

Can we measure the gravitational quantum states of Positronium?

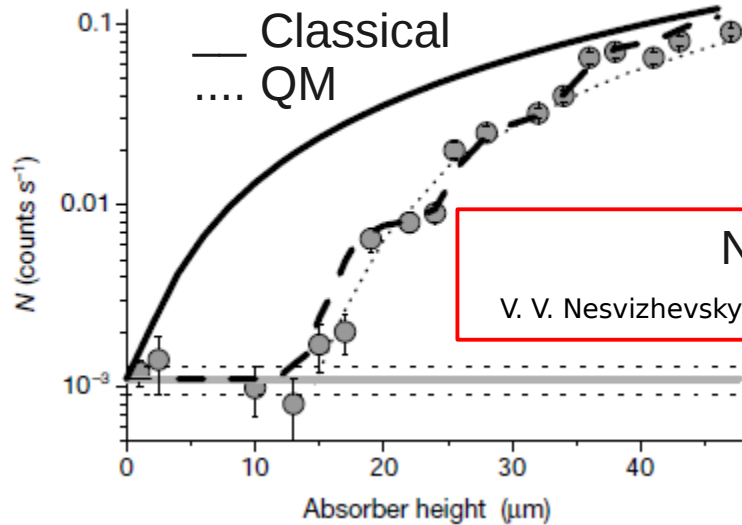
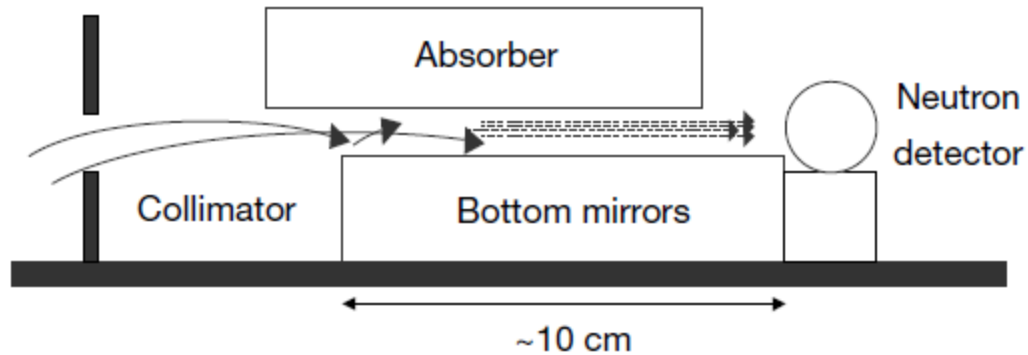
Paolo Crivelli

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GRANIT Workhop, Les Houches, 5th of February 2014

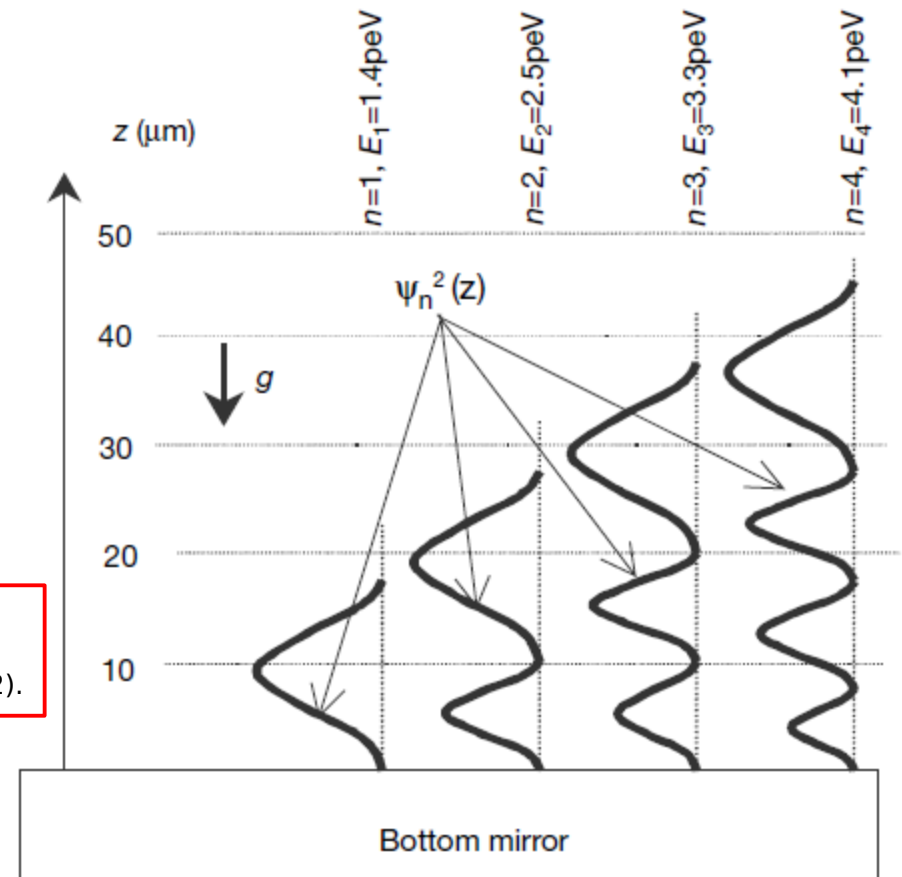
My work is supported by the Swiss National Science Foundation
(Ambizione grant) and the ETH research grant ETH-47-12-1

Gravitational QM states



Neutrons

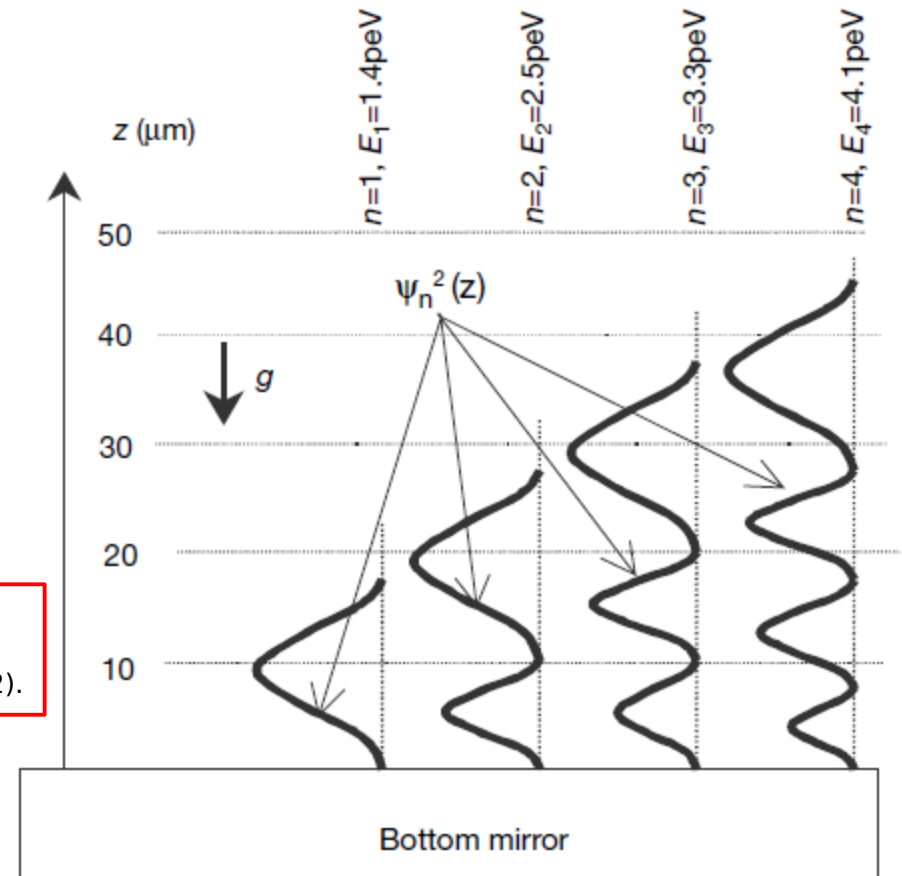
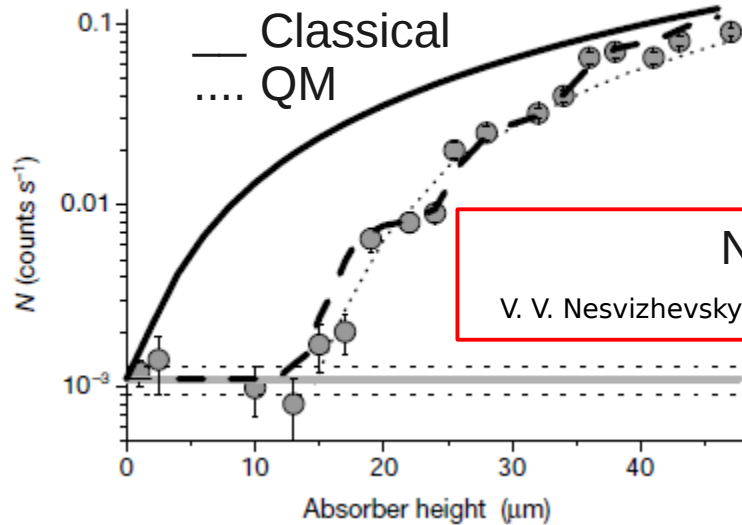
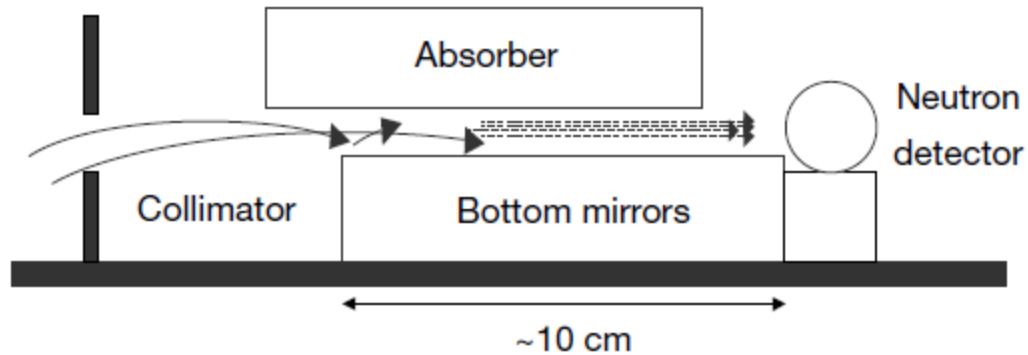
V. V. Nesvizhevsky et al., Nature 415, 297 (2002).



$$E_n \cong \sqrt[3]{\left(\frac{9 \cdot m}{8}\right) \cdot \left(\pi \cdot \hbar \cdot g \cdot \left(n - \frac{1}{4}\right)\right)^2}$$

$$z_0 = \sqrt[3]{\frac{\hbar^2}{2 \cdot m^2 \cdot g}}$$

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Positronium: $z_1 = 1.2 \text{ mm}$, $E_1 = 0.13 \text{ peV}$

Requirements for observation with Ps

To observe QM behavior:

1) Ps vertical velocity so that $E_{\text{kin}_y} \sim E_1 \sim 0.1 \text{ peV}$

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Is it possible to get a some Ps with vertical velocities $v_y < 0.15 \text{ m/s}$ and horizontal velocities $v_{x,z} < 50 \text{ m/s}$ (to keep the size of the apparatus to 0.5 m in 10 ms)?

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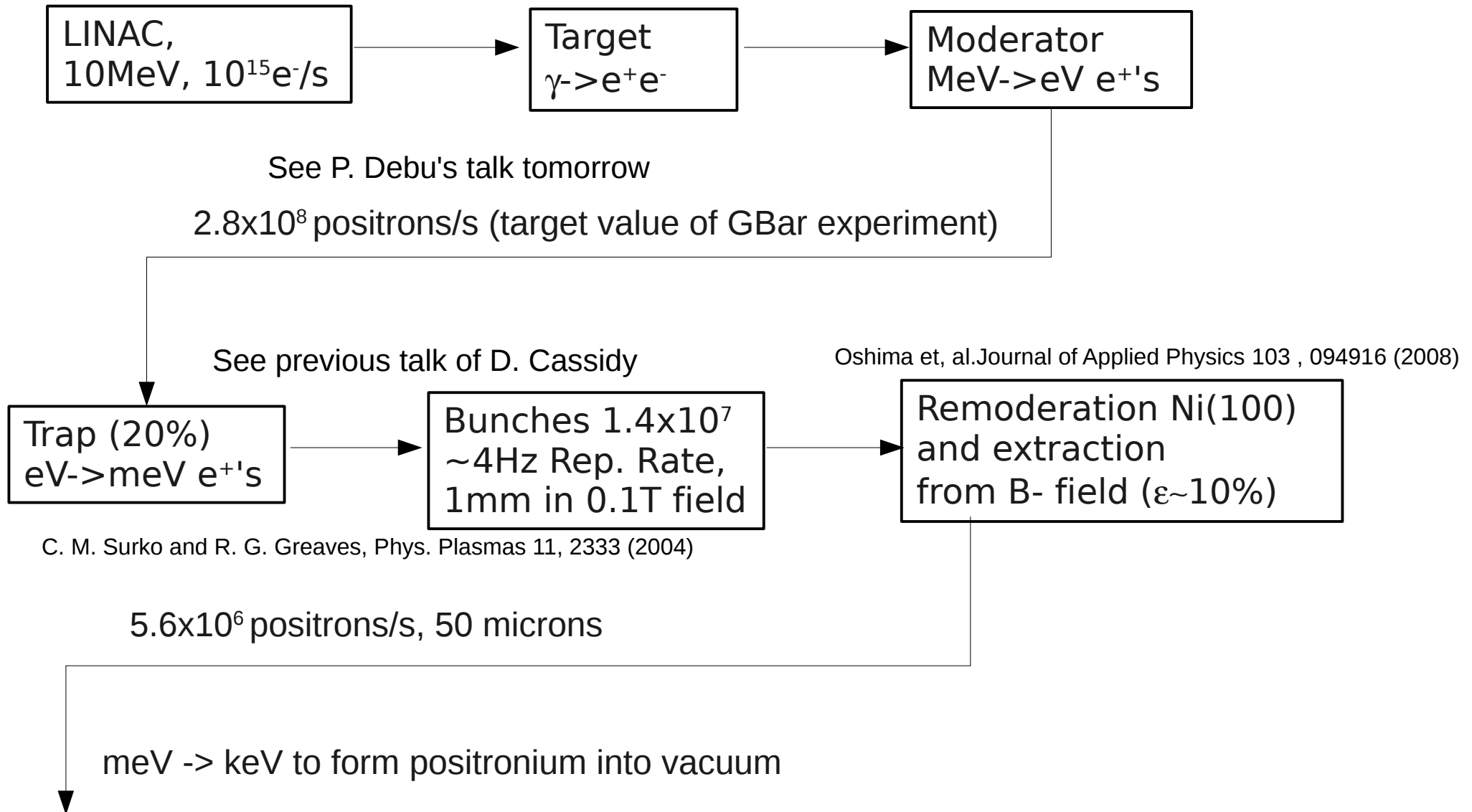
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Ps (triplet) lifetime in ground state 142 ns

-> excitation to a Rydberg state $n=33$ and $l = 32$ mean time to emit the first photon will be about 10 ms.

(Mills & Leventhal, NIMB 192 (2002) 102–106)

Positron source



Positron implantation

By varying the positron implantation energy one can tune the mean implantation depth in a given material.

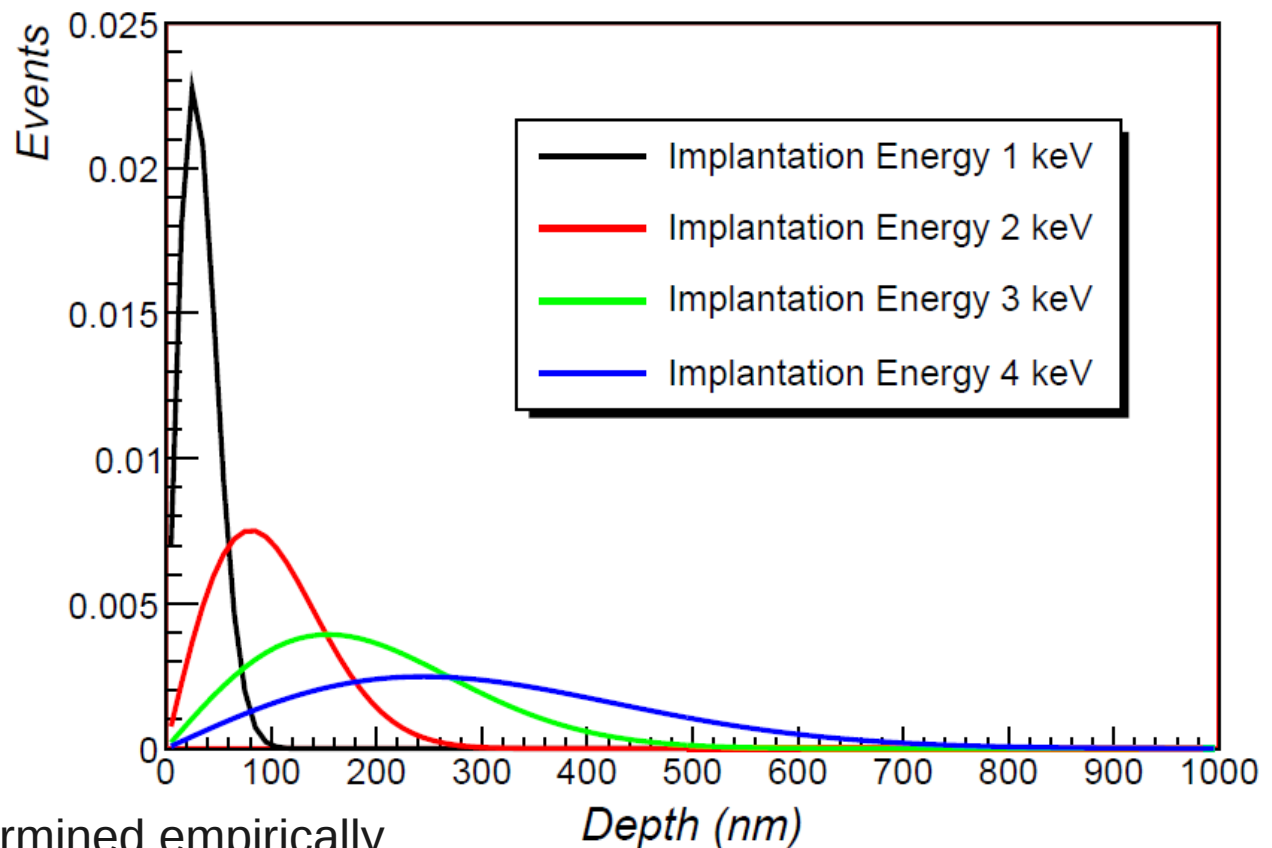
Makhovian
profile

$$P(x, E) = \frac{mx^{m-1}}{x_0^m} e^{-(x/x_0)^m},$$

$$x_0 = \frac{x_m}{\Gamma((1/m) + 1)},$$

$$x_m = \frac{40}{\rho} E^{1.6},$$

Dependent on the material, determined empirically



Positronium production

Very good conversion of about 30 % can be achieved with different methods.

1) Thermal desorption from Al(111) kept at 600K -> 30%

Lower temperature reported with oxygen monolayer at the surface
but very unstable.

A. P. Mills, Jr, Phys. Rev. Lett. 41, 1828 (1978);
K. G. Lynn, Phys. Rev. Lett. 43, 391 (1979);

2) Nano-channeled silicon (5-8 nm) -> 7% at 150 K

Brusa et al., PRL 104, 243401 (2010)

3) Emission from Si(100) surface via exciton-like formation, 0.16 eV

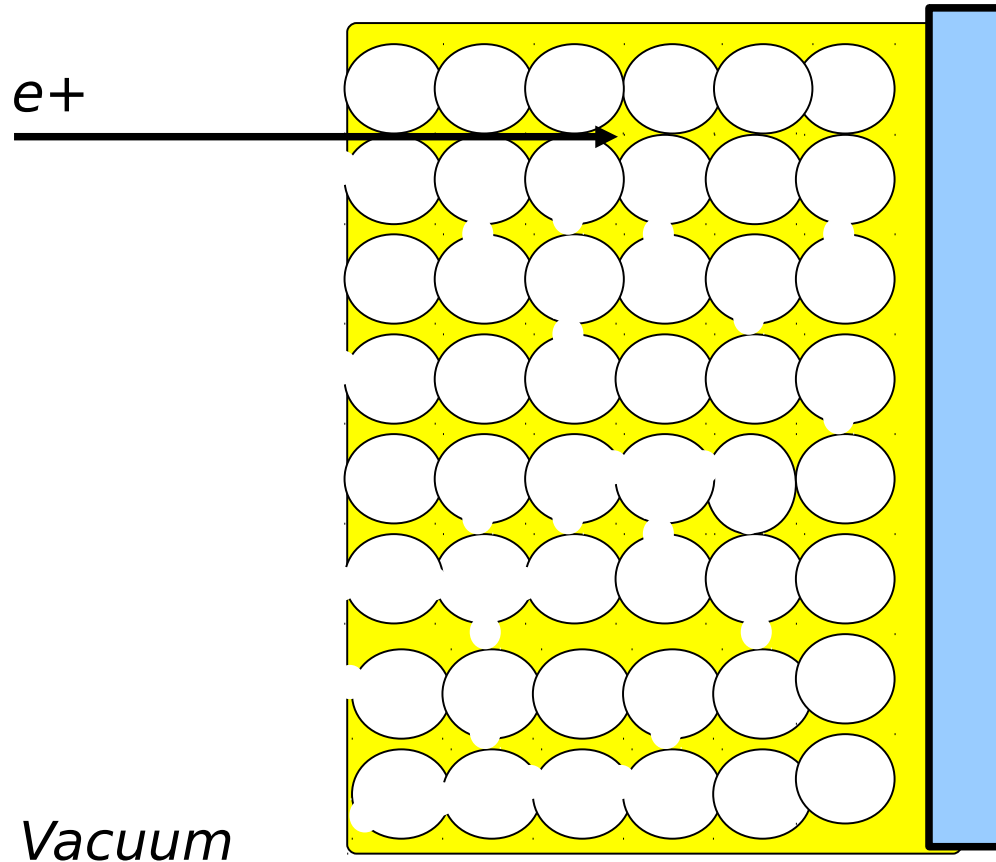
D. Cassidy et al., PRL 106, 133401 (2011)

4) Porous silica films, 30%, 40 meV

P. Crivelli et al., Phys. Rev. A. 81, 052703 (2010)

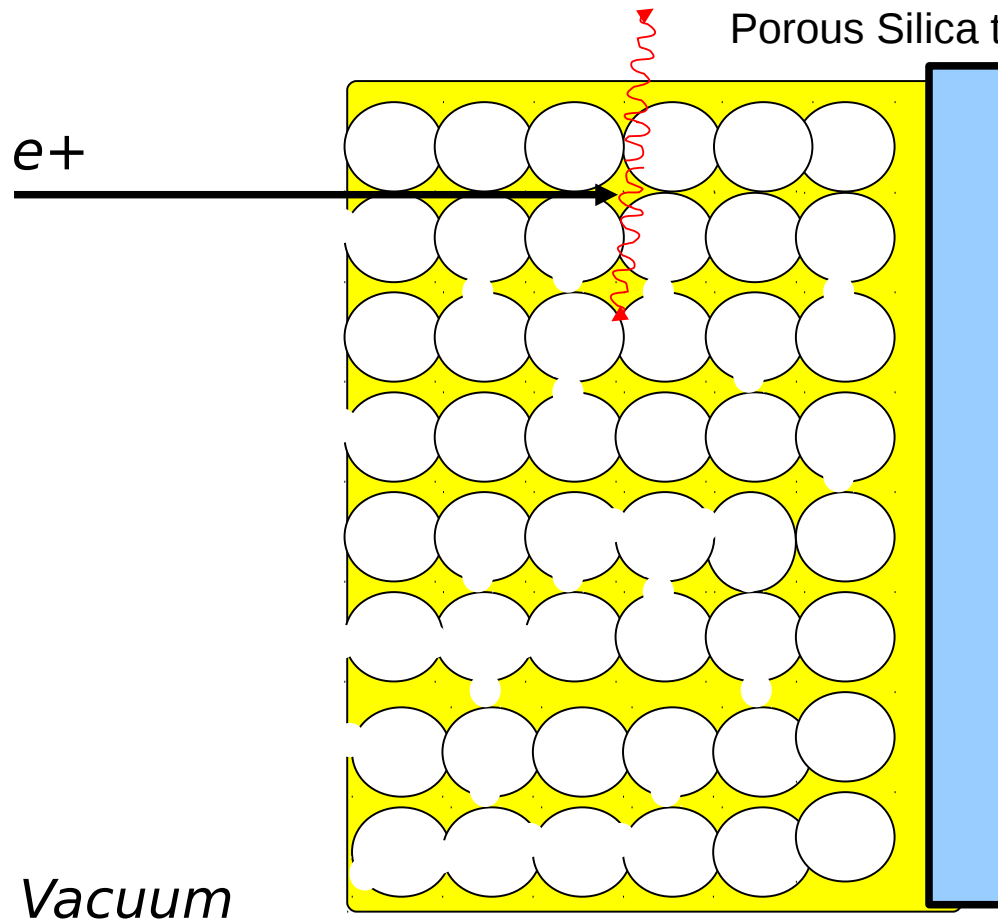
Positronium formation

Porous Silica thin film ~1000nm 3-4 nm pore size



- Positron implanted with keV energies
- Rapidly thermalizes in the bulk (~ps)

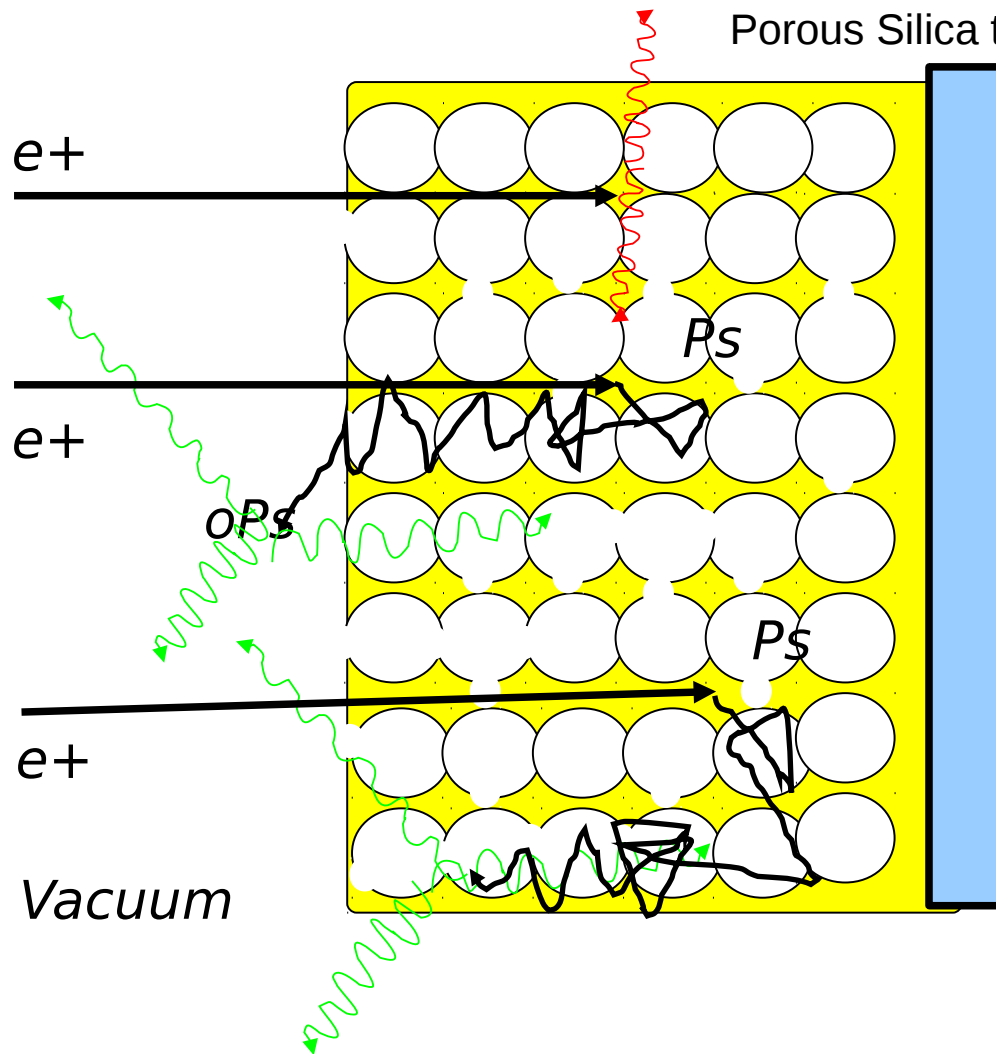
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- Positron diffusion and annihilation

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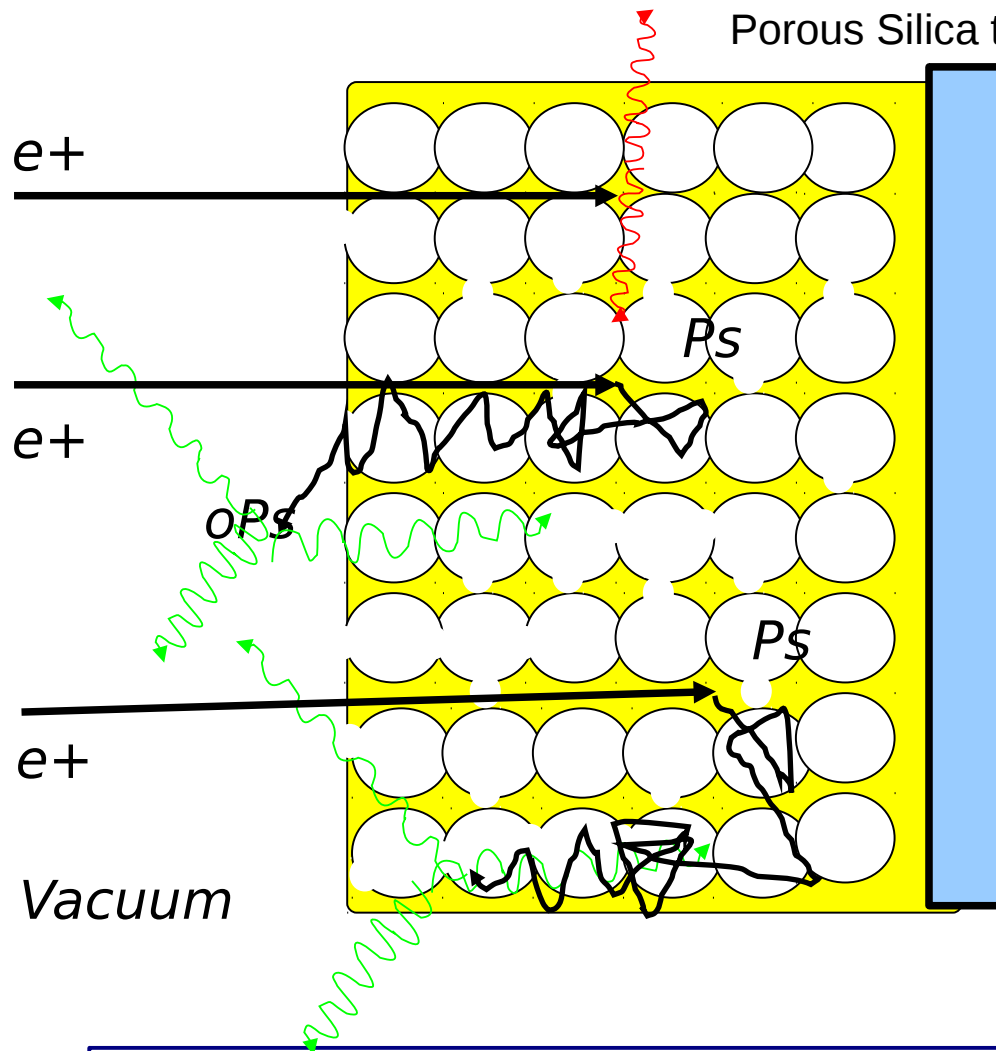


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- Positron diffusion and annihilation

- Positronium formation ($1/4$ pPs, $3/4$ oPs) in SiO_2 by capturing 1 ionized electron
- Diffusion to the pore surface and emission in the pores:
- $W_{Ps} = \mu_{Ps} + E_B - 6.8 \text{ eV} = -1 \text{ eV}$
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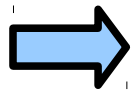
30% of the incident positrons are converted in positronium at 40 meV.

Ps as a particle in a box

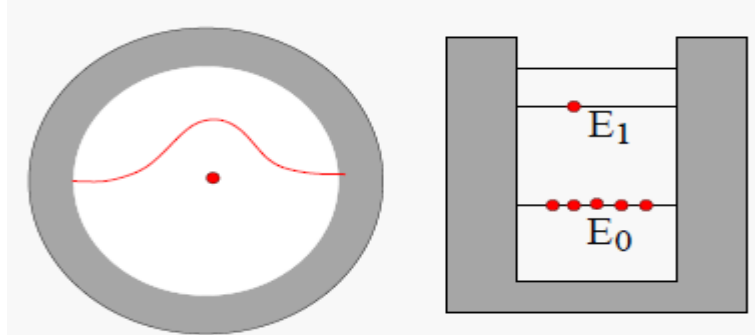
Ps de Broglie wavelength at kinetic energy E_{Ps} ,

$$\lambda_{Ps} = h(2m_{Ps}E_{Ps})^{-1/2} \sim 0.9 \text{ nm}(1 \text{ eV}/E_{Ps})^{1/2},$$

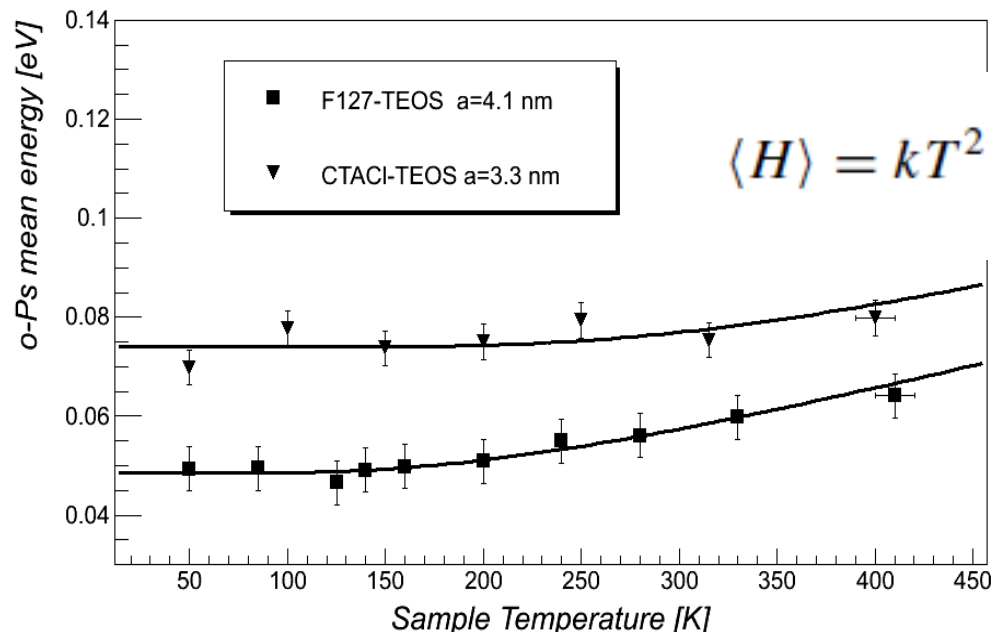
For Ps with 100 meV, λ_{Ps} is comparable with the pore size.



QM effects!



$$E_{Ps} = \frac{h^2}{2m d^2} \approx 0.8 \text{ eV}(1 \text{ nm}/d)^2$$



$$\langle H \rangle = kT^2 \left(\frac{1}{Z(a)} \frac{dZ(a)}{dT} + \frac{1}{Z(b)} \frac{dZ(b)}{dT} + \frac{1}{Z(c)} \frac{dZ(c)}{dT} \right)$$

Z is the partition function defined as

$$Z(a) = \sum_{n=1}^{\infty} e^{-\frac{h^2 n^2}{8ma^2} / kT},$$

Colder Ps from porous materials?

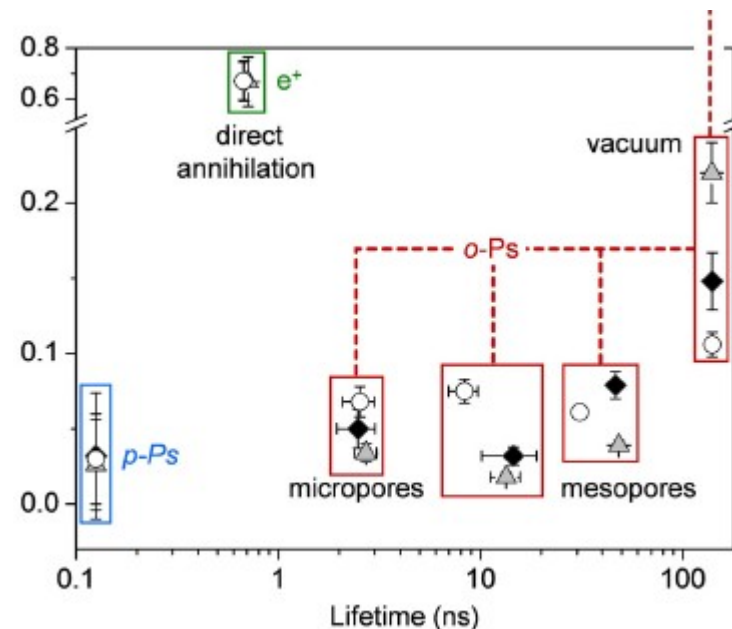
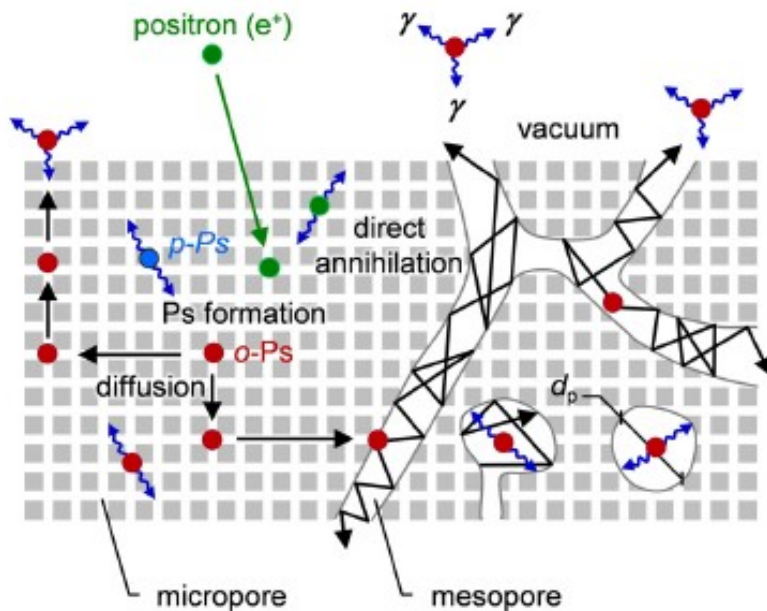
In principle it should be easy: use larger pores of 8-10 nm confinement energy ~ 100 K (for muonium we did reach 100 K with 4 nm since de Broglie wavelength much smaller)

A, Antognini, P. Crivelli et al., PRL 108, 143401 (2012)

In practice: not easy to find the right recipe...work in progress
(in collaboration with CEA Saclay).

Testing different porous materials: hierarchical zeolites and metal organic frameworks

<http://arxiv.org/abs/1309.6497>



Collimation

Ps from silica monoenergetic (Ps velocity 10^5 m/s determined by pore size) with an angular distribution proportional to $\cos(\theta)$ (from random orientation of the pores at the surface):

Collimation: $v_y < 0.15$ m/s and $v_x < 50$ m/s \rightarrow **2×10^{-9}**

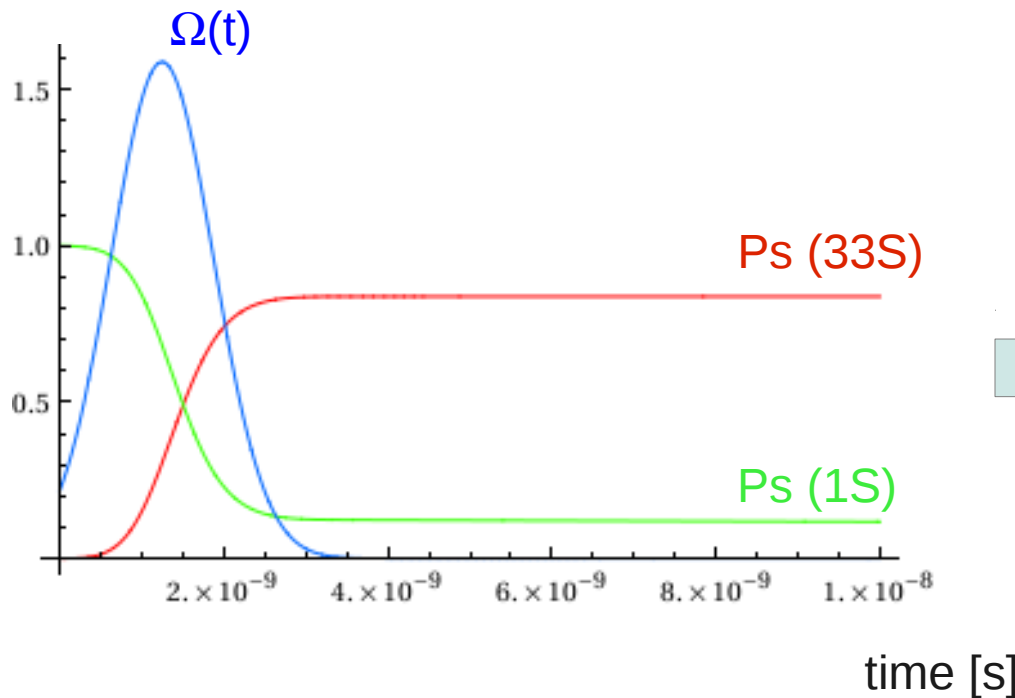
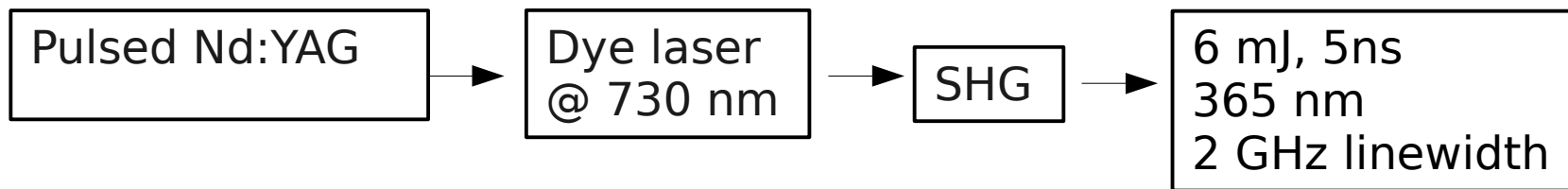
What about v_z (velocity normal to the surface)?

In 10 ms a Ps atom would have traveled (~ 1 km) \rightarrow not very practical for the experiment...

Two photons Rydberg excitation

To resolve the qm gravitational state Ps has to live at least 4.5 ms
-->Ps in Rydberg state

Ziock et al., Phys.Rev. Lett 64, 2366 (1990), Cassidy et al.Phys. Rev. Lett. 108, 043401 (2012)



An efficiency of more than 50% of Ps in $n=33$ atoms seems achievable. Circularly polarized MW radiation to spun up the Ps to $l=32$ (estimated efficiency of 10%).

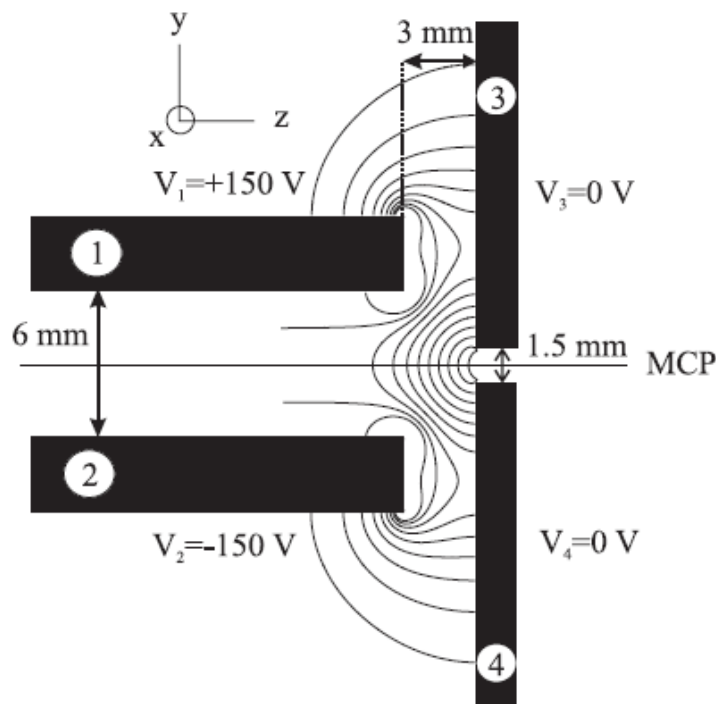
Rydberg Ps deceleration

Large dipole moment $3/2 a_0 nk$ of Rydberg atoms \rightarrow deceleration with inhomogeneous electric fields.

(where a_0 , e , and k are the Bohr radius, the elementary charge, and the difference between the parabolic quantum numbers n_1 and n_2)

Hydrogen stark deceleration and trapping demonstrated in this way by Merkt's group (ETHZ).

see, e.g. S. Hogan, F. Merkt, PRL 100, 043001 (2008)



$$a(\text{m s}^{-2}) = 76 \nabla F (\text{V cm}^{-2}) \frac{1}{m(u)} nk$$

For H ($n=25$) $a = 2 \times 10^8 \text{ m/s}^2$

From 700 m/s \rightarrow at rest for a 3 mm gap with an efficiency of 80%

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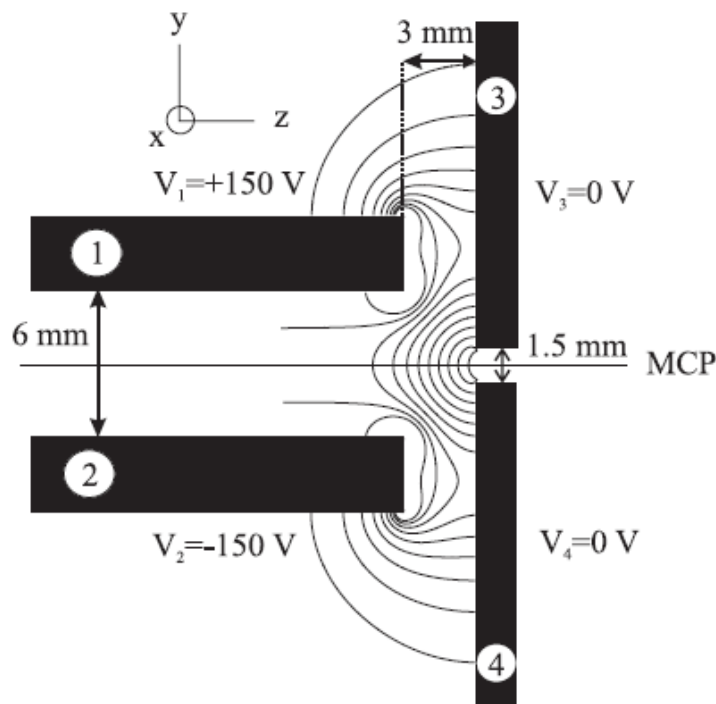
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From 700 m/s \rightarrow at rest for a 3 mm gap with an efficiency of 80%

For Ps ($n=25$) $a = 4 \times 10^{11} \text{ m/s}^2$

From $10^5 \text{ m/s} \rightarrow v_z < 50 \text{ m/s}$

Seems achievable with an efficiency of 20%
 \rightarrow simulation is in progress...)



Proposed scheme

Pulsed e^+ Beam
($50\text{ }\mu\text{m}$, $5.6 \times 10^6\text{ e}^+/\text{s}$)

Thermally (4K) and e-m shielded region

SiO₂
target,
30% Ps

Ps*

collimator

Absorber plate 25 cm

~1. mm
variable

Mirror plate

BGO detector

Rectangular
MCPs

BGO detector

MWs

Stark
decelerator

Mirror:

- Use wires with constant current.
- Gradient of magnetic fields: $\mathbf{F} = \mu \mathbf{grad} B$
- To equate Ps $E_{\text{kin}} = 0.1\text{ peV} \rightarrow \text{grad } B \sim 17\text{ mG}$

Estimated signal rate

Expected signal rate is given by:

$$R_{\text{signal}} = R_{e^+} \epsilon_{\text{Ps}} \epsilon_{n33} \epsilon_{l32} \epsilon_{\text{tails}} \epsilon_{\text{tau}} \epsilon_{\text{det}} \epsilon_{\text{dec}} = 4 \times 10^{-5} \text{ events/s}$$

where

$$R_{e^+} = 5. \times 10^6 \text{ positrons/s,}$$

$$\epsilon_{\text{Ps}} = 0.3 \text{ the fraction of Ps emitted into vacuum,}$$

$$\epsilon_{n33} = 0.5 \text{ the excitation probability in the } n=33 \text{ state,}$$

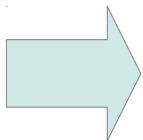
$$\epsilon_{l32} = 0.1 \text{ the efficiency to spin the Ps in the } l=32 \text{ state,}$$

$$\epsilon_{\text{tails}} = 2 \times 10^{-9} \text{ fraction of Ps atoms after collimator (} v_x < 50 \text{ and } v_y < 0.15),$$

$$\epsilon_{\text{tau}} = 0.6 \text{ fraction of atoms surviving the 5 ms time-of-flight,}$$

$$\epsilon_{\text{det}} = 0.5 \text{ efficiency of signal detection.}$$

$$\epsilon_{\text{dec}} = 0.2 \text{ efficiency of stark deceleration}$$



This means that 3 events/day with a background of 0.06 bkg/day (requirement of MCP signal + TOF + back-2-back 511 keV photons) should be detected.

Concluding remarks

- A measurement of the Ps gravitational QM state looks very challenging indeed but maybe not impossible.
- This experiment would profit greatly if a source of cold Ps could be produced.
- Further simulations/calculations are needed to understand all the efficiencies and sources of systematic effects.
- Preliminary experiments to understand if cold enough Ps atoms could be produced and the best way to Ps mirror seems to be justified.
- Any advice/remark or comment is more than welcome!

Thank you for your attention!

Specials thanks to V. Nesvizhevsky and A. Voronin for the very useful discussions.