# What can we learn from modified jets in Heavy Ion collisions?

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## Introduction

- Jets & Hadrons are suppressed at RHIC and the LHC
- Jets gain lose energy and gain (some) transverse momentum from the medium
- Jet interactions are split into elastic and radiative, the strength of each is quantified by a transport coefficient: e-hat and q-hat.
- Measuring or calculating these is a major goal of heavy-ion jet physics
- What does the current zoo of jet observables actually tell us about the physics of quenching?



## VNI/BMS Jet Simulation Methodology

## VNI/BMS - 2.0, a simple JET transport model

- Partonic transport via the Boltzmann equal footing.
   s is a simplistig model, with the LPM interferences only also between the potential alculation [P] the gluon is alcowed to bescatter after emission and inalytic treatment (see §5 in [12]).
- Interactions are tree level 2->2 scatterings and final-state radiation. Radiation includes leading order (BDMPS-Z) LPM effect.
- Medium is a box of thermal partonic QGP at a fixed temperature. No expansion!
- Cross sections are screened by Debye mass, computed using the box temperature
- A generated jet is injected, cascade of interacting partons are tracked. Evolution of entire jet is recorded.



## Zapp and Wiedemann, LPM Algorithm



- Quark and gluon propagate freely
- Simulates coherent emission from multiple centers

Zapp K, Wiedemann U. *Phys Rev Lett,* 103 (2009) JEWEL CCS, S.A.Bass, D.K.Srivastava, *hep-ph/1101.48*95

## Jet Simulation Method: Box Mode



- Insert all partons from each jet into parton cascade box and evolve for a fixed path length.
- Medium is partonic and static, temperature can be fixed for the entire evolution.
- Medium and jet interact on an equal footing, track all resulting partons

A fully controllable brick of QCD matter

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## Hadronization In VNI/BMS - 2.0



## Hadronization In VNI/BMS

100 GeV quark evolved for 4fm in a box at T=350 MeV



## Jet Observables - VNI/BMS - Box Mode

## Dijet Simulation Method: Circle Mode



- Vacuum shower evolution takes place in Pythia.
- Sample production vertices uniformly within a circular medium of some radius R. Sample jet paths as length of random chord generated from vertices to edge of medium.
- Insert all partons from each jet into parton cascade box and evolve for sampled path length.
- Medium is partonic and static, temperature is fixed for duration of the evolution.

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## Process Dependence





## Di-jet Asymmetry at RHIC

- Use VNI/BMS to try and understand dependence of Aj on:
  - Qhat, determined by the medium temperature: T = 250, 350 MeV
  - Distance travelled by the both jets
  - Cuts on leading jet energy and cone-radius: Elead > 20, 35, 50 GeV,  $R = \{0.2, 0.3, 0.4\}$
  - Interaction mechanism, elastic or elastic+rad
- Setup:
  - Di-jets generated by p+p at 200GeV, with cuts / jet definition applied
  - Cuts and jet definition are applied in post-processing

C.C-S, B.Müller, "What can we learn from Dijet suppression at RHIC", Phys.Rev.C 86 (2012)

## Dijet Asymmetry - Varying Medium Temperature



## Dijet Asymmetry - Varying Jet Cone Radius



## RHIC Dijet Asymmetry - Varying Strong Coupling



C.C-S, B.Müller, "What can we learn from Dijet suppression at RHIC", Phys.Rev.C 86 (2012)

## Reconstruct jets with Anti-Kt at successively larger cone radii

## RHIC - Jet Shape



### Jet Shape - 2, varying strong coupling



#### RHIC-Fragmentation - Longitudinal $z = E_T / E_{T,Jet} \cos \Delta R$



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#### Jet Fragmentation - Transverse

 $J_T = E_T \sin \Delta R$ 



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$$z = E_T / E_{T,Jet} \cos \Delta R$$
$$J_T = E_T \sin \Delta R$$

## 2d Fragmentation





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#### What do jet transport coefficients tell us about QGP?

- Jet transport coefficients encode medium density, and kinematic information
- Knowing q-hat and e-hat would allow for true jet tomography, given a set of observed jets from an event one could reconstruct the density profile.
- q-hat measures transverse momentum transfer, e-hat measures longitudinal energy loss
- In the infinite probe momentum limit, radiation off static charges (q-hat) dominates energy loss
- In the low probe momentum limit, recoil of medium from elastic scattering (e-hat) dominates energy loss
- The ratio of the two coefficients is proportional to the mass of the medium constituents seen by the jet.





## Thermal Masses for medium partons

- Previous results from VNI/BMS derived from a medium of massless partons
- Interaction cross-sections are always screened by the Debye mass  $m_D^2 = \frac{1}{6}(2N_c + N_f)g^2T^2$
- Introduce asymptotic HTL masses for medium partons

$$m^2 = kg^2T^2, \quad k = \begin{cases} \frac{1}{6}N_c + \frac{1}{12}N_f & \text{gluon} \\ \frac{1}{4}C_F & \text{quark} \end{cases}$$

 Introduce a dimensionless scaling parameter µs to 'dial' medium masses

$$m^2 = k\mu_s^2 g^2 T^2.$$

NB: Medium number density now scales with masses

C.C-S, B.Müller, "Constituent Mass Dependence of transport coefficients in a quark gluon plasma", hep-ph/1209.3328



 $\frac{N}{V} = \frac{gTm^2}{2\pi^2} K_2\left(\frac{m}{T}\right)$ 

## Measuring Transport Coefficients in VNI/BMS

- Fix medium temperature T=350 MeV, run code without radiation or hadronization.
- Run events with quark probes at fixed energies
- Extract q-hat and e-hat from transverse momentum and energy loss accumulated by the probe

$$\hat{q} = \frac{1}{L} \sum_{i=1}^{N_{\text{coll}}} \Delta p_{T,i}^2$$

• For light probes in a massless medium we expect from pQCD calculations

$$\hat{e} = \frac{1}{L} \sum_{i=1}^{N_{\text{coll}}} \Delta E_i,$$
  

$$\hat{q}(T) = 4\pi C_R \alpha_s^2 \mathcal{N}(T) \ln\left(\frac{q_{\text{max}}^2}{m_D^2} + 1\right),$$
  

$$\mathcal{N}(T) = \frac{\zeta(3)}{\pi^2} \left(2N_c + \frac{3}{2}N_f\right) T^3,$$
  

$$\hat{e}(T) = \frac{4\pi \alpha^2 T^2}{3} \left(1 + \frac{N_f}{6}\right) \ln\left(\frac{ET}{m_D^2}\right)$$

## Elastic Energy Loss



Confirms kinematic intuition, higher medium constituent masses recoil less resulting in a lower energy loss

T=350 MeV,Alpha = 0.3

## Q-hat / E-hat Ratio



- Ratio scales linearly with the medium mass scale µs.
- Experimental measurements of qhat and e-hat could provide insight into the nature of the QGP as seen by jets.
- Measurements made at different jet scales may reveal structure in quasi-particle mass spectrum.
- A possible precision measurement of hard probes and the QGP

T=350 MeV, Et > 20GeV

## Dijet Response at RHIC - Aj



### Elastic Interactions Only

## Dijet Response at RHIC - Subleading Jet Shape

Varying mu-s

Varying alpha-s



Jet shape is sensitive to elastic energy loss

## Conclusions

- RHIC Dijet Asymmetry is sensitive to: strong coupling, medium radius, medium temperature and cone radius. Sometimes this is subtle.
- Modified jets have a softened radial profile (jet shape), partons are scattered transverse to the jet axis, transverse fragmentation profile softened.
- Fragmentation distributions look similar to vacuum for narrow jets, these jets are still strongly modified. Fragmentation of broad jets shows strong modifications over vacuum
- Modified jets retain a hard core that looks like a vacuum jet, surrounded by a soft cloud of radiated and liberated particles.
- Measurements of qhat/ehat ratio could provide insight into the nature of the medium as seen by jets.

## How can we use what we know?

Unknown - modification process

Initial state - pQCD (pythia)

$$P(E^{f}, p_{t}^{f}) = \int P(E^{f}, p_{t}^{f} | E^{I}, p_{t}^{I}, Q^{2}, \rho, L, \ldots) P(\rho, L | b) P(L | b) P(b) P(E^{I}, p_{t}^{I}, Q^{2}) d \ldots$$

'measured' distributions of modified jets

Hydro - path length and density encountered

- Experimental data gives a window into the after effects of the modification process
- How can we extract the maximum information about the modification process from the data?
- Typically people propose a form for the kernel and then attempt to reproduce observables
- Instead propose the most general meta-model for the modification kernel and then restrict the allowed form by requiring consistency with successively more difficult observables. A non-parametric model selection problem

## Extras

## Leading Parton Energy Loss



Baier et al (BDMPS) *Nucl Phys.B478* (1996), B.Zakharov, *JETP Lett.63* (1996),

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## VNI/BMS - 2.0 + Hydro

Keep jetty partons at each time step. Discard 'old' medium partons



## VNI/BMS - 2.0 Current + Future Features

- Elastic and Radiative energy loss models, with BDMPS-Z LPM effect
- Lund Stringy Hadronization with full color tracking from the initial generator.
- Fixed medium temperature in simple box mode.
- Event by event hydro background.
- Variable medium constituent masses give control of qhat/ehat ratio.
- Integration with Pythia for hard process generation
- Jet level data analysis built in (and single hard probe)
- Relatively simple user-options
- Modern build system (CMAKE)

## Leading Parton Energy Distribution, E=20 GeV



## Leading Parton Pt Distribution, E=20 GeV



gluon E=20, T=0.2, MS=0

## **VNI/BMS** Cross-Sections

g g → g g	$\frac{9}{2}\left(3-\frac{tu}{s^2}-\frac{su}{t^2}-\frac{st}{u^2}\right)$	$q q' \rightarrow q q'$	$\frac{4}{9}\frac{s^2+u^2}{t^2}$
q g→ q g	$-\frac{4}{9}\left(\frac{s}{u}+\frac{u}{s}\right)+\frac{s^2+u^2}{t^2}$	q qbar→ q' qbar'	$\frac{4}{9}\frac{t^2+u^2}{s^2}$
g g → q qbar	$\frac{1}{6}\left(\frac{t}{u}+\frac{u}{t}\right)-\frac{3}{8}\frac{t^2+u^2}{s^2}$	q g →q γ	$-\frac{e_q^2}{3}\left(\frac{u}{s}+\frac{s}{u}\right)$
$\mathbf{q} \ \mathbf{q}  ightarrow \mathbf{q} \ \mathbf{q}$	$\frac{4}{9} \left( \frac{s^2 + u^2}{t^2} + \frac{s^2 + t^2}{u^2} \right) - \frac{8}{27} \frac{s^2}{tu}$	q qbar → g γ	$\frac{8}{9}e_q^2\left(\frac{u}{t}+\frac{t}{u}\right)$
q qbar → q qbar	$\frac{4}{9} \left( \frac{s^2 + u^2}{t^2} + \frac{u^2 + t^2}{s^2} \right) - \frac{8}{27} \frac{u^2}{st}$	q qbar → γ γ	$\frac{2}{3}e_q^4\left(\frac{u}{t}+\frac{t}{u}\right)$
q qbar → g g	$\frac{32}{27} \left( \frac{t}{u} + \frac{u}{t} \right) - \frac{8}{3} \frac{t^2 + u^2}{s^2}$		

Dominant elastic interactions
 including screening mass

$$\mu_D = \sqrt{(2N_c + N_f)/6gT}$$

$$\begin{split} \frac{d\sigma^{gg \to gg}}{dq_{\perp}^2} &= 2\pi \alpha_s^2 \frac{9}{4} \frac{1}{(q_{\perp}^2 + \mu_D^2)^2},\\ \frac{d\sigma^{gq \to gq}}{dq_{\perp}^2} &= 2\pi \alpha_s^2 \frac{1}{(q_{\perp}^2 + \mu_D^2)^2},\\ \frac{d\sigma^{qq \to qq}}{dq_{\perp}^2} &= 2\pi \alpha_s^2 \frac{4}{9} \frac{1}{(q_{\perp}^2 + \mu_D^2)^2}, \end{split}$$