Starobinsky-like models of inflation

Impact of BICEP2

Effect of HDOs in inflation

Starobinsky-like models

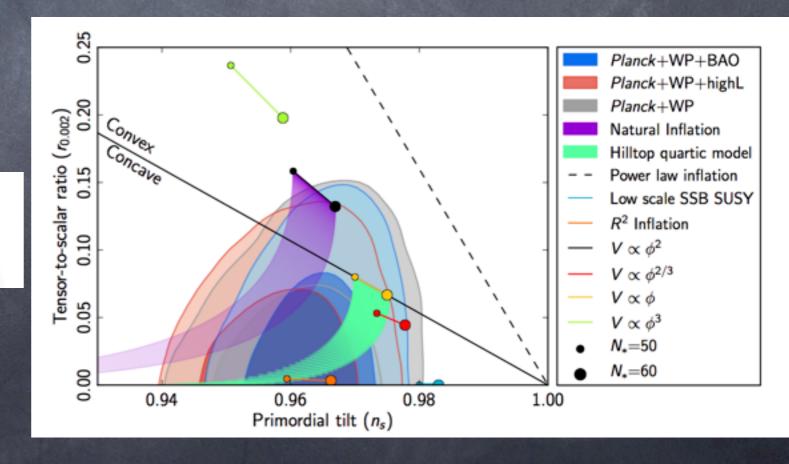
The model

$$S_{\rm S} = \frac{1}{2} \int d^4x \sqrt{-g} \left(M_{\rm p}^2 R + \frac{1}{6M^2} R^2 \right)$$

Kehagias, Dizgah & Riotto 1312.1155 Giudice & Lee 1402.2129

leads to

$$n_s - 1 \approx -\frac{2}{N}, \quad r \approx \frac{12}{N^2}$$



Starobinsky-like models

Not a single model, but a class

Example: Higgs inflation

$$S_{\rm HI} = \int \mathrm{d}^4 x \sqrt{-g} \left(\frac{M_{\rm p}^2}{2} R + \frac{1}{2} \xi h^2 R - \frac{1}{2} \partial_\mu k \partial^\mu h - \frac{\lambda}{4} h^4 \right)$$

integrate out the scalar

$$\xi hR - \lambda h^3 = 0,$$

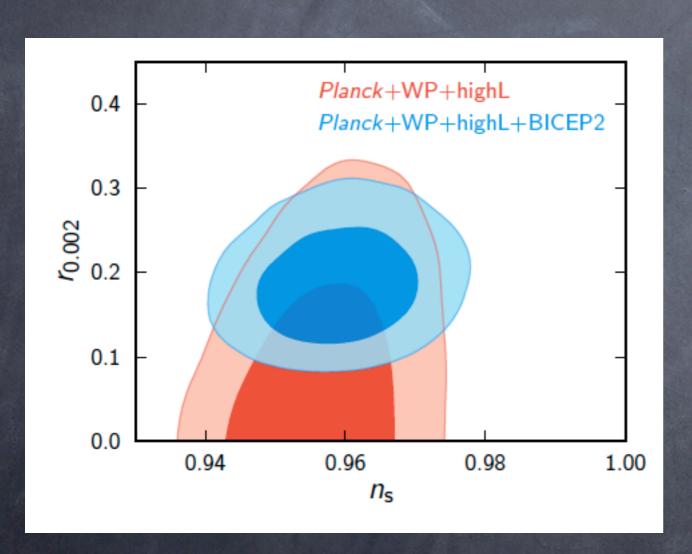
$$h^2 = \frac{\xi R}{\lambda}$$
.

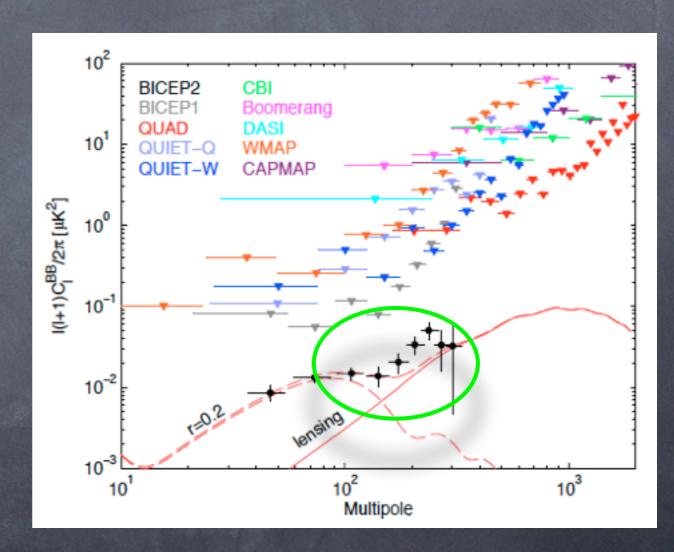
$$S_{\rm HI} = \int \mathrm{d}^4 x \sqrt{-g} \left(\frac{M_{\rm p}^2}{2} R + \frac{\xi^2}{4\lambda} R^2 \right)$$

this is the case for many models:
S-inflation, gravity induced...
Same inflation, but unitarity and other properties can be different

$$A_T = 8 \left(\frac{H_I}{2\pi M_{pl}} \right)^2$$

The impact of BICEP2





The impact of BICEP2

chaotic inflation is a winner

BUT

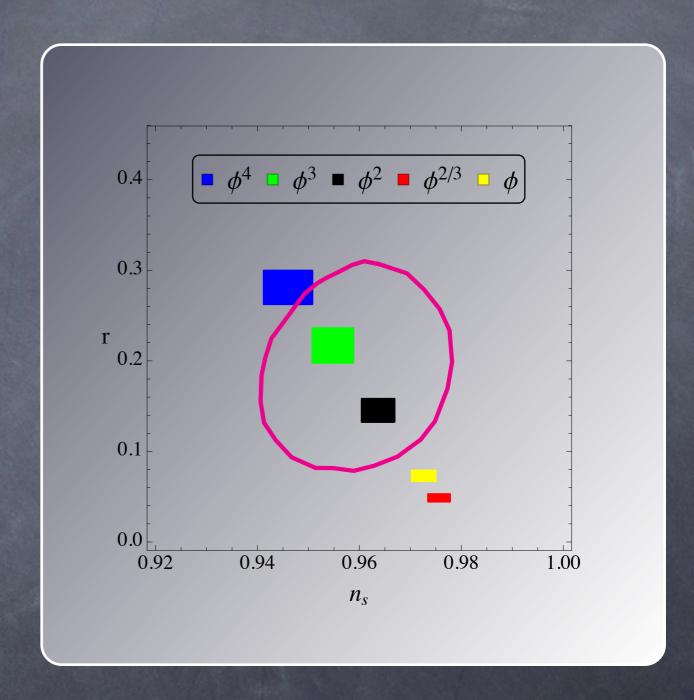
needs trans-Planckian excursions of the field

$$V(\tilde{\phi})_{CI} = \bar{M}_P^4 \tilde{m}^2 \tilde{\phi}^2$$

$$\tilde{\phi} = \phi/M_{Pl}$$

$$\tilde{m} = m/M_{Pl}$$

$$\tilde{\phi}_N = N/2\pi$$



$$N \in [50, 60]$$

Calmet, VS 1403.5100

These trans-Planckian excursions must come at a price

Quantify these effects?

BHs, graviton exchange, more exotic stuff...

Higher-order operators

$$V(\phi) = V_{ren}(\phi) + \sum_{n=5}^{\infty} c_n \frac{\phi^n}{\bar{M}_P^{n-4}}$$

Example: Impact on chaotic inflation

$$V(\tilde{\phi}) = \bar{M}_P^4 \left(\tilde{m}^2 \tilde{\phi}^2 + c_n \tilde{\phi}^n \right)$$

Example: Impact on chaotic inflation

$$V(\tilde{\phi}) = \bar{M}_P^4 \left(\tilde{m}^2 \tilde{\phi}^2 + c_n \tilde{\phi}^n \right)$$

$$c_6 = \alpha_m \tilde{m}^2 \to V(\tilde{\phi}) = \bar{M}_P^4 \tilde{m}^2 \tilde{\phi}^2 \left(1 + \alpha_m \tilde{\phi}^4 \right)$$

$$\begin{split} c_6 &= \alpha_m \tilde{m}^2 \to V(\tilde{\phi}) = \bar{M}_P^4 \tilde{m}^2 \tilde{\phi}^2 \left(1 + \alpha_m \tilde{\phi}^4\right) \quad \begin{bmatrix} \text{naively} \\ |\alpha_m| \tilde{\phi}^4 < 1 \end{bmatrix} \\ \epsilon &= \frac{1}{16\pi} \left(\frac{V'(\tilde{\phi})}{V(\tilde{\phi})}\right)^2 = \frac{1}{4\pi} \frac{1}{\tilde{\phi}^2} \left(\frac{1 + 3\alpha_m \tilde{\phi}^4}{1 + \alpha_m \tilde{\phi}^4}\right)^2 = \epsilon_{CI} + \frac{\alpha_m \tilde{\phi}^4}{\pi \tilde{\phi}^2} + \mathcal{O}(\alpha_m \tilde{\phi}^4)^3 \end{split}$$

the largest value of phi is at the beginning of inflation

Example: Impact on chaotic inflation

$$V(\tilde{\phi}) = \bar{M}_P^4 \left(\tilde{m}^2 \tilde{\phi}^2 + c_n \tilde{\phi}^n \right)$$

$$c_6=lpha_m ilde{m}^2 o V(ilde{\phi})=ar{M}_P^4 ilde{m}^2 ilde{\phi}^2\left(1+lpha_m ilde{\phi}^4
ight)$$
 naively $|lpha_m| ilde{\phi}^4<1$

$$\epsilon = \frac{1}{16\pi} \left(\frac{V'(\tilde{\phi})}{V(\tilde{\phi})} \right)^2 = \frac{1}{4\pi} \frac{1}{\tilde{\phi}^2} \left(\frac{1 + 3\alpha_m \tilde{\phi}^4}{1 + \alpha_m \tilde{\phi}^4} \right)^2 = \epsilon_{CI} + \frac{\alpha_m \tilde{\phi}^4}{\pi \tilde{\phi}^2} + \mathcal{O}(\alpha_m \tilde{\phi}^4)^3$$

the largest value of phi is at the beginning of inflation

$$\tilde{\phi}_N^2 \simeq \tilde{\phi}_{N,CI}^2 + \frac{N^3}{12\pi^3} \alpha_m \simeq \frac{N}{2\pi} \left(1 + \frac{N^2 \alpha_m}{6\pi^2} \right)$$

$$|\alpha_m| \tilde{\phi}_N^4 \simeq \frac{N^2 |\alpha_m|}{4\pi^2} \lesssim 1 \rightarrow |\alpha_m|^{EFT} \lesssim 2 \times 10^{-2}$$

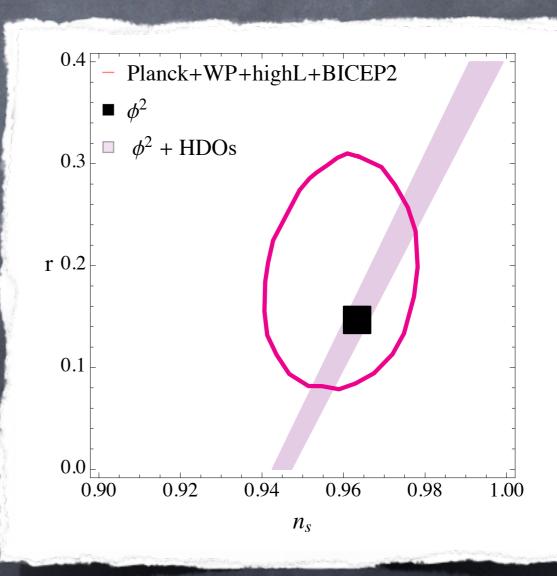
$$n_s - 1 = (n_s - 1)_{CI} \left(1 - \frac{5}{3} \frac{N^2 \alpha_m}{4\pi^2} \right)$$

$$r = r_{CI} \left(1 + \frac{10}{3} \frac{N^2 \alpha_m}{4\pi^2} \right)$$

$$\alpha_m^{BICEP2} \in [-2, 3] \times 10^{-3}$$

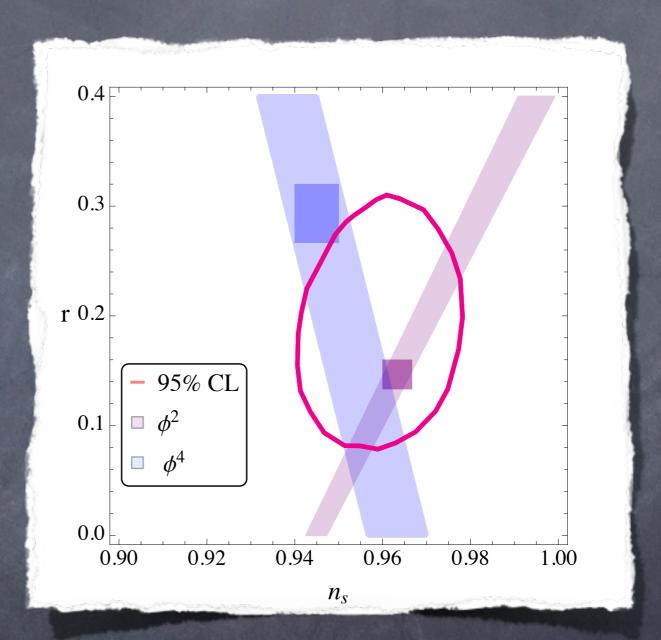
$$c_6 \lesssim 10^{-9}$$

If MPI->MGUT, even stronger constraints



models can be 'rescued'

with HDOs well within the EFT validity region



$$\alpha_{\lambda}^{BICEP} \in [-0.06, 0] \to c_6 < 10^{-15}$$