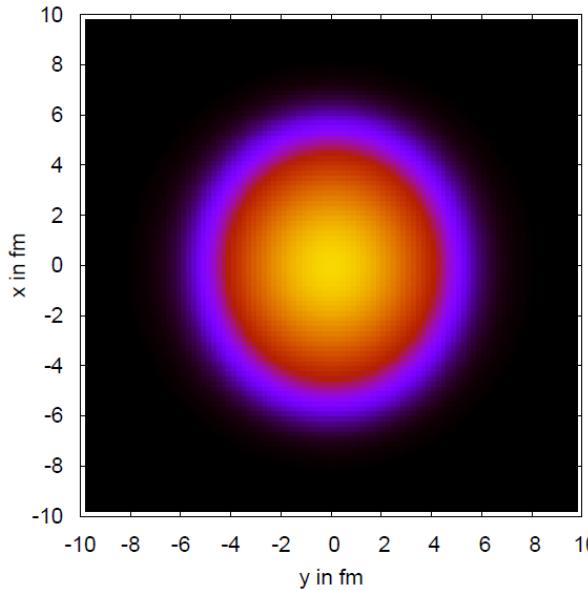


Selected values; good  
agreement with recent  
RHIC measurement

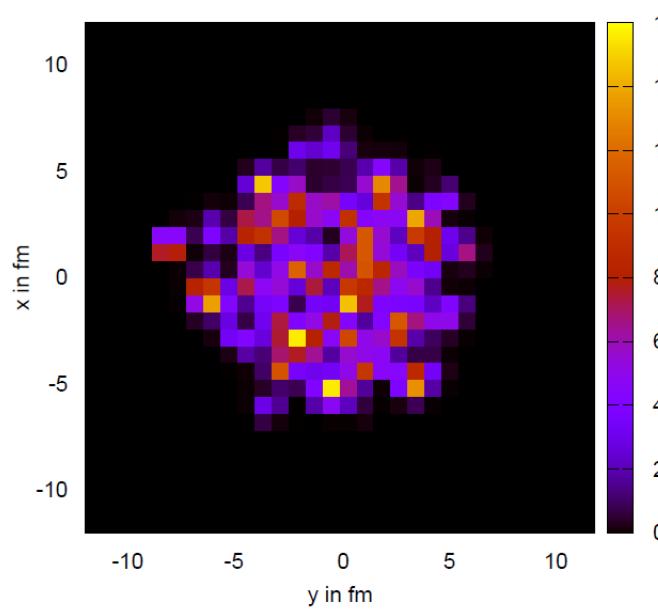
# EPOS as a background for MC@sHQ

EPOS: state of the art framework that encompass pp, pA and AA collisions

Initial energy density @ RHIC (central Au-Au)



Kolb Heinz (used previously)



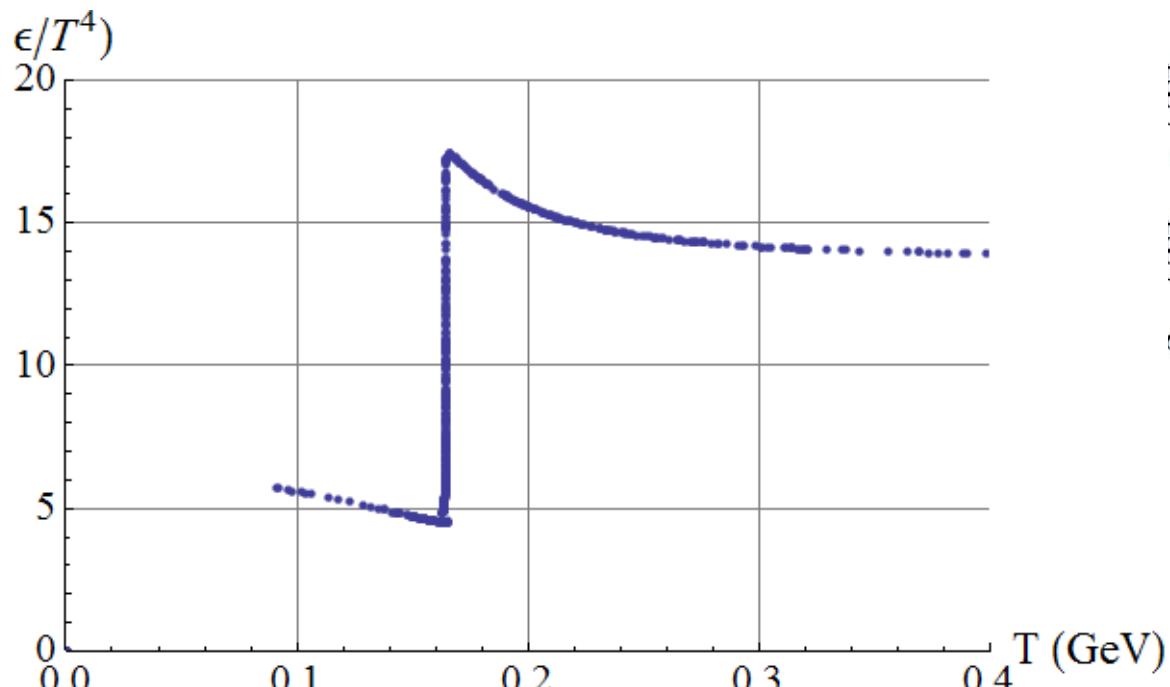
EPOS

Beware:  $\neq$  color scales

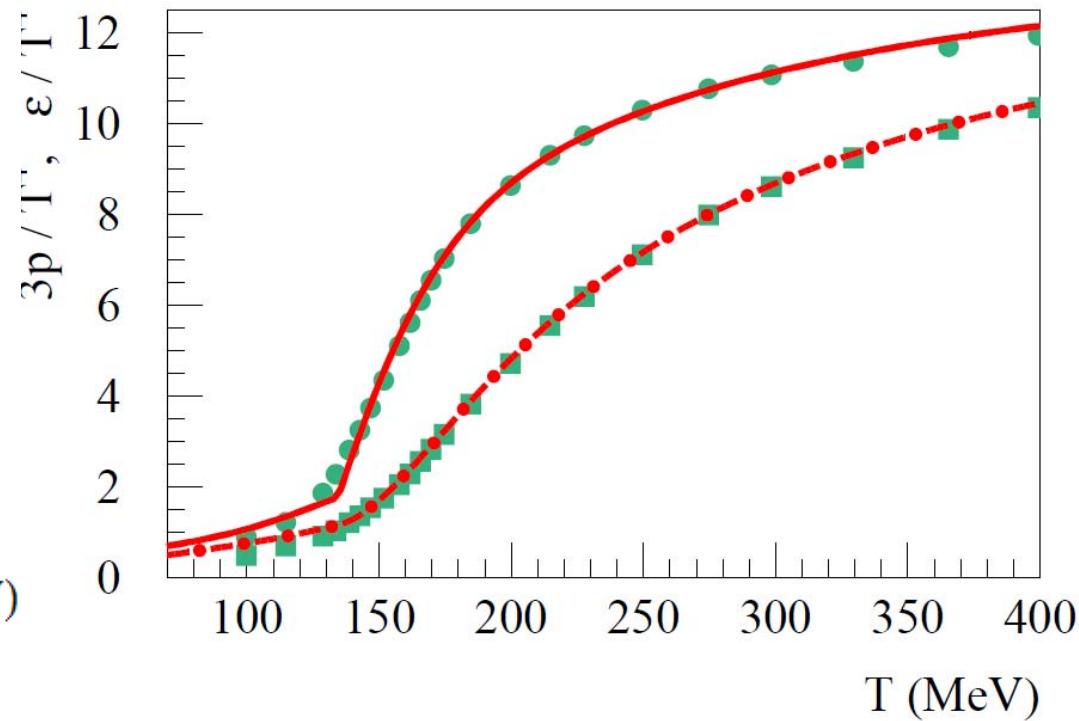
More realistic hydro and initial conditions => original HQ studies such as:

- 1) fluctuations in HQ observables (some HQ might « leak » through the « holes » in the QGP)
- 2) correlations between HF and light hadrons

# Large differences in the EOS !

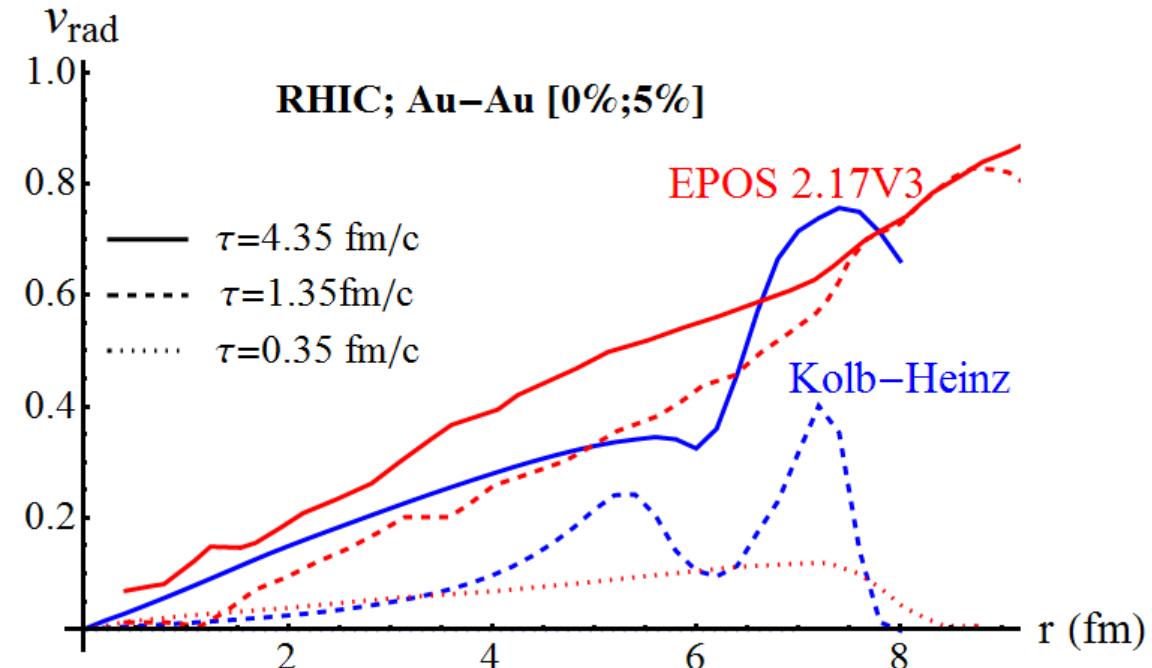
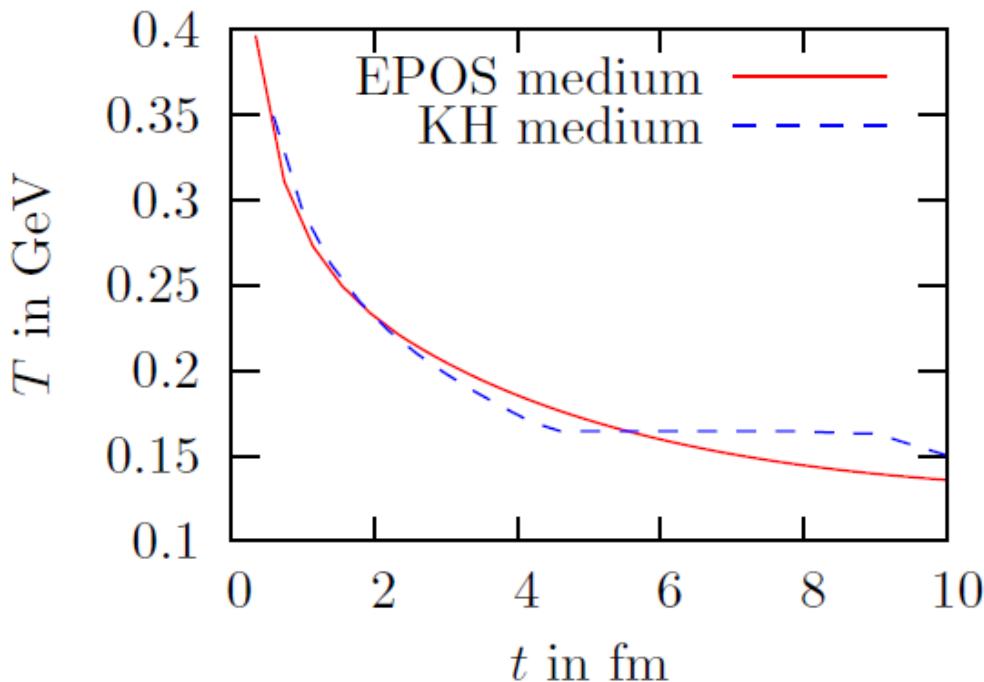


Kolb Heinz: bag model  
(1rst order transition  
btwn hadronic phase  
and massless partons)



EPOS2: fitted on the  
lattice data from the  
Wuppertal-Budapest  
collaboration

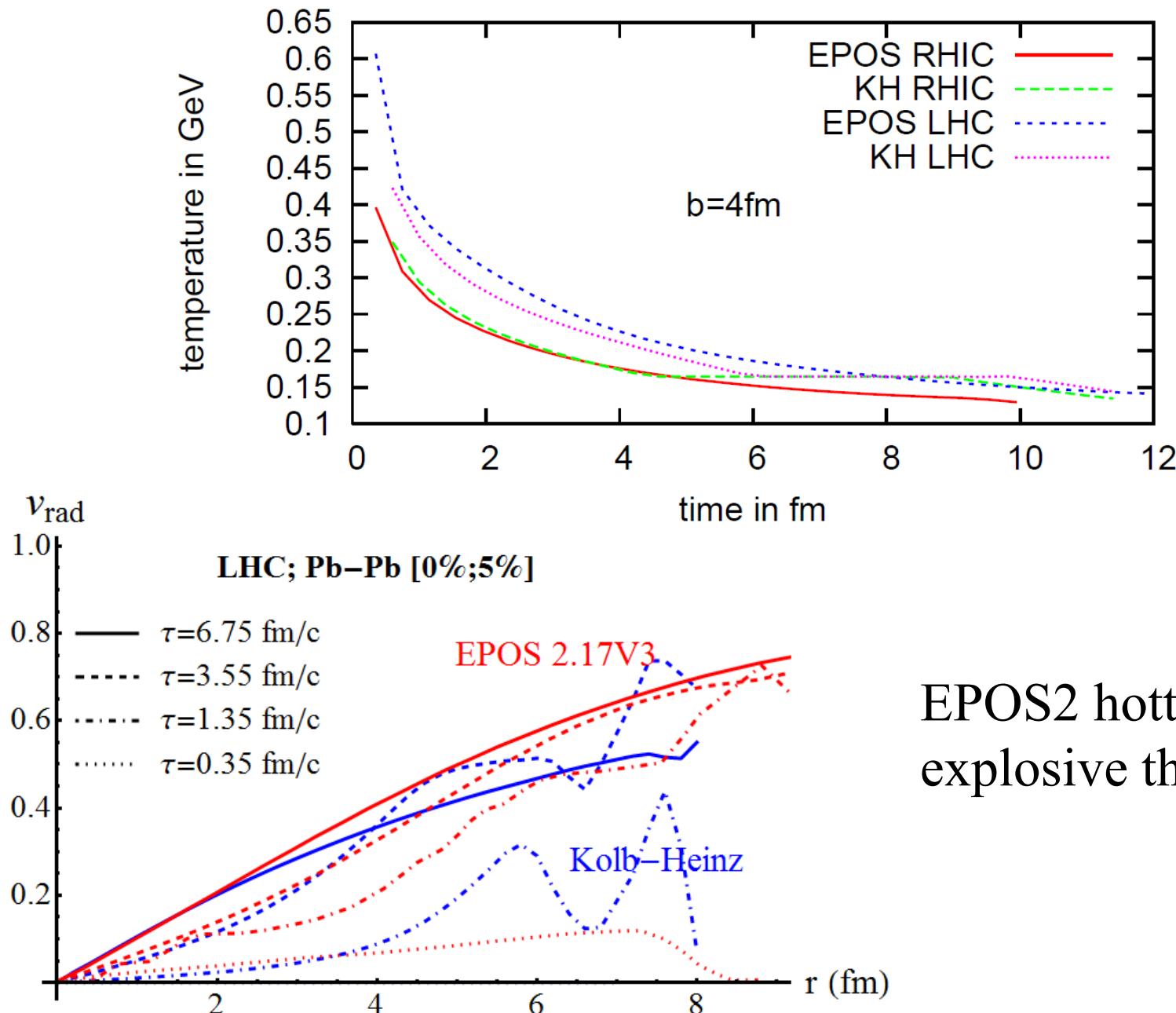
# Medium comparison at RHIC



Gross features of T-evolution  
are identical in the  
« plasma » phase ( $T > 200$   
MeV)

Radial velocities differ  
significantly, starting from the  
earliest times in the evolution

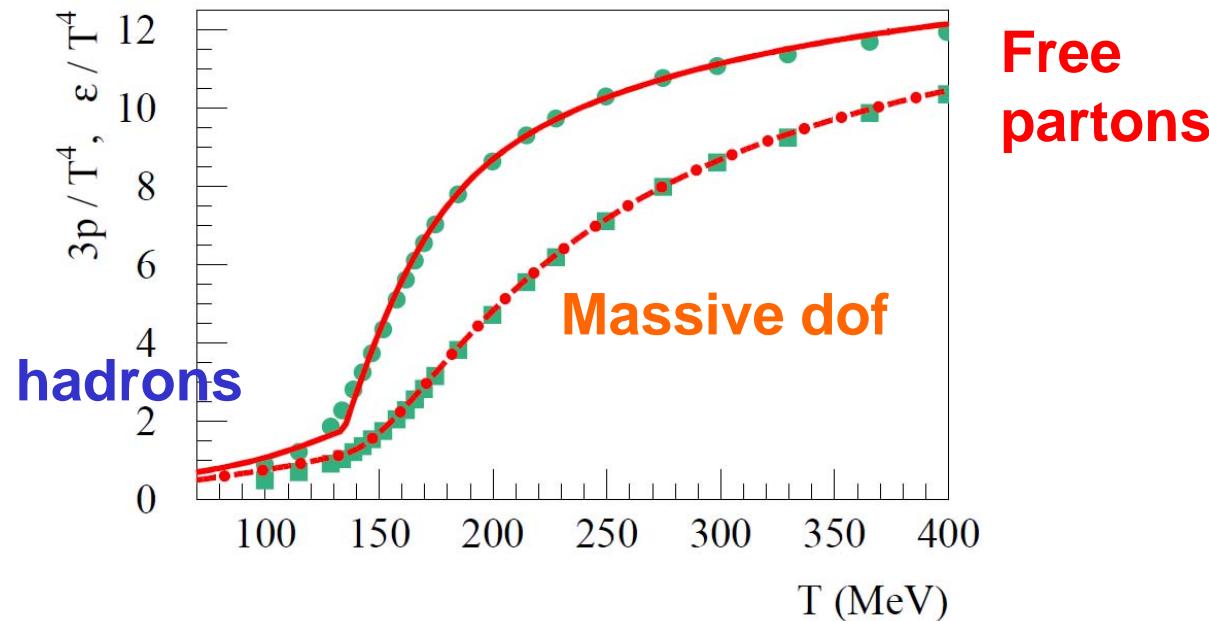
# Medium comparison at LHC



# Coupling EPOS and MC@<sub>s</sub>HQ

Two main (physical) issues:

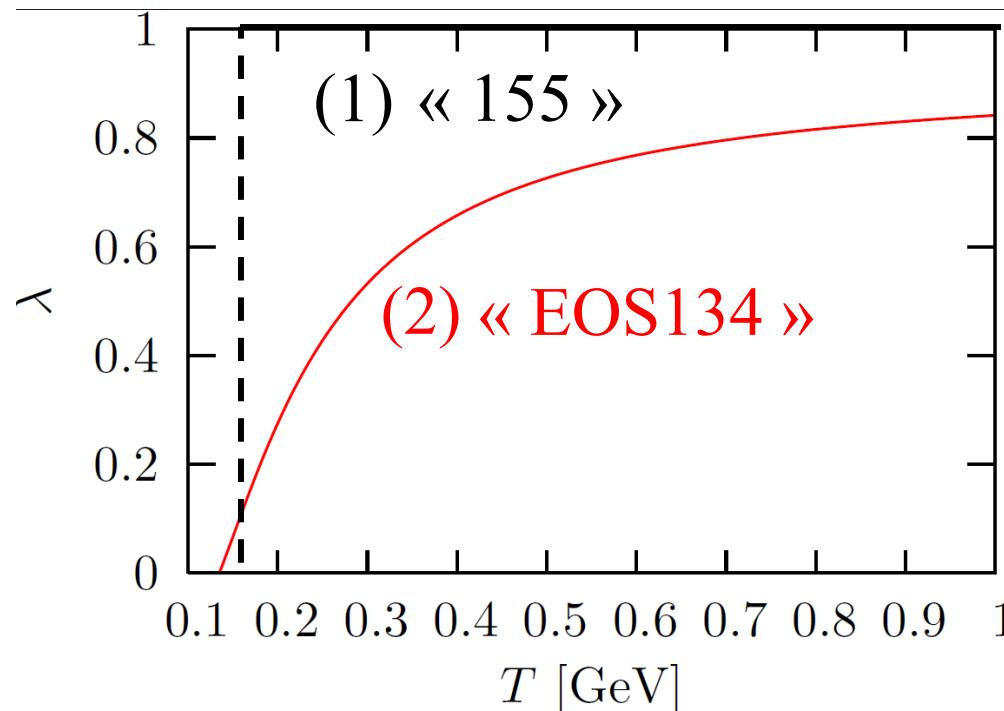
- 1) Generating initial HQ consistently with the multipartonic approach in EPOS (ongoing project)
- 2) Dealing properly with the underlying degrees of freedom in a crossover evolution between hadronic phase and QGP.



# Coupling EPOS and MC@<sub>s</sub>HQ

For the time, 2 prescriptions:

- 1) Interactions as in KH medium (evaluated with massless partons) down to  $T_c=155\text{MeV}$  (in the bulk of the range for the transition temperatures given from lattice)... **most conservative**
- 2) Reduction of effective dof ( $1 \rightarrow \lambda$ ) using the EPOS parametrization of the EOS in terms of partonic and hadronic dofs... down to  $T_c=134\text{MeV}$  (value at which  $\lambda=0$ )

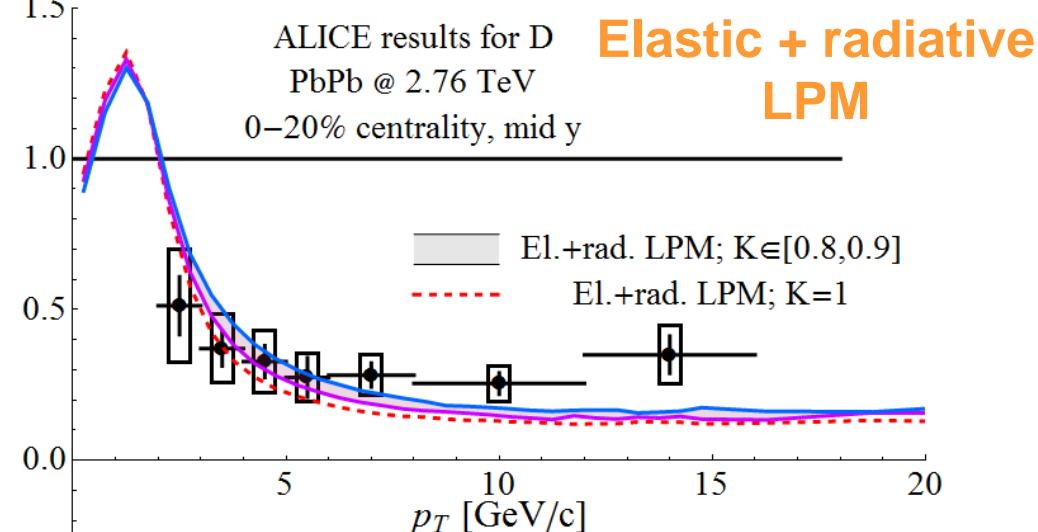
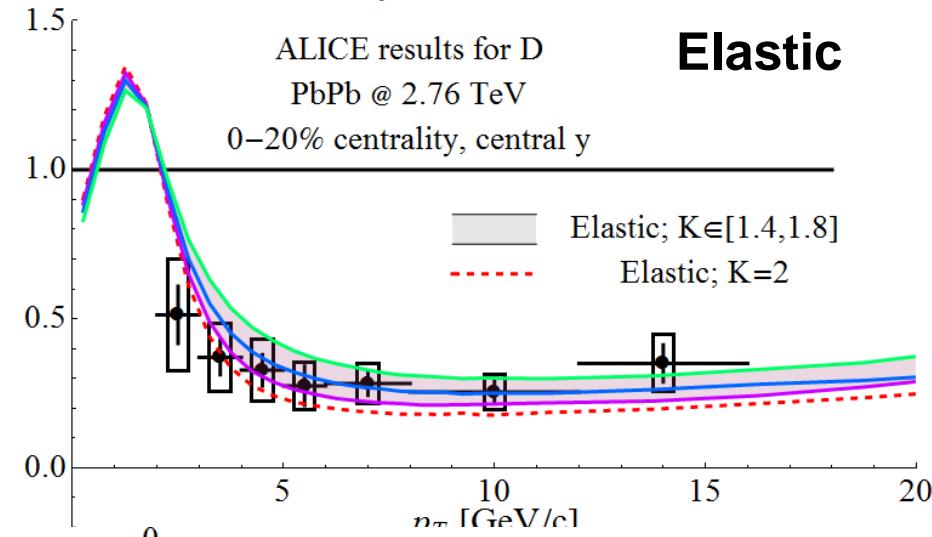


# Some EPOS+MC@<sub>s</sub>HQ results at LHC

NO SHADOWING (yet)

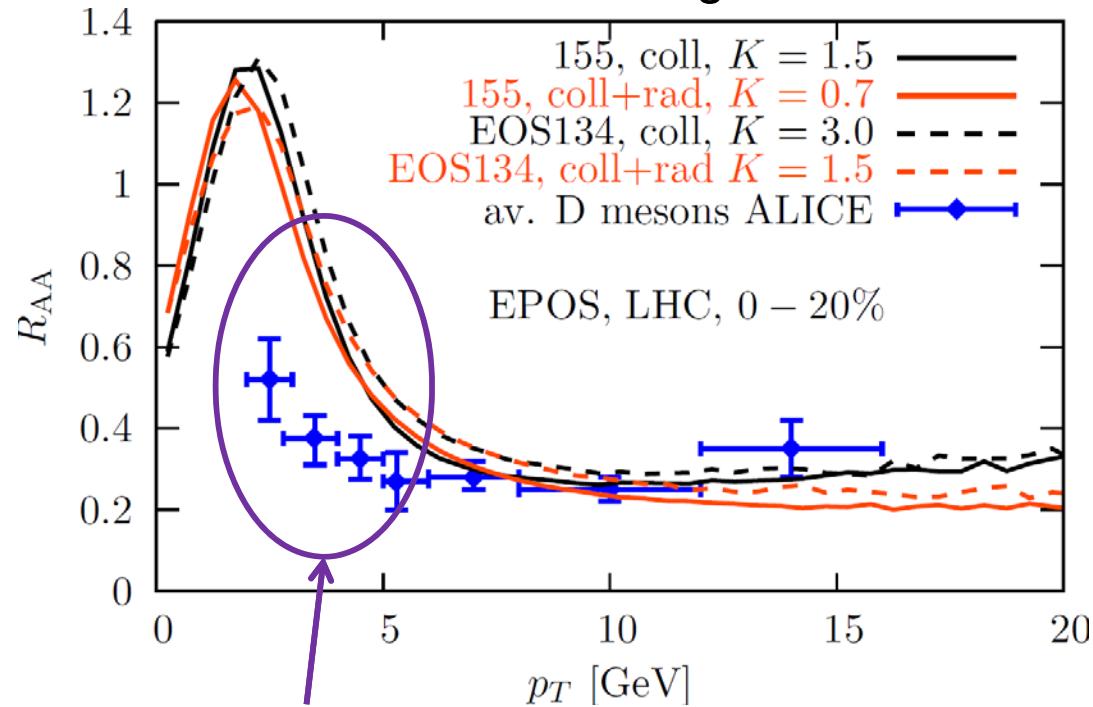
Kolb-Heinz Hydro ( $dN_{ch}/dy = 1600$ )

K close to unity if rad + col considered



EPOS background

K values fixed at  $p_T=10$  GeV/c, x2 if reduction of dof according to EOS134 !



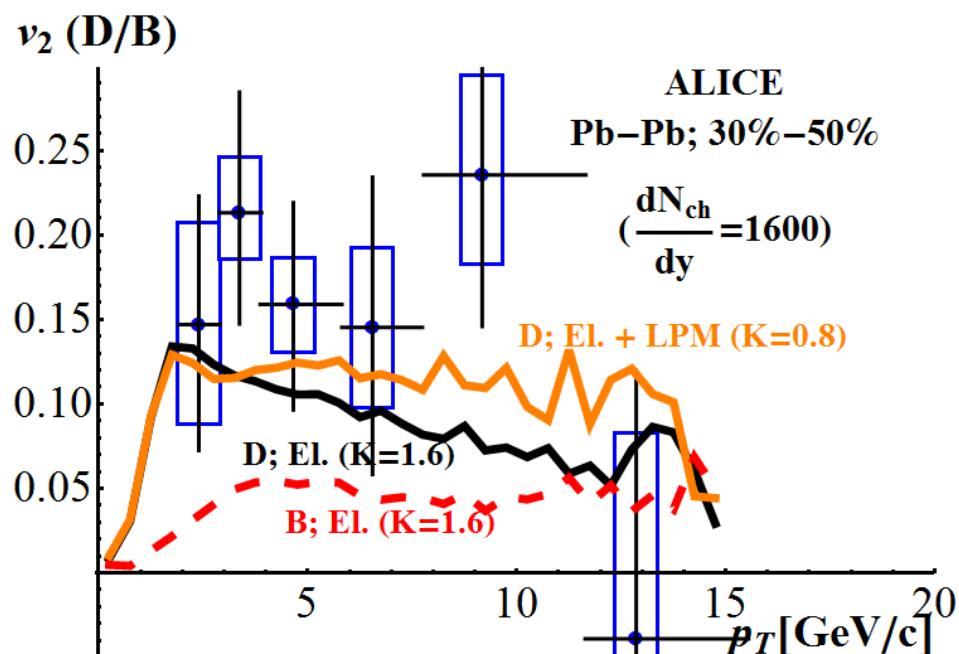
Large push from radial flow; discrepancy unlikely to be explained by shadowing alone.

Concern: Need to revisit the model for small  $p_T$  ?

# Some EPOS+MC@<sub>s</sub>HQ results at LHC

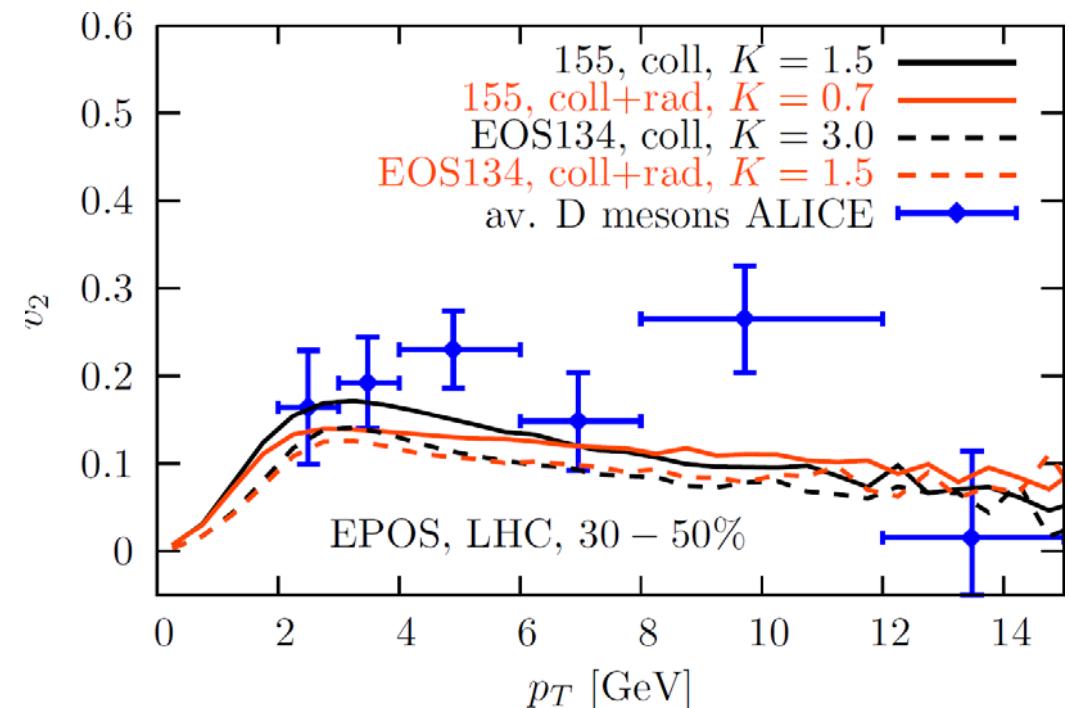
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Kolb-Heinz Hydro ( $dN_{ch}/dy = 1600$ )  
K close to unity if rad + col considered



EPOS background

K values fixed at  $p_T = 10 \text{ GeV}/c$ , x2 if reduction of dof according to EOS134 !

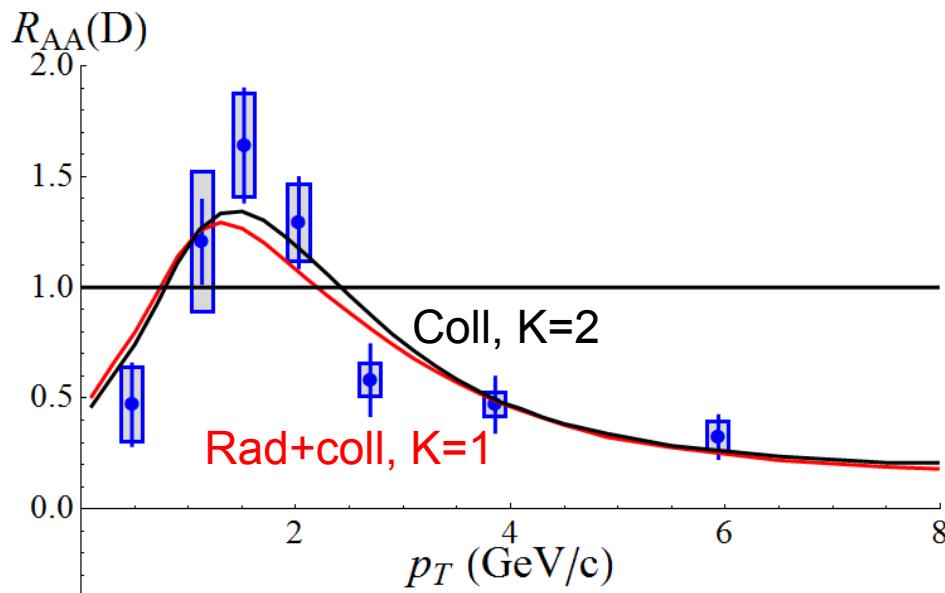


Concerns: Need to revisit the model for small  $p_T$ ?... (Bad) consequences for  $v_2$ ?

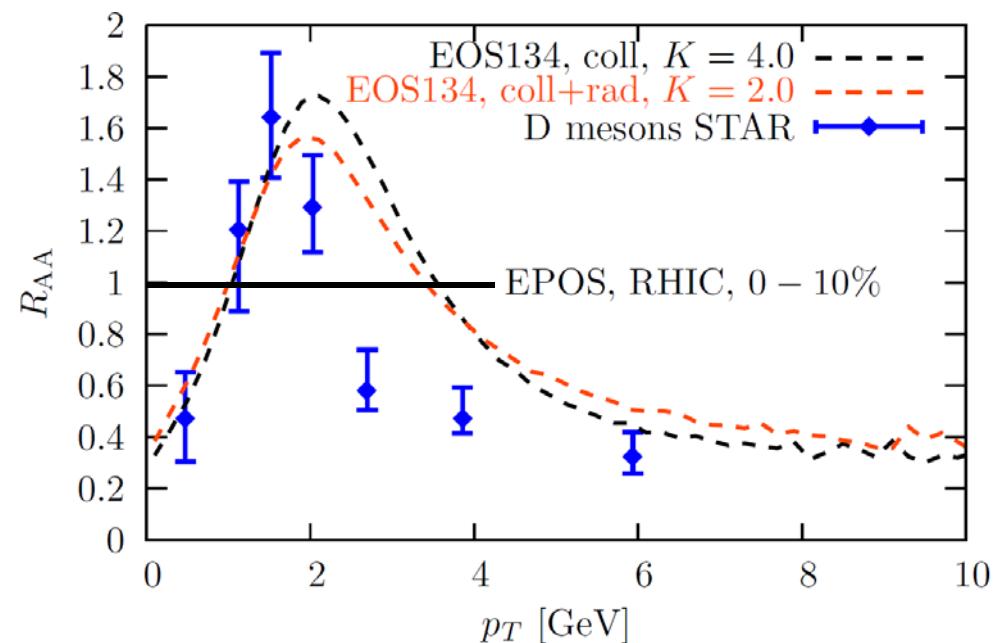
Main message: the models of HF energy loss and the background medium (including its microscopic content) are bound together

# Some EPOS+MC@<sub>s</sub>HQ results at RHIC

KH background



EPOS background + reduction of dof



Slightly larger radial flow in EPOS

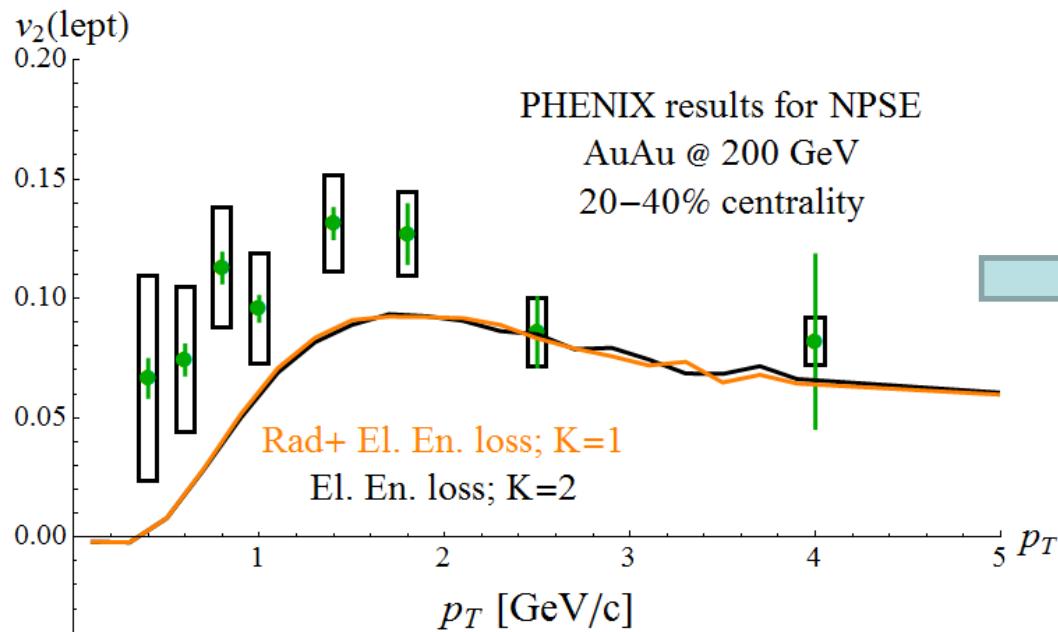
Both « cocktails » (HF energy loss + background + K factor) provide a fair agreement with the data

Data at larger  $p_T$  would help a lot !

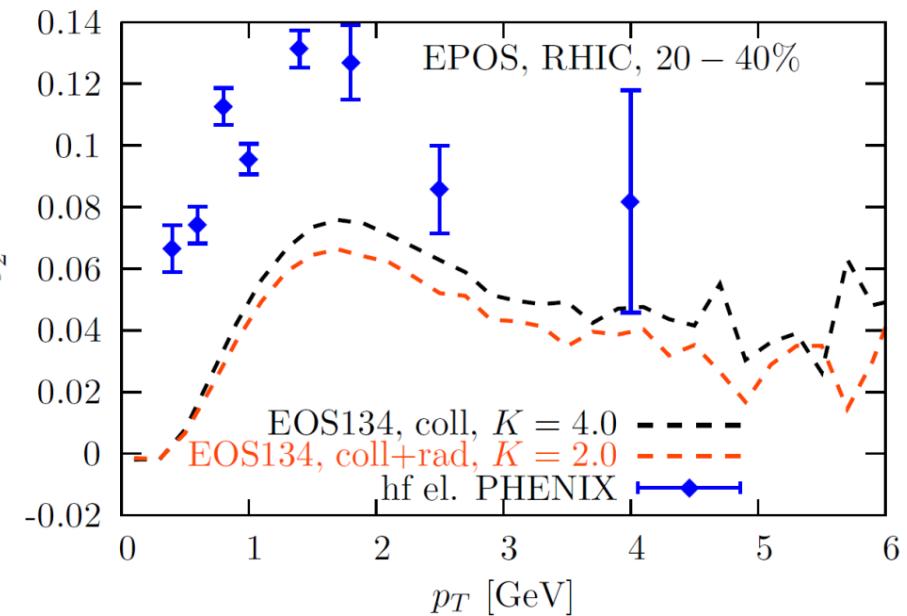
**Main message: the models of HF energy loss and the background medium (including its microscopic content) are bound together**

# Some EPOS+MC@<sub>s</sub>HQ results at RHIC

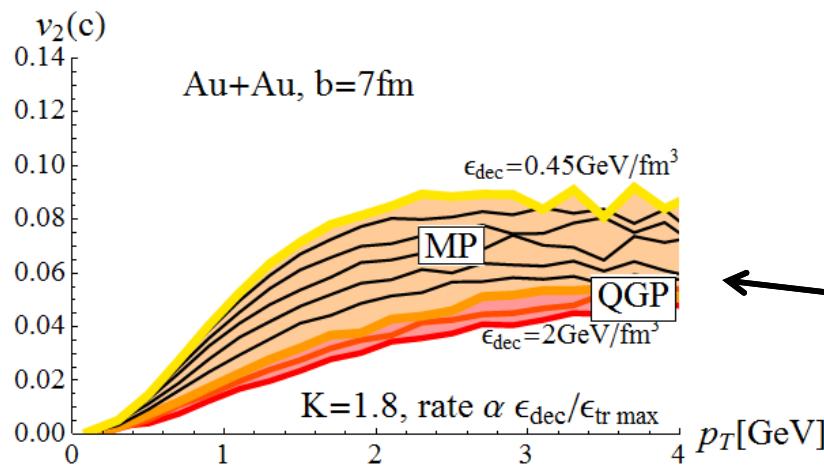
KH background



EPOS background + reduction of dof



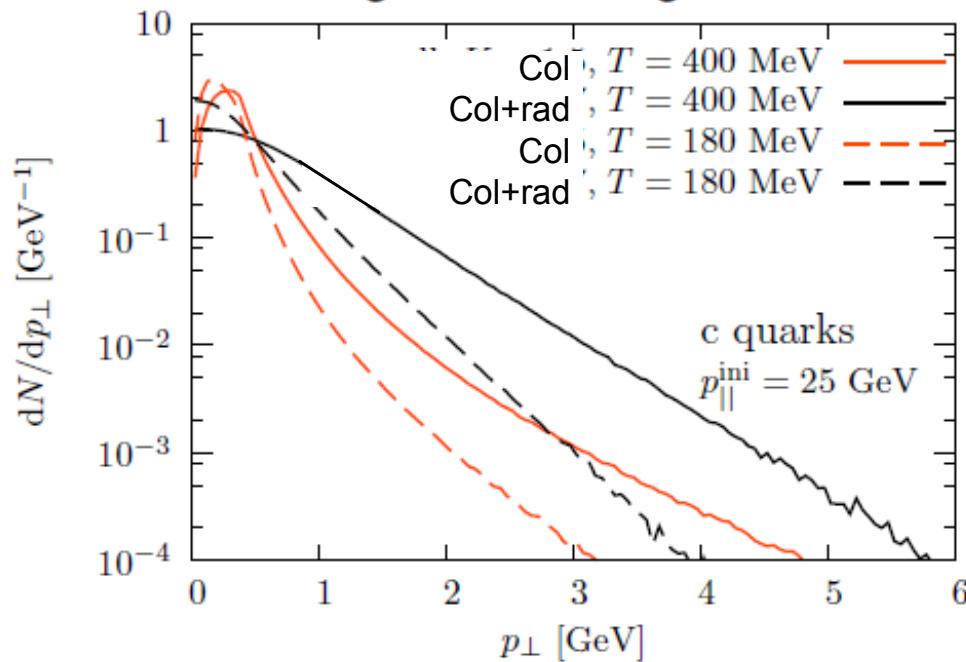
Elliptic flow is reduced by  $\approx 1\%$



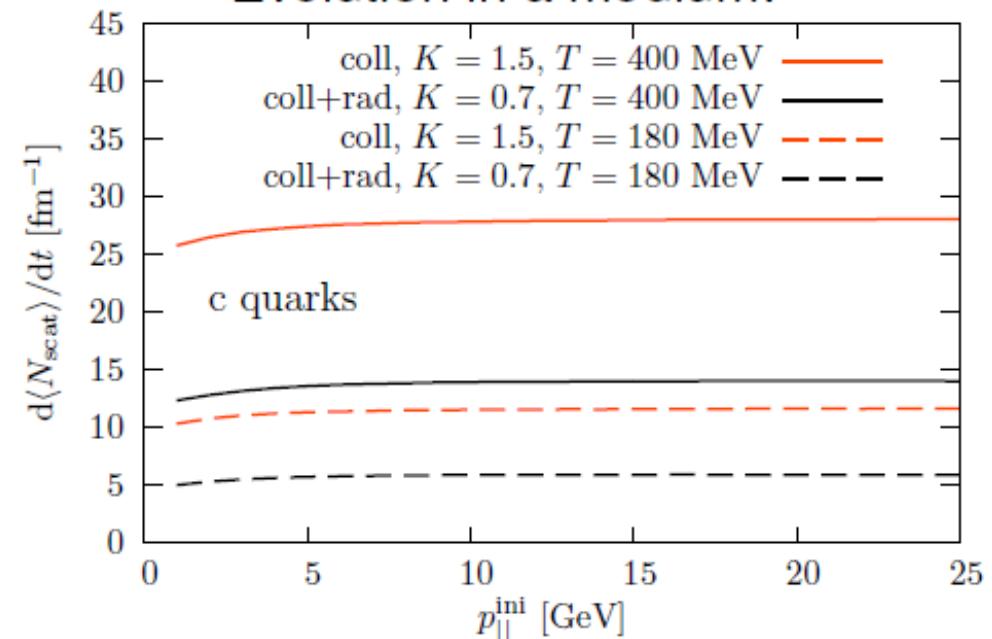
Systematic underprediction of  $v_2$ , ... which develops continuously in time. Probably need for hadronic afterburner in order to reach experimental values.

# Properties of the interactions

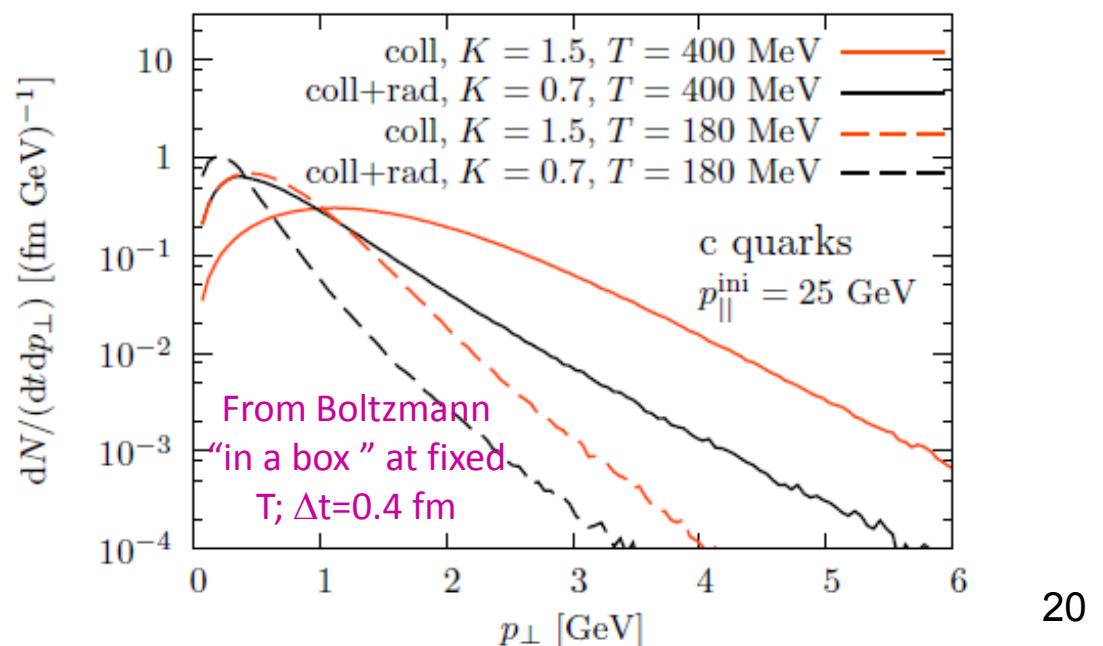
Single scattering:



Evolution in a medium:

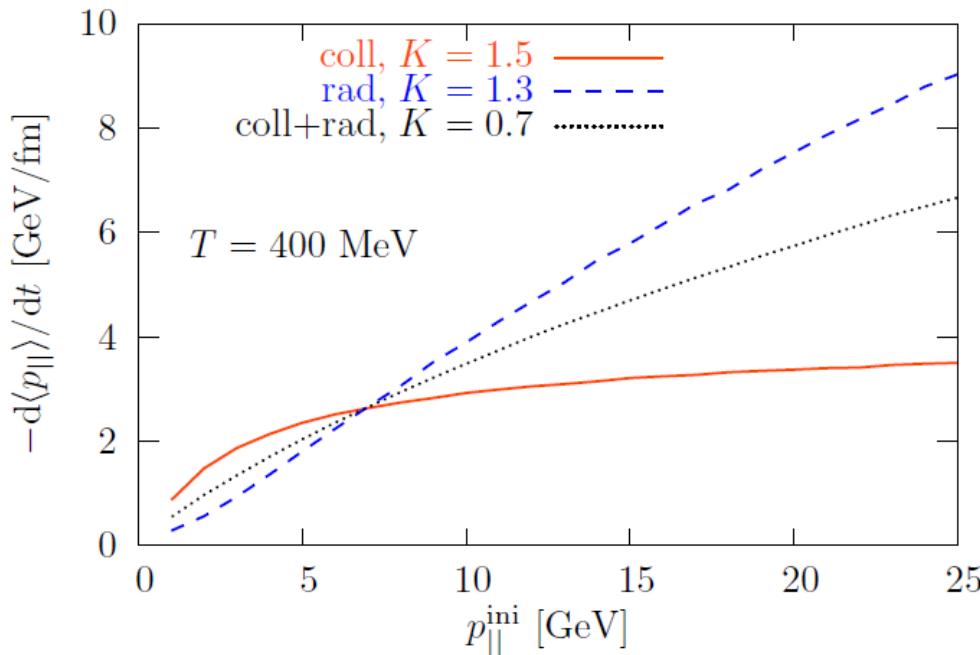


- $p_T$ -distribution in a single scattering: larger  $\langle p_T \rangle$  for **coll+rad** than for **coll** alone
- Scattering rates are **much larger** for **coll** ( $K=1.5$ )
- Differential broadening per unit time: larger for **coll** ( $K=1.5$ )
- Similar trends for b-quarks

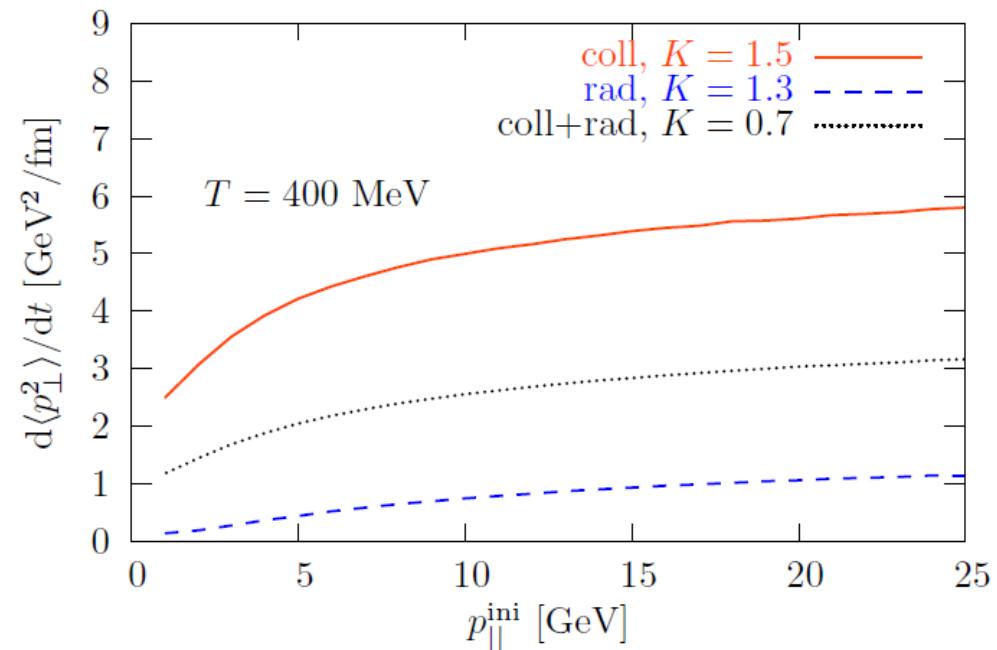


# Properties of the interactions

Drag force

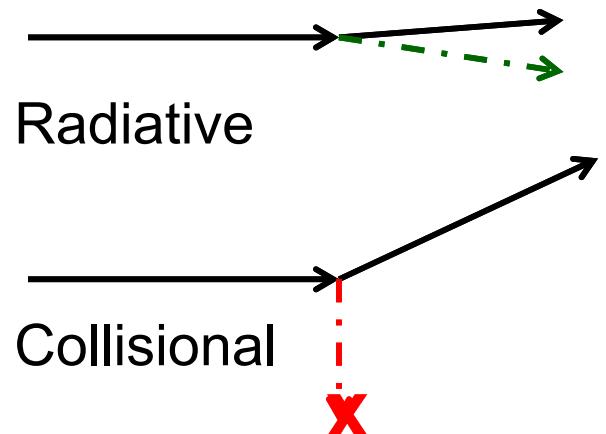


Average perpendicular broadening

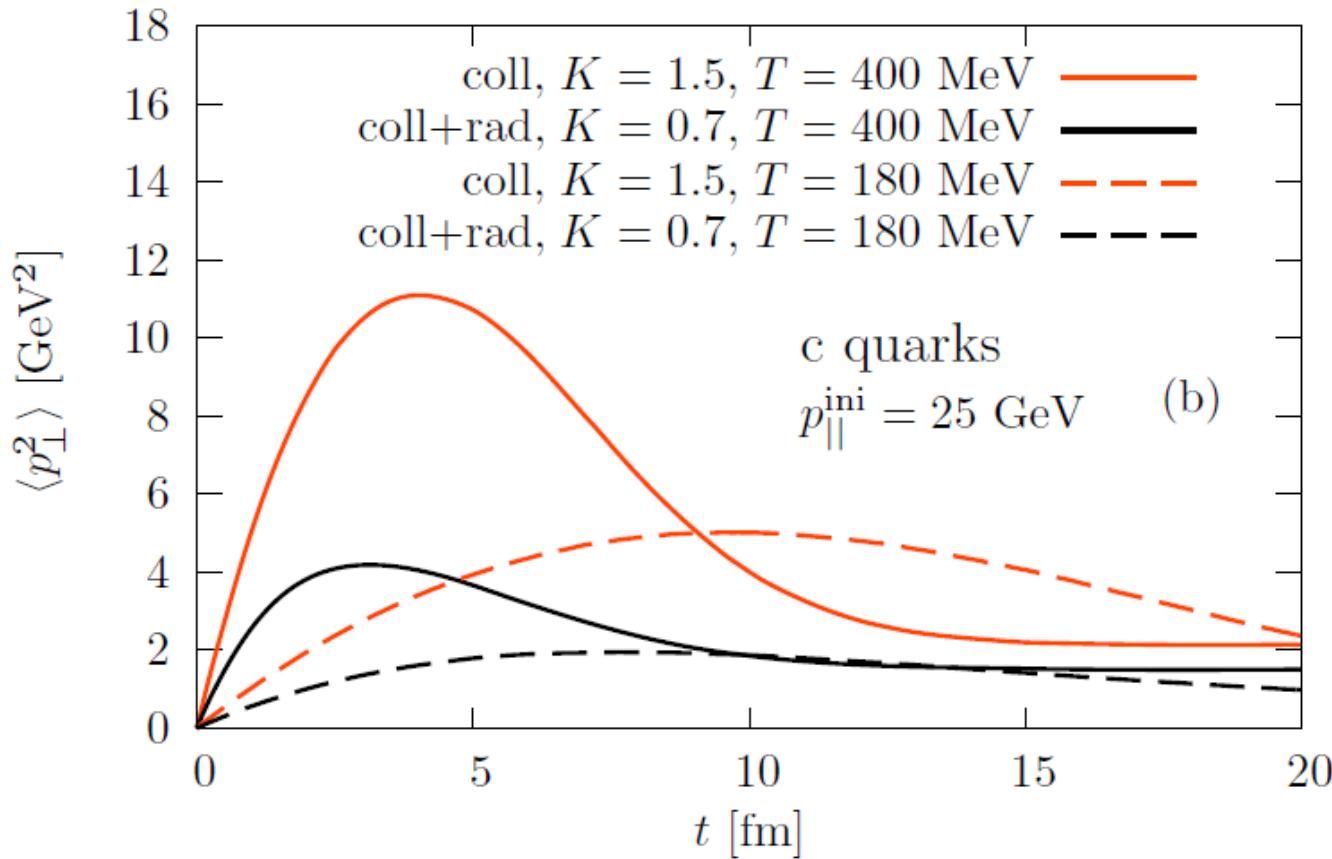


- The  $d\langle p_{||} \rangle/dt$  (drag force) differ between both models, although correct agreement with the data at intermediate  $p_T$ . larger increase of the model with some radiative component.
- The purely collisional scatterings lead to a larger average  $\langle p_t^2 \rangle$  than the radiative “corrections” (need for large scattering to be efficient)
- Expectation: Initial correlations will be broadened more effectively in a purely collisional model.

For BK: acceptable description of the data with moderate  $q\hat{}$



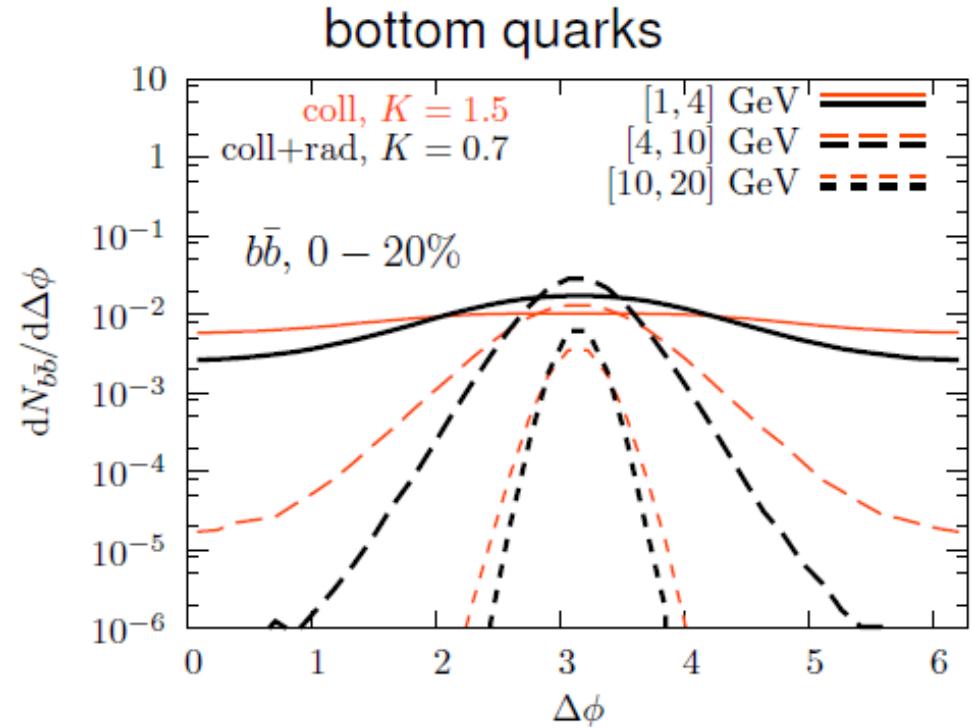
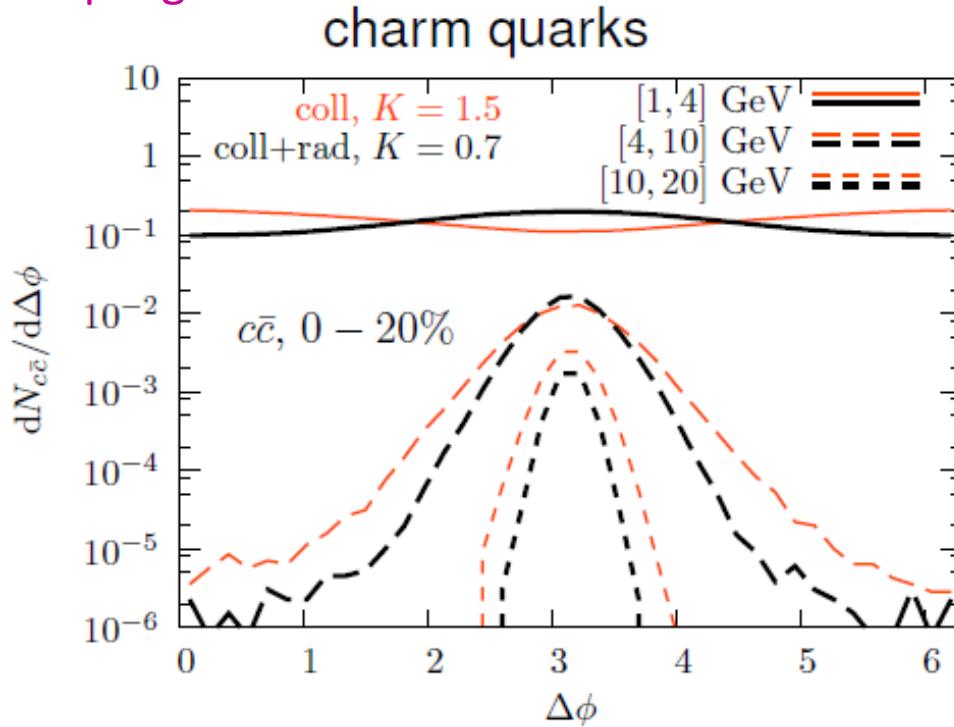
# Evolution in a box



- Linear initial increase, followed by thermalisation at late stage; overshoot at intermediate times... sign of the propagation through hot medium
- Maximum reached at longer times for smaller temperatures (smaller increase, smaller energy loss => smaller thermalisation rate).

# Heavy quarks azimuthal correlations: Back-to-back

Pb-Pb at LHC, HQ initialized back-to-back, no background from uncorrelated pairs, eff.deg=1; decoupling at T=155 MeV



- Stronger broadening in a purely **collisional** than in a **collisional+radiative** interaction mechanism
- At low pT, initial correlations are almost washed out. Some collectivity seen in the purely **collisional** scenario
- Variances in the intermediate pT range (4 GeV-10 GeV): **0.18** vs **0.094** (charm) and **0.28** vs **0.12** (bottom)
- At higher pT, initial correlations survive the propagation in the medium

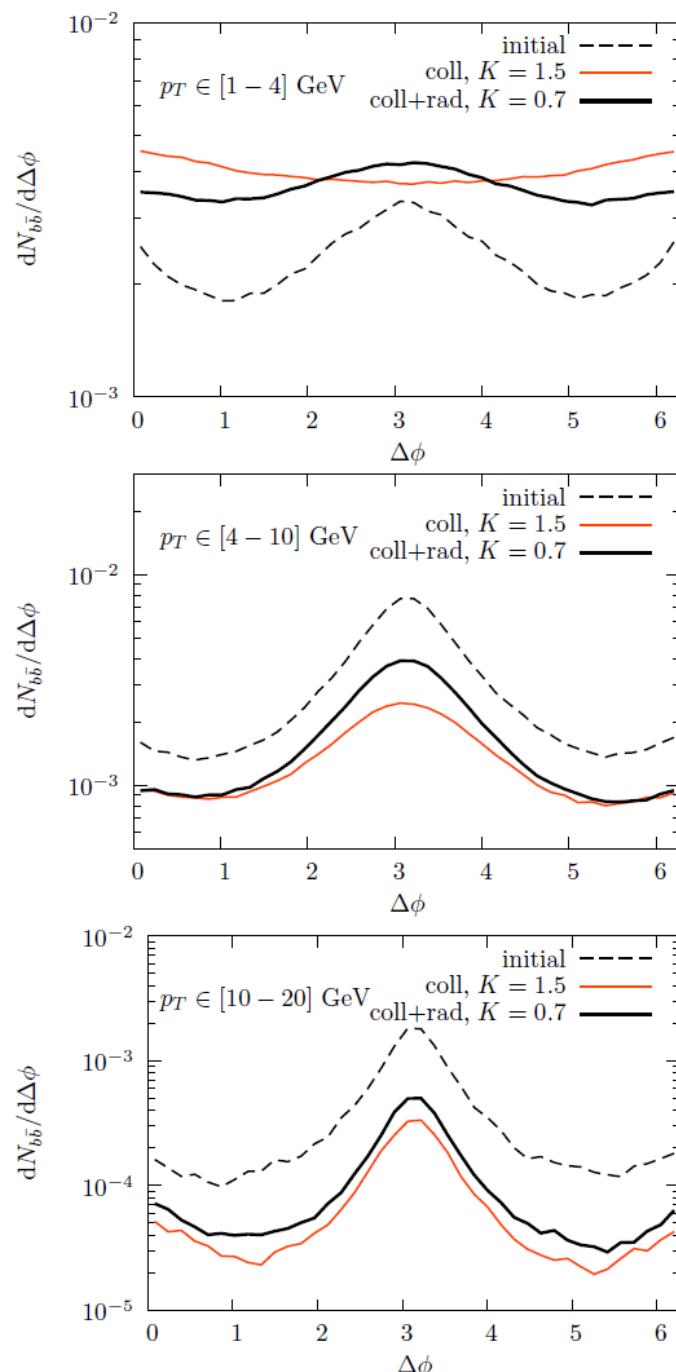
# ... and with Realistic initial distributions: MC@NLO

Next-to-leading order QCD matrix elements coupled to parton shower (HERWIG) evolution: MC@NLO

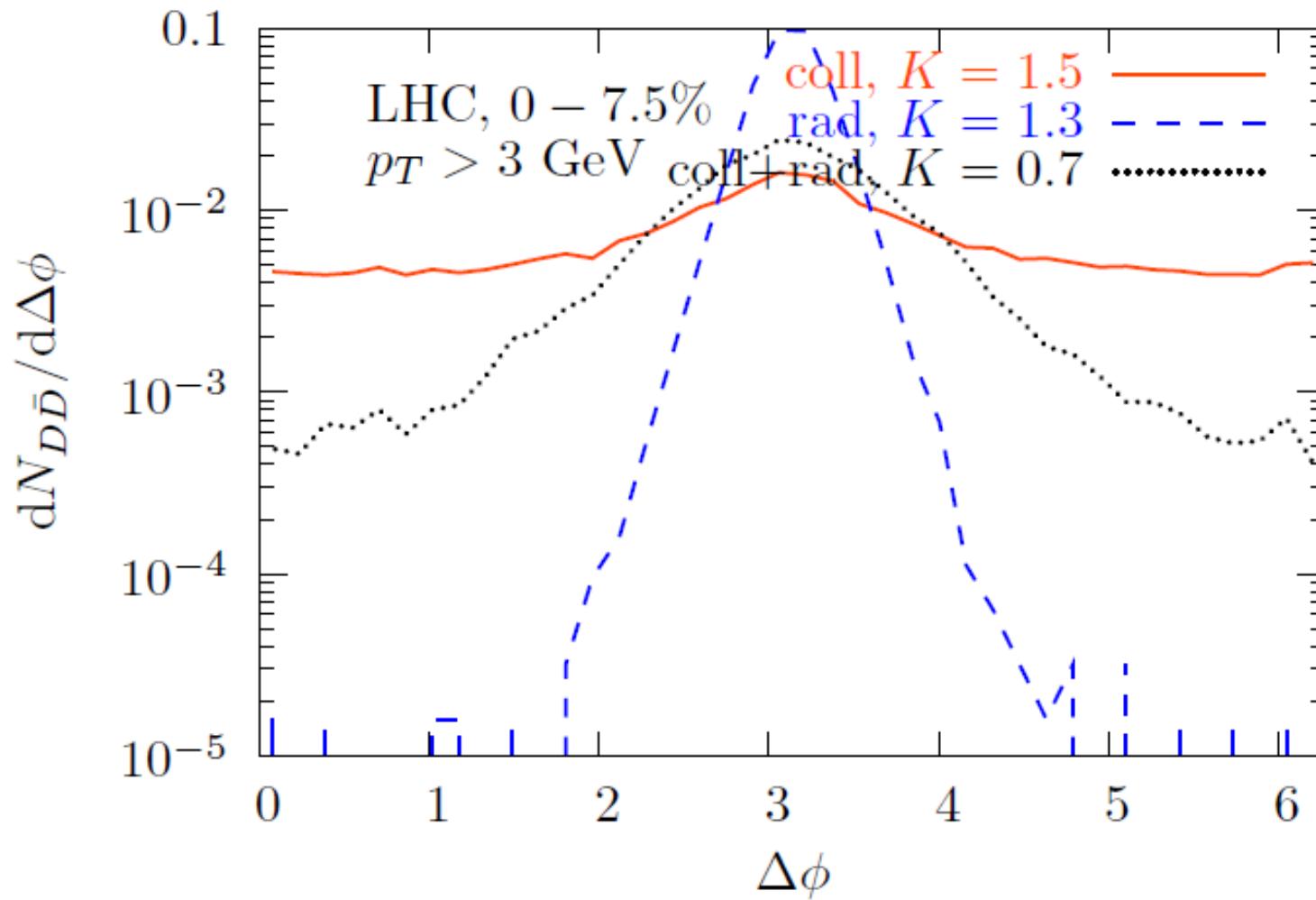
S. Frixione and B. R. Webber, JHEP **0206** (2002)

S. Frixione, P. Nason and B. R. Webber, JHEP **0308** (2003)

- Gluon splitting processes lead to an initial enhancement of the correlations at  $\Delta\phi \approx 0$ .
- For intermediate pT : increase of the variances from 0.43 (initial NLO) to 0.51 ( $\approx 20\%$ ) for the purely **collisional** mechanisms and to 0.47 ( $\approx 10\%$ ) for the interaction including **radiative** corrections (**no additivity with initial width**).
- At larger pT, the deviations from back to back correlations are mostly due to initial NLO corrections.
- Different NLO+parton shower approaches agree on bottom quark production, differences remain for charm quark production!

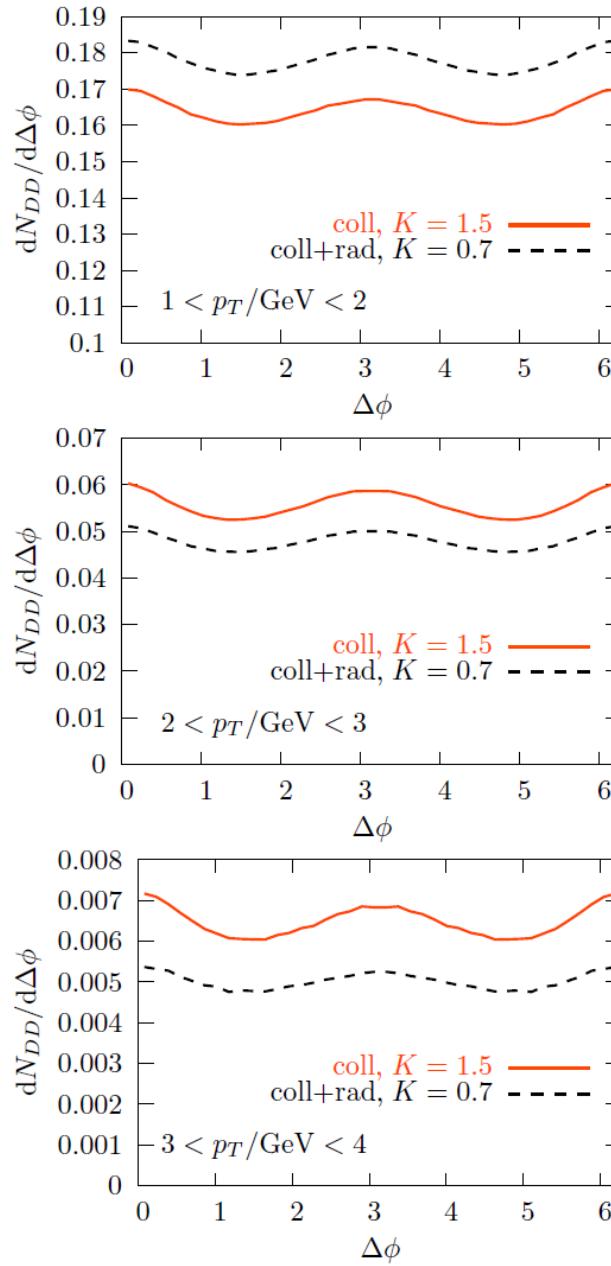


## ... until then: D mesons with back-to-back ccbar



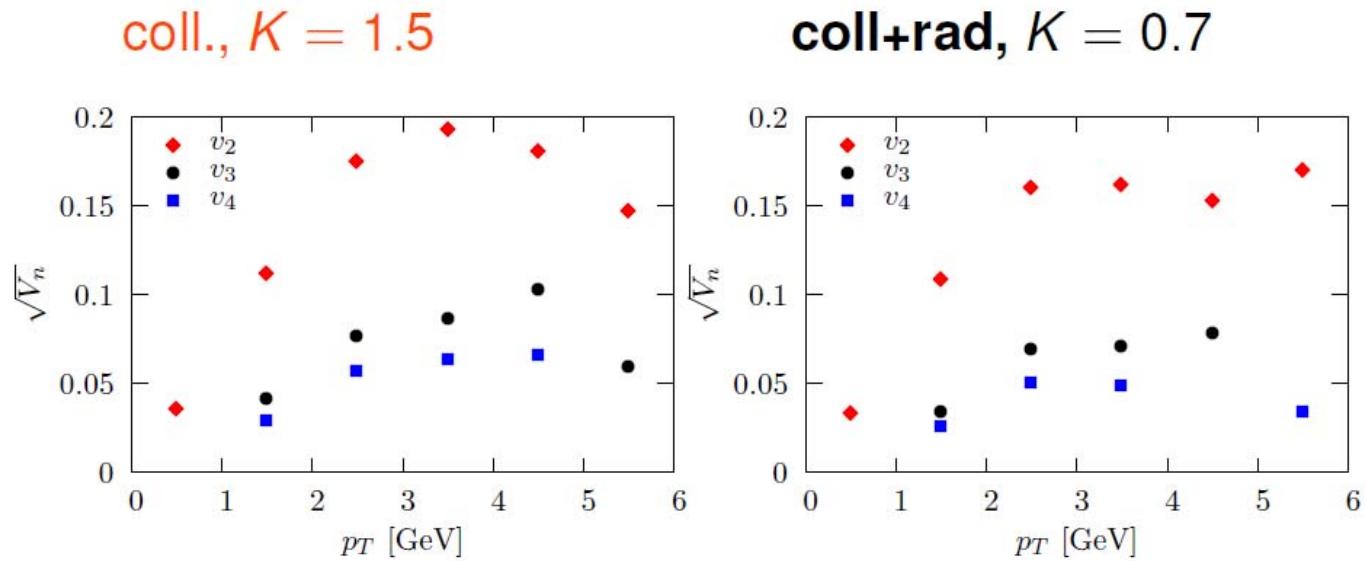
Provide at least some hope that azimuthal correlations could survive the QGP at intermediate  $pT$  and help us understand the basic mechanisms at play

# Peripheral PbPb: Azimuthal correlations and flow



- 30-50% centrality, DD correlations
- Flow harmonics from 2-particle correlation functions

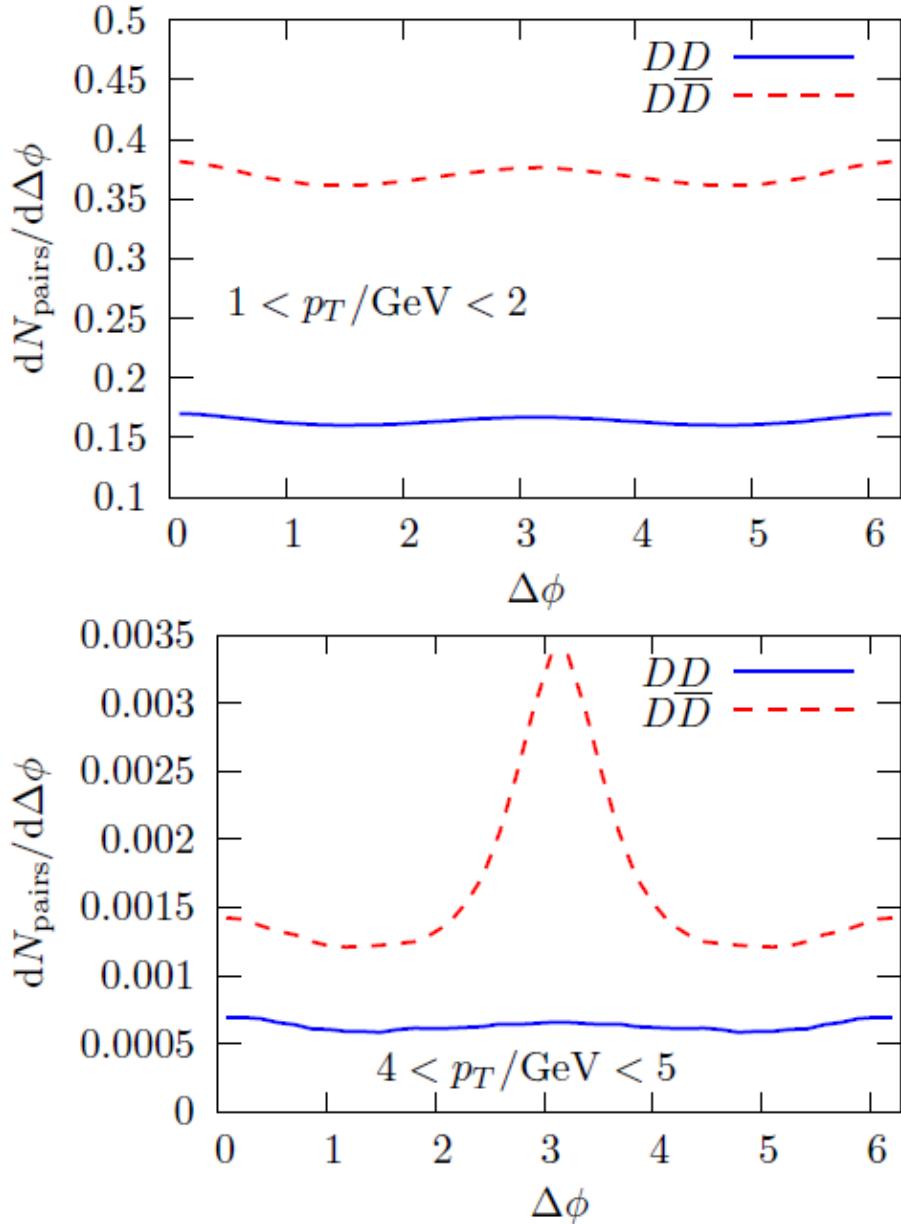
$$\propto \frac{N}{2\pi} (1 + 2 \sum V_n \cos(n\Delta\phi))$$



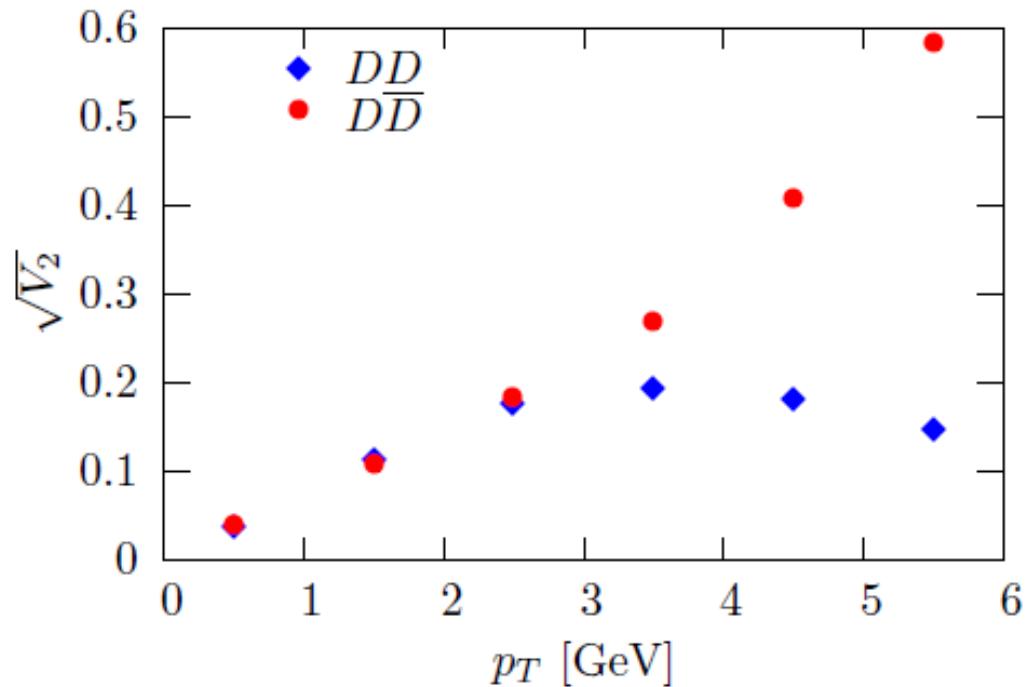
- Similar  $v_n$  for both interaction mechanisms at low  $pT$
- Nonvanishing higher flow coefficients

# Peripheral PbPb: Azimuthal correlations and flow

as an example collisional,  $K = 1.5$



- Compare DD correlations to DDbars correlations to learn about the flow contribution and the degree of isotropization of DDbars pairs



- Similar  $v_2$  for both DD and DDbars at small  $p_T$
- Dominant initial back-to-back correlation for DDbars at large  $p_T$

## Conclusions & Perspectives

- Data at intermediate  $p_T$  can be reproduced by several model(s).
- D suppression at Large  $p_T$  favors collisional energy loss... or suggests improvements are in order for our treatment of radiative energy loss (finite path length, finite gluon width,...)
- Discrepancy at small  $p_T$  is increased using a more realistic model of the medium (need to be confirmed using EPOS with viscous hydro).
- Azimuthal correlations of HF mesons at intermediate pT may help constraining the models, but initial NLO distributions need to be better controlled for c-quarks, as well as possible CNM effects (not included here)
- To do: predictions for correlations including the decay chains.

# Based on

- *Towards an understanding of the single electron data measured at the BNL Relativistic Heavy Ion Collider (RHIC)*, P.B. Gossiaux & J. Aichelin, Phys. Rev. C **78**, 014904 (2008); [[arXiv:0802.2525](#) ]
- *Tomography of quark gluon plasma at energies available at the BNL Relativistic Heavy Ion Collider (RHIC) and the CERN Large Hadron Collider (LHC)*, P.B. Gossiaux, R. Bierkandt & J. Aichelin, Physical Review C **79** (2009) 044906; [[arXiv:0901.0946](#)]
- *Tomography of the Quark Gluon Plasma by Heavy Quarks*, P.-B. Gossiaux & J. Aichelin, J. Phys. G **36** (2009) 064028; [[arXiv:0901.2462](#)]
- *Energy Loss of Heavy Quarks in a QGP with a Running Coupling Constant Approach*, P.B. Gossiaux & J. Aichelin, Nucl. Phys. A **830** (2009), 203; [[arXiv:0907.4329](#)]
- *Competition of Heavy Quark Radiative and Collisional Energy Loss in Deconfined Matter*, P.B. Gossiaux, J. Aichelin, T. Gousset & V. Guiho, J. Phys. G: Nucl. Part. Phys. **37** (2010) 094019; [[arXiv:1001.4166](#)]
- *Plasma damping effects on the radiative energy loss of relativistic particles*, M. Bluhm, P. B. Gossiaux, & J. Aichelin, Phys. Rev. Lett. **107** (2011) 265004 [[arXiv:1106.2856](#)]
- *Theory of heavy quark energy loss*, P.B. Gossiaux, J. Aichelin, T. Gousset, [[arXiv:1201.4038v1](#)]

# Based on

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- ... other recent publications all available on arxiv
- **Azimuthal correlations of heavy quarks in Pb+Pb collisions at LHC ( $\sqrt{s} = 2.76$  TeV)**, Marlène Nahrgang, Jorg Aichelin, Pol Bernard Gossiaux, and Klaus Werner [arxiv 1305.6544]