# Higher twist from the isolated particle yield 

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## Introduction: $x_{\mathrm{T}}$ Scaling

Factorization of inv. cross section by dimensional analysis

$$
E \frac{\mathrm{~d}^{3} \sigma}{\mathrm{~d} p^{3}}=\frac{1}{p_{\mathrm{T}}^{n}} F\left(\frac{p_{\mathrm{T}}}{\sqrt{s}}\right)=\frac{1}{\sqrt{s}^{n}} G\left(x_{\mathrm{T}}\right), n++=4
$$


$x_{T}=\frac{2 p_{\mathrm{T}}}{\sqrt{s}}$
$F\left(\frac{p_{\mathrm{T}}}{\sqrt{s}}\right) \& G\left(x_{\mathrm{T}}\right)$ dimensionless
scaling functions
(Left fig. [JHEP 1108:086,2011])
QCD has $\Lambda_{\mathrm{QCD}}, \alpha_{s}$, etc $\Longrightarrow n>4$.
At $\sqrt{s_{\text {LHC }}}$ pQCD predicts $n \approx 5 \ldots 6$.
$n$ OBTAINED $\sigma^{\text {inv }}$ WITH TWO DIFFERENT $\sqrt{s}$

$$
n\left(x_{\mathrm{T}}, \sqrt{s}_{1}, \sqrt{s_{2}}\right)=\frac{\ln \left(\sigma^{\mathrm{inv}}\left(x_{\mathrm{T}}, \sqrt{s_{2}}\right) / \sigma^{\mathrm{inv}}\left(x_{\mathrm{T}}, \sqrt{s_{1}}\right)\right)}{\ln \left(\sqrt{s}_{1} / \sqrt{s}_{2}\right)}
$$

## SCALING EXPONENT \& "TWIST"

Leading-twist (LT) processes = particle produced by fragmentation $\Longrightarrow n \approx 5 \ldots 6$ (NLO pQCD at LHC energies)

Higher-twist (HT) processes = particle produced directly in the subprocess
$\Longrightarrow n$ significantly larger (low- $p_{\mathrm{T}}$ phenomenon)


Leading-twist process


Left: Leading-twist, Right: Direct Higher-twist

## HT prediction for RHIC and LHC

Prediction bands by F.Arleo et al. [PRL. 105 (2010) 062002]
$\Delta^{\text {fit }}=n_{\text {measured }}-n_{\mathrm{NLO}}="(\mathrm{LT}+\mathrm{HT})-\mathrm{LT} "$


FIGURE : Prediction of $\Delta^{\text {fit }}$ based on a global fit of RHIC and Tevatron data with preliminary data points from PHENIX. [F.Arleo, Moriond 2010]

Non-zero $\Delta^{\text {fit }}$ data points $\Longrightarrow$ A hint of HT at RHIC data?

## ALICE DATA vs. PYTHIA8

Inclusive charged hadrons at 7 TeV and 2.76 TeV



Left: Data \& PYTHIA8 inclusive invariant $p_{\mathrm{T}}$ distributions of charged hadrons
Right: Data/PYTHIA8 ratio. PYTHIA describes the inclusive spectra within 20\%

* finite tracking efficiency included in PYTHIA8


## ALICE DATA vs. NLO

ALICE inv. cross sections with NLO (DSS FF) [arXiv:1307.1093], NLO overpredicts the data by a factor two (See backup).

Here NLO with Kretzer FF is used, describes the data within $25 \%$.
Calculation by H.Paukkunen, Jyväskylä.


## Scaling exponents vs. predictions <br> Left: $n$ from NLO, ALICE data and PYTHIA8

Right: $\Delta n$ to NLO.


## REMARKS:

1) Data points seem to agree with LHC prediction
2) Non-zero $\Delta n$ even from PYTHIA8 which has no Higher Twist processes.

## HT Prediction for Isolated Particles

It is suggested by F.Arleo et al. [PRL. 105 (2010) 062002] that isolation cuts should enhance the fraction of HT processes

Hadrons from direct production $(\mathrm{HT})$ are expected to be accompanied with less hadronic activity than jets.

I aim to verify this using PYTHIA8.


Isolation cut:

- Calculate $\sum p_{\mathrm{T}}$ of associated particles inside the cone
- If the sum is smaller than pre-defined limit, the particle is isolated
- I used cone with radius $\mathrm{R}=0.4$ in $\eta-\phi$
- The $\sum p_{\mathrm{T}}$ was required to be less than $10 \%$ of the trigger $p_{T}$


## ISOLATED $n$ FROM PYTHIA8

Below is the $n$ of inclusive (black) and isolated spectra from PYTHIA8.



The isolation cut increases $n$, even in PYTHIA8 without HT. This suggests a kinematical bias caused by the isolation itself.

## HT IMPLEMENTATION FOR PYTHIA8

With help of T.Sjöstrand, I added a HT process [Phys.Rev.D23,99] to PYTHIA8.

$$
\begin{equation*}
(\mathrm{qg} \rightarrow \mathrm{q} \pi) \quad \frac{\mathrm{d} \hat{\sigma}}{\mathrm{~d} \hat{t}}=\frac{\pi\left(\alpha_{s}\left(p_{\mathrm{T}}\right)\right)^{3}}{\hat{s}^{2}} \frac{1}{4 \pi} \frac{s_{0}}{\hat{s}} F(\cos \theta) \tag{1}
\end{equation*}
$$

with HT scale $s_{0}=16 \pi^{2} f_{\pi}^{2}, f_{\pi}=93 \mathrm{MeV}$ and $F(z=\cos \theta)=\frac{2}{27}(1-z)\left(1+\frac{4}{(1+z)^{2}}\right)$



Left: HT as in [Phys.Rev.D23,99] Right: HT multiplied by 1000 to magnify effect
With this implementation, the effect of HT to $n$ is very weak, even when enhanced by factor of 1000 .

## CONCLUSION

- The $x_{T}$ scaling exponents $n$ for inclusive charged hadrons were extracted from ALICE pp data, PYTHIA and NLO at $\sqrt{s}=7 \mathrm{TeV}$ and 2.76 TeV .
- The $\Delta n$ results are in agreement with HT predictions from F.Arleo et al.
- Isolation cuts increase $n$ due to a kinematical bias.
- Even when multiplied by 1000 , the direct $\pi^{ \pm}$production as described in [Phys.Rev.D23,99] is barely visible on $n$ at LHC energies in PYTHIA8.
- The HT PYTHIA implementation suggests that the direct pion production HT contribution is very small at LHC energies


## Backup slides

## The ALICE Experiment

Part of the Large Hadron Collider (LHC) at CERN, Switzerland.
Designed for heavy-ion collisions, but has also a pp program.
Excellent particle identification cabability, down to even $p_{\mathrm{T}}=100 \mathrm{MeV} / c$.


- For the isolation checking in PYTHIA, we wanted to mimick the low $p_{T}$ performance of a real detector.
- This was done by removing particles from the PYTHIA event randomly according to the efficiency curve.
- When obtaining the $p_{\mathrm{T}}$ spectra, the inverse of the efficiency was used to recover the lost particles, like in the real data.
- After this, the isolated spectra were in better agreement. Figures below are 7 TeV data/PYTHIA ratios before (left) and after (right) the effi was applied in PYTHIA.


Efficiency not taken into account


Efficiency taken into account

## BACKGROUND ACTIVITY PLOTS <br> PYTHIA8 ISR=FSR=MPI=kT=off <br> Activity in a $\mathrm{R}=0.4$ cone for HT pions and all hadrons



Very hard to distinguish HT from all hadrons with isolation cut

## HT IN PYTHIA8


$p_{\mathrm{T}}$ analytical \& PYTHIA8


HT pion correlation function

## HT IN PYTHIA8

## Ratio of HT pions / all charged hadrons



Left: HT as in [Phys.Rev.D23,99] Right: HT multiplied by 1000 to magnify effect

## ALICE AND NLO $p_{T}$



Left: Fig. from [arXiv:1307.1093]


Right: From H.Paukkunen, Jyväskylä

## FF AND $n$




Left: $n$ with DSS and Kretzer FF
Right: Ratio Kretzer/DSS NLO spectra
from H.Paukkunen, Jyväskylä

## Toy Monte Carlo model

To study the role of kinematics further, a toy Monte Carlo model was made in the spirit of Bjorken's Parent-Child-Relationship [Phys.Rev.D8,4098 (1973)]

$$
\frac{1}{p_{\mathrm{T}}} \frac{\mathrm{~d} N}{\mathrm{~d} p_{\mathrm{T}}}=\int f_{q}\left(p_{\mathrm{Tq}}\right) \cdot D(z) \frac{1}{z^{2}} \mathrm{~d} z \sim \frac{1}{p_{\mathrm{T}}^{\mathrm{a}}} \int_{x_{\mathrm{T}}}^{1} z^{a-2} D(z) \mathrm{d} z
$$

where $z=p_{\mathrm{T}} / p_{\mathrm{Tq}}, D(z)$ is the fragmentation function $(\mathrm{FF})$ and $f_{q}\left(p_{\mathrm{Tq}}\right)$ is the final parton spectrum.

The spectra of isolated hadrons was needed, so a Monte Carlo cascade process was used instead of integration:

- Parton momentum $p_{\text {Tq }}$ is sampled from power law
- $f_{q}\left(p_{\mathrm{Tq}}\right) \sim p_{\mathrm{Tq}}^{-6}$ for " 2.76 TeV "
- $f_{q}\left(p_{\mathrm{Tq}}\right) \sim p_{\mathrm{Tq}}^{-5}$ for " 7 TeV "
- $z$ is sampled from the approximation of $\mathrm{FF}, \exp (-8.2 \cdot z)$
- $p_{\mathrm{T}}=z \cdot p_{\mathrm{Tq}}$ is saved and subtracted from the $p_{\mathrm{Tq}}$ Now the scaling exponent $n$ could be extracted from the hadron spectra.


## Toy MC - $\Delta n$ Prediction (RELATIVE ISOLATION)





Fractions $f$ of isolated particles with the relative cut:
$f_{1}=0.0257(1)$, green points at 7 TeV
$f_{2}=0.0478(1)$, blue points at 2.76 TeV

$$
\Delta n_{\text {pred. }}=\frac{\ln \left(f_{1} \sigma_{1} / f_{2} \sigma_{2}\right)}{\ln \left(\sqrt{s_{2}} / \sqrt{s}_{1}\right)}-\frac{\ln \left(\sigma_{1} / \sigma_{2}\right)}{\ln \left(\sqrt{s_{2}} / \sqrt{s}_{1}\right)}=\frac{\ln \left(f_{1} / f_{2}\right)}{\ln \left(\sqrt{s_{2}} / \sqrt{s_{1}}\right)}
$$

$\Longrightarrow$ non-zero $\Delta n$ just from the isolation probabilities!

